LHC data confront Monte Carlo predictions

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ICHEP 2010







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Outline

- Introduction: why minimum bias events?
- 2 Models for soft particle production
- 3 A potpourri of results from Tevatron to LHC
- A snapshot of one new development
- 5 Conclusions



Introduction: why minimum bias events?

The bird's eye view

• Name suggests: most complete view of physics

(at the LHC or any other experiment)

- As such: intellectual challenge to grasp complete picture: Up to now no complete model based on first principles including all facets elastic scattering, diffractive events & hard jets - on the same footing; instead: phenomenological models with many choices & parameters.
- First day physics at the LHC new energy regime to challenge our understanding of soft particle production and the corresponding tunes of the event generators.
- Intimate connection of minimum bias physics to underlying event

(Although underlying event \neq minimum bias event)

Image: Image:

 Therefore: Immediate impact on some searches for new physics (see below), jet physics (e.g. jet energy scale, relation of hadrons ↔ partons, ...) etc..

Introduction		

Importance of particle production: Rapidity gaps

As an example consider Higgs physics at the LHC.

- Important channels for Higgs searches: VBF
- Signature: two forward jets + rapidity gap, filled by Higgs decay
- Essential for S/B: rapidity gap (and its survival rate)
- Typical manifestation of gap: central jet veto.
- Also discussed: track based rapidity gaps.
- Typically, the requirement of central gap effectively suppresses perturbative QCD-driven backgrounds (hadrons produced along colour exchange between pp)
- Obvious: Underlying event/soft particle production beyond perturbation theory may spoil this picture.



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Rapidity gaps, once again

- Another daring idea: elastic Higgs production
- Signature: two intact protons in the forward direction, only Higgs in central detector.
- Again: rapidity gap (suppresses background)
- Rare process: very few events.
- If possible to trigger on forward protons: super clean signature, full control over FS kinematics.
- Most likely not a discovery channel, but due to spin-selection (only $J^{PS} = 0^{++}$ possible) a great chance to measure the quantum numbers of H.
- Need a straw-man to check calculations: central diffractive production.





Models		

Models for soft particle production

Reminder: Eikonals

• Optical theorem relates total cross section σ_{tot} with elastic scattering amplitude $\mathcal{A}(s,t)$ through

$$\sigma_{ ext{tot}}(s) = rac{1}{s} \operatorname{Im}[\mathcal{A}(s, t = 0)]$$

• Rewrite $\mathcal{A}(s,t)$ as $a(s,b_{\perp})$ in impact parameter space

$$\mathcal{A}(s,t=-ec{q}_{\perp}^2)=2s\int\mathrm{d}^2b_{\perp}e^{iec{q}_{\perp}\cdotec{b}_{\perp}}a(s,ec{b}_{\perp})$$

and introduce eikonal Ω

$$a(s,\,ec{b}_\perp)=rac{e^{-\Omega(s,ec{b}_\perp)}-1}{2i}.$$

• Total cross section now reads:

$$\sigma_{
m tot}(s) = 2 \int {
m d}^2 b_\perp [1 - e^{-\Omega(s, ec{b}_\perp)}]$$

and similar expressions for elastic & inelastic cross sections.

(guarantees unitarity of cross sections)

Models		

Formulation of models based on eikonals

• Typical procedure: Write eikonals as sum of soft and hard part:

$$\Omega(s,ec{b}_{\perp}) = \Omega_{\mathcal{S}}(s,ec{b}_{\perp}) + \Omega_{ ext{H}}(s,ec{b}_{\perp})$$

• Use perturbative QCD for hard part:

$$\Omega_{\rm H}(s, \vec{b}_{\perp}) = \frac{1}{2} \rho(\vec{b}_{\perp}) \hat{\sigma}_{2 \rightarrow 2}(s)$$

Here:

- $\rho(\vec{b}_{\perp}) =$ spatial distribution of partons, parametrised with form factors;
- *^ˆ*_{2→2}(s) parton-parton cross section from QCD, typically collinear factorisation (see next slide).
- Can fix Ω_S as constant to reproduce anticipated total *pp*-xsec from fits, set it to 0, or give some dynamics based on Regge-physics or similar.
- Produce hard and soft interactions according to their eikonals, typically as a Poissonian distribution with argument like Ω/σ_{tot} .



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	Results	

MC vs. Tevatron, 630 & 1800 GeV, inclusive data



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	Results	

MC vs. underlying event at Tevatron, Run I



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	Results	

MC vs. ATLAS, 900 GeV

- Charged particles only, analysis track-based, corrected to particle level.
- \bullet Trigger: at least one particle with $p_\perp > 500$ MeV, $|\eta| < 2.5$



	Results	

MC vs. underlying event at ATLAS, 7 TeV



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	Results	

Why I did not show any CMS/ALICE data

- I wanted to show some results of tuned SHERPA and PYTHIA and not just regurgitate the plots shown by the collaborations.
- Main problem of CMS data: Not corrected for detector effects this makes it nearly impossible to run your MC and draw any conclusion from the comparison with data.
- ALICE is an even sadder story:
 - In their first publication, ALICE corrected (with Monte Carlo) on "non-single diffractive" events. This selection is based on rapidity gaps, which however, may also occur as fluctuation, due to hadronisation etc., of "proper" inelastic events.
 - In addition, they filled the region outside their acceptance (especially the low p_{\perp} -region) with MC.

Therefore it is hard to see how you can test MCs with this data in an unbiased fashion.

	Results	

A remark on tuning

- Tuning can help a model to describe a limited subset of data but deficiencies of the model will show up at other places.
- Always try to check for the overall picture.
- Measure R₃₂ the three-to-two jet ratio.
- Fairly independent of PDFs, allows to test running of α_S at hadron colliders.
- Receives contributions from hadronisation and underlying event activity.
- Two parameters: $p_{\perp}^{\max} = \text{minimal } p_{\perp} \text{ of 1st jet}$ $p_{\perp}^{\min} = \text{scale of other jets}$



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	Results	

Rapidity gaps from fluctuations - a remark on diffraction

- Often, diffractive events are identified by a rapidity gap of a certain size.
- Nearly always the effect of fluctuations, where perturbative events produce rapidity gaps due to hadronisation effects, are ignored.
- It is hard to estimate the probability for this to happen from first principles or a Monte Carlo simulation, see below.
- Figure to the right:
 - SHERPA simulation with native cluster and Lund fragmentation, both tuned to LEP data.
 - Figure shows probability to find a rapidity gap of a given size, with different p_{\perp} -thresholds (bottom to top: 100 MeV, 500 MeV, 1000 MeV) at a c.m.-energy of 7 TeV.



A new model for Minimum Bias (and the underlying event)

Underlying ideas

- Multi-channel eikonal approach (decompose proton in more than one diffractive eigenstate, one eikonal for each pair of states): allows for natural description of low-mass diffraction.
- Rooted in unitarisation by exponentiating eikonals.
- BFKL-inspired interpretation: exchange of "ladders" (cut pomerons) between hadrons yields eikonal.
- Ladders described by evolution equations in rapidity y, with form factor as starting condition at $y = \pm \infty''$; evolution for both hadrons coupled through rescattering/absorption in effective description.
- $\bullet\,$ Number of ladders $\propto\,$ eikonal, partons emitted by ladder according to differential equation.
- Naturally incorporates diffraction/diffractive parts in ladder dynamics.
- \bullet Work in progress, expect model in \approx 4 weeks.

	New developments	

Some appetisers



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			Conclusions
Conclusio	าร		

- There are two classes of models currently used for describing minimum bias/ soft particle production data: Regge-based and PQCD-based.
- The former (implemented e.g. in PHOJET) have some difficulties describing Tevatron data, energy extrapolation, and perturbative QCD I did not discuss it here.
- The latter (implemented in standard MCs such as PYTHIA and SHERPA) do somewhat better they're still far from being perfect, and very susceptible to tuning. PDF effects also play a substantial role.
- None of the models manages to describe all data satisfactory.
- In my opinion this shows that we have not understood minimum bias physics.