

# Electroweak stars

*Dejan Stojkovic*

SUNY at Buffalo

$$SU(2)_L \times U(1)_Y$$

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## Based on:

*Electroweak stars: how nature may capitalize on the standard model's ultimate fuel.*

D. Dai<sup>1</sup>, A. Lue<sup>2</sup>, G. Starkman<sup>3</sup>, D. Stojkovic<sup>1</sup>

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<sup>1</sup> SUNY at Buffalo

<sup>2</sup> MIT

<sup>3</sup> CWRU



# Motivation

- **Stellar evolution** is one of the most interesting questions in modern astrophysics
- Despite tremendous progress in this field, it looks like it is far from being closed

**Conventional evolution stops with neutron stars**

New stages proposed: quark stars, preon stars, dark stars...

**Electroweak stars!**

# Outline

## • ~~Brief overview~~

- ~~Stellar evolution: protostar – neutron star~~
- ~~Quark star~~

## • Question:

*Can Electroweak Stars powered by baryon number violating processes exist?*

- Can a star's core be compressed to EW densities without being within its own  $R_S$ ?
- Can this new phase last long enough to be called a star?

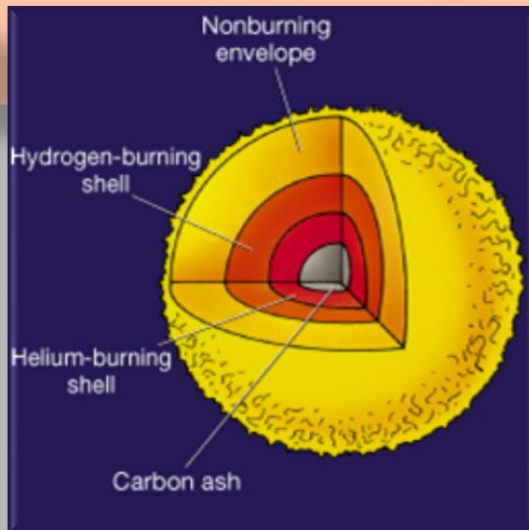
Perhaps YES!



# The hierarchy in the central core density

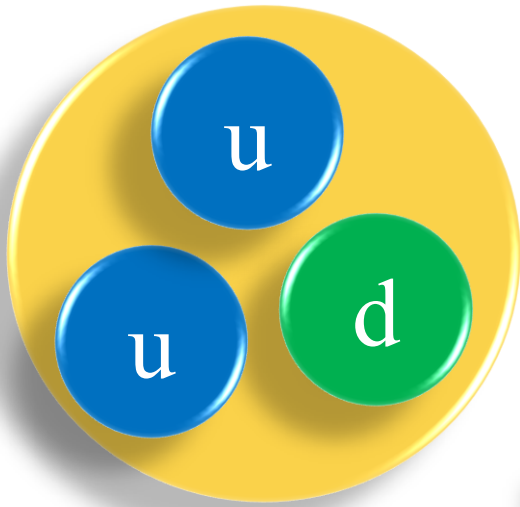
- Protostar
- Main Sequence Star
- White Dwarf
- Neutron Star
- Quark Star ?
- ???

- Central core density grows

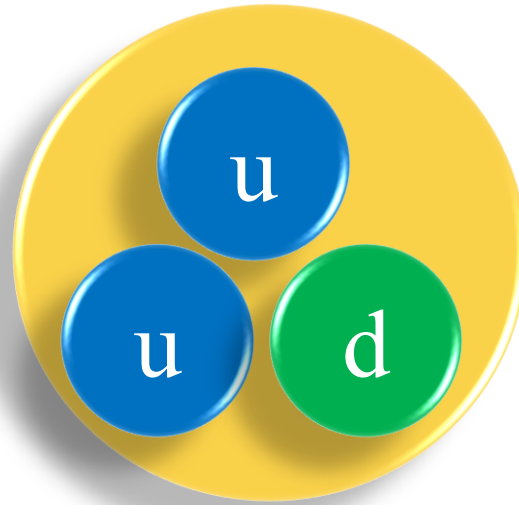


# Quark Star

- Fermi pressure can support only neutron stars lighter than about  $2.1 M_{\text{sun}}$  → **Tolman-Oppenheimer-Volkoff** limit
- Cores of stars heavier than this have densities comparable with QCD phase transition densities  $\sim (100 \text{ MeV})^4$
- This may lead to a **quark star**:  
state where one can not distinguish between the nucleons



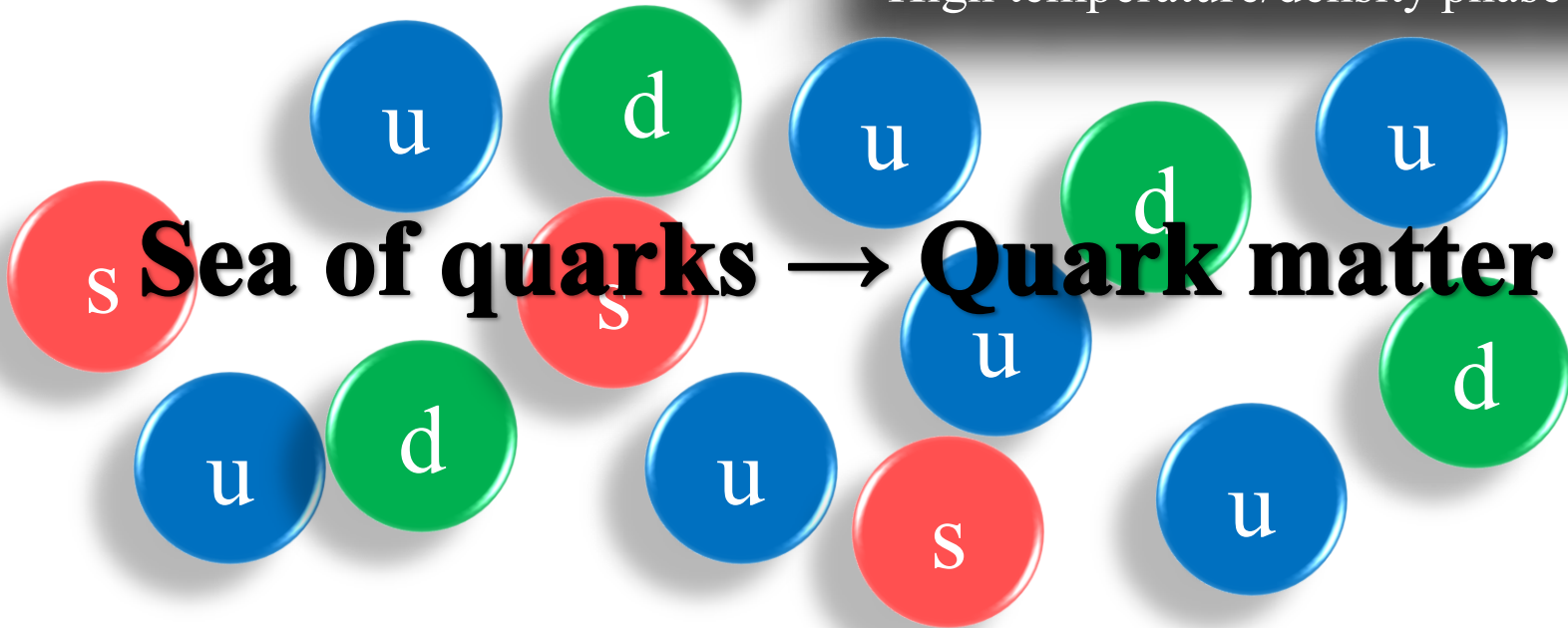
**Neutron**



**Neutron**



High temperature/density phase transition



**Sea of quarks → Quark matter**



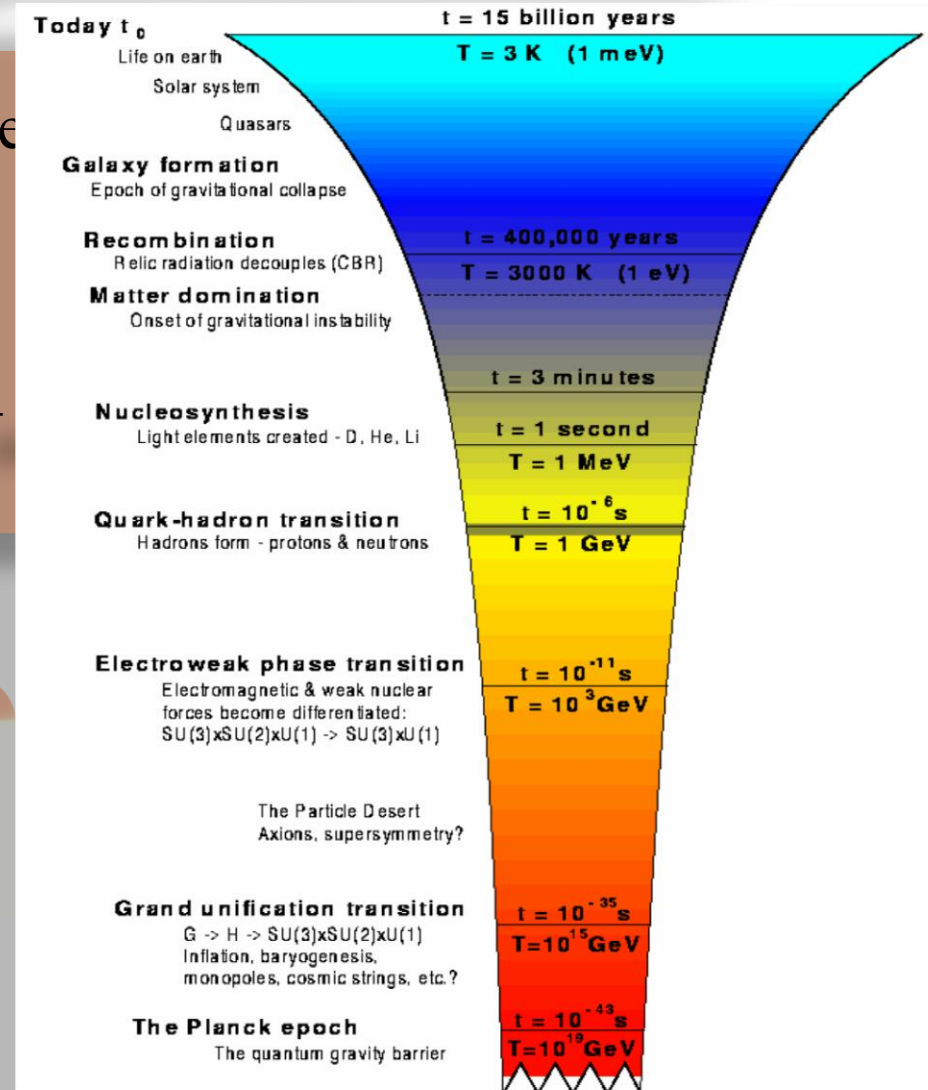
*“Og discovered fire, and Thorak invented the wheel. There's nothing left for us.”*



# What happens next?

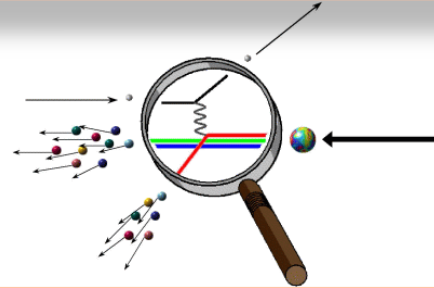
- Gravitational collapse continues
- Inverse Big Bang
- Densities reach  $\sim (100 \text{ GeV})^4$

Electroweak phase transition happens!



# Standard Model's ultimate fuel

- Why is that important for the evolution of the star?
- We can take advantage of the Standard Model's ultimate fuel



Non-perturbative baryon number violating electroweak processes

**Conversion of quarks to leptons**

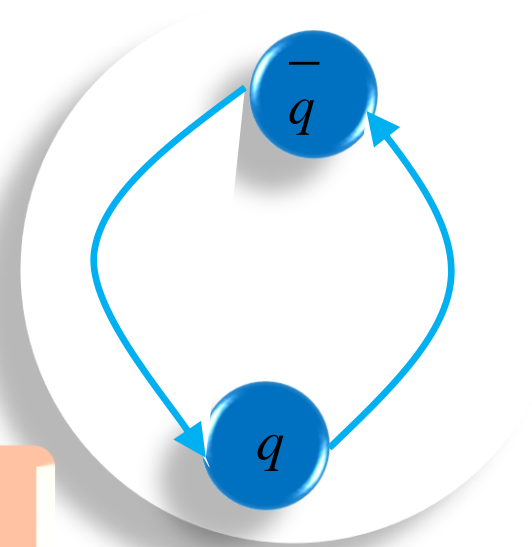
# Baryon Number

- Baryon number is a global symmetry of the SM Lagrangian

$$\partial_\mu J_B^\mu = \sum_i \partial_\mu \bar{q}_i \gamma^\mu q_i \neq 0$$

- Baryonic current is perturbatively conserved

- Quarks can appear and disappear only in pairs



# Baryonic current anomaly

Quantum corrections destroy conservation of **baryon number**

$$\partial_{\mu} J_B^{\mu} = \frac{g^2 \text{const}}{16\pi^2} F_{\mu\nu}^a \tilde{F}^{a\mu\nu}$$

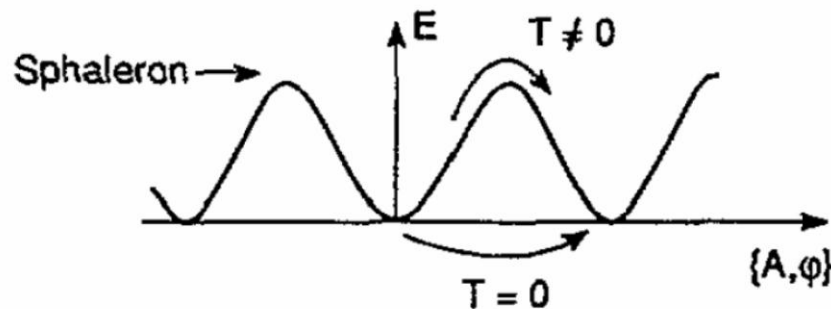
$F_{\mu\nu}^a \rightarrow$  gauge field strength

Known as Adler-Bell-Jackiw anomaly

Baryon number violating processes are inherently quantum (instanton) processes

# Sphaleron

At zero temperature such processes can occur only by quantum tunneling and are exponentially suppressed by  $e^{-8\pi\alpha} \sim e^{-3000}$



At finite temperature one may classically pass over the barrier with the Boltzmann factor  $e^{-E/T}$

**Sphaleron** - an unstable solution to the equations of motion

# Sphaleron rate

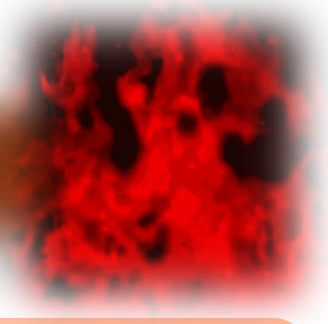
At finite temperature, the sphaleron rate is  $\frac{\Gamma_{sph}}{V} \approx \frac{1}{4} \alpha_w T^4 e^{-E_{sph}/T}$

$$E_{sph} \approx \frac{M_W}{\alpha_w} \approx 10 \text{ TeV} \quad \text{for } T < T_c$$

$$E_{sph} \approx T \quad \text{for } T > T_c$$

So, above EW symmetry-breaking scale ( $T_c \approx 100 \text{ GeV}$ ), baryon number violating processes are essentially unsuppressed

# Electroweak Burning



Quarks can then be effectively converted into leptons. In this **electroweak burning** huge amounts of energy can be released.

B-L preserving interaction can convert 9 quarks into 3 anti-leptons

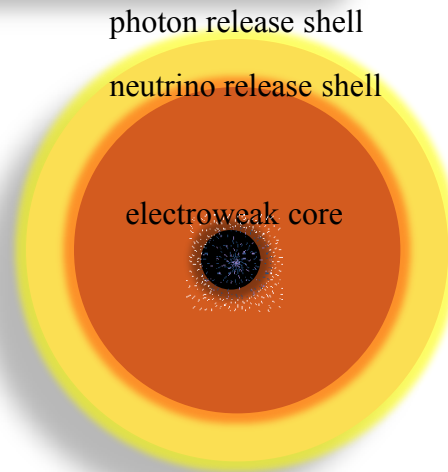
$$\begin{array}{l} udd \\ css \\ tbb \end{array} \rightarrow \bar{\nu}_e, \bar{\nu}_\mu, \bar{\nu}_\tau$$

At these temperatures each particle carries about 100 GeV of energy, so this process can release about 300 GeV per neutrino

# The structure of the star

## *Electroweak core*

- T above EW-breaking scale:  $T > 100 \text{ GeV}$



## *Photon release shell*

- Effective radius of the star

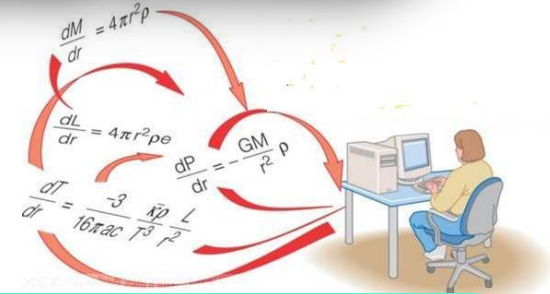
## *Neutrino release shell*

- Inside the shell neutrinos are trapped
- Outside they can freely stream



# Finding the solution

## Tolman-Oppenheimer-Volkoff equations



$$\frac{dP}{dr} = - \frac{(\varepsilon + P)(M + 4\pi pr^3)}{r^2 (1 - 2M/r)}$$

$$\frac{dM}{dr} = 4\pi\varepsilon r^2$$

$$g_{tt} = \left( 1 - \frac{2M_{star}}{R_{surface}} \right) \exp \left( - 2 \int_0^{P(r)} \frac{dP}{P + \varepsilon} \right)$$
$$g_{rr} = \frac{1}{1 - \frac{2M(r)}{r}}$$

$\varepsilon$  = total energy density

$P$  = total pressure

Metric coefficients  $g_{tt}$  and  $g_{rr}$  very important!

$P(r) = ?$

$\varepsilon(r) = ?$

$M(r) = ?$

# Variables and parameters

The pressure, energy density, and number density of particles can be well approximated from an **ideal gas distribution**

$$p_i = \frac{g_i}{6\pi^2} \int_{m_i}^{\infty} dE \sqrt{E^2 - m_i^2}^3 f_i(E)$$
$$n_i = \frac{g_i}{2\pi^2} \int_{m_i}^{\infty} dE \sqrt{E^2 - m_i^2}^2 E f_i(E)$$
$$\varepsilon_i = \frac{g_i}{2\pi^2} \int_{m_i}^{\infty} dE \sqrt{E^2 - m_i^2}^2 E^2 f_i(E)$$

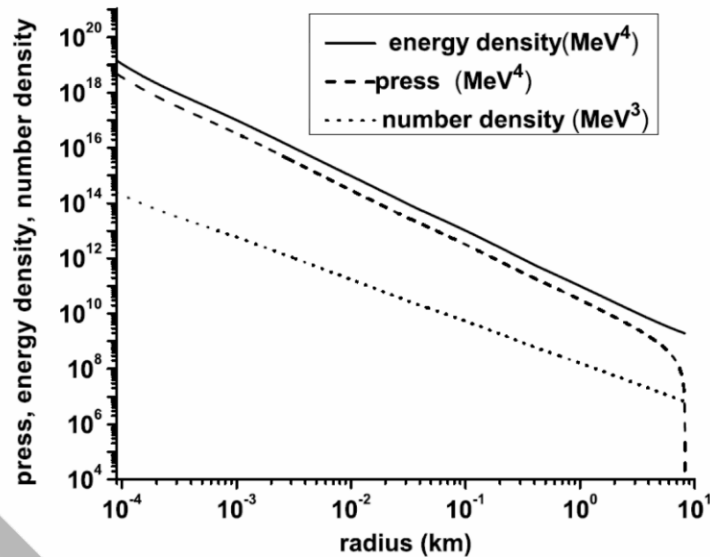
$$f_i(E) = \frac{1}{1 \pm e^{(E - \mu_i)/T}}$$

$$P = \sum_i p_i - B$$
$$\varepsilon = \sum_i \varepsilon_i + B$$



B = “bag” energy = 145 MeV  
from the “bag” model of nucleons

# Star Parameters



Radius of the star is where P and  $\epsilon$  drop to zero



$$R_{\text{star}} = 8.2 \text{ km}$$
$$M_{\text{star}} = 1.3 M_{\text{sun}}$$

The pressure, energy density and particle number density dependence of the radius of the star

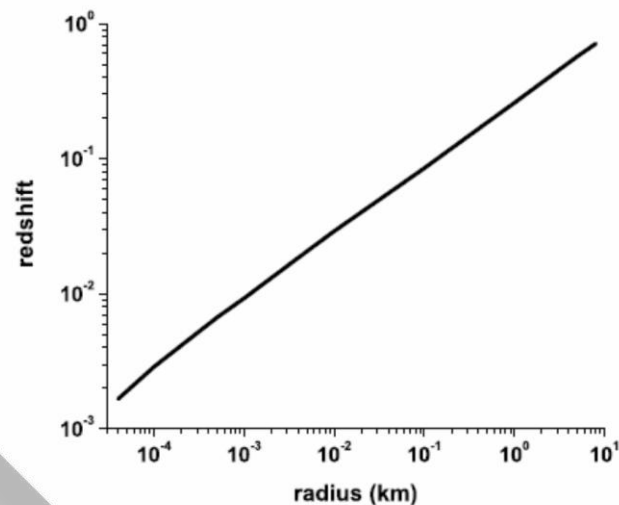
- Core ends where  $\epsilon$  drops below  $(100\text{GeV})^4$
- Several cm in size

The solution is non-singular at the center – not a black hole

# Neutrino energy redshift

- Unlike ordinary stars, particles propagating through the electroweak stars suffer large **gravitational redshift**
- Energy changes as:

$$E_{\nu}(r) = \frac{\sqrt{g_{tt}(r_0)}}{\sqrt{g_{tt}(r)}} E_{\nu}(r_0)$$



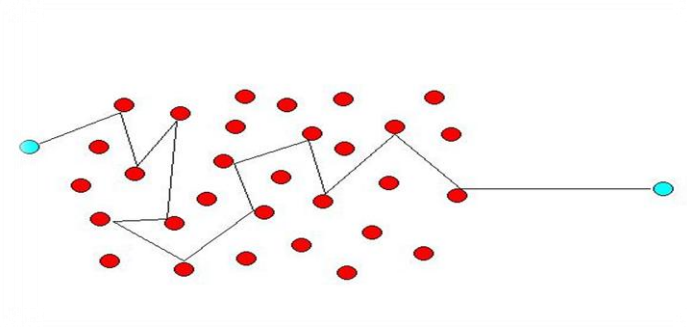
The redshift factor  $\sqrt{g_{tt}(r)}$  inside the star. A particle with the original energy of 100GeV near the center carries away only 100MeV as it leaves the surface

# Neutrino mean free path

The neutrino mean free path inside the star



$$\frac{1}{\lambda} = \sum_i \sigma_i n_i$$



with

$$\sigma_i \approx \frac{G_F^2 E_\nu E_i}{\pi}$$

Near the core, the mean free path is  $\lambda \sim 10^{-14} \text{m}$

Neutrinos interact many times before they leave the star

# Estimated luminosity

Neutrino luminosity



$$L_{\nu} \approx \sigma T_{core}^4 4\pi r_{core}^2$$

$$\approx 10^{41} \text{ MeV}^2 \approx 10^{53} \text{ erg /sec}$$

At this rate it would take less than a second to release  $M_{\text{sun}}$

***However, this is a severe over-estimate!***

**Not taken into account:**

- GR effects
- Luminosity depends not just on the  $T$  and  $\epsilon$  but also on their gradients: *the net outward flux of energy*

# Maximal energy release rate

Free fall time of the quark shell into the EW-burning core

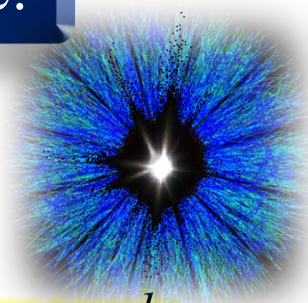


upper bound on  $dE/dt$

$$\left(\frac{dE}{dt}\right)_{\max} = 4\pi r_{ew}^2 \sqrt{\frac{2\lambda M(r_{ew})}{r_{ew}^2}} \varepsilon_{\nu}(r_{ew}) g_{tt}(r_{ew}) \approx 10^{27} \text{ MeV}^2$$

Compare this with EW baryogenesis/baryodestruction rate:

$$\left(\frac{dE}{dt}\right)_{EWbd} = 0.1 \cdot 4\pi r_{ew}^2 \alpha_{ew}^4 T^4 \approx 10^{34} \text{ MeV}^2$$



As quarks reach the EW core, they are converted into neutrinos instantaneously. Otherwise, the infalling matter would pile up and form a black hole.

# Modeling the energy transport through the star

Relativistic transport of energy can be described by

$$\frac{dE_{\infty}}{dt} = -4\pi r^2 \lambda \frac{d[S(r) g_{tt}(r)]}{\sqrt{-g_{rr}(r)} dr}$$

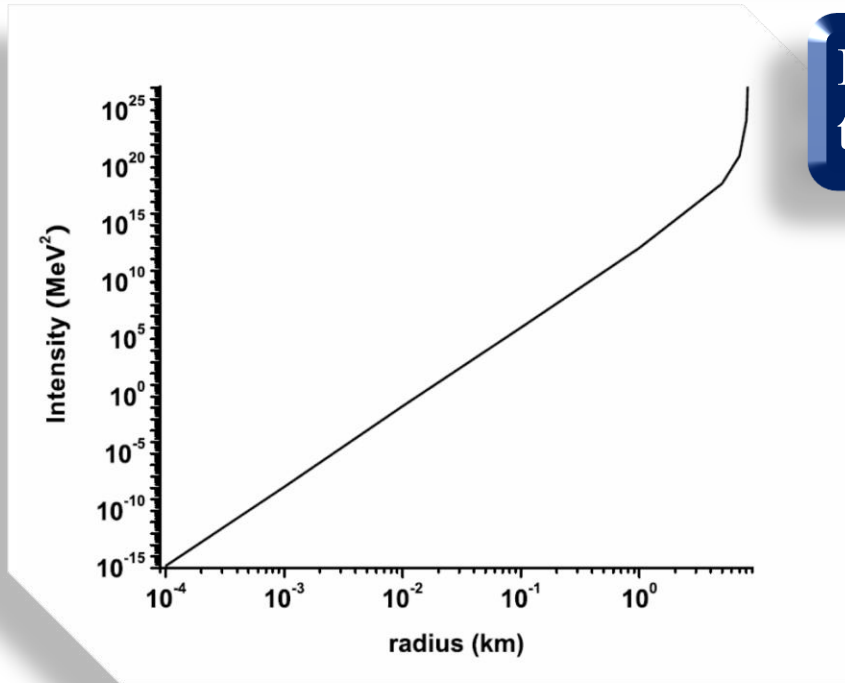
with

$$S(r) = \frac{dE(r)}{dt dA}$$

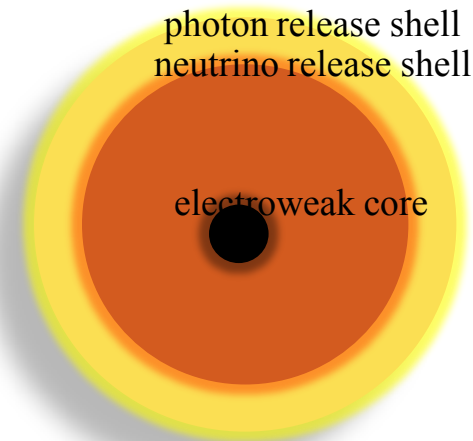
- Metric coeff.  $g_{tt}(r)$  describes both the redshift and time delay
- Energy flux  $S(r)$  can be modeled by the energy density  $\varepsilon$  knowing the energy (i.e. chemical potential) of neutrinos.



# The energy release rate vs. neutrino release radius:



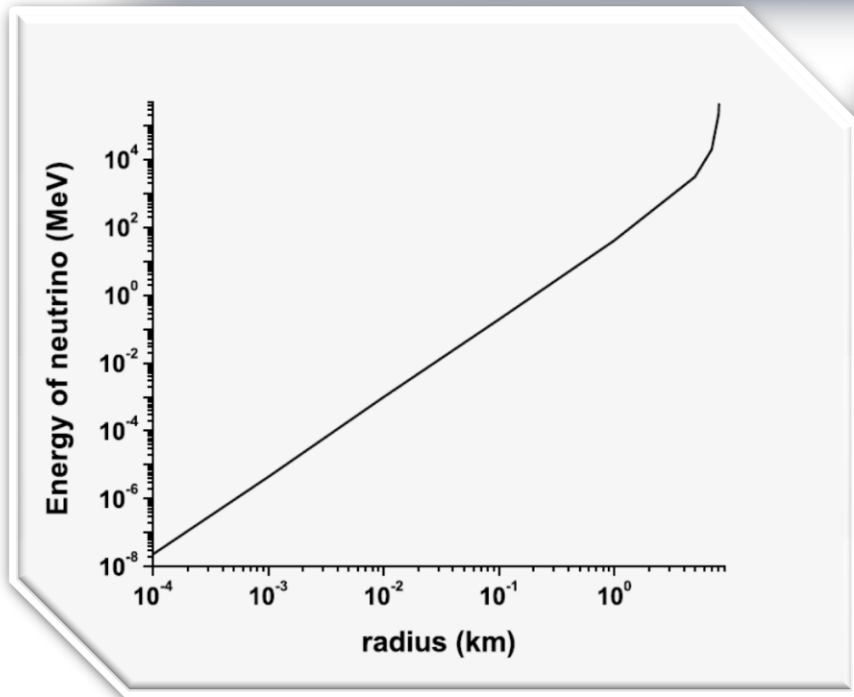
Energy release rate increases with the radius of the neutrino release shell



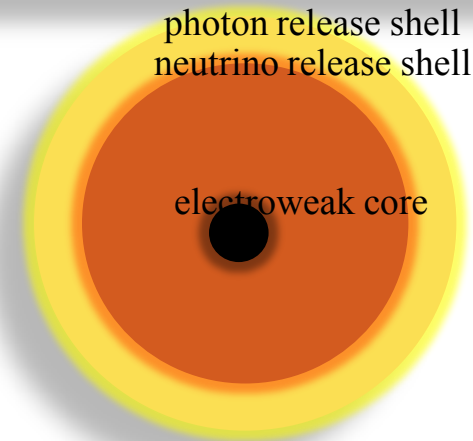
The maximum energy release rate is on the surface, but there it exceeds the limit from the quark shell free fall into the electroweak core  $\approx 10^{27} \text{ MeV}^2$

This implies that the neutrino release shell must be inside the star

# Neutrino energy at the surface of the core vs. neutrino release radius:



Higher energy release rate needs the source of higher energy. Energy of neutrinos therefore increases with the radius of release shell.



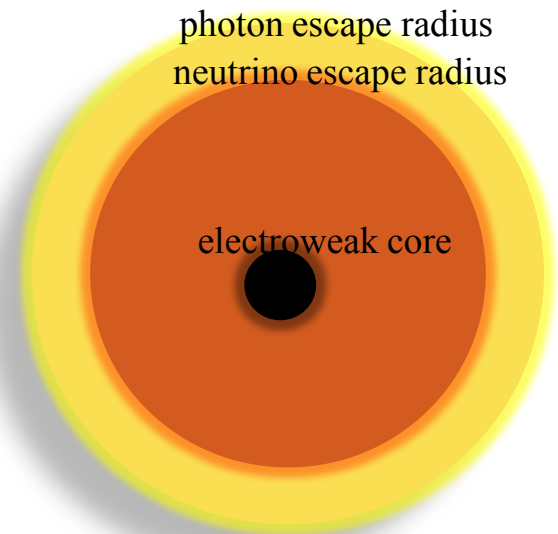
However, the energy is already larger than  $300\text{GeV}$  at  $8.1 \text{ km}$

*The star cannot support this amount of energy implying that the radius of the neutrino-sphere must be smaller  $8.1 \text{ km}$*

# The life-time of the electroweak star

If the neutrino escape radius is **8.1 km**,  
the energy release rate is  **$10^{24} \text{ MeV}^2$**

This implies that it takes about  
**10 million years** to release  **$1 M_{\text{Sun}}$**

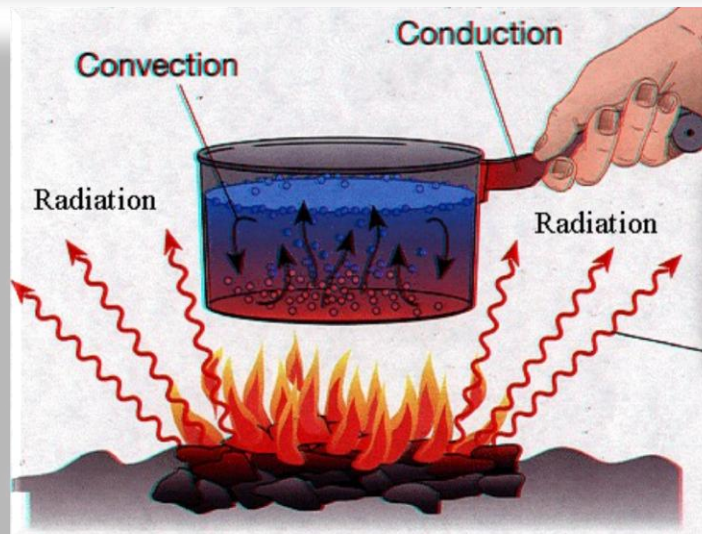


**This is the minimal life-time of the electroweak star**  
*(provided that all the available quark fuel burns)*

# What did we neglect?

We neglected the fact that some fraction of energy is carried away by photons, since they have shorter mean free path

We also ignored the effects of energy transport due to convection



Order of magnitude estimate will not change

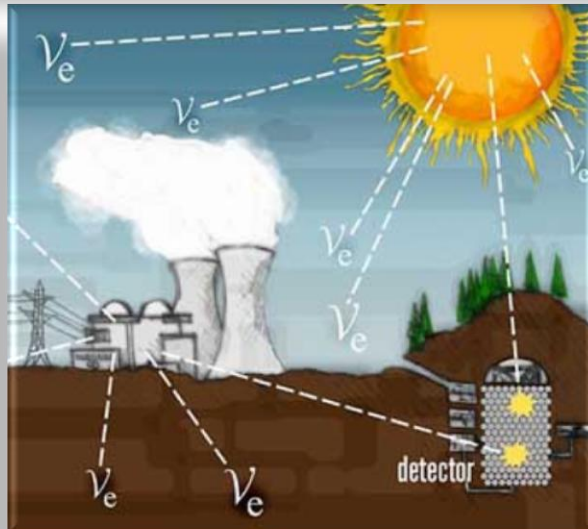
# Observational signature?

Most of the energy is emitted in form of neutrinos

When quarks are transformed into leptons, we get mostly *antineutrinos*

Most of the astrophysical object emit (**electron**) neutrinos

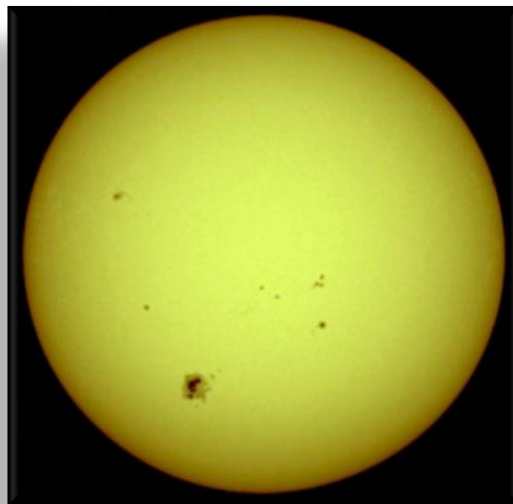
Antineutrino emitters (**of all families**) might be electroweak stars!



# Observational signature?

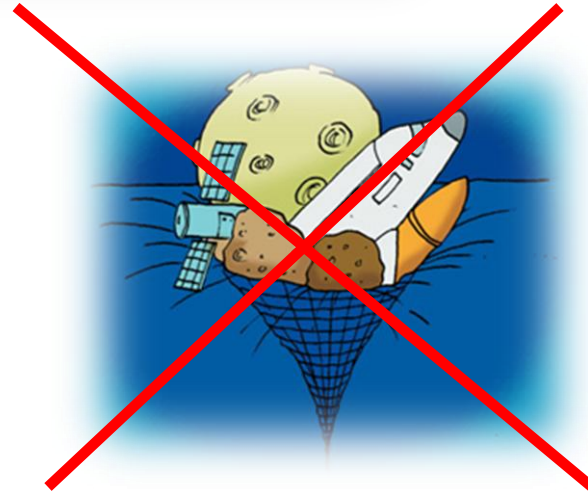
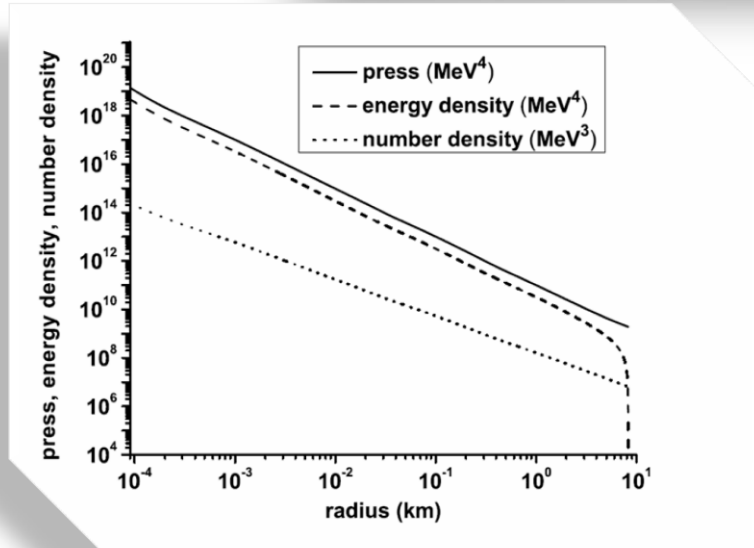
Energy released in photons very important for the signature

- Most of the energy emitted in neutrinos
- But significant percentage might be photons



Important to calculate  
photo-luminosity!

# Potential problems



- We found a non-singular solution to TOVE → not a black hole
  - However, TOVE assume quasi-equilibrium
  - We have not solved a full time-dependent evolution of the system
- Will EW density be reached before the object crosses its own Schwarzschild radius?

# Full time-dependent analysis



The core of the star:  
Earth mass in a region of the size of an apple

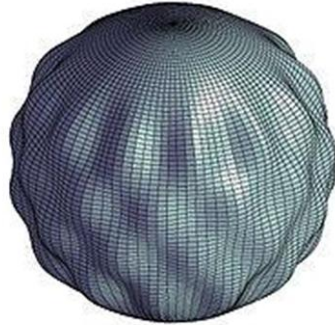
It is almost a black hole!

In this regime most of the approximations  
(e.g. homologous collapse, PPN) fail

Must use full GR with back reaction included  
Very challenging task!



# Full stability analysis



The electroweak star can't be in a static equilibrium like a neutron star or a black hole

Gravity is balanced by the gas pressure and also radiation pressure

Dynamic rather than static equilibrium

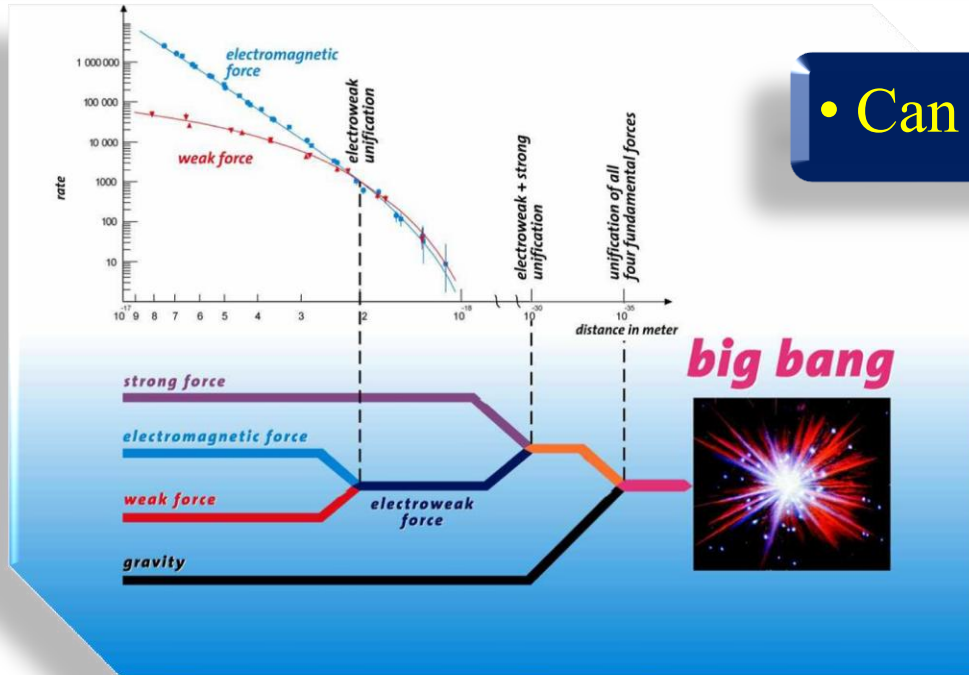
Stability is much more difficult to check!

# Important questions:

- Should a star spend at least  $10^7$  years in the EW stage before it becomes a black hole?
- When does electroweak burning stop?
- Can electroweak burning spend all of the quark fuel (i.e. black hole never forms) ?

Requires much more detailed modeling of the star!

# Next? GUT phase transition?



• Can GUT stars exist?

• Core of the GUT star must be microscopic

- Unlike to support the whole star
- Unlike to happen before the star crosses its own  $R_S$

# Conclusions

**Electroweak star** is an interesting new phase in stellar evolution

## We found the solution and basic properties:

- Enormous energy is released at the core
- Energy release rate is moderate at the surface
- Life-time can be as long as 10 million years



## Remains to be done:

- Photo-luminosity and other observational signature
- Full time-dependent analysis



A black oval button with a slight shadow on the left side, containing the text "THANK YOU" in a glowing yellow font.

THANK YOU