MATHEMATICS OF THE UNIVERSE

Probing Variant Axion Models at LHC

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Chen, Frampton, FT, Yanagida, arXiv:1005.1185, to appear in JHEP.

<u>Strong CP problem:</u>

$$\mathcal{L}_{\theta} = \theta \frac{g_s^2}{32\pi^2} G^{\mu\nu} \tilde{G}_{\mu\nu}$$

Experimental bound (from neutron EDM) reads $|\theta| < 10^{-(9-10)} \equiv \theta^{(\exp)}$

Why is θ so small??

The Peccei-Quinn mechanism: Peccei and Quinn, `77

We introduce a global chiral $U(1)_{PQ}$ symmetry, which has a QCD anomaly.

$$\begin{split} \mathcal{L} &= \bar{q}_L q_R \Phi + \text{h.c.} \\ \mathbf{P}_{\mathbf{Q}} \text{ trans.:} \quad \begin{array}{l} q \to e^{i\alpha\gamma_5} q \\ \Phi \to e^{-2i\alpha} \Phi \end{array} \quad \begin{array}{l} \text{results in} \quad \Delta \mathcal{L} = 2\alpha \frac{g_s^2}{32\pi^2} G^{\mu\nu} \tilde{G}_{\mu\nu} \end{split}$$

If U(1)_{PQ} is spontaneously broken by $\langle\Phi
angle=f_a
eq 0$, there appears the associated NG boson, a, the axion.

$$\mathcal{L} = \left(\frac{a}{f_a} + \theta\right) \frac{g_s^2}{32\pi^2} G^{\mu\nu} \tilde{G}_{\mu\nu}$$

<u>Astrophysical/cosmological constraints</u>

invisible axion $f_a \gg v_{\rm EW} \quad \Longrightarrow \quad \text{Axion may play an important role in cosmology.}$ Star cooling/supernovae: $f_a \gtrsim 10^9 \, \text{GeV}$ Dark matter abundance: $f_a \lesssim 10^{12} \, \text{GeV}$ Moreover,

if U(1)_{PQ} is spontaneously broken during inflation:
 DM isocurvature fluc. w/ non-Gaussianity

if U(1)_{PQ} is restored during/after inflation:
 axionic strings/walls.
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This talk

If the $U(1)_{PQ}$ symmetry is restored during/ after inflation, the axionic strings and domain walls may be produced.



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Cosmic strings are formed at $T \approx f_a$



Allen, Shellard

Moreover, domain walls may be formed, if there are multiple degenerate vacua. Sikvie, 82





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Solutions to the domain wall problem

Add explicit breaking of U(1)_{PQ}
 Low inflation/reheating
 N_{DW} = 1

Dine-Fischler-Srednicki-Zhitnitsky (DFSZ) axion model has $N_{DW} = 3$.

(This is because there are three generations of quarks which carry PQ charges)

Kim-Shifman-Vainshtein-Zakharov axion model has $N_{DW} = 1$, if there is only one heavy quark with a PQ charge.

Variant axion models have $N_{DW} = 1$.

Variant Axion Models

[Peccei, Wu, Yanagida `86 Krauss, Wilczek `86]

Two Higgs doublets: $\Phi_1, \ \Phi_2$ (+ PQ singlet: σ)

In DFSZ model, Φ_1 is coupled to the down-type quarks and Φ_2 is to the up-type quarks. $\longrightarrow N_{\rm DW} = 3$

If we couple Φ_2 to only the t (or u or c) quark, we can avoid the domain wall problem!



<u>PQ charge assignment</u> [Model T]

	Φ_1	Φ_2	σ	t_R	others
PQ charge	0	-1	1	-1	0

$$V(\Phi_{1}, \Phi_{2}, \sigma) = \lambda_{1} \left(|\Phi_{1}|^{2} - \frac{v_{1}^{2}}{2} \right)^{2} + \lambda_{2} \left(|\Phi_{2}|^{2} - \frac{v_{2}^{2}}{2} \right)^{2} + \lambda \left(|\sigma|^{2} - \frac{v^{2}}{2} \right)^{2} + a |\Phi_{1}|^{2} |\sigma|^{2} + b |\Phi_{2}|^{2} |\sigma|^{2} + \left(m \Phi_{1}^{\dagger} \Phi_{2} \sigma + \text{h.c.} \right) + d |\Phi_{1}^{\dagger} \Phi_{2}|^{2} + e |\Phi_{1}|^{2} |\Phi_{2}|^{2}.$$

Yukawa interactions

E

$$-\mathcal{L}_{ ext{Yukawa}}$$

$$y_{ij}^{(d)} \bar{Q}_{Li} \Phi_1 d_{Rj} + y_i^{(t)} \bar{Q}_{Li} \tilde{\Phi}_2 t_R + y_i^{(u)} \bar{Q}_{Li} \tilde{\Phi}_1 u_R + y_i^{(c)} \bar{Q}_{Li} \tilde{\Phi}_1 c_R$$

Light and heavy Higgs, h and H:

$$\Phi_1 = \begin{pmatrix} \phi_1^+ \\ \frac{1}{\sqrt{2}}(v_1 + h_1 + ig_1) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} \phi_2^+ \\ \frac{1}{\sqrt{2}}(v_2 + h_2 + ig_2) \end{pmatrix}$$

 $\begin{pmatrix} H \\ h \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} h_1 \\ h_2 \end{pmatrix}, \quad \tan \beta = \frac{v_2}{v_1}$

1.Where is the light higgs?

 $\sin \alpha$ In extreme cases, h_1 or h_2

2. The couplings to gauge bosons?





<u>Couplings to gauge bosons:</u>

V = W or Z

➡ SM like

$\frac{\alpha \sim 0 \ \beta \sim \frac{\pi}{2}}{\text{Couplings to fermions: (Model T)}}$

HVV : $\cos(\beta - \alpha) g_{\rm SM}^{hVV}$,

 $hVV : \sin(\beta - \alpha) g_{\rm SM}^{hVV},$

$$\begin{aligned} hcc &: -\frac{\sin \alpha}{\cos \beta} g_{\rm SM}^{hcc}, \\ hbb &: -\frac{\sin \alpha}{\cos \beta} g_{\rm SM}^{hbb}, \\ htt &: -\frac{\cos \alpha}{\sin \beta} g_{\rm SM}^{htt}. \end{aligned}$$

Suppressed, if $|\sin \alpha| \ll \cot \beta$

⇒ SM like

The SM Higgs boson decay branching ratios



If h→bb (and h→gg) is suppressed, h→2 γ can be enhanced.







W⁺W⁻, ZZ





Production processes

Model T: gluon-gluon fusion (GGF)



+ the sub-dominant processes below.

Models U and C: Vector boson fusion (VBF)



Associated production (VH)



<u>Prospect for discovery of the light Higgs</u> <u>through the two-photon decay</u>

SM Higgs search at LHC:

Will be discovered through $h \rightarrow \gamma \gamma$ at ATLAS/ CMS with 30fb⁻¹ and 14TeV c.m. energy for M_h < 130GeV in the inclusive search.

Settimated significance for $h \rightarrow \gamma \gamma$ in VBF and VH is 2.2 sigma with 30fb⁻¹ at CMS.

In the case of VAM:

Model T
 ✓Enhanced by 4(3) at M_h = 120GeV for sin α = 0(-0.05)
 ✓Will be discovered with 3fb⁻¹(4fb⁻¹).

✓ Production c.s. will be reduced by a factor of (3-4) with 7TeV compared to 14TeV. The significance will be about 2 sigma with 1 fb⁻¹.

- Models U and C

 \checkmark Will be discovered only with 3fb⁻¹ (10fb⁻¹).

Conclusions

The PQ mechanism predicts the presence of a light particle, axion.

If the PQ symmetry is restored during/after inflation, domain walls may be produced. One way-out is the variant axion model.

The special Yukawa structure of the VAM may lead to the enhancement of $h \rightarrow 2\gamma$.



Back-up slides





