

CMB and Secondaries: the Cold Spot

Isabella Masina

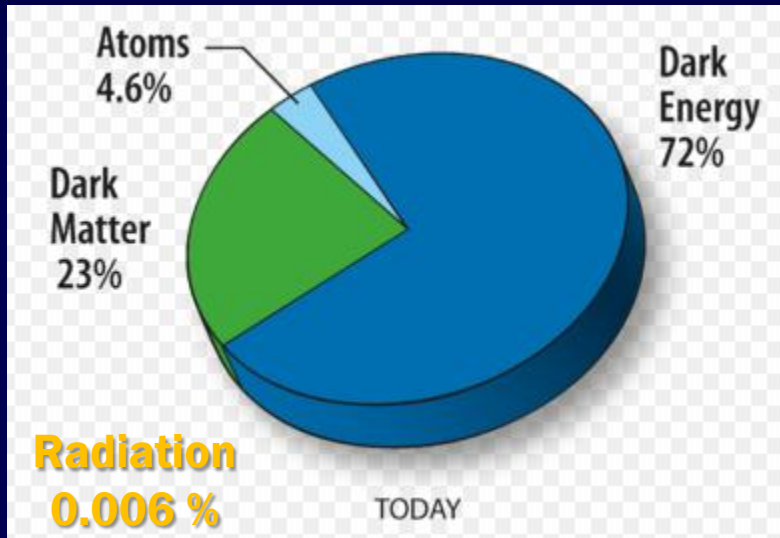
(Ferrara U., Italy & CP³-SDU, Denmark)

Based on: I.M, A.Notari, JCAP 0902:019,2009. [arXiv:0808.1811](#)
JCAP 0907:035,2009. [arXiv:0905.1073](#)

Plan

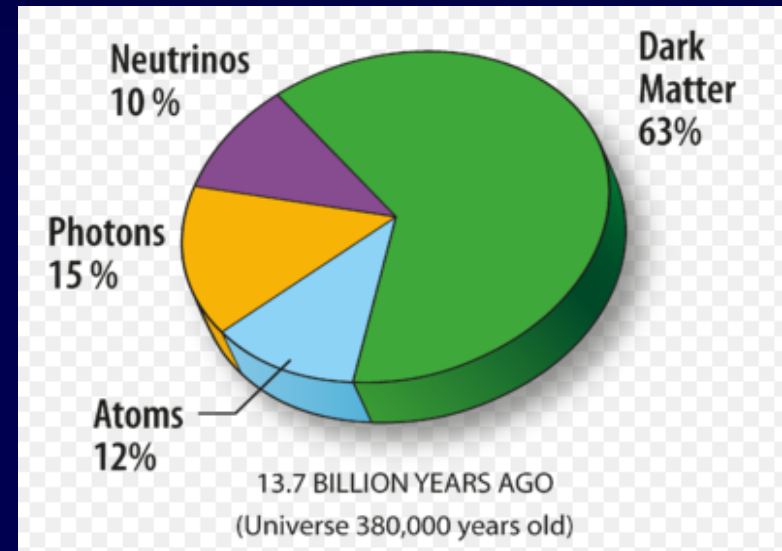
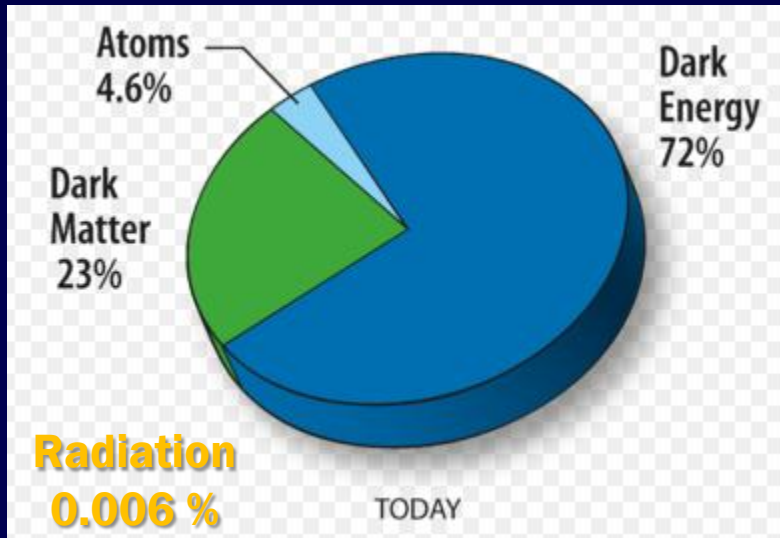
- Inflationary BB and CMB
- Unexpected features: the Cold Spot
- The Cold Spot as a Void on the line of sight:
secondary anisotropies associated to
the Rees-Sciama & Lensing effects
- Conclusions and perspectives

Content of Universe



radiation density scales as a^{-4}
WHILE
matter density as a^{-3} & dark energy stays constant

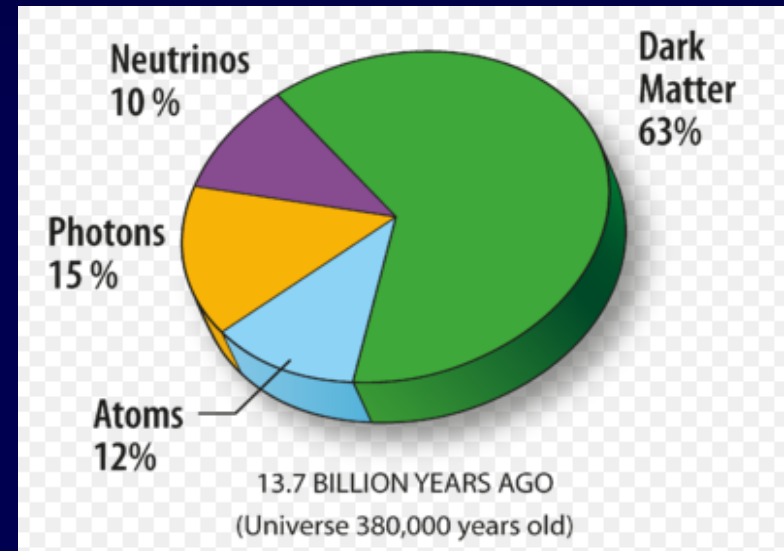
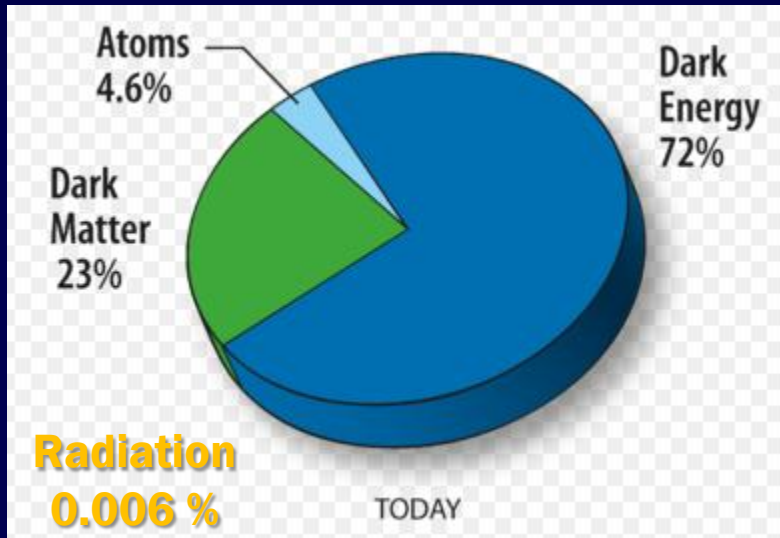
Content of Universe



but there has been a time when radiation and matter densities were comparable:

13.7 Gyr ago

Content of Universe



but there has been a time when radiation and matter densities were comparable:

According to BB theory,
the CMB gives a snapshot of the universe at that time

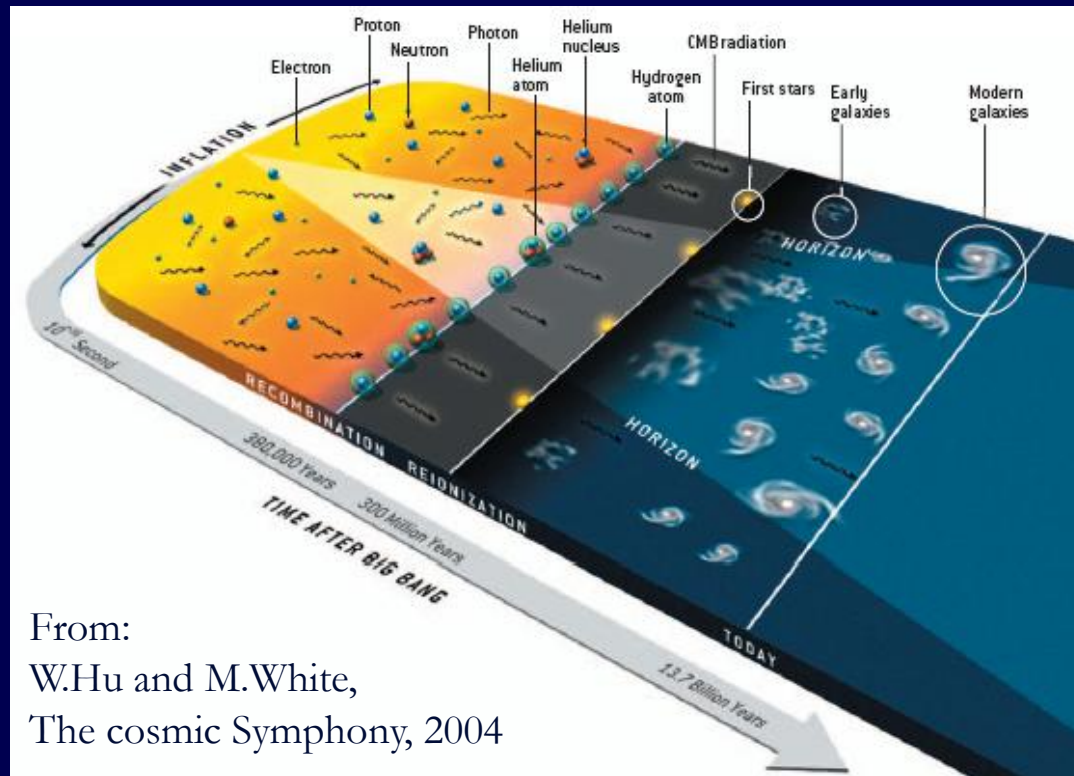
Fiat lux

Because it was just at that time that T dropped enough to allow e and p to form H atoms, thus making the **universe transparent to radiation** (since then in the eq with matter, including the dark one).



Sistine Chapel – Separation of Light and Darkness

This “last scattering” or “recombination” or “decoupling” period occurred
in between **0.38 – 0.48 Myr** after BB
when the universe was **about 1000 times smaller** ($z=1100$)
and had temperature **T about 3000 K** (kT about 0.25 eV).



From:
W.Hu and M.White,
The cosmic Symphony, 2004

Anisotropies and Inflationary BB

No model other than the inflationary BB has yet explained the temperature fluctuations or ANISOTROPIES, which are

- about 10^{-5} (rms variation $18\mu\text{K}$)
- more pronounced on 1° (twice full moon).

Contributions to CMB anisotropy

PRIMARY due to the physics at the LSS and before

→ dominant and linked to fundamental cosmological parameters

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SECONDARY induced on photons in their travel from the LSS to us

Those associated with a gravitational potential are:

- ❖ Integrated Sachs-Wolfe
- ❖ Rees-Sciama
- ❖ Gravitational lensing

→ generically small but could plague extraction of fundamental parameters.

Spherical harmonics expansion

$$T(\theta, \phi) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\theta, \phi)$$

represents a very useful tool to study anisotropies

BECAUSE

theoretical models for inflation generically predict the PRIMARY $a_{\ell m}$ to be nearly Gaussian random fields.

Correlation functions

Brackets stand for a statistical average over an ensemble of possible realizations of the Universe

- 2-point correlation function or
POWER SPECTRUM
- 3-point correlation function or
BISPECTRUM
- 4-point correlation function or
TRISPECTRUM
- ...

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$$\langle a_{\ell_1 m_1}^{(P)} a_{\ell_2 m_2}^{(P)*} \rangle = \delta_{\ell_1 \ell_2} \delta_{m_1 m_2} \langle C_{\ell_1}^{(P)} \rangle$$

$$\langle a_{\ell_1 m_1}^{(P)} a_{\ell_2 m_2}^{(P)} a_{\ell_3 m_3}^{(P)} \rangle = 0$$

Higher statistics correlation functions are fully determined by the power spectrum

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HOWEVER

secondary anisotropies are not random gaussian!

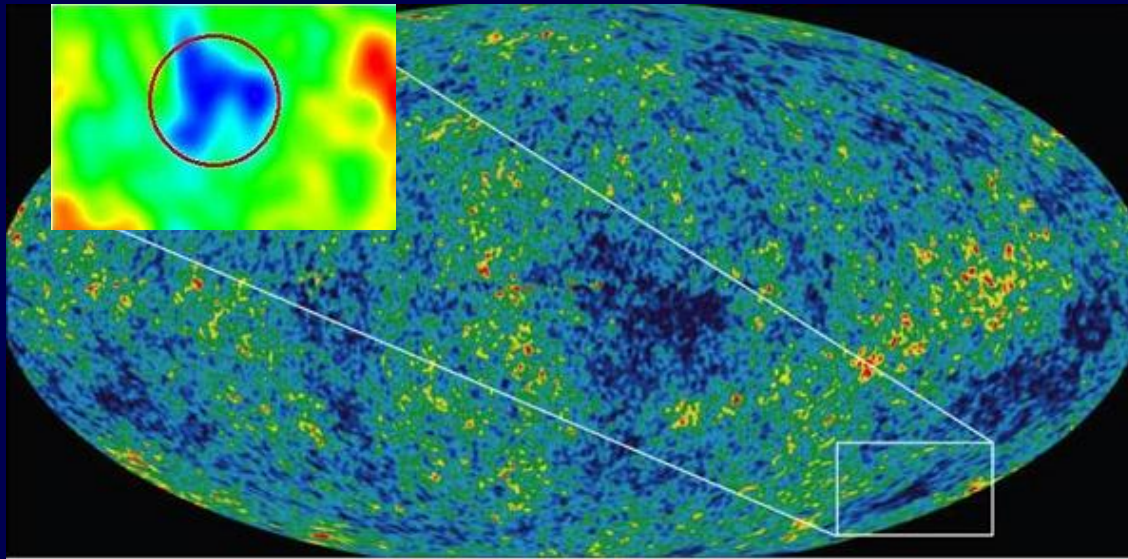
They could be studied by looking at deviations from the predictions above

**Let's apply all this
to a concrete case:**

The Cold Spot

The Cold Spot

Large circular region on about 10° angular scale which is anomalously cold: $\Delta T = 190 \pm 80 \mu\text{K}$ [A = $(7 \pm 3) \times 10^{-5}$]



Probability of this spot to come from Gaussian fluctuations has been estimated to be $< 2\%$ [Cruz et al., astro-ph/0603859]

→ explore other possibilities, e.g. secondaries

The Cold Spot as a “Void”

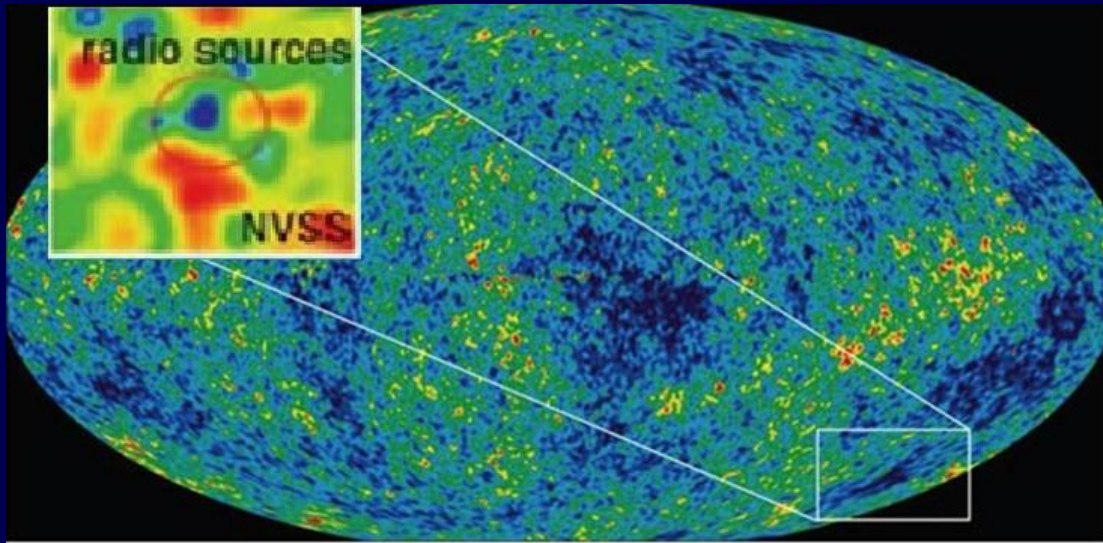
K. Tomita, Phys. Rev. D 72, 103506 (2005) [Erratum-ibid. D 73, 029901 (2006)] [astro-ph/0509518].

K. T. Inoue and J. Silk, Astrophys. J. 648, 23 (2006) [astro-ph/0602478]; Astrophys. J. 664, 650 (2007) [astro-ph/0612347].

suggested it could be due to a
large spherical underdense region (of some unknown origin),
on the line of sight between us and the LSS.

Further support ?

McEwen et al. (2006) & Rudnick et al. (2007) claimed that looking at the direction of the Cold Spot in the Extragalactic Radio Sources (NVSS survey), an underdense region is visible at $z \sim 1$



Smith et al. (2008) challenge this claim.

Granett et al. (2009) find no underdense region at $z < 1$.

Bremer et al. (2010) confirm Granett et al.

The Cold Spot as a “Void”

[I.M, A.Notari, JCAP 0902:019,2009. arXiv:0808.1811 , JCAP 0907:035,2009.arXiv:0905.1073]

Modelling it through an inhomogeneous LTB metric (requires an overdense compensating shell),

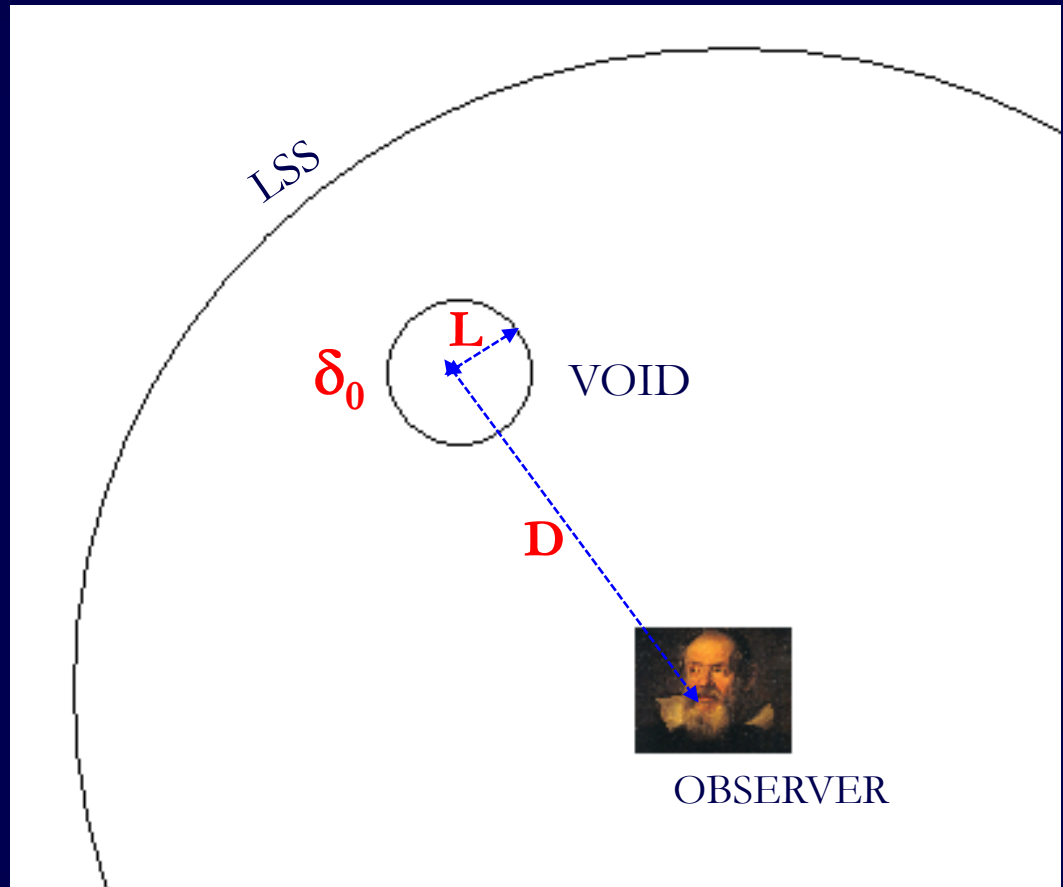
we computed the secondary effects occurring to the CMB photons that travel through the Void:

Rees-Sciama & Lensing

L comoving radius,

D comoving distance,

δ_0 density contrast at centre today



Radius-Redshift relation

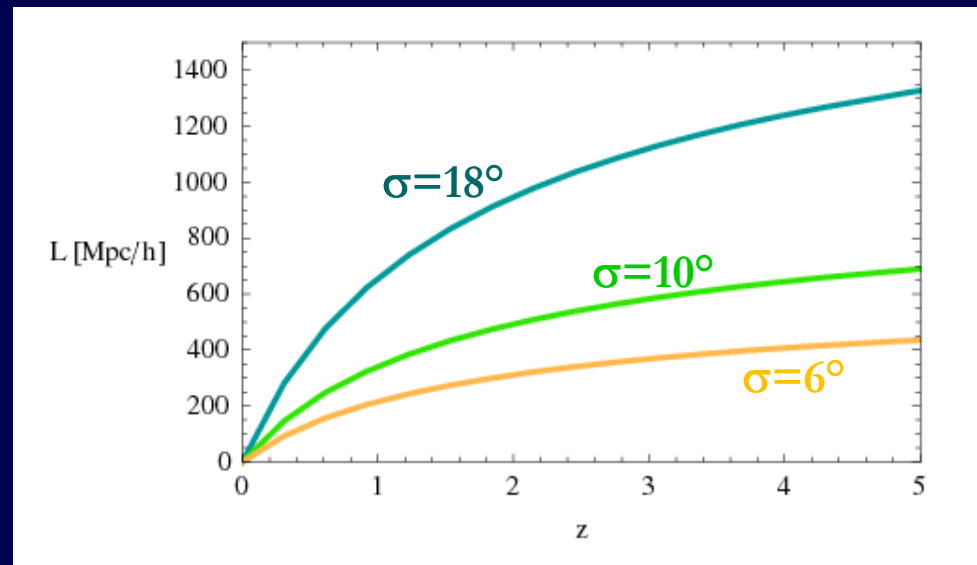
$$L = \frac{2 \tan \theta_L}{H_0} \left(1 - \frac{1}{\sqrt{1+z}} \right)$$

$$\tan \theta_L = L/D$$

$z > 1$ means indeed large Void:

$L > 300 \text{ Mpc}/h$ for cold
spot angular size **$\sigma = 10^\circ$**

Such big underdense region should
clearly come from a different mechanism
than standard inflation, maybe bubble
nucleation due to phase transitions or ...



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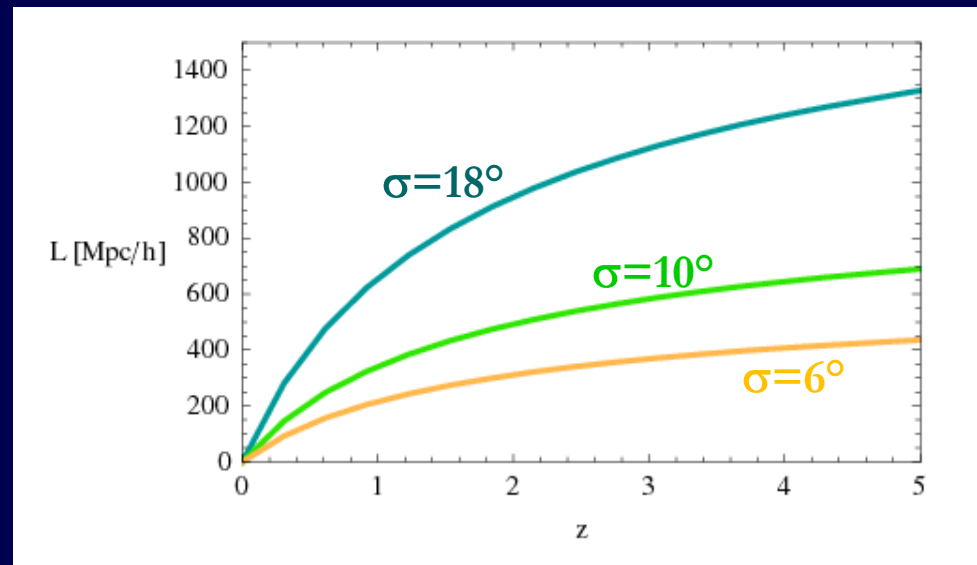
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HERE we show that **Planck** will be able to reject or confirm
this hypothesis via the study of secondary effects

The procedure

1. Find the spherical harmonics decomposition of the i -th ($i=RS, L$) temperature anisotropy profile (\hat{n} is the direction of observation)

$$a_{\ell m}^{(i)} \equiv \int d\hat{n} \frac{\Delta T^{(i)}(\hat{n})}{T} Y_{\ell m}^*(\hat{n})$$

2. Add them all

$$a_{\ell m} = a_{\ell m}^{(P)} + a_{\ell m}^{(RS)} + a_{\ell m}^{(L)}$$

3. Estimate 2 and 3 point correlation functions
4. Look for quantities that would vanish in case of random Gaussian alm

Rees-Sciama

Passing through a Void, photons suffer some blue-shift due to the fact that the gravitational potential ϕ is not exactly constant in time – the so-called Integrated Sachs-Wolfe (1966)

$$\frac{\Delta T}{T}(\vec{n}) = \frac{\Delta T}{T}^{(P)}(\vec{n}) + 2 \int_{\tau_{rec}}^{\tau_0} \dot{\phi} d\tau$$

line-of-sight integral over the conformal time from recombination to present time

The **RS** effect is the ISW part associated to the variation of ϕ due to **non-linear** effects.

Actually, just the linear level effect is the one usually called “ISW” effect.

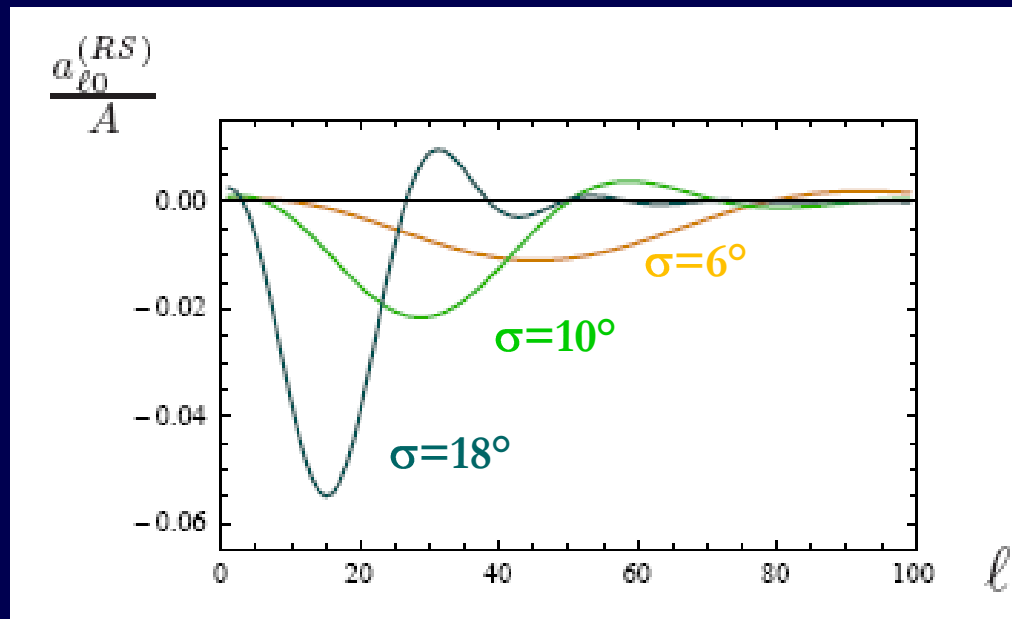
It would vanish in a matter dominated flat Universe. It is significant only when a Dark Energy component becomes dominant with respect to matter ($z < 1$).

→ here focus attention on RS effect, which is always present
(we checked that ISW is not bigger than RS)

For RS, due to spherical symmetry, the only non-vanishing alm are those with $m=0$
 (axis z pointing from observer to the center of the Void)

A is the amplitude of T fluctuation
 at Void centre (fitted experimentally)

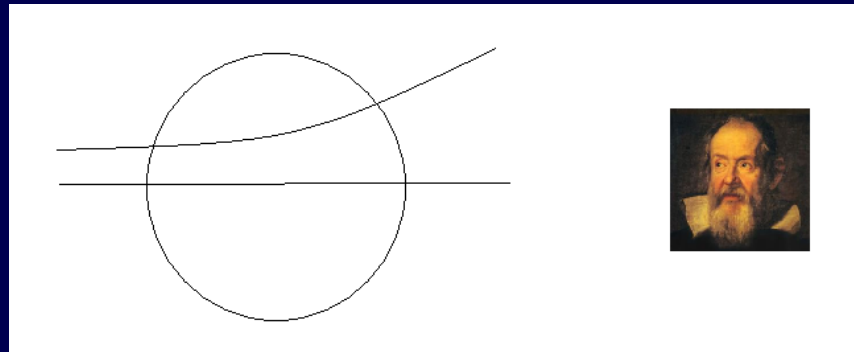
$$A \approx 0.5 \delta_0^2 \frac{(LH_0)^3}{\sqrt{1+z}}$$



Lensing

of the primordial perturbations.

See also: Das, Spergel, arXiv: 0809.4704



For Lensing we need the so-called Lensing potential, related to the gravitational potential by

$$\nabla_{\perp} \Theta = -2 \int_{\tau_{LSS}}^{\tau_0} d\tau \frac{\tau_{LSS} - \tau}{\tau_{LSS}} \nabla_{\perp} \Phi$$

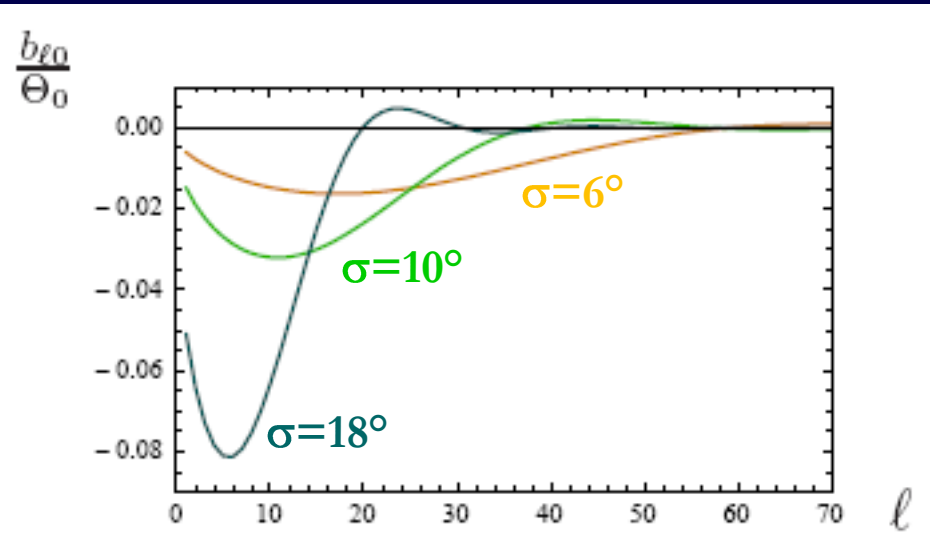
Now define

$$b_{\ell m} \equiv \int d\hat{n} \Theta(\hat{n}) Y_{\ell m}^*(\hat{n})$$

due to spherical symmetry, the only non-vanishing blm are those with $m=0$

Θ_0 is the amplitude of the lensing potential at Void centre

$$\Theta_0 \approx \left(\frac{A L H_0 \tan^2 \theta_L}{1 - \frac{L H_0}{2 \tan \theta_L}} \right)^{1/2}$$



At 1st order

$$a_{\ell m}^{(L)(1)} = \sum_{\ell', \ell''} G_{\ell' \ell''}^{-mm0} \frac{\ell'(\ell'+1) - \ell(\ell+1) + \ell''(\ell''+1)}{2} a_{\ell' - m}^{(P)*} b_{\ell'' 0}$$

Diagonal 2-p function

(power spectrum)

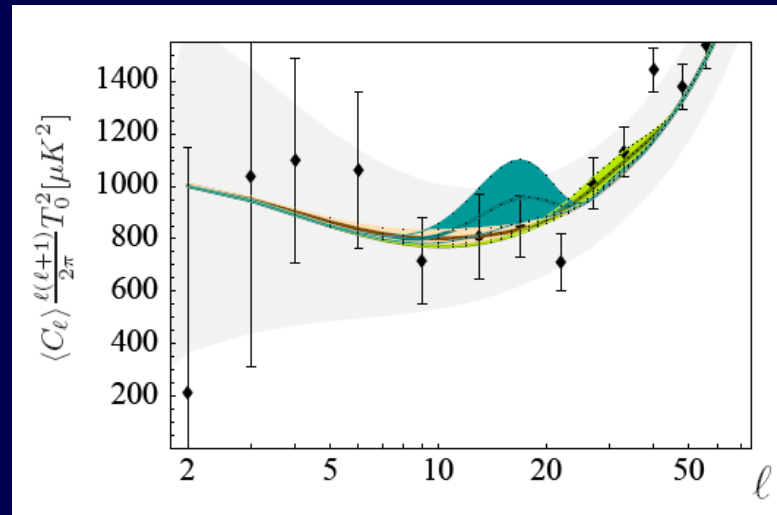
$$\langle a a \rangle = \langle P P \rangle + \cancel{\langle P R S \rangle} + \langle R S R S \rangle + \underbrace{\langle P L \rangle + \langle L L \rangle}_{\text{negligible}}$$

vanishing
(P&RS uncorrelated)

negligible

$\langle P L_1 \rangle = 0 \rightarrow \langle L_1 L_1 \rangle$ and $\langle P L_2 \rangle$

$$\langle C_\ell \rangle = \langle C_\ell^{(P)} \rangle + C_\ell^{(RS)}$$



There is a slight -negligible- increasement in chi-square.

Bispectrum

The basic quantities are the B coefficients

$$B_{l_1 l_2 l_3} = \sum_{m_1, m_2, m_3} \begin{pmatrix} l_1 & l_2 & l_3 \\ m_1 & m_2 & m_3 \end{pmatrix} a_{l_1 m_1} a_{l_2 m_2} a_{l_3 m_3}$$

once we calculate the statistical average,

$$\begin{aligned} \langle a a a \rangle = & \langle \cancel{P P P} \rangle + \langle \cancel{P P RS} \rangle + \langle \cancel{P P L} \rangle + \langle \cancel{P RS RS} \rangle + \langle \cancel{P L L} \rangle \\ & + \langle \cancel{RS L L} \rangle + \langle \cancel{RS RS L} \rangle + \langle P L RS \rangle + \langle RS RS RS \rangle + \langle \cancel{L L L} \rangle \end{aligned}$$

non-trivial

negligible

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non-trivial
negligible

the only contributions left are

1. RS^3

$$\langle B_{\ell_1 \ell_2 \ell_3}^{(RS)} \rangle = \sum_{m_1, m_2, m_3} \begin{pmatrix} \ell_1 & \ell_2 & \ell_3 \\ m_1 & m_2 & m_3 \end{pmatrix} \langle a_{\ell_1 m_1}^{(RS)} a_{\ell_2 m_2}^{(RS)} a_{\ell_3 m_3}^{(RS)} \rangle,$$

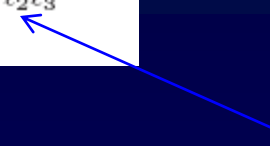
2. $P L RS$

$$\langle B_{\ell_1 \ell_2 \ell_3}^{(PLRS)} \rangle = \sum_{m_1, m_2, m_3} \begin{pmatrix} \ell_1 & \ell_2 & \ell_3 \\ m_1 & m_2 & m_3 \end{pmatrix} \langle a_{\ell_1 m_1}^{(P)} a_{\ell_2 m_2}^{(L)} a_{\ell_3 m_3}^{(RS)} \rangle + (5 \text{ permutations}).$$

For a signal labeled by i , the **SIGNAL TO NOISE** ratio is

$$(S/N)_i = \frac{1}{\sqrt{F_{ii}^{-1}}} , \quad F_{ii} = \sum_{2 \leq l_1 \leq l_2 \leq l_3 \leq l_{\max}} \frac{(B_{l_1 l_2 l_3}^{(i)})^2}{\sigma_{l_1 l_2 l_3}^2} ,$$

cosmic variance
of bispectrum

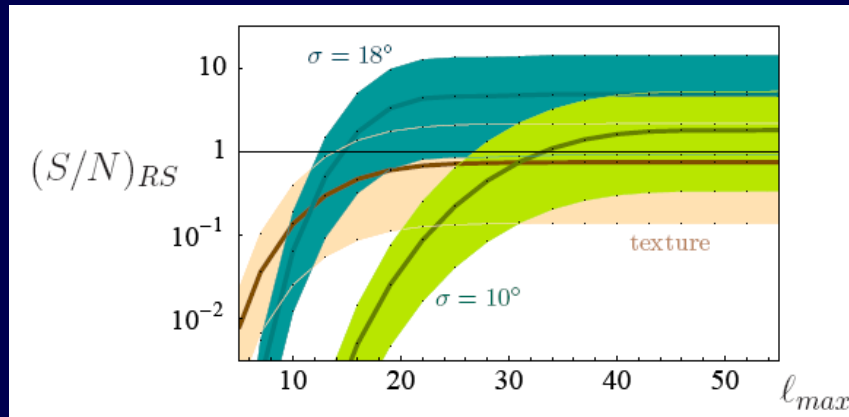


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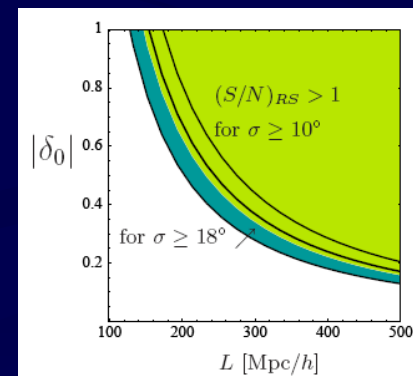
cosmic variance
of bispectrum

1. RS³



exceeds 1 for $l > 20-40$,
according to the void size

If no signal found, grt constraints
on Void parameter space

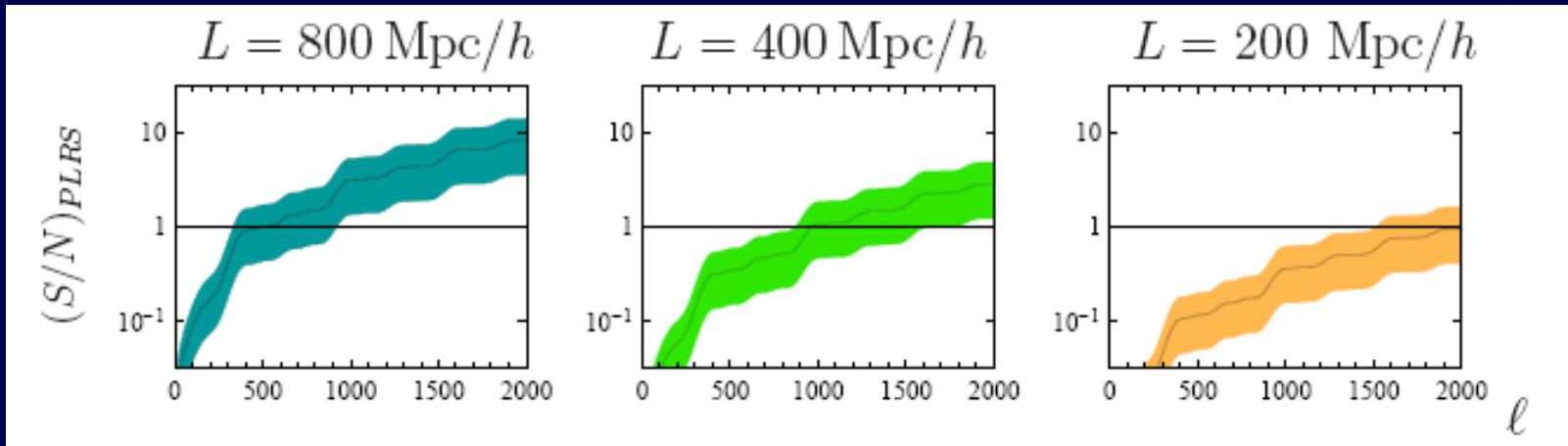


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cosmic variance
of bispectrum

2. PLRS



Detectable by Planck if Void radius $L > 400 \text{ Mpc}/h$

$L < 300 \text{ Mpc}/h$ means $z < 1$, which has been excluded by galaxy surveys

Signal at low multipoles from RS3 + high multipoles from P L RS
is a **UNIQUE SIGNATURE** of a Void

→ **confirm of reject** the Void explanation of the Cold Spot

Contamination of f_{NL}

The primordial non-gaussianity parameter f_{NL} is defined parametrizing the primordial perturbation as

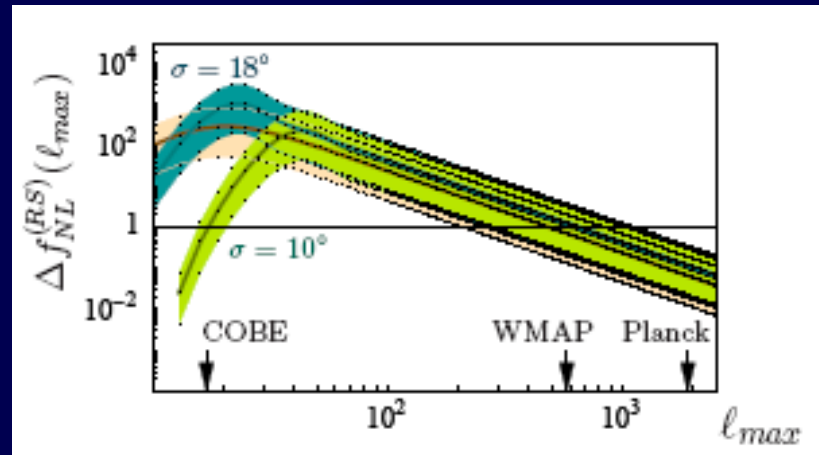
$$\phi(x) = \phi_L(x) + f_{NL}(\phi_L^2(x) - \langle \phi_L^2(x) \rangle)$$

linear gaussian part of the perturbation

Notice that inflationary model generically predict $f_{NL} = O(0.1)$; models exist with $O(1)$ f_{NL} .

The result is
(lensing negligible)

$$\Delta f_{NL}^{(RS)} \approx \begin{cases} 1 & \text{for WMAP} \\ 0.1 & \text{for Planck} \end{cases}$$



Non-diagonal 2-p function

$$\langle a a \rangle = \langle \cancel{P P} \rangle + \langle \cancel{P RS} \rangle + \langle P L \rangle + \underbrace{\langle RS RS \rangle + \langle L L \rangle}_{\text{negligible}}$$

vanishing
(for gaussian and/or isotropous fields)

vanishing
(P&RS uncorrelated)

PROMISING!

(depends on axis orientation)

Non-Diagonal S/N
should exceed 1 at $l > 1000$
(preliminary)

Conclusions and perspectives

- ✓ Correct interpretation of new data from Planck requires better understanding of CMB secondaries
- ✓ Cold Spot: Planck could confirm or discard the Void explanation by looking at bispectrum
- ✓ Non-diagonal 2-point function seems also promising
- ✓ Polarization effects also deserve study

[Vielva et al.,
arXiv:1002.4029]