

FROM SOFT WALLS TO INFRARED BRANES

Gero von Gersdorff (École Polytechnique)
Planck 2010

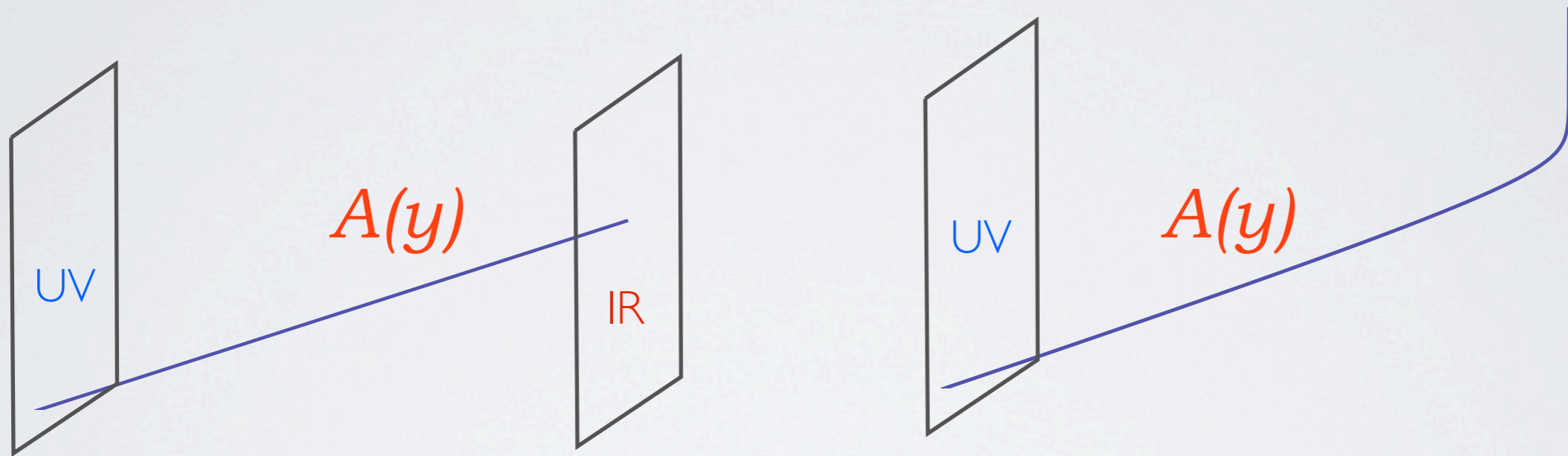
ArXiv 1005.5134 (hep-ph)

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Soft Wall models are generalizations of RS that do not possess an IR brane

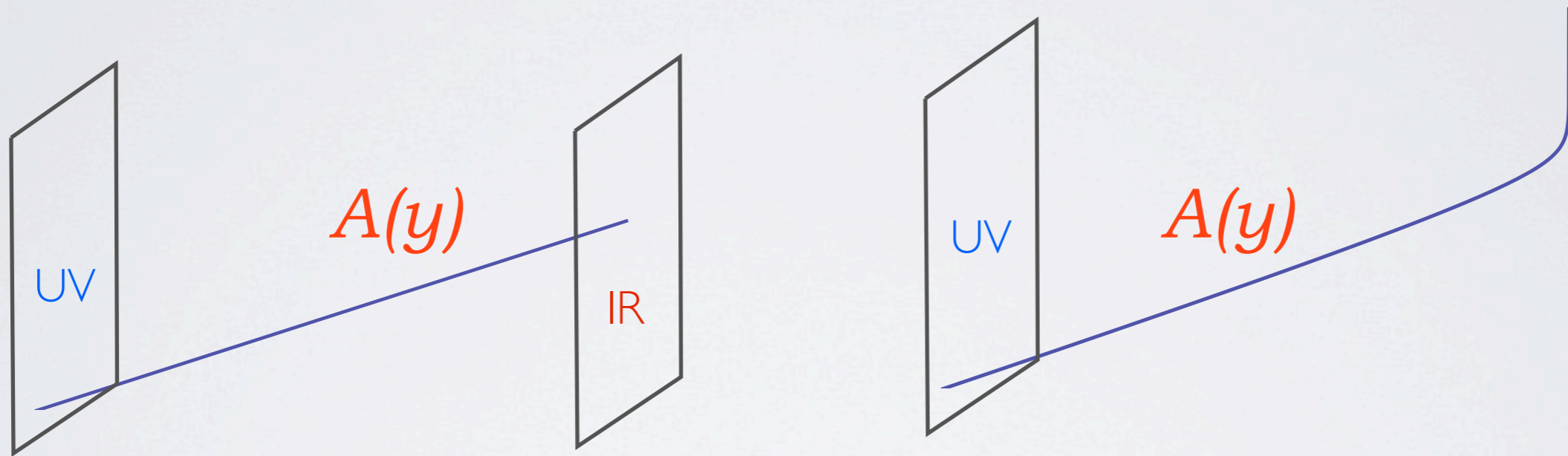
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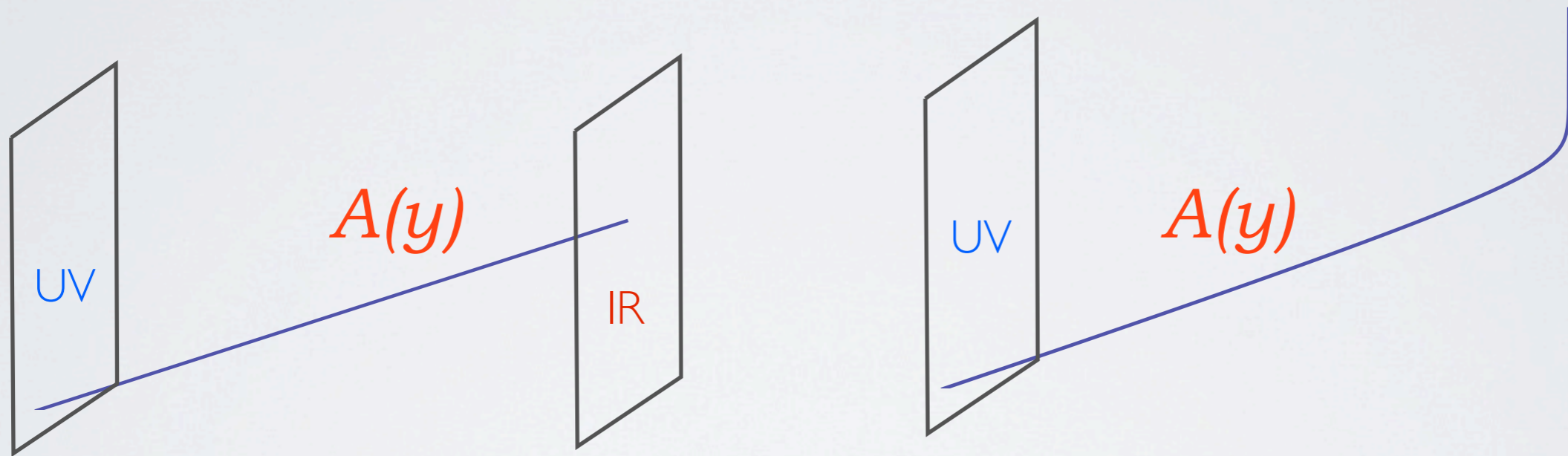


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Singularities are “good” if

$$e^{-A(y)} \rightarrow 0 \text{ faster than } (y_s - y)^{\frac{1}{4}}$$

Gubser'02, Forste et al '02
Cabrer, GG & Quirós '09

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 - Need to specify infinitely many operators?
- Soft Walls give “**resolution**” of IR brane

CLASSES OF SOFT WALLS

Soft Walls are rather restricted by demanding:

- ◆ A Gap in the spectrum should exist
- ◆ The singularity should be a “good”

Then, asymptotics near singularity fall into 2 classes

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Proper Length coordinates

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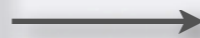
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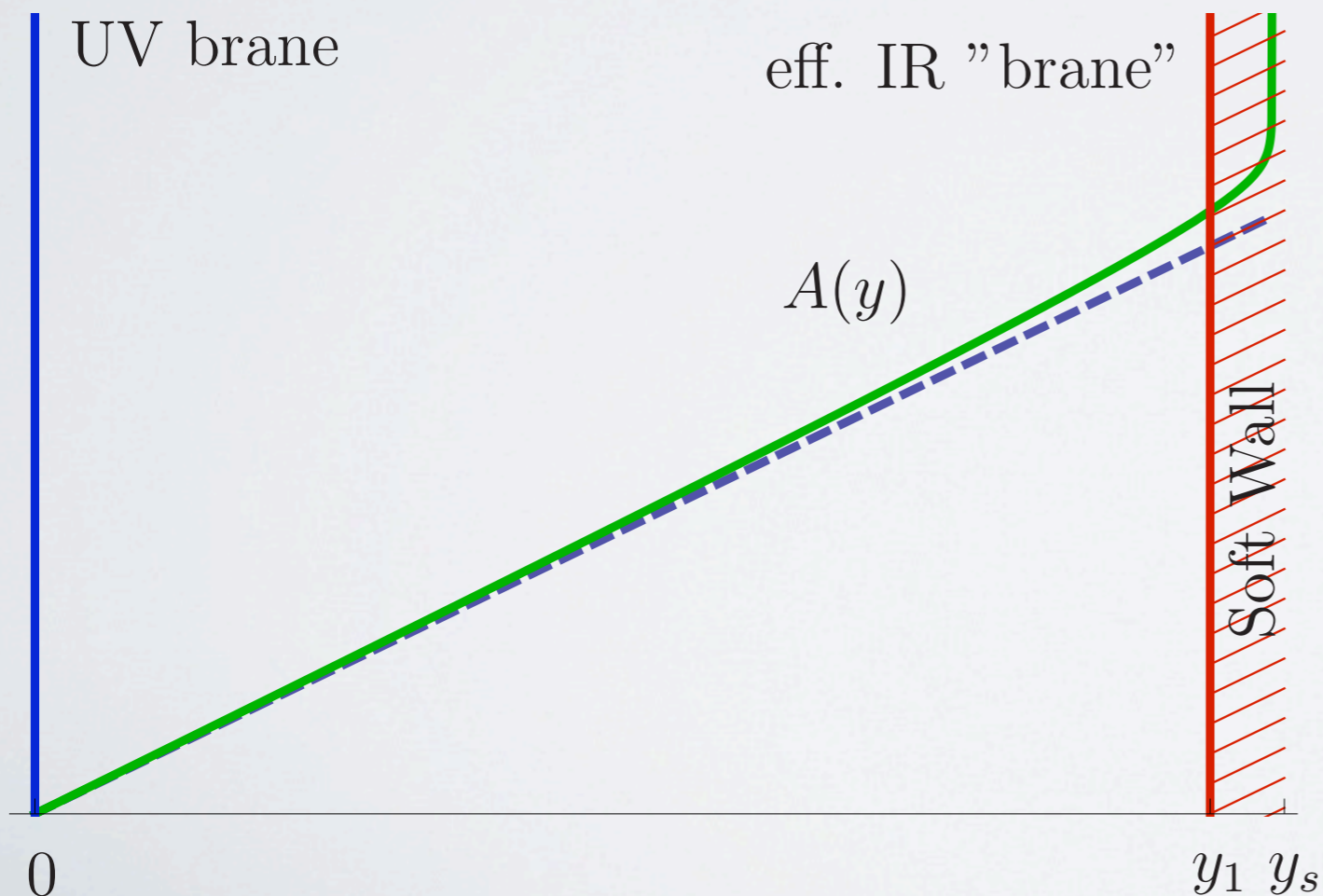
PLAN

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Integrate over the region near the singularity

$$S = \int_0^{y_s} d^5x \mathcal{L}_{\text{bulk}} \longrightarrow S = \int_0^{y_1} d^5x [\mathcal{L}_{\text{bulk}} + \mathcal{L}_{\text{SW}} \delta(y - y_1)]$$

Equivalent description of SW in terms of IR brane

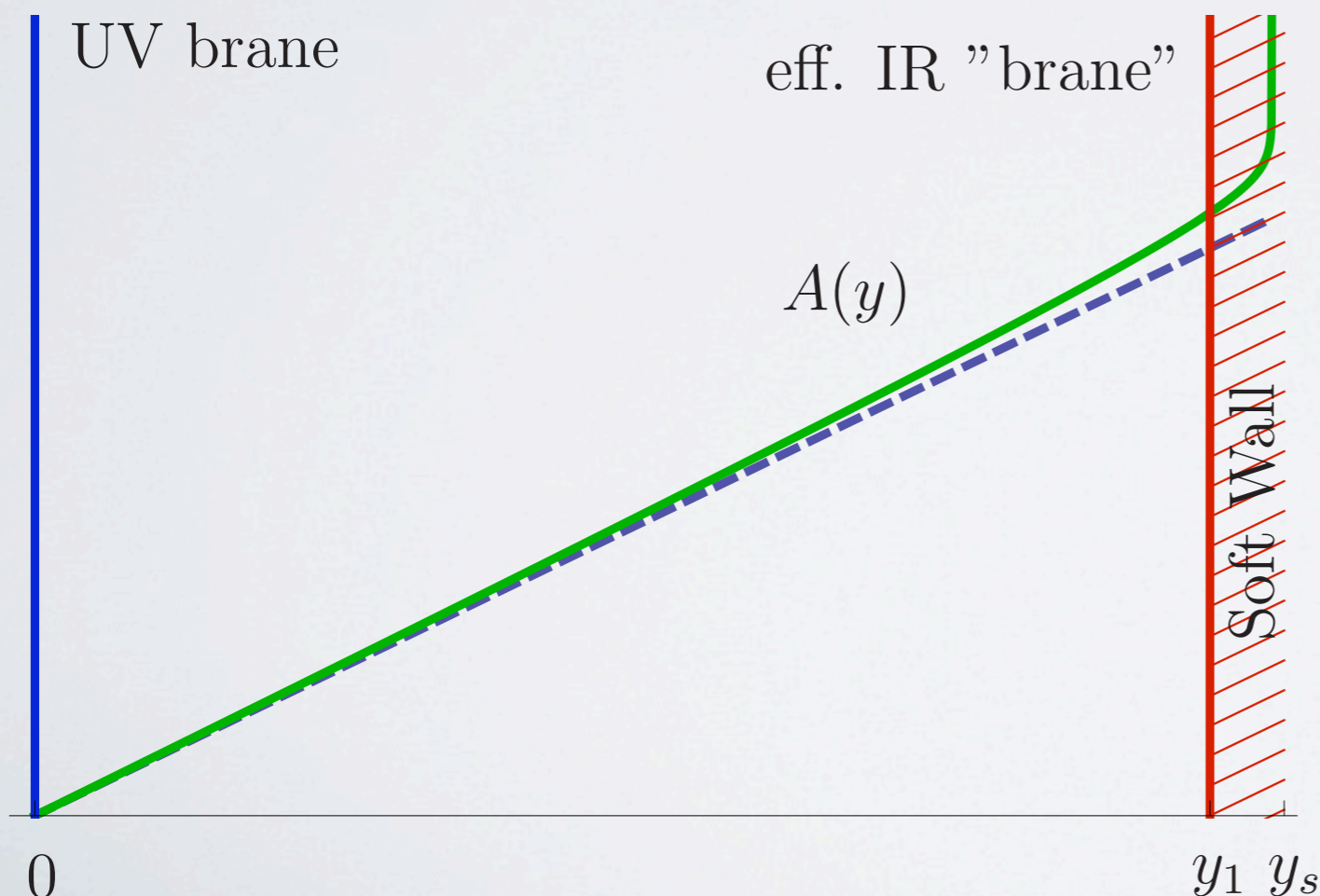


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Equivalent description of SW in terms of IR brane



- ◆ For y_1 close to y_s close IR Lagrangian "universal"
- ◆ Facilitates comparison with standard 2 brane compactif.
- ◆ IR Lagrangian makes sense even for $p > \text{TeV}$
- ◆ Useful approximation scheme: Approximate "new" bulk by RS metric

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$$(e^{2A(y)} \partial_\mu \partial^\mu + \partial_y^2 - 4A'(y) \partial_y - M^2) K(x, x'; y) = 0 \quad \text{Witten '98}$$

$$K(x, x'; y_1) = \delta(x - x') \quad \psi(x, y) = \int d^4 x' K(x, x'; y) \psi(x', y_1)$$

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Plug back into action

$$\mathcal{F}(p) = K'(p, y_1)$$
$$\mathcal{S}_{\text{SW}} = \frac{1}{2} \int \frac{d^4 p}{(2\pi)^4} \sqrt{-g(y)} \mathcal{F}(p) \psi(-p, y) \psi(p, y) \delta(y - y_1)$$

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$$\mathcal{F}(p) = e^{A(z_1)} \frac{\sqrt{-p^2} J_{\alpha+1}(\sqrt{-p^2} \Delta_z)}{J_{\alpha}(\sqrt{-p^2} \Delta_z)} \quad \alpha = \frac{4 - \nu^2}{2(\nu^2 - 1)} ;$$

RESULTS FOR SW1, SW2

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SW2

$$\mathcal{F}(p^2) = e^{A(z_1)} \sqrt{-p^2} \tan \left(c_{\beta} \left[-\frac{p^2}{\rho^2} \right]^{\frac{1}{4\beta}} \right) \quad \beta = \frac{\sigma - 1}{2\sigma}$$

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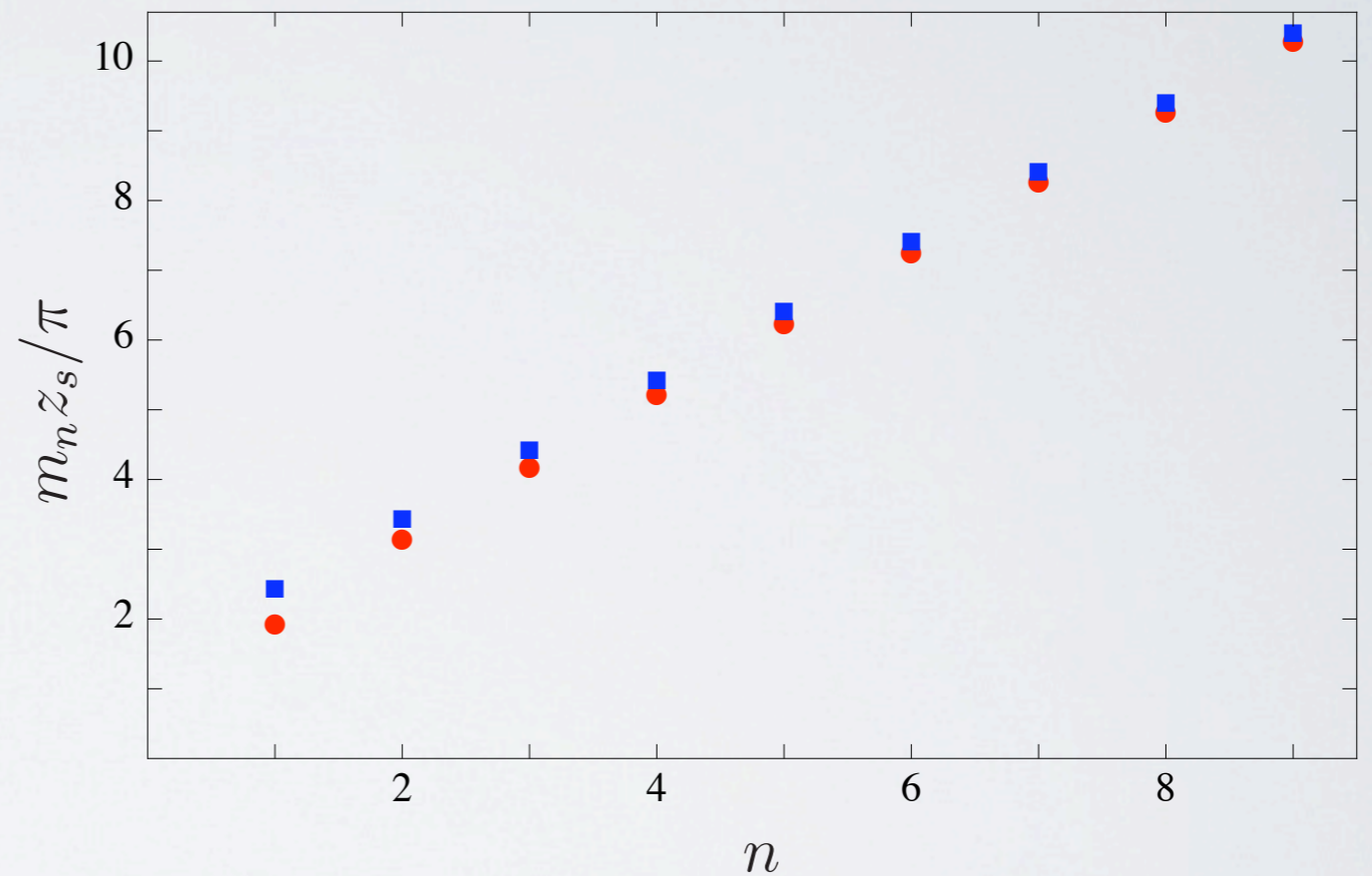
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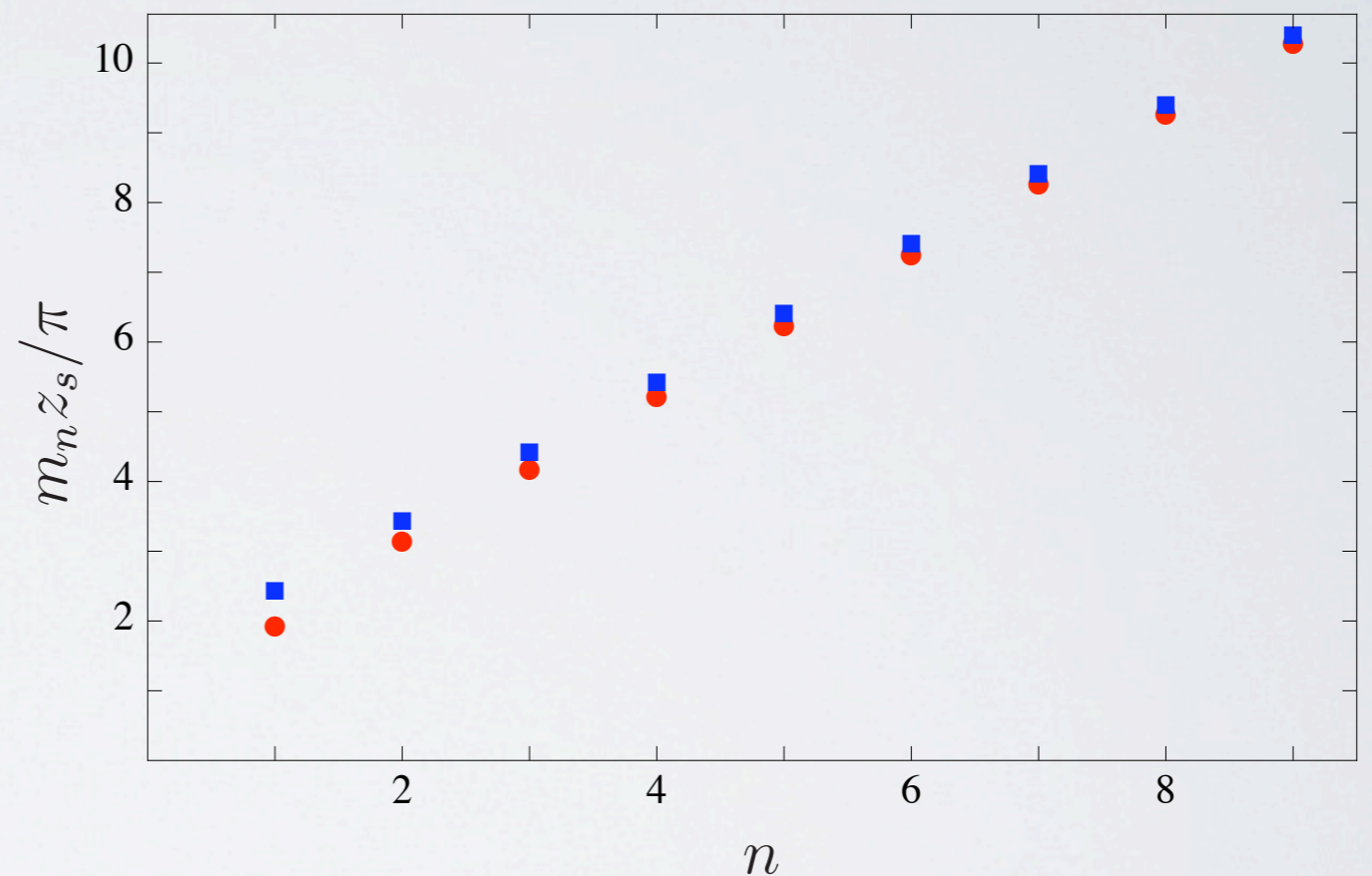


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Asymptotic Spectrum (n large)

Approximation $m_n z_s / \pi = n + 0.34$

Numerical $m_n z_s / \pi = (10.40, 50.37, 100.36, 500.35)$.

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Compare with **exact** solution (analytic only for $\beta = \frac{1}{4}$)

Exact spectrum is $m_n = \sqrt{12 n} \rho$

Concides with previous method since $c_{\frac{1}{4}} = \frac{\pi}{12}$

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- Future work: higher spin, symmetry breaking...