$V{+jets \ with \ Sherpa} \\ {\rm Validating \ ME+PS \ merging \ against \ Tevatron \ data} \\$

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MC4LHC readiness CERN 01/04/10



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- The Sherpa approach (\geq v1.2.0)
 - tree-level matrix elements and truncated showers
 - inherent systematics
- Validation/Application for DY & prompt-photons at Tevatron

Simulating Hadron–Hadron Collisions

The Monte Carlo approach

• Hard interaction

exact matrix elements $|\mathcal{M}|^2$

• QCD bremsstrahlung

parton showers in the initial and final state

• Multiple Interactions

beyond factorisation: modelling

• Hadronisation

non perturbative QCD: modelling

• Hadron Decays

phase space or effective theories



\Rightarrow fully exclusive hadronic final states \Rightarrow direct comparison with experimental data

Herwig, Pythia, Sherpa

modulo detector simulation

Catani-Seymour dipole factorisation based Shower

Catani-Seymour local subtraction term

$$\int_{m+1} \mathrm{d}\sigma^{\mathrm{A}} = \sum_{\mathrm{dipoles}} \int_{m} \mathrm{d}\sigma^{\mathrm{B}} \otimes \int_{1} \mathrm{d}V_{\mathrm{dipole}}$$

$$\sup_{\mathrm{spin-\& color correlation}} \longleftrightarrow \qquad \longleftrightarrow \qquad \text{universal dipole term}$$

[Catani, Seymour Nucl. Phys. B 485 (1997) 291]

[Catani, Dittmaier, Seymour, Trocsanyi Nucl. Phys. B 627 (2002) 189]

Ansatz: complete factorisation through

- projection onto leading term in $1/N_c$
- spin averaged dipole terms $V_{
 m dipole}
 ightarrow \langle V_{
 m dipole}
 angle$

Shower Algorithm [Krauss, S. JHEP 0803 (2008) 038]

- color connected emitter-spectator 'dipoles'
- subsequent branchings of type II, IF, FI, FF
- exact momentum mappings invertable
- emissions ordered in k_{\perp}^2



Matrix Elements and Truncated Showers: ME⊕TS

Describe few hardest emissions through ME

- emission phase space sliced IR sensible measure Q
 - Soft/collinear emissions from shower $Q < Q_{
 m cut}$
 - Hard emissions from matrix element $Q > Q_{
 m cut}$
- → Sudakov form factor factorises

$$\begin{split} & \Delta_{\mathfrak{a}}(\mu^{2},k_{\perp}^{2}) = \Delta_{\mathfrak{a}}^{\mathrm{PS}}(\mu^{2},k_{\perp}^{2}) \, \Delta_{\mathfrak{a}}^{\mathrm{ME}}(\mu^{2},k_{\perp}^{2}) \\ & \nleftrightarrow \quad \mathcal{K}_{\mathfrak{b}\mathfrak{a}}^{\mathrm{ME}}(z,k_{\perp}^{2}) \to \frac{1}{\mathrm{d}\hat{\sigma}_{\mathfrak{a}}^{(N)}(\Phi_{N})} \, \frac{\mathrm{d}\hat{\sigma}_{\mathfrak{b}}^{(N+1)}(z,k_{\perp}^{2};\Phi_{N})}{\mathrm{d}\log(k_{\perp}^{2}/\mu^{2})\,\mathrm{d}z} \end{split}$$

A new merging algorithm [Höche, Krauss, S., Siegert JHEP 0905 (2009) 053]

- \rightarrow pseudo shower history for MEs cluster algorithm inverse to the shower
- → Truncated Shower

PS starts at $2\to 2$ core and can radiate partons off intermediate lines i.e. "between" ME $_{\text{partons}}$

ME branchings must be respected evolution-, splitting- & angular variables $\{k_{\perp}^2,z,\phi\}$ preserved



Matrix Elements and Truncated Showers: ME⊕TS

The proposed measure: final-state splitting $(ij) \rightarrow ij$

$$Q_{ij}^2 = 2 p_i p_j \min_{k \neq i, j} \frac{2}{C_{i,j}^k + C_{j,i}^k}; \qquad C_{i,j}^k = \begin{cases} \frac{p_i p_k}{(p_i + p_k)p_j} - \frac{m_i^2}{2 p_i p_j} & \text{if } j = g \\ 1 & \text{else} \end{cases}$$

 \hookrightarrow minimize over color partners k

The proposed measure: initial-state splitting $a \rightarrow (aj)j$

$$Q_{aj}^{2} = 2 p_{a} p_{j} \min_{k \neq a, j} \frac{2}{C_{(aj), j}^{k} + C_{j, (aj)}^{k}}; \qquad C_{(aj), j}^{k} = \begin{cases} \frac{p_{aj} p_{k}}{(p_{aj} + p_{k}) p_{j}} & \text{if } j = g \\ 1 & \text{else} \end{cases}$$

- measure correctly identifies IR enhanced phase-space regions
- regulator for matrix elements & veto measure for the shower
- no additional free parameter is introduced as e.g. D for the Run II k_T algorithm

Systematics

Sources of Systematics

merging related we try hard to minimize that

- $\bullet\,$ value of phase-space separation cut, ${\it Q}_{\rm cut}\,$ [cancels to the log-accuracy of the shower]
- $\bullet\,$ maximum number of jets from hard ME's, $\textit{N}_{\rm max}$
- choice of internal separation measure

pQCD related dynamical scale choice inherent

- scale uncertainties from ME's $_{[e.g.\ \mu_R\ of\ the\ 2\ \rightarrow\ 2\ core]}$
- scale uncertainties from PS's [e.g. rescaled factorisation/start scale]

pQCD/npQCD transition

- parton-shower IR cut-off / intrinsic transverse momentum [tuned at LEP & low-pT DY]
- PDF + α_S from the fit [enter in ME and PS]

npQCD related need to be tuned

- hadronisation parameters [Professor tune to LEP data]
- Underlying Event parameters [so far hand-tuned only]

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 Q_{cut} and/or N_{max} variation should affect σ_{tot} only beyond (N)LL

Example: DY-pair production $\sigma_{\rm tot}$ @ Tevatron

		Nmax						
		0	1	2	3	4	5	6
$Q_{ m cut}$	20 GeV	192.6(1)	191.0(3)	190.5(4)	189.0(5)	189.4(7)	188.2(8)	189.9(10)
	30 GeV		192.3(2)	192.7(2)	192.6(3)	192.9(3)	192.7(3)	193.2(3)
	45 GeV		193.6(1)	194.4(1)	194.3(1)	194.4(1)	194.6(2)	194.4(1)

 \rightsquigarrow "merging systematics" of $\sigma_{\rm tot} < \pm 3\%$

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Z^0 +jets at Tevatron: jet multiplicities

Jet rates and -spectra improved compared to pure PS simulation due to exact real emission ME's

 $\label{eq:constraint} \begin{array}{l} \mbox{Example: DY-pair production $\sigma_{e^+e^-+N_{jet}}$ CDF Data: PRL 100 (2008) 102001 $$ [Sherpa normalized to $\sigma_{e^+e^-+1_{jet}}$]$ \end{array}$



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Variation of Q_{cut} should affect distributions only beyond (N)LL But Q_{cut} must be in range where PS approximation is valid!

Example: All-jets p_T 's in DY-pair production CDF Data: PRL 100 (2008) 102001 [Sherpa normalized to $\sigma_{e^+e^-+1_{jet}}$]



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V+jets with Sherpa

Variation of Q_{cut} should affect distributions only beyond (N)LL But Q_{cut} must be in range where PS approximation is valid!

Example: Differential \mathbf{k}_{T} jet rates



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Example: Differential \mathbf{k}_{T} jet rates

⇒ 'old' Sherpa CKKW for W+jets [Alwall et. al Eur. Phys. J. C 53 (2008) 473]



Z^0 +jets at Tevatron: Z/γ^* transverse momentum

Comparison to Sherpa's CKKW implementation in v1.1.3

Example: DY- p_T in Z/γ^* +jet events DØ Data: Phys. Lett. B 669 (2008) 278



Sherpa v1.1.3

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Sherpa v1.1.3

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Prompt-Photons

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The perturbative QCD approach

Direct production



- fixed-order calculations
- $\gamma+jet$ @ NLO (JetPhox) [Catani et. al]
- $\gamma\gamma$ @ NLO (DiPhox) [Binoth et. al]
- $\gamma\gamma$ +jet @ NLO [Del Duca et. al]
- $gg
 ightarrow \gamma \gamma g$ [de Florian et. al]

Pre-requisits

• IR safe cross-section definition based on photon-isolation criterium

[cone, smooth isolation, democratic approach]

- input for the non-perturbative component of $D_{q,g}^{\gamma}(z,Q^2)$
- non-prompt component, e.g. $\pi^0 \to \gamma \gamma$, $\eta \to \gamma \gamma$, experimentally separable

Fragmentation component



- QED γq collinear singularity
- resummation to all orders α_s
- fragmentation function D^γ_{q,g}
 [Llewellyn Smith Phys. Lett. B 79 (1978) 83]

A new Monte Carlo model [Höche, S., Siegert PRD 81 (2010) 034026]

Democratic Monte Carlo model

- treat photons and QCD partons fully democratically [Glover, Morgan Z. Phys. C 62 (1994) 311]
- combine matrix elements of different parton/photon multiplicity with
- interleaved QCD \oplus QED evolution and hadronization \rightsquigarrow models $D_{q,g}^{\gamma}(z, Q^2)$

Generalized merging formalism

• emission probabilities factorise trivially

$$\Delta_{a}(\mu^{2},k_{\perp}^{2}) = \Delta_{a}^{\text{QCD}}(\mu^{2},k_{\perp}^{2})\Delta_{a}^{\text{QED}}(\mu^{2},k_{\perp}^{2})$$

ullet implemented by adding splitting functions ${\it q} \rightarrow {\it q} \gamma$

$${\cal K}^{
m QED}_{(ij)i}(z,{
m k}^2_\perp) \sim \; {lpha({
m k}^2_\perp)\over 2\pi} \; \sum_k \langle {
m V}^{
m QED}_{(ij)i,k}({
m k}^2_\perp,z)
angle$$

- different than large-N_C QCD: spectators all particles with opposite charge
- neglect (negative) interference with same-sign charges [Dittmaier Nucl. Phys. B 565 (2000) 69]

Measuring the photon fragmentation function in $e^+e^- \rightarrow$ Hadrons

- validation of the shower/hadronization component
- perform jet finding including final-state photons
- study photon-energy fraction wrt its containing jet: $z_{\gamma} \equiv E_{\gamma}/E_{\rm jet}$
- measured by ALEPH for 2-jet, 3-jet, \geq 4-jet events at varying y_{cut}^{Durham}

[Z. Phys. C 69 (1996) 365]





Isolated prompt-photon production at Tevatron

• CDF data [Phys. Rev. D 80 (2009) 111106]

cuts: 30 $< E_T^{\gamma} <$ 400 GeV, $|\eta^{\gamma}| <$ 1, ~ isolation: $E_T^{R=0.4} - E_T^{\gamma} <$ 2 GeV

• Sherpa: $jj \rightarrow jj/\gamma j/\gamma \gamma$ + shower vs. ME \oplus TS with $N_{max} = 2$



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Di-Photon production at Tevatron

DØ data [arXiv:1002.4917]

cuts: $E_T^{\gamma_1} > 21$ GeV, $E_T^{\gamma_2} > 20$ GeV, $|\eta^{\gamma_{1,2}}| < 0.9$ isolation: $E_T^{R=0.4} - E_T^{\gamma} < 2.5$ GeV ResBos, DiPhox, Pythia

• Sherpa: merged 2 \rightarrow {2, 3, 4} plus $gg \rightarrow \gamma\gamma$ box



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Summary/Outlook

Summary

- Multijet ME-PS merging sustainable approach to describe multijet events
- Virtues of two complementary approaches combined
 - hard emissions through exact tree-level matrix elements
 - (intra) jet evolution through truncated QCD parton showers
- New formalism available within Sherpa-1.2 [Höche, Krauss, S., Siegert JHEP 0905 (2009) 053]
 - MEs from Comix [Gleisberg, Höche JHEP 0812 (2008) 039]
 - dipole subtraction based shower [S., Krauss JHEP 0803 (2008) 038]
 - interleaved QCD ⊕QED evolution [Höche, S., Siegert PRD 81 (2010) 034026]
 - Deep Inelastic Scattering [Carli, Gehrmann, Höche arXiv:0912:3715]

Validation

- improved/reduced systematics wrt 'old' CKKW implementation
- good agreement with Tevatron data & pQCD NLO calculations
- Tevatron data extremely useful to develop/validate our MC models
- Rivet proves to be the perfect tool for direct comparison.