

BEYOND THE MINIMAL
COMPOSITE HIGGS MODEL

2 HIGGSSES AS COMPOSITE PSEUDO-NGB'S

Javi Serra

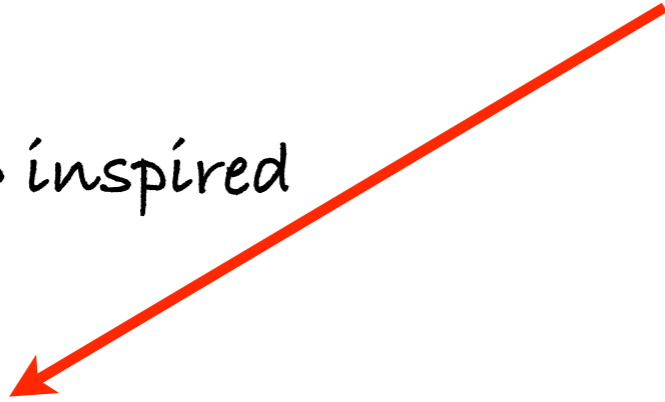
(Universitat Autònoma de Barcelona)

work (in progress) in collaboration with: J.Mrazek, A.Pomarol, R.Rattazzi, M.Redi, A.Wulzer

WHY A COMPOSITE-NGB-HIGGS ?

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QCD inspired



**NO Naturalness or Hierarchy
problems**



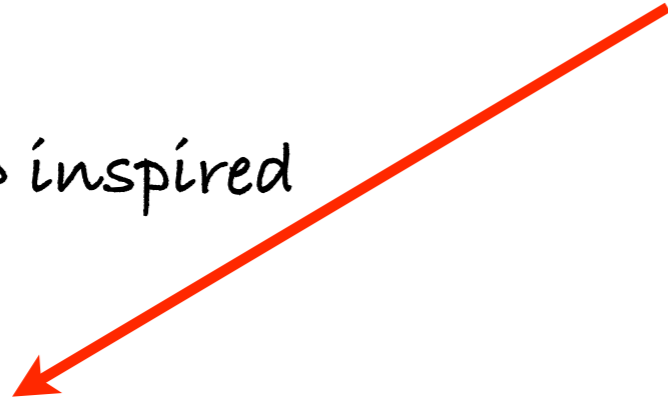
dynamical

EW SCALE

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consistency with EWPT (S, T)

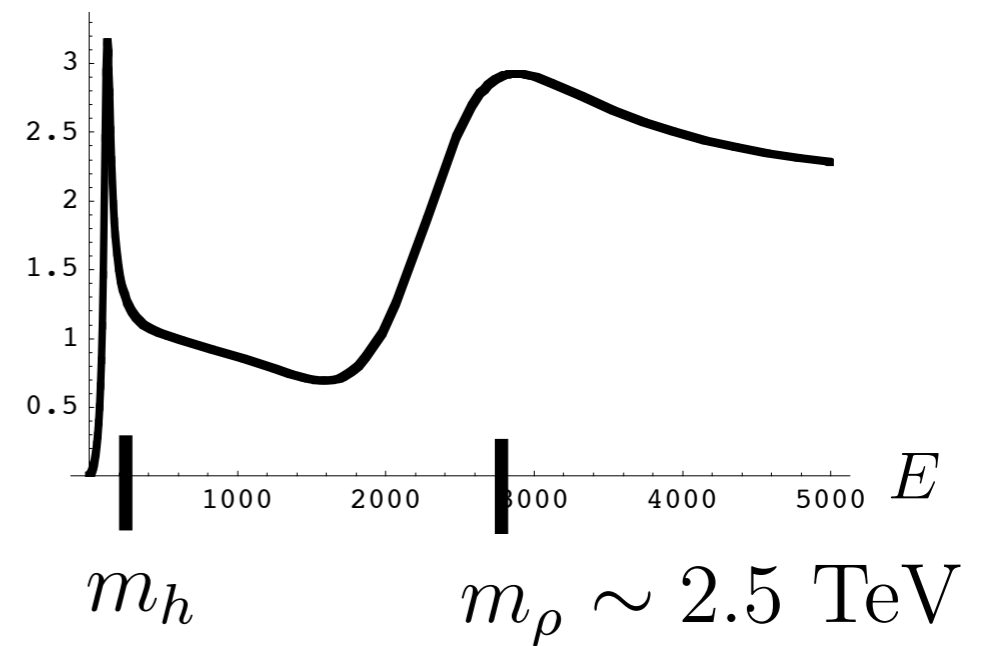
STRONG SECTOR



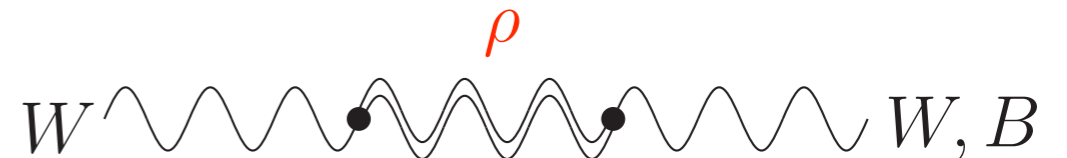
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$\mathcal{M}(W_L W_L \rightarrow W_L W_L)$



resonances pushed up in energy



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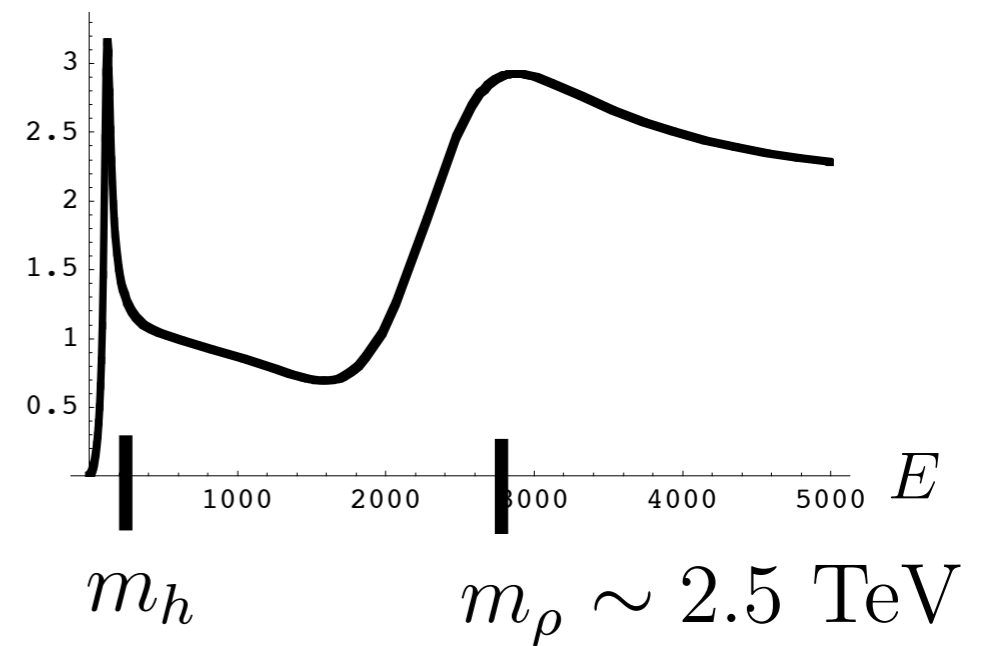
STRONG SECTOR

dynamical

EW SCALE

like QCD pions

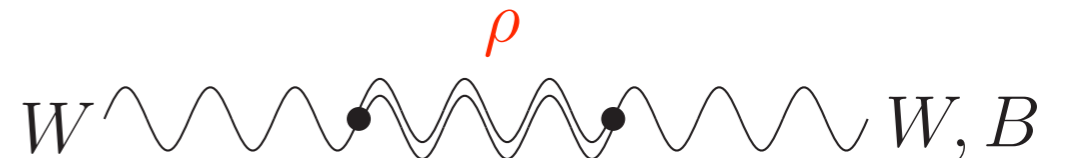
$\mathcal{M}(W_L W_L \rightarrow W_L W_L)$



resonances pushed up in energy

Higgs lighter than other resonances

$$m_h \ll m_\rho$$



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= G/H sigma model scale

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ii) $H \supset O(4) \cong SU(2)_A \times SU(2)_B \times P_{AB}$ → $\delta g_{Zbb} = 0$

↙
 $\langle h \rangle = \left. \begin{pmatrix} 0 \\ 0 \\ 0 \\ v \end{pmatrix} \right\} SO(3) \cong SU(2)_{A+B}$
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h , Higgs complex doublet

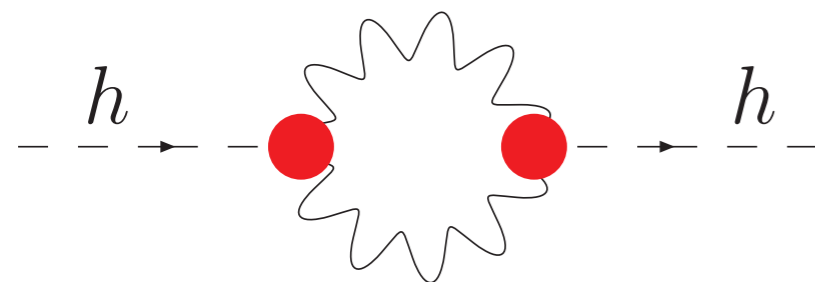
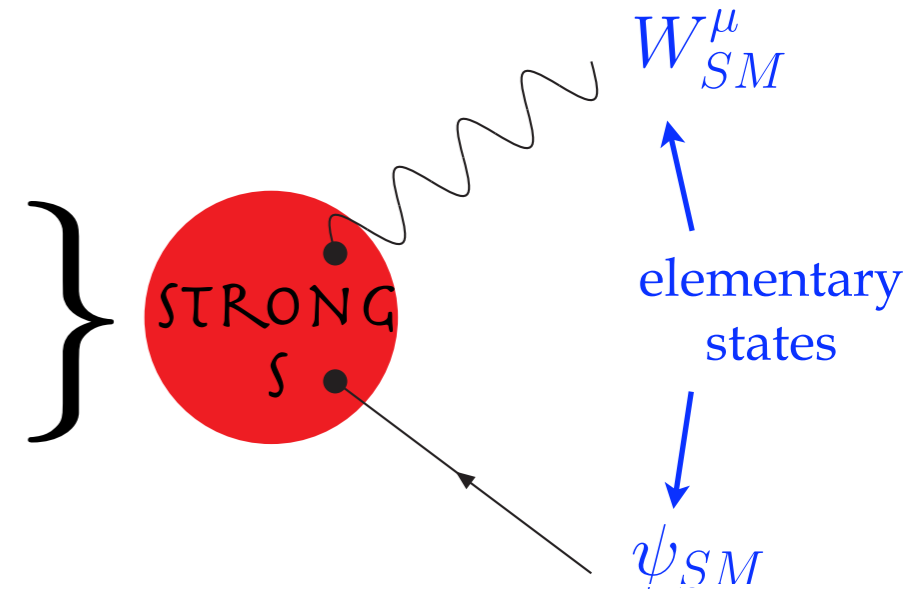
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Explicit Symmetry Breaking:

from the Strong Sector alone: $V(h) = 0$

from interactions with SM fields:

- SM gauge interactions (subgroup K of G is gauged)
- SM fermion interactions (proto-Yukawas, mostly the top)



$$V(h) \neq 0$$



$$\langle h \rangle = v < f$$

$$\xi = \frac{v^2}{f^2} \sim \frac{1}{4}$$

- other sources of explicit breaking? like “quark” masses,...

$$f \gtrsim 500 \text{ GeV}$$

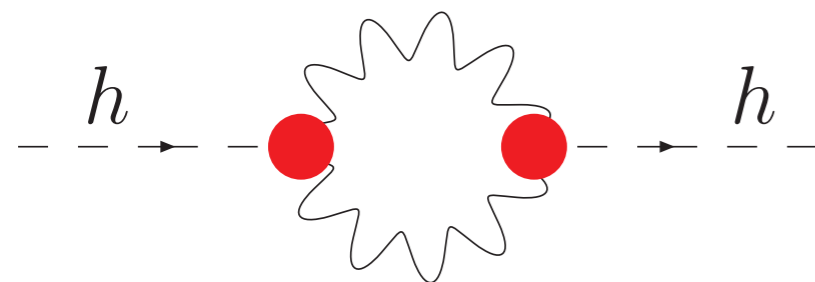
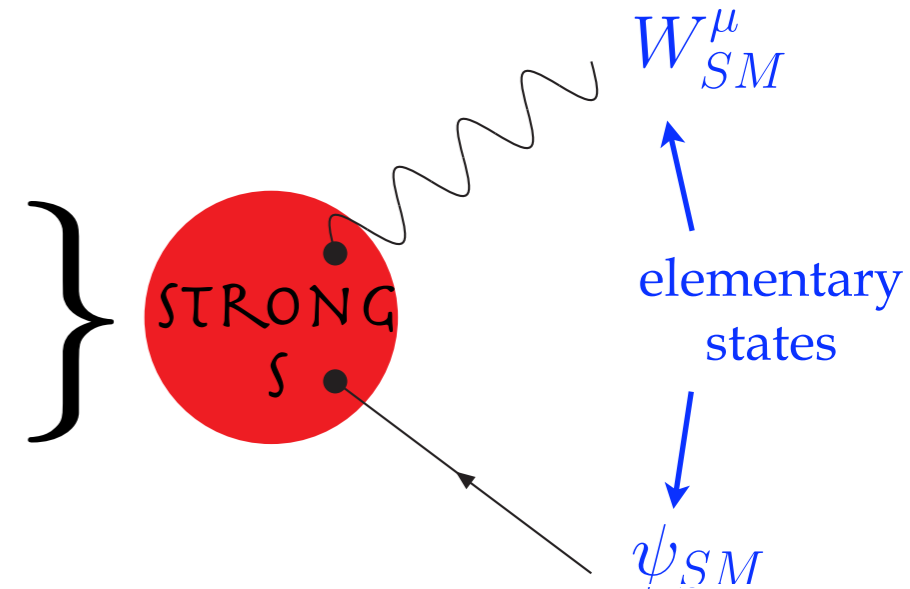
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Calculability:

Effective low-energy “chiral” Lagrangian

+

Holography and Warped extra-d

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Example of *minimal* cosets from constituent "quarks":

$$\text{SU}(4) \longrightarrow \text{Sp}(4) \quad G/H = 5 = 4 + 1 = (2,2) + (1,1)$$

extra singlet

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 **models with 2 Higgs doublets**

2 COMPOSITE-NGB-HIGGS: IS IT VIABLE?

$$h_i = (\mathbf{2}, \mathbf{2}), \quad i = 1, 2$$

➔ Large contributions to $\delta\rho$:

$$\left. \langle h_1 \rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ v_1 \end{pmatrix} \langle h_2 \rangle = \begin{pmatrix} 0 \\ 0 \\ v_2 \\ 0 \end{pmatrix} \right\} \text{SO}(2) \cong \text{U}(1)_Q \quad \text{custodial SO}(3) \text{ breaking}$$

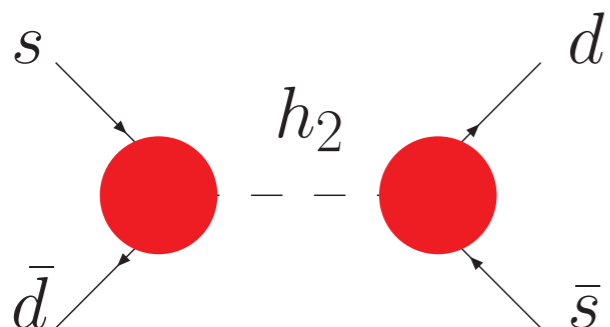
misaligned VEV's

$$\frac{c_T}{f^2} (h_1^T D_\mu h_2)^2 \longrightarrow \delta\rho \sim c_T \frac{v^2}{f^2}$$

$$v^2/f^2 \lesssim 1 \times 10^{-3}$$

➔ Dangerous Flavor transitions:

$$\bar{q}_L \left(y_1^u \tilde{h}_1 + y_2^u \tilde{h}_2 \right) u_R + \bar{q}_L \left(y_1^d h_1 + y_2^d h_2 \right) d_R \quad \text{misaligned Yukawa's}$$



$$\sim \frac{m_s^2 V_{us}^2}{2m_{h_2}^2} (\bar{s}_L d_R)^2 \longrightarrow m_{h_2} \gtrsim 2 \text{ TeV}$$

INERT COMPOSITE-HIGGS MODEL

defined by extra Z_2 :

$$\begin{aligned} h_1 &\rightarrow + h_1 \\ h_2 &\rightarrow - h_2 \quad + \quad \langle h_2 \rangle = 0 \\ \text{SM} &\rightarrow + \text{SM} \end{aligned}$$

Z_2 exactly conserved
NO $(h_2)^{2n}$ couplings

immediate consequences:

- ➔ Automatic solution of $\delta\rho$ and Flavor problems.
- ➔ Stability of lightest h_2 component (DM candidate).
- ➔ Double h_2 production @ colliders.

EXAMPLE: $G = O(6) \longrightarrow H = O(4) \times O(2)$

INERT HIGGS: $O(6)/O(4) \times O(2)$

Breakings:

$(SU(2)_L, SU(2)_R)_{Z_2}$

spontaneous, $\Sigma(h_1, h_2) = O(6)/O(4) \times O(2) \in \mathbf{20}'$
 $= (\mathbf{2}, \mathbf{2})_+ + (\mathbf{2}, \mathbf{2})_-$

explicit, $\psi_{SM} \in \mathbf{6} \quad \mathbf{6} = (\mathbf{2}, \mathbf{2})_+ + (\mathbf{1}, \mathbf{1})_+ + (\mathbf{1}, \mathbf{1})_-$

$q_L = (u_L, d_L) \quad u_R, d_R$

$W_{SM}^\mu \in \mathbf{15}$

$\mathbf{15} = (\mathbf{1}, \mathbf{1})_+ + (\mathbf{3}, \mathbf{1})_+ + (\mathbf{1}, \mathbf{3})_+ + (\mathbf{2}, \mathbf{2})_+ + (\mathbf{2}, \mathbf{2})_-$

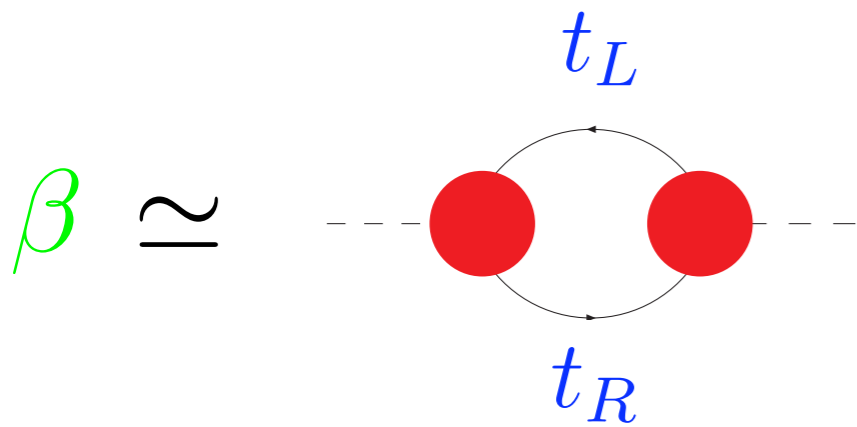
W_L

B_Y

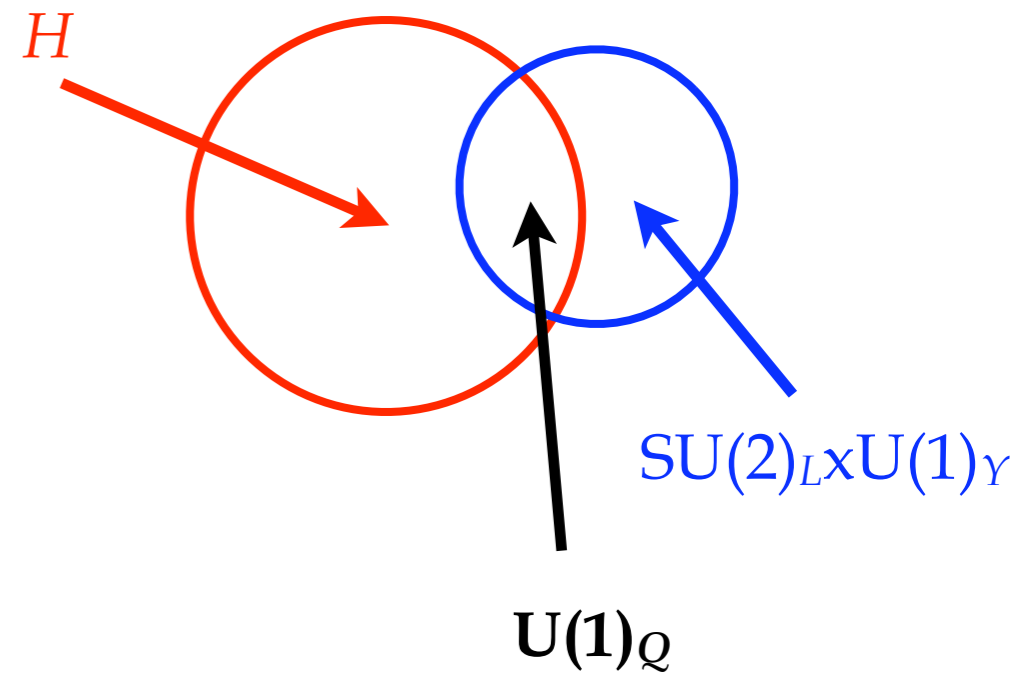
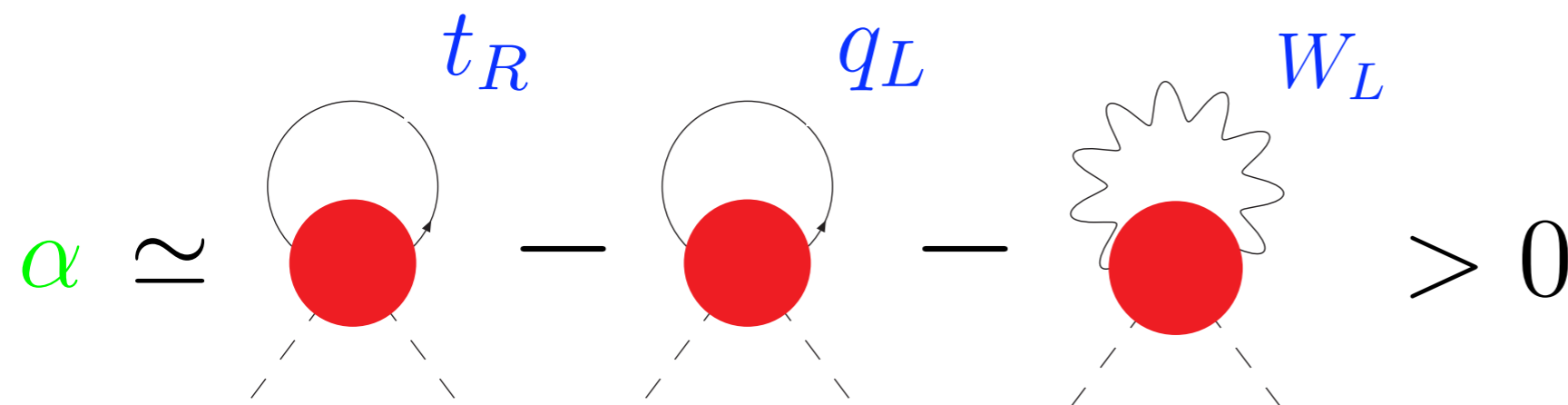
INERT HIGGS: $O(6)/O(4) \times O(2)$

EW Symmetry Breaking: $\langle h_2 \rangle = 0$

$$V(h_1) = \alpha \cos(h_1/f) - \beta \sin^2(h_1/f) \longrightarrow \cos(h_1/f) = -\alpha/2\beta$$

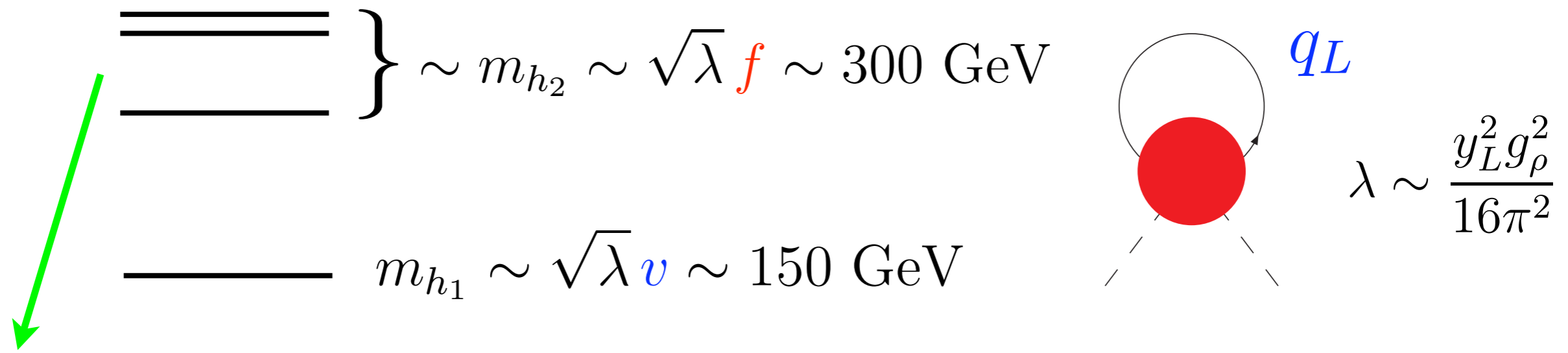


thanks to large top-loops



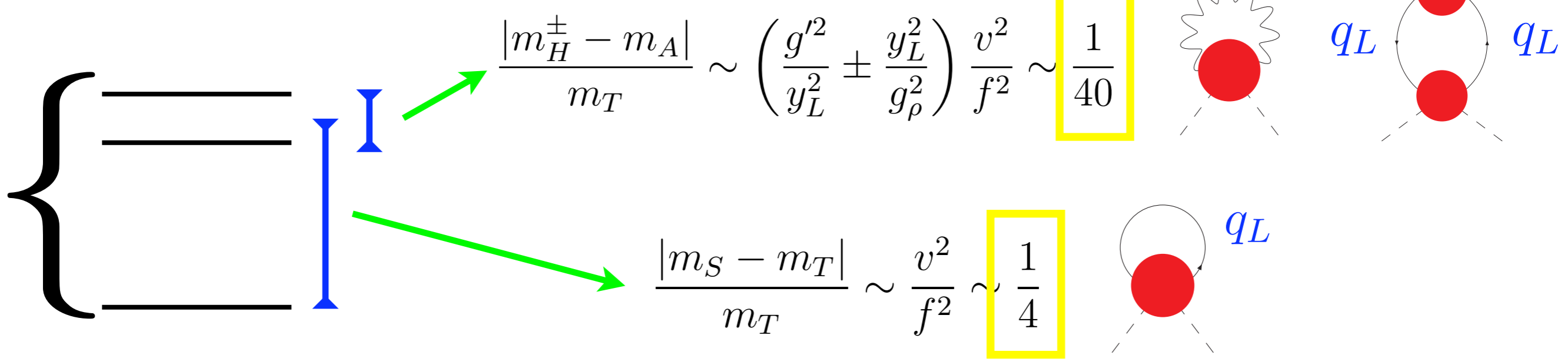
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Spectrum:



splittings, $h_2 \rightarrow T = (H^\pm, A) \oplus S$

$SO(4) \quad 4 = 3 + 1 \quad SO(3)$



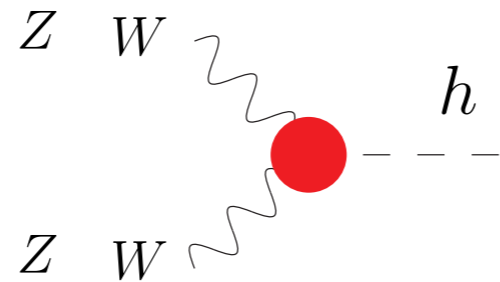
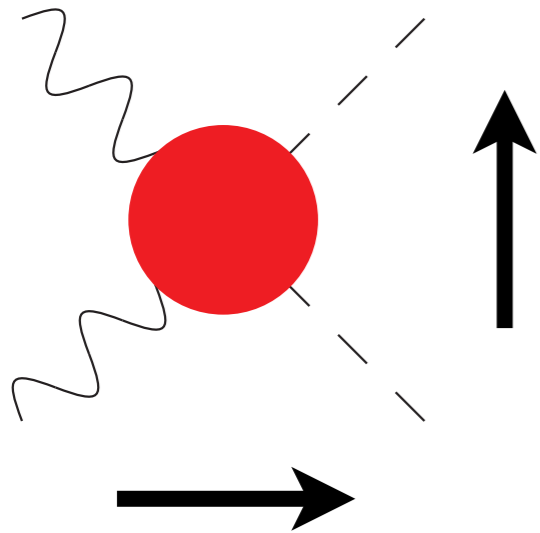
SO(6)/SO(4) x SO(2)

Signatures of composite-NGB: $\mathcal{L}_{\text{coset}}[h_1/f, h_2/f] = \frac{f^2}{2} \text{Tr}[|D_\mu \Sigma|^2]$

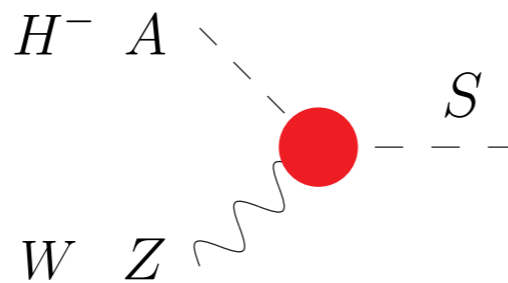
$$= \frac{1}{2} (D_\mu h_j)^2 + \frac{c_\partial}{f^2} (h_i \partial_\mu h_i)^2 + \frac{c_W}{f^2} (h_i W_\mu h_j)^2 + \frac{c_{\partial W}}{f^2} (h_i \partial_\mu h_i) (h_j W^\mu h_j)$$

$$\langle h_1 \rangle = v, \quad \langle h_2 \rangle = 0$$

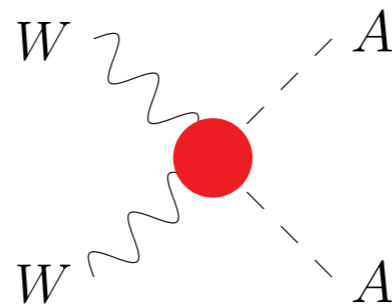
non-unitarity of
amplitudes,



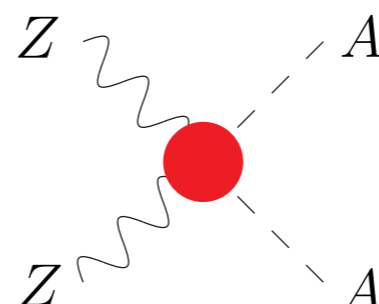
$$\sum_i g_{h_i^0 WW}^2 = g^2 m_W^2 \left(1 - \frac{2v^2}{3f^2} \right)$$



$$\sum_i g_{h_i^0 AZ}^2 = \frac{g^2}{\cos^2 \theta_W} \left(1 - \frac{1v^2}{6f^2} \right)$$



÷



$$= \cos^2 \theta_W \left(1 + \frac{v^2}{f^2} \right)$$

CONCLUSIONS

If the hierarchy problem is solved by **strong dynamics**, we can expect rich phenomenology of pseudo-NGB's.



motivated framework for extended Higgs sectors

example, **Inert composite-Higgs model** (not near MSSM)

→ Completely **viable**

→ “predictions”:
large h_2 masses, **$m_{h_2} \sim 300\text{GeV}$**
small custodial splittings, **$\Delta m_T \sim 10\text{GeV}$**
 $\mathcal{O}(v^2/f^2) \sim 20\%$ deviations in couplings