

Characterizing the underlying event in hadron-hadron collisions

Sebastian Sapeta

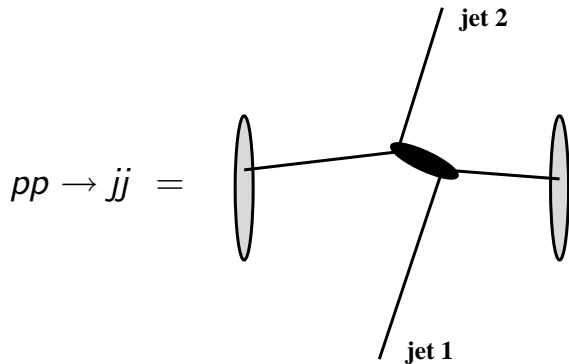
LPTHE, UPMC, CNRS, Paris

in collaboration with Matteo Cacciari and Gavin Salam[§]

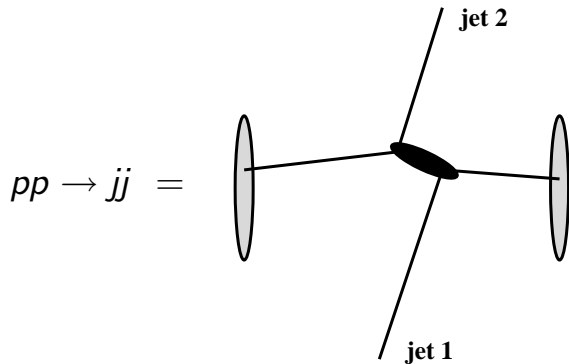
*Workshop on Hadron-Hadron & Cosmic-Ray Interactions at multi-TeV Energies
ECT* - Trento, Nov 29th - Dec 3rd, 2010*

[§]M.Cacciari, G.P.Salam and SS, JHEP 1004 (2010) 065

What is the underlying event?

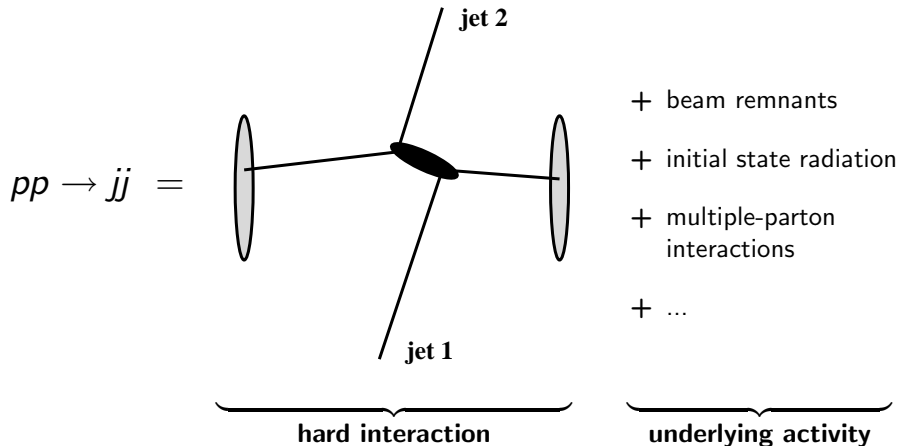


What is the underlying event?



- + beam remnants
- + initial state radiation
- + multiple-parton interactions
- + ...

What is the underlying event?



- ▶ these are ingredients of present Monte Carlo models

Problems and questions

Definition of underlying event (UE) is ambiguous ...

- ▶ there is only one event with no clear bound between hard part and UE

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This leads us to the following two questions:

1. what do we really measure with existing methods of UE determination?
→ test the methods with toy model
2. which observables are interesting to measure?
→ study UE from Monte Carlo models

What can we measure about UE?

Relevant characteristics of energy flow of UE

- ▶ ρ – level of transverse momentum per unit area
- ▶ rapidity dependence of ρ
- ▶ point-to-point fluctuations within a single event ($\equiv \sigma$)
- ▶ fluctuations from event to event
- ▶ point-to-point correlations

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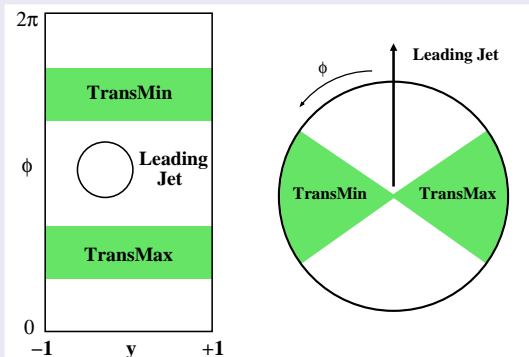
Two existing methods for measuring UE

- ▶ traditional approach
- ▶ jet area/median based approach

Traditional approach [Marchesini & Webber (1988), UA1 (1988), Field et al.]

For each event

1. take charged particles with $p_t > 0.5$ GeV and $|y| < y_{\max}$
2. cluster with cone jet algorithm with $R = 0.5 - 0.7$ to find the leading jet
3. define typical p_t of UE as $\langle p_t \rangle$ in TransMin, TransMax or TransAv regions



$$\text{TransAv: } \mathcal{O}(\alpha_s)$$

$$\text{TransMax: } \mathcal{O}(\alpha_s)$$

$$\text{TransMin: } \mathcal{O}(\alpha_s^2)$$

- **topological** separation: UE defined as particles entering certain region of (y, ϕ) space

Parenthesis: Sequential recombination jet algorithms

Distance measure between pair of particles d_{ij} and between particle and beam d_{iB}

$$d_{ij} = \min(k_{ti}^{2p}, k_{tj}^{2p}) \frac{\Delta_{ij}^2}{R^2}, \quad d_{iB} = k_{ti}^{2p},$$

where $\Delta_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$ is a geometrical distance in the (y, ϕ) plane

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- ▶ identify the smallest of d_{ij} and d_{iB}
- ▶ recombine particles i and j (if d_{iB} is the smallest call i a jet and remove it)
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$p = 1$: k_t algorithm

[Catani, Dokshitzer, Seymour, Webber (1993); Ellis, Soper (1993)]

$p = 0$: Cambridge/Aachen algorithm

[Dokshitzer, Leder, Moretti, Webber (1997); Wobisch, Wengler (1999)]

$p = -1$: anti- k_t algorithm

[Cacciari, Salam, Soyez (2008)]

Parenthesis: Jets have areas [Cacciari, Salam, Soyez (2008)]

To determine the **active area** of a jet

- ▶ supplement a set of physical particles $\{p_i\}$ with an ensemble of dense, infinitely soft, randomly distributed *ghost particles* $\{g_i\}$
- ▶ cluster the set $\{p_i, g_i\}$
- ▶ compute the active area of a jet J for this specific ensemble of ghosts $\{g_i\}$

$$A(J | \{g_i\}) = \frac{\mathcal{N}(J)}{\nu_g},$$

where $\mathcal{N}(J)$ is the number of ghosts contained in the jet J and ν_g is the number of ghosts per unit area

- ▶ average over many ghost ensembles

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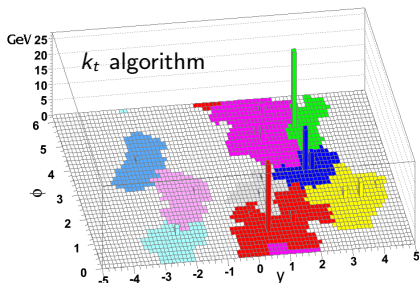
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<http://fastjet.fr>



Area/median approach [Cacciari, Salam, Soyez (2008), <http://fastjet.fr>]

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1. cluster particles with an infrared safe jet finding algorithm (all particles are clustered so we have set of jets ranging from hard to soft)

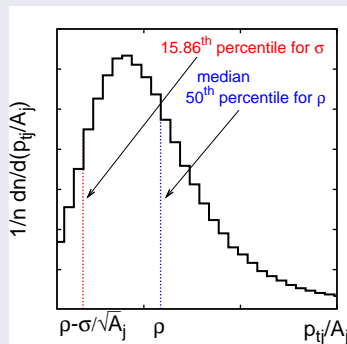
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2. from the list of all jets (no cuts required!) determine

$$\rho = \text{median} \left[\left\{ \frac{p_{t,j}}{A_j} \right\} \right]$$

and its uncertainty σ

- ▶ median gives a typical value of p_t/A for a given event
- ▶ using median is a way to **dynamically** separate hard and soft parts of the event



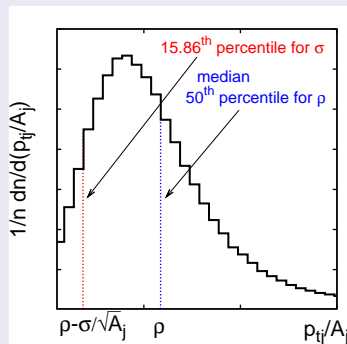
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- ▶ ρ may be used e.g. to correct hard jet transverse momentum

$$p_{t,j}^{(\text{sub})} = p_{t,j} - \rho A_j \pm \sigma \sqrt{A_j}$$

since jet area measures the jet susceptibility to the soft radiation

Understanding the methods – a toy model study

Two component model: soft UE + hard contamination

soft component (UE)

- ▶ take the region of area A in (y, ϕ) space \rightarrow transverse region (traditional approach) or jet area (area/median approach)
- ▶ number of particles in this region, n , given by Poisson distribution with the average $\langle n \rangle$
- ▶ single-particle p_t distribution given by

$$\frac{dpt_1}{dp_t} = \frac{1}{\mu} e^{-p_t/\mu}$$

- ▶ parameters:
 - μ – average p_t of particle,
 - $\nu = \frac{\langle n \rangle}{A}$ – density of particles
- ▶ in this model $\rho = \mu\nu$ is the true value of p_t/A of UE

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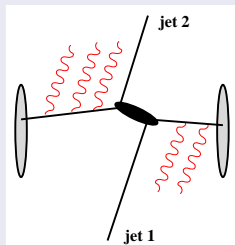
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hard component (ISR)

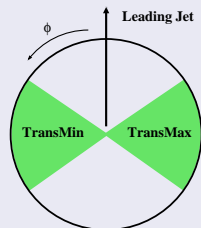


- ▶ soft and collinear partons from primary emissions:
$$\frac{dn}{dp_t dy d\phi} \simeq \frac{C_i}{\pi^2} \frac{\alpha_s(p_t)}{p_t}$$
- ▶ hard scale cut $Q = \frac{1}{2} p_t = 50$ GeV
- ▶ partons distributed uniformly in angle and rapidity

Two component model: biases

Traditional approach

- ▶ TransAv and TransMin variants



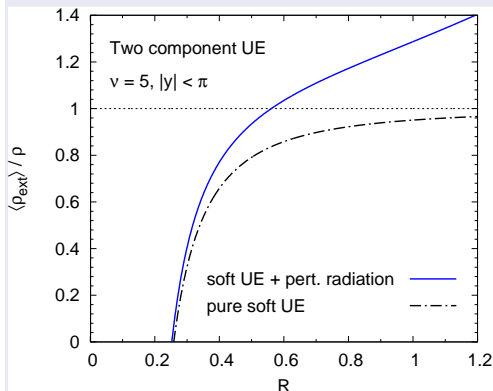
$$\langle \rho_{\text{ext,Av}} \rangle = \rho + \frac{C_i \alpha_s}{\pi^2} Q$$

$$\langle \rho_{\text{ext,Min}} \rangle \simeq \rho - \frac{\sigma \mathcal{P}}{\sqrt{\pi A_{\text{Trans}}}} + 2 \left(\frac{C_i \alpha_s}{\pi^2} \right)^2 A_{\text{Trans}} Q$$

\mathcal{P} – fraction of events with perturbative radiation smaller than soft fluctuations

Two component model: biases

Area/median approach



$$\langle \rho_{\text{ext}} \rangle \simeq \langle \rho_{\text{ext}}^{(\text{soft})} \rangle + \sqrt{\frac{\pi C_J}{2}} \sigma R \frac{\langle n_h \rangle}{A_{\text{tot}}}$$

$\langle n_h \rangle$ – number of perturbative part.

σ – measure of fluctuations

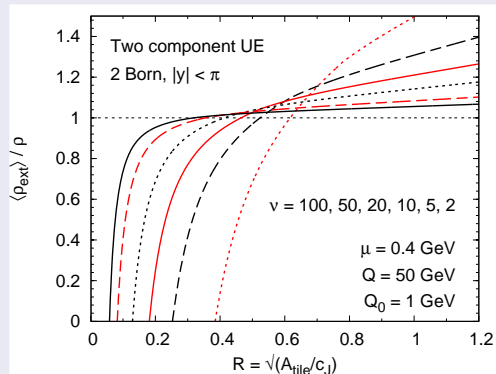
ρ – true value of p_t/A

$$\frac{\langle n_h \rangle}{A_{\text{tot}}} \simeq \frac{n_b}{A_{\text{tot}}} + \frac{C_i}{\pi^2} \frac{1}{2b_0} \ln \frac{\alpha_s(Q_0)}{\alpha_s(Q)}$$

- ▶ the two terms bias $\langle \rho_{\text{ext}} \rangle$ in opposite directions
- ▶ for $R \simeq 0.5 - 0.6$ (used in most MC analysis of UE) the biases largely cancel
- ▶ similar picture and conclusions for σ

Two component model: biases

Area/median approach



Turn-on point:

$$R_{\text{crit}} \simeq 0.41 \cdot \frac{\sigma}{\rho} = 0.41 \cdot \sqrt{\frac{2}{\nu}}$$

Point of zero bias:

$$R_{\text{zero-bias}} \simeq 0.87 R_{\text{crit}}^{\frac{1}{3}} \left(\frac{C_A}{C_i} \right)^{\frac{1}{3}}$$

- ▶ for $R \simeq 0.5 - 0.6$ the uncertainty always stays below 20% for a broad range of particle densities ν

Fluctuations in estimation of ρ

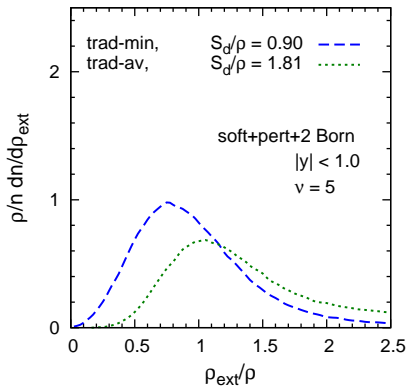
In the toy model: the same ρ distribution used to generate all events

- ▶ however, there are event-to-event fluctuations of ρ due to restricted area
- ▶ this sets the lower limit for the uncertainty of ρ determination

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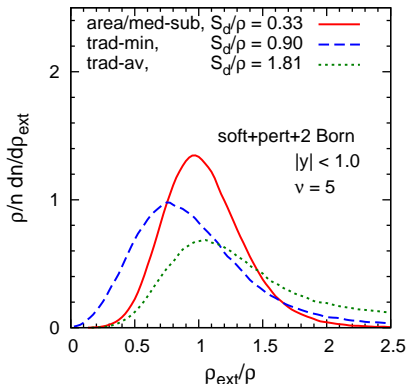
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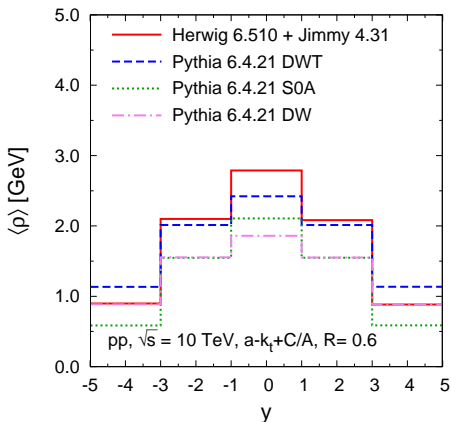


- ▶ traditional approach suffers significantly more from the hard contamination
 $S_d \sim Q$

Approaching real life – Monte Carlo study

Average ρ as a function of y

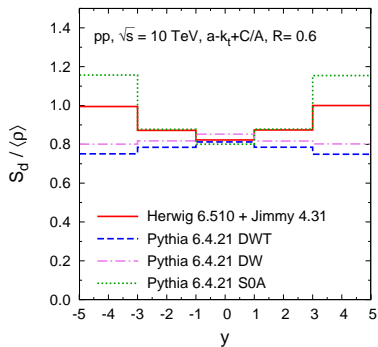
- ▶ dijets at the LHC, $\sqrt{s} = 10$ TeV, $p_t > 100$ GeV, $|y| < 4$



- ▶ significant y dependence
- ▶ strips of $\Delta y = 2$ sufficient for robust ρ determination

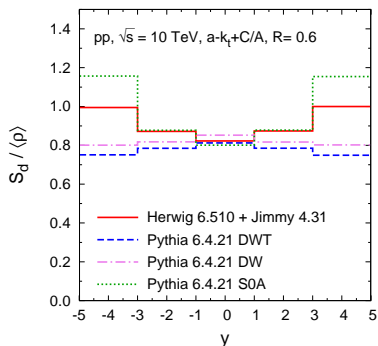
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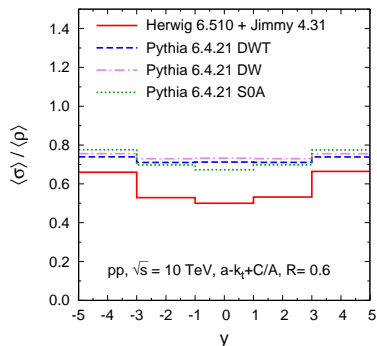


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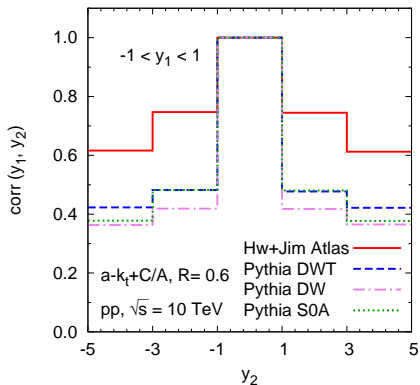


- ▶ within an event



- ▶ large inter-event and intra-event
- ▶ two patterns of rapidity dependence
- ▶ sizable difference between Herwig+Jimmy and Pythia

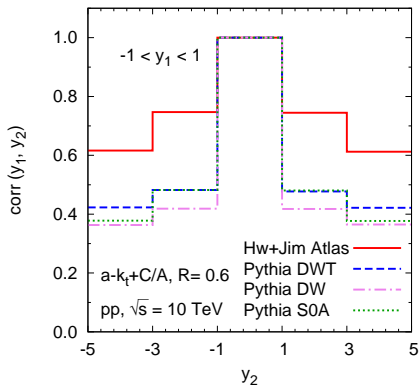
Correlations



$$\text{corr}(y_1, y_2) = \frac{\langle \rho(y_1)\rho(y_2) \rangle - \langle \rho(y_1) \rangle \langle \rho(y_2) \rangle}{S_d(y_1)S_d(y_2)}$$

- ▶ y_1, y_2 – rapidity bins of width $\Delta y = 2$
- ▶ $\langle \dots \rangle$ – average over many events

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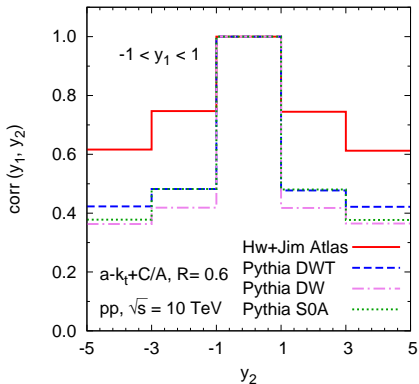


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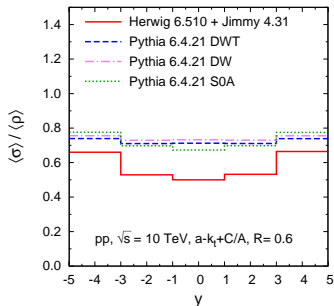
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- ▶ significant difference between Herwig + Jimmy and Pythia
- ▶ qualitatively consistent with $\langle \sigma \rangle / \langle \rho \rangle$: smaller fluctuations within event \Leftrightarrow larger correlations



CMS analysis at $\sqrt{s} = 0.9$ TeV

Experimental conditions

- ▶ only charged tracks
- ▶ $p_t > 0.3$ GeV
- ▶ $|\eta| < 2.3$

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Modification of $\rho \rightarrow \rho'$

$$\rho' = \text{median}_{j \in \text{physical jets}} \left[\left\{ \frac{p_{t,j}}{A_j} \right\} \right] \cdot \mathcal{C}$$

$$\mathcal{C} = \frac{\sum_{j \in \text{physical jets}} A_j}{A_{\text{tot}}}$$

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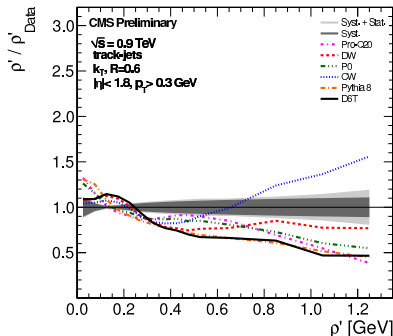
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- ▶ can discriminate between UE tunes even at extreme conditions
- ▶ none of the tunes describes data
- ▶ looking forward to the 7 TeV results (also for other observables)

Summary

Measurement of UE is difficult both in principle and in practice

- ▶ we have considered a simple toy model to better understand the methods
- ▶ both traditional and area/based approach perform comparably well in measuring average quantities
- ▶ for event-to-event measurements traditional approach suffers significantly from hard radiation

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The study of UE from MC with the area/median method suggests the set of observables deserving dedicated measurements

- ▶ dependence of ρ on rapidity
- ▶ fluctuations from event to event (large for all generators/tunes)
- ▶ fluctuations within an event, σ , (significant differences between Herwig+Jimmy and Pythia)
- ▶ correlations (large differences between Herwig+Jimmy and Pythia)