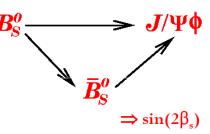
# New Measurement of B<sub>s</sub> Mixing Phase at CDF

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#### CP Violation in $B_s \rightarrow J/\Psi\Phi$ Decays

- Analogously to the neutral  $B^0$  system, CP violation in  $B_s$  system is accessible through interference of decays with and without mixing

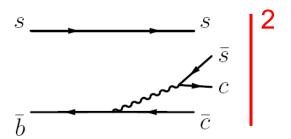


- In SM, *CP* violation phase  $\beta_s$  is predicted to be very small  $\sim sin^2(\theta_{Cabibbo})$ 

$$\beta_s^{\text{SM}} = \arg(-V_{ts}V_{tb}^*/V_{cs}V_{cb}^*) \approx 0.02$$

dominant contribution from top quark

$$\begin{pmatrix}
V_{ud} & V_{us} & V_{ub} \\
V_{cd} & V_{cs} & V_{cb} \\
V_{td} & V_{ts} & V_{tb}
\end{pmatrix}$$

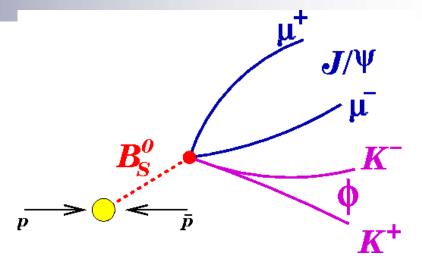


New Physics particles?

- New physics particles running in the mixing diagram may enhance  $\beta_{s}$ 
  - large  $\beta_s$   $\rightarrow$  clear indication of New Physics!

#### $B_s \rightarrow J/\Psi\Phi$ Decays

- Measure:
  - $B_s$  lifetime  $\tau_s$
  - $B_{sH}$ ,  $B_{sL}$  decay width difference  $\Delta \Gamma_s$
  - *CP* violating phase  $\beta_s$

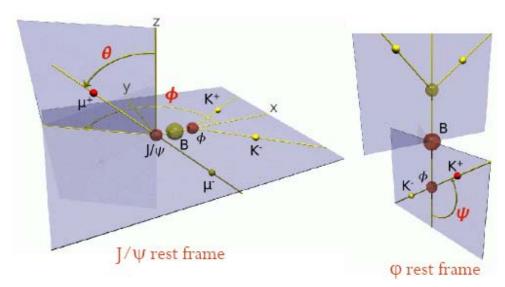


- Decay of  $B_s$  (spin 0) to  $J/\Psi$  (spin 1) and  $\Phi$  (spin 1) leads to three different angular momentum final states:

$$L = 0$$
 (s-wave), 2 (d-wave)  $\rightarrow CP$  even ( = short lived or light  $B_s$  if no  $CPV$ )

$$L = 1$$
 (p-wave)

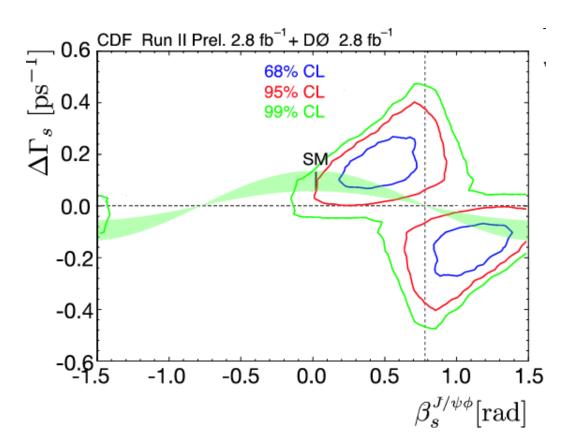
$$\rightarrow$$
 CP odd ( = long lived or heavy  $B_s$  if no CPV)



- Three decay angles  $\overrightarrow{\rho} = (\theta, \phi, \psi)$  describe directions of final decay products  $\mu^+ \mu^- K^+ K^-$ 

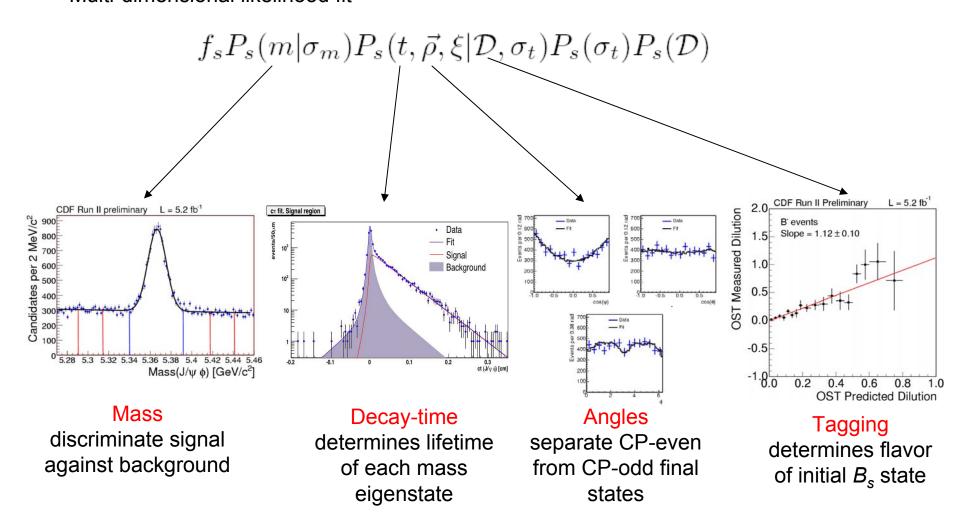
#### Status Before This Update: CDF + DØ Combination

- CDF + DØ combination done by the Tevatron B Working Group: http://tevbwg.fnal.gov/
- Shows intriguing 2.1σ deviation from SM expectation (CDF note 9787)



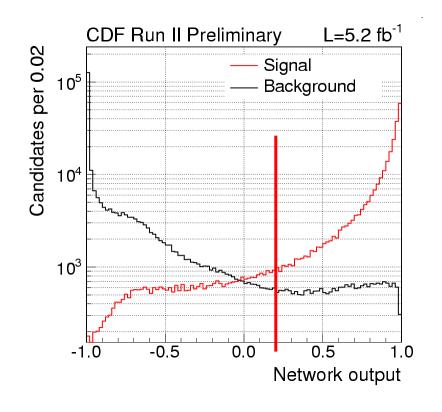
#### **Analysis Components**

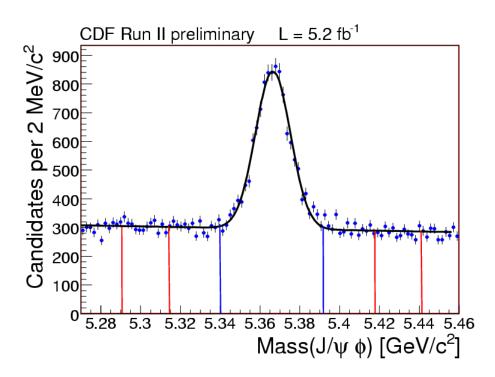
- Multi-dimensional likelihood fit



#### Signal Reconstruction

- Reconstruct  $B^0_s \rightarrow J/\psi \Phi$  in 5.2 fb<sup>-1</sup> of data from sample selected by di-muon trigger
- Combine kinematic variables with particle ID information (dE/dx, TOF) in neural network to discriminate signal from background
- Yield of ~6500 signal  $B_s$  events with  $S/B \sim 1$  (compared to ~3150 in 2.8 fb<sup>-1</sup>)

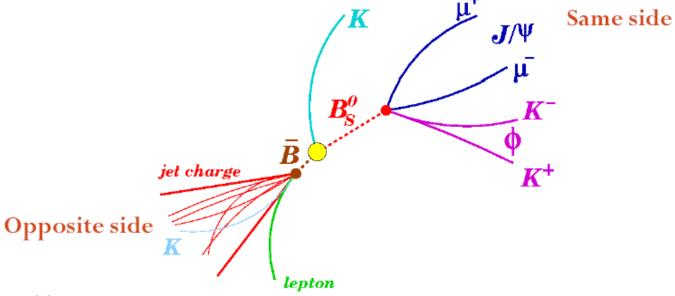




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#### Flavor Tagging

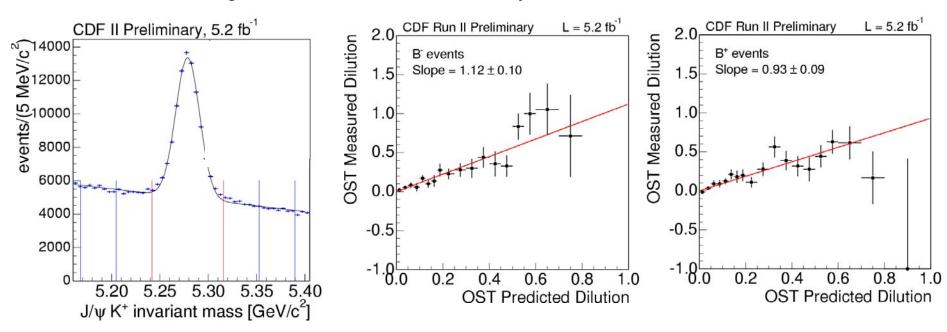
- Tevatron: b-quarks mainly produced in pairs of bottom anti-bottom
  - → flavor of the *B* meson at production inferred with:
- Opposite Side Tagger (OST): exploits decay products of other *b*-hadron in the event
- Same Side Kaon Tagger (SSKT): exploits correlations with particles produced in fragmentation



- Output of flavor tagger:
  - flavor decision (*b*-quark or anti-*b*-quark)
  - probability that the decision is correct: P = (1 + Dilution) / 2

#### Opposite Side Tagging Calibration and Performance

- OST combines in a NN opposite side lepton and jet charge information
- Initially calibrated using a sample of inclusive semileptonic B decays
  - predicts tagging probability on event-by-event basis
- Re-calibrated using  $\approx$  52,000  $B^{+/-} \rightarrow J/\Psi K^{+/-}$  decays



- OST efficiency = 94.2 +/- 0.4%, OST dilution = 11.5 +/- 0.2 % (correct tag probability ~56%)
- Total tagging power = 1.2%

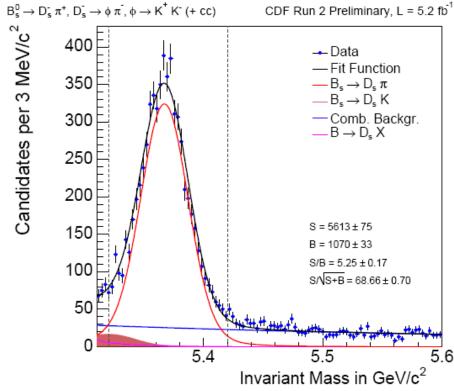
#### Same Side Tagging Calibration

- Event-by-event predicted dilution based on simulation
- Calibrated with 5.2 fb<sup>-1</sup> of data
- Simultaneously measuring the  $B_s$  mixing frequency  $\Delta m_s$  and the dilution scale factor A

$$P_{Sig}(ct|\sigma_{ct}, \xi = \xi_D \cdot \xi_P, D) = \frac{1}{N} \cdot \left[ \frac{1}{\tau} e^{-\tilde{t}/\tau} \cdot (1 + \xi \mathcal{A}D \cdot \cos(\Delta m_s \tilde{t})) \right] \otimes \mathcal{G}(c\tilde{t}|\sigma_{ct}) \cdot \epsilon(ct|\sigma_{ct})$$

- *D* event by event predicted dilution
- $\frac{-\xi}{B_s}$  and un-tagged events
- Fully reconstructed  $B_s$  decays selected by displaced track trigger

Decay Channel	S
$B_s^0 \to D_s^- \pi^+, \ D_s^- \to \phi \pi^-$	$5613 \pm 75$
$B_s^0 \to D_s^- \pi^+, \ D_s^- \to K^* K^-$	$2761 \pm 53$
$B_s^0 \to D_s^- \pi^+, \ D_s^- \to (3\pi)^-$	$2652 \pm 52$
$B_s^0 \to D_s^-(3\pi)^+, \ D_s^- \to \phi\pi^-$	$1852 \pm 43$
Sum	$\boxed{12877 \pm 113}$





- $B_s$  oscillation frequency measured  $\Delta m_s = (17.79 \pm 0.07) \; \mathrm{ps}^{-1}$  (statistical error only)
- In good agreement with the published CDF measurement with 1 fb<sup>-1</sup> PRL 97, 242003 2006, PRL 97, 062003 2006

$$\Delta m_s = 17.77 \pm 0.10 \; (\mathrm{stat}) \pm 0.07 \; (\mathrm{syst}) \; \mathrm{ps}^{-1}$$
 used as external constraint in  $\beta_s$  measurement

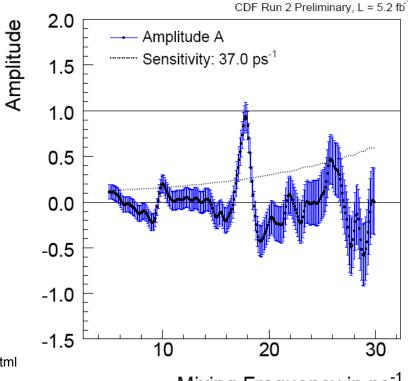
 Dilution scale factor (amplitude) in good agreement with 1:

$$A = 0.94 \pm 0.15 \ (stat.) \pm 0.13 \ (syst.)$$

- Largest systematic uncertainty from decay time resolution modeling
- Total SSKT tagging power:

$$\varepsilon \mathcal{A}^2 D^2 = (3.2 \pm 1.4) \%$$

http://www-cdf.fnal.gov/physics/new/bottom/100204.blessed-sskt-calibration/index.html CDF public note 10108

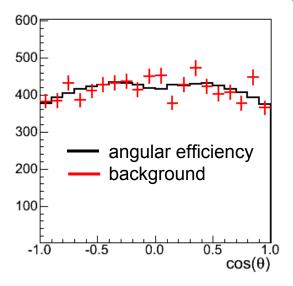


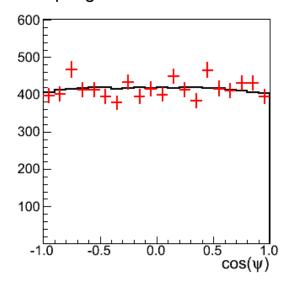
Mixing Frequency in ps<sup>-1</sup>

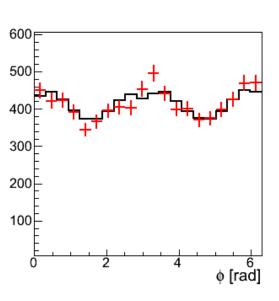


- CP even and CP odd final states have different angular distributions
  - $\rightarrow$  use angles  $\rho = (\theta, \phi, \psi)$  to statistically separate *CP* even and *CP* odd components
- Detector acceptance distorts the angular distributions
  - → determine 3D angular efficiency function from simulation and account for this effect in the fit
- Cross check angular efficiency by comparing with background angular distributions
  - good agreement indicates good modeling of angular efficiency

#### CDF Simulation of Detector Angular Sculpting







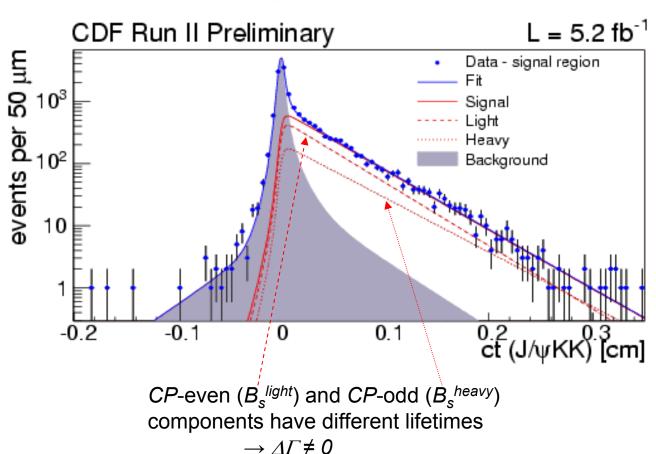
### B<sub>s</sub> Lifetime and Decay Width Difference

- Assuming no *CP* violation ( $\beta_s = 0$ ) obtain most precise measurements of lifetime  $\tau_s$  and decay width difference  $\Delta \Gamma_s$ :

$$au_s = 1.53 \pm 0.025 \text{ (stat.)} \pm 0.012 \text{ (syst.) ps}$$
  
 $\Delta \Gamma = 0.075 \pm 0.035 \text{ (stat.)} \pm 0.01 \text{ (syst.)} \ ps^{-1}$ 

compared to PDG 2009 averages:

$$\tau_{\rm s} = 1.472^{+0.024}_{-0.026} \, {\rm ps}$$
  
 $\Delta \Gamma_{\rm s} = 0.062^{+0.034}_{-0.037} \, {\rm ps}^{-1}$ 



#### Polarization Amplitudes

Most precise measurement of polarization amplitudes

100E

-0.5

0.5

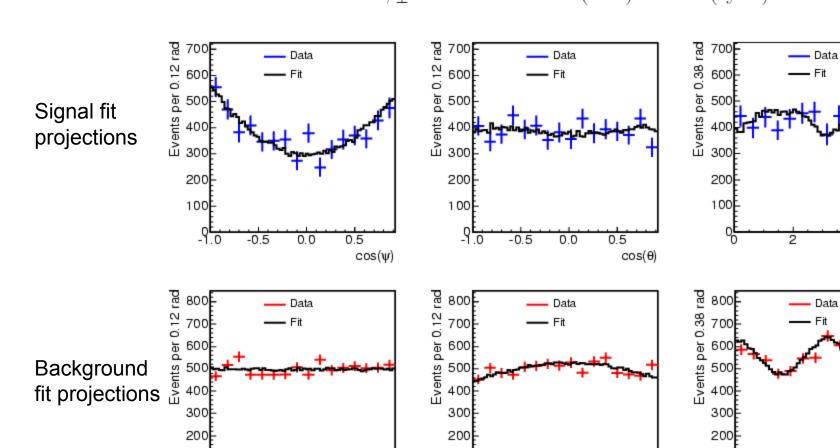
cos(ψ)

0.0

$$|A_{\parallel}(0)|^2 = 0.231 \pm 0.014 \text{ (stat)} \pm 0.015 \text{ (syst.)}$$
  
 $|A_0(0)|^2 = 0.524 \pm 0.013 \text{ (stat)} \pm 0.015 \text{ (syst.)}$   
 $\phi_{\perp} = 2.95 \pm 0.64 \text{ (stat)} \pm 0.07 \text{ (syst.)}.$ 

100

2



100

-0.5

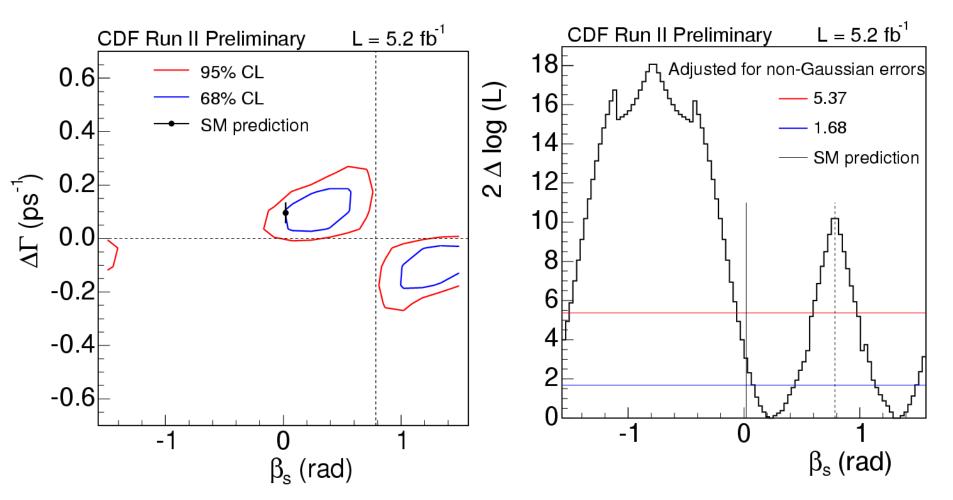
0.0

0.5

 $cos(\theta)$ 

#### CP Violation Phase $\beta_s$ with 5.2 fb<sup>-1</sup> at CDF

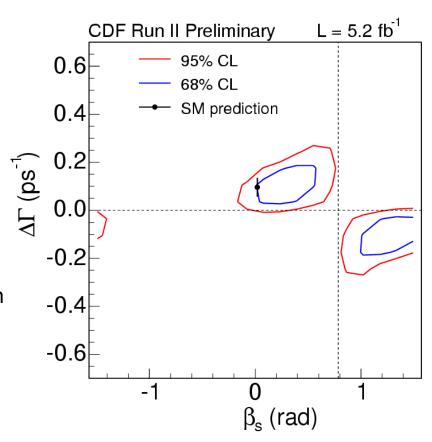
- Final confidence regions in  $\beta_s$ - $\Delta\Gamma_s$  space: [0.02, 0.52] U [1.08, 1.55] at 68% C.L.
- Agreement with SM at ~1σ level





#### **Conclusions**

- Measurement of CP violation in  $B_s$  system updated by CDF with 5.2 fb<sup>-1</sup>
- Tightened constraints in  $\beta_s$  space: [0.02, 0.52] U [1.08, 1.55] at 68% C.L.
- Improved agreement with SM expectation at ~1σ level
- Best measurements of  $B_s$  lifetime, decay width difference  $\Delta \Gamma_s$  and polarization amplitudes





#### **Prospects**

- Possible analysis improvements:
  - Improve statistics by ~25-30% by adding  $B_s \to J/\Psi\Phi$  decays from displaced track trigger (difficult due to trigger effects on decay time )
  - Addition of new decay modes:

 $B_s \rightarrow J/\Psi f^0$ , with  $f^0 \rightarrow \pi\pi$  (less statistics but no angular analysis needed since final state is CP eigen-state)

$$B_s \rightarrow \psi(2s)\Phi$$

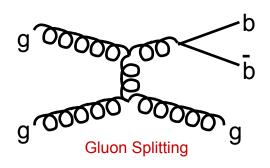
- Add more data!
  - 7 fb<sup>-1</sup> already recorded
  - expect to double sample size (~10 fb-1) by end of Tevatron running in 2011

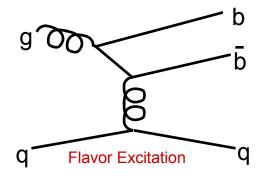
## Backup Slides

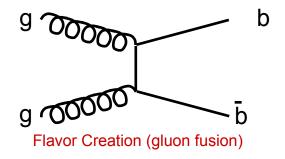


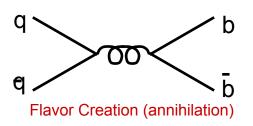
#### B Physics at the Tevatron

- Mechanisms for b production in pp collisions at 1.96 TeV









- At Tevatron, b production cross section is much larger compared to B-factories

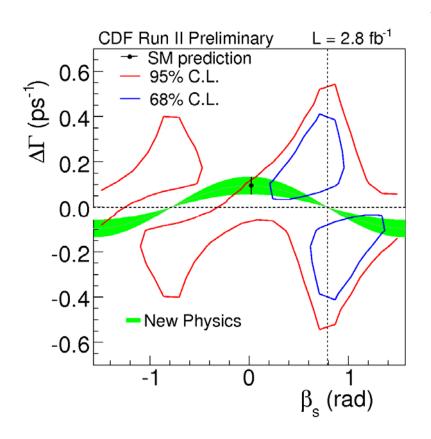
  → Tevatron experiments CDF and DØ enjoy rich B Physics program
- Plethora of states accessible only at Tevatron:  $B_s$ ,  $B_c$ ,  $\Lambda_b$ ,  $\Xi_b$ ,  $\Sigma_b$ ...  $\to$  complement the B factories physics program
- Total inelastic cross section at Tevatron is ~1000 larger than b cross section

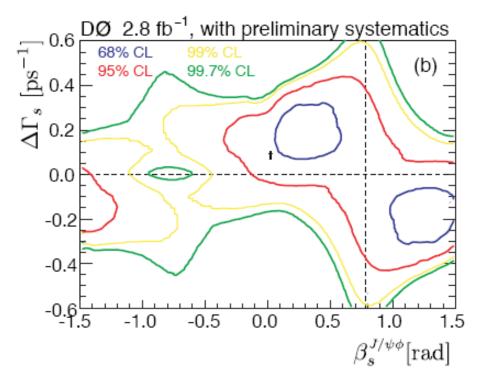
  → large backgrounds suppressed by triggers that target specific decays



#### Status Before This Update

- Both CDF (public note 9458) and DØ (conference Note 5933-CONF) showed ~1.5σ deviations from SM in the same direction







$$\beta_{\rm S}$$
 VS  $\phi_{\rm S}$ 

- Up to now, introduced two different phases:

$$\phi_{\rm s}^{\rm SM} = \arg\left(-\frac{M_{12}}{\Gamma_{12}}\right) \approx 4 \times 10^{-3} \qquad \text{ and } \qquad \beta_s^{\rm SM} = \arg(-V_{ts}V_{tb}^*/V_{cs}V_{cb}^*) \approx 0.02$$

- New Physics affects both phases by same quantity  $\phi_s^{\mathrm{NP}}$  (arxiv:0705.3802v2):

$$2\beta_s = 2\beta_s^{\text{SM}} - \phi_s^{\text{NP}}$$
$$\phi_s = \phi_s^{\text{SM}} + \phi_s^{\text{NP}}$$

- If the new physics phase  $\phi_s^{\rm NP}$  dominates over the SM phases  $2\beta_s^{\rm SM}$  and  $\phi_s^{\rm SM}$   $\to$  neglect SM phases and obtain:

$$2\beta_s = -\phi_s^{\text{NP}} = -\phi_s$$

#### Decay Rate

-  $B_s \rightarrow J/\Psi\Phi$  decay rate (A.S. Dighe *et al.*, Phys. Lett. B **36**9 144 (1996)) :

$$P_B(\theta, \phi, \psi, t) = \frac{9}{16\pi} |\mathbf{A}(t) \times \hat{n}|^2$$

where: 
$$\mathbf{A}(t) = (\mathcal{A}_0(t)\cos\psi, -\frac{\mathcal{A}_{\parallel}(t)\sin\psi}{\sqrt{2}}, i\frac{\mathcal{A}_{\perp}(t)\sin\psi}{\sqrt{2}})$$
 and  $\hat{n} = (\sin\theta\cos\phi, \sin\theta\sin\phi, \cos\theta)$ 

- Time evolution of transversity amplitudes  $A_0$ ,  $A_{||}$ ,  $A_{\perp}$ :

$$\mathcal{A}_i = \frac{e^{-\Gamma t/2}}{\sqrt{\tau_H + \tau_L \pm \cos 2\beta_s \left(\tau_L - \tau_H\right)}} \left[ E_+(t) \pm e^{2i\beta_s} E_-(t) \right] a_i$$

where  $\pm$  corresponds to CP-even and CP-odd final states,  $\sum_i |a_i|^2 = 1$  and

$$E_{\pm}(t) \equiv \frac{1}{2} \left[ e^{+\left(\frac{-\Delta\Gamma}{4} + i\frac{\Delta m}{2}\right)t} \pm e^{-\left(\frac{-\Delta\Gamma}{4} + i\frac{\Delta m}{2}\right)t} \right]$$

- Finally:

$$P_{B}(\theta, \psi, \phi, t) = \frac{9}{16\pi} \left\{ |\mathbf{A}_{+}(\mathbf{t}) \times \hat{n}|^{2} + |\mathbf{A}_{-}(\mathbf{t}) \times \hat{n}|^{2} + 2Re((\mathbf{A}_{+}(\mathbf{t}) \times \hat{n}) \cdot (\mathbf{A}_{-}^{*}(\mathbf{t}) \times \hat{n})) \right\}$$

$$= \frac{9}{16\pi} \left\{ |\mathbf{A}_{+} \times \hat{n}|^{2} |f_{+}(t)|^{2} + |\mathbf{A}_{-} \times \hat{n}|^{2} |f_{-}(t)|^{2} + 2Re((\mathbf{A}_{+} \times \hat{n}) \cdot (\mathbf{A}_{-}^{*} \times \hat{n}) \cdot f_{+}(t) \cdot f_{-}^{*}(t)) \right\}$$

$$|f_{\pm}(t)|^2 = \frac{1}{2} \frac{(1 \pm \cos 2\beta_s)e^{-\Gamma_L t} + (1 \mp \cos 2\beta_s)e^{-\Gamma_H t} \mp 2\sin 2\beta_s e^{-\Gamma t} \sin \Delta mt}{\tau_L (1 \pm \cos 2\beta_s) + \tau_H (1 \mp \cos 2\beta_s)} \qquad f_{\pm}(t)f_{\pm}^*(t) = \frac{e^{-\Gamma t} \cos \Delta mt + i \cos 2\beta_s e^{-\Gamma t} \sin \Delta mt + i \sin 2\beta_s (e^{-\Gamma_L t} - e^{-\Gamma_H t})/2}{\sqrt{[(\tau_L - \tau_H)\sin 2\beta_s]^2 + 4\tau_L \tau_H}}$$

#### Decay Rate with S-Wave Included

- Including the s-wave contribution the probability density function becomes:

$$\rho_B(\theta,\phi,\psi,t,\mu) = \frac{9}{16\pi} \left| \left[ \sqrt{1-F_s} g(\mu) \mathbf{A}(t) + e^{i\delta_s} \sqrt{F_s} \frac{h(\mu)}{\sqrt{3}} \mathbf{B}(t) \right] \times \hat{n} \right|^2$$
 where: 
$$\mathbf{B}(t) = (\mathcal{B}(t),0,0) \quad \text{and} \quad \mathcal{B}(t) = \frac{e^{-\Gamma t/2}}{\sqrt{\tau_H + \tau_L - \cos 2\beta_s \left(\tau_L - \tau_H\right)}} \left[ E_+(t) - e^{2i\beta_s} E_-(t) \right]$$

 $g(\mu)$  is relativistic Breit-Wigner to describe asymmetric  $\Phi$  mass shape and  $h(\mu)$  is constant

- Integrating out the dependence on the KK mass:

$$\rho_B(\theta, \psi, \phi, t) = (1 - F_s) \cdot P_B(\theta, \psi, \phi, t) + F_s Q_B(\theta, \psi, \phi, t)$$

$$+ 2 \frac{\sqrt{27}}{16\pi} Re \left[ \mathcal{I}_{\mu} \left( (\mathbf{A}_{-} \times \hat{n}) \cdot (\mathbf{B} \times \hat{n}) \cdot |f_{-}(t)|^2 + (\mathbf{A}_{+} \times \hat{n}) \cdot (\mathbf{B} \times \hat{n}) \cdot f_{+}(t) \cdot f_{-}^*(t) \right) \right]$$

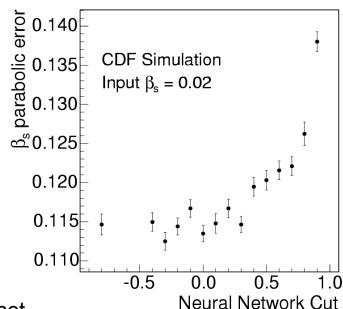
where:  $I(\mu)$  is a function of the s-wave phase and  $Q_B(\theta, \phi, \psi, t) = \frac{3}{16\pi} |\mathbf{B}(t) \times \hat{n}|^2$ 



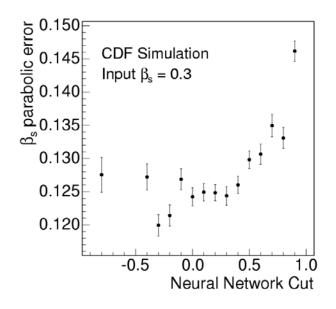
#### Analysis Improvements with Respect To 2008 Update

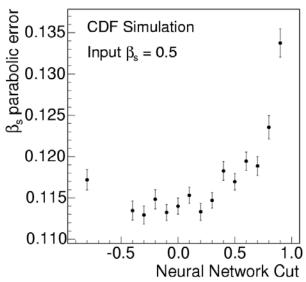
- Almost doubled data sample (from 2.8 fb<sup>-1</sup> in 2008 to 5.2 fb<sup>-1</sup> now)

- Improved signal selection:
  - use particle ID (dE/dx and TOF) for full dataset
  - use pseudo-experiments to optimize neural network selection to minimize  $\beta_s$  statistical uncertainty (previously used  $S/(S+B)^{1/2}$  as figure of merit)



- Same side kaon tagger (SSKT) used for the full dataset
  - re-calibrated by measuring B<sub>s</sub> mixing frequency with 5.2 fb<sup>-1</sup>
- Inclusion of S-wave contamination in the likelihood fit







$$c au = 458.64 \pm 7.54 \; ({
m stat.}) \; \mu m$$
  $c au = 459.1 \pm 7.7 \; ({
m stat.}) \; \mu m$   $\Delta\Gamma = 0.075 \pm 0.035 \; ({
m stat.}) \; ps^{-1}$   $\Delta\Gamma = 0.073 \pm 0.03 \; ({
m stat.}) \; ps^{-1}$   $|A_{\parallel}|^2 = 0.231 \pm 0.014 \; ({
m stat.})$   $|A_{\parallel}|^2 = 0.232 \pm 0.014 \; ({
m stat.})$   $|A_{\parallel}|^2 = 0.523 \pm 0.012 \; ({
m stat.})$   $\phi_{\perp} = 2.95 \pm 0.64 \; ({
m stat.})$   $\phi_{\perp} = 2.80 \pm 0.56$ 

Tagged, with S-wave

Untagged, with S-wave

$$c au = 456.93 \pm 7.69 \; ({
m stat.}) \; \mu m$$
  $c au = 457.2 \pm 7.9 \; ({
m stat.}) \; \mu m$   $\Delta\Gamma = 0.071 \pm 0.036 \; ({
m stat.}) \; ps^{-1}$   $\Delta\Gamma = 0.070 \pm 0.04 \; ({
m stat.}) \; ps^{-1}$   $|A_{\parallel}|^2 = 0.233 \pm 0.015 \; ({
m stat.})$   $|A_{\parallel}|^2 = 0.233 \pm 0.016 \; ({
m stat.})$   $|A_{0}|^2 = 0.521 \pm 0.013 \; ({
m stat.})$ 

 $c\tau = 458.64 \pm 7.54 \text{ (stat.) } \mu m$   $c\tau = 459.1 \pm 7.7 \text{ (stat.) } \mu m$  $|A_0|^2 = 0.523 \pm 0.012$  (stat.)  $\phi_{\perp} = 2.80 \pm 0.56$ 

Tagged, no S-wave

Untagged, no S-wave

$$c au = 456.93 \pm 7.69 \; ext{(stat.)} \; \mu m$$
  $c au = 457.2 \pm 7.9 \; ext{(stat.)} \; \mu m$   $\Delta\Gamma = 0.071 \pm 0.036 \; ext{(stat.)} \; ps^{-1}$   $\Delta\Gamma = 0.070 \pm 0.04 \; ext{(stat.)} \; ps^{-1}$   $|A_{\parallel}|^2 = 0.233 \pm 0.015 \; ext{(stat.)}$   $|A_{\parallel}|^2 = 0.233 \pm 0.016 \; ext{(stat.)}$   $|A_{0}|^2 = 0.520 \pm 0.013 \; ext{(stat.)}$ 

# 

## Systematic Uncertainties

Systematic	ΔΓ	$c au_s$	$ A_{\ }(0) ^2$	$ A_0(0) ^2$	$\overline{\phi_{\perp}}$
Signal efficiency:					
Parameterisation	0.0024	0.96	0.0076	0.008	0.016
MC reweighting	0.0008	0.94	0.0129	0.0129	0.022
Signal mass model	0.0013	0.26	0.0009	0.0011	0.009
Background mass model	0.0009	1.4	0.0004	0.0005	0.004
Resolution model	0.0004	0.69	0.0002	0.0003	0.022
Background lifetime model	0.0036	2.0	0.0007	0.0011	0.058
Background angular distribution:					
Parameterisation	0.0002	0.02	0.0001	0.0001	0.001
$\sigma(c\tau)$ correlation	0.0002	0.14	0.0007	0.0007	0.006
Non-factorisation	0.0001	0.06	0.0004	0.0004	0.003
$B^0 \to J \psi K^*$ crossfeed	0.0014	0.24	0.0007	0.0010	0.006
SVX alignment	0.0006	2.0	0.0001	0.0002	0.002
Mass error	0.0001	0.58	0.0004	0.0004	0.002
c au error	0.0012	0.17	0.0005	0.0007	0.013
Pull bias	0.0028		0.0013	0.0021	
Totals	0.01	3.6	0.015	0.015	0.07

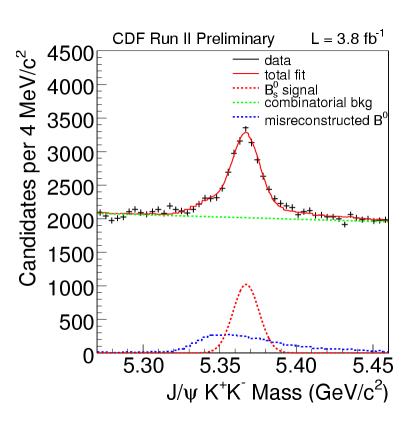
# м

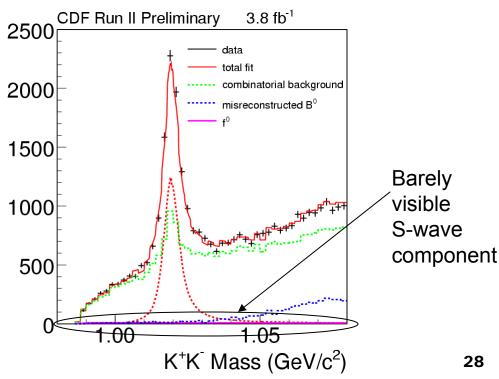
## Dilution Scale Factor Systematic Uncertainties

Modification	Systematic Uncertainty	
Proper decay time resolution scaling	0.11	
Resolution model	0.06	
Cabibbo reflection	0.03	
Cabibbo fraction	negligible	
Mass window	negligible	
Selection of upper side band	negligible	
$\Lambda_b$ template	negligible	
$\Delta\Gamma/\Gamma$	negligible	
Mean Lifetime	negligible	
Trigger Composition	negligible	
Signal Mass Model	negligible	
Total 0.13		

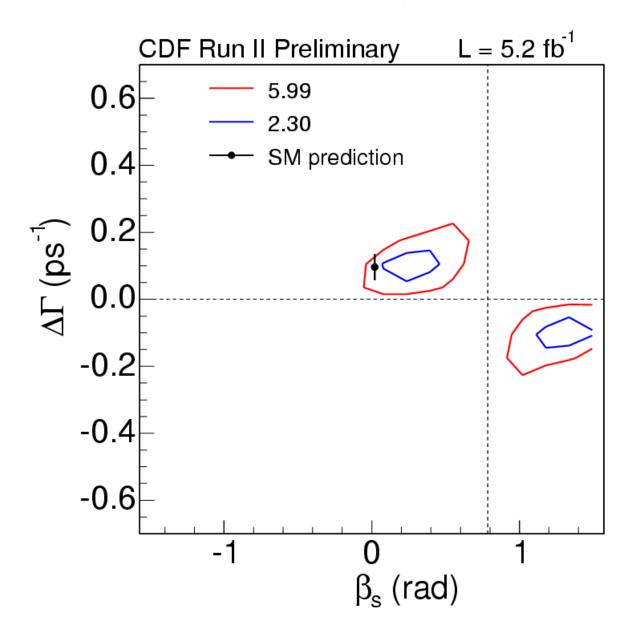
#### S-Wave Cross Check Using KK Mass Spectrum

- Cross check the result from angular fit by fitting the KK invariant mass spectrum
- From a fit to the  $B_s$  mass distribution with wide KK mass range selection (0.980,1.080 GeV), determine contributions of combinatorial background, mis-reconstructed  $B^0$ , and  $B_s$  events
- Good fit of the KK mass spectrum with 2% f<sup>0</sup> contributions

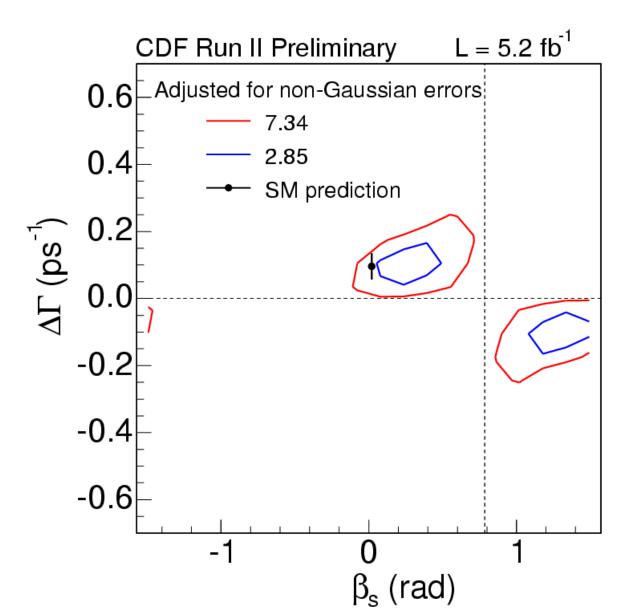




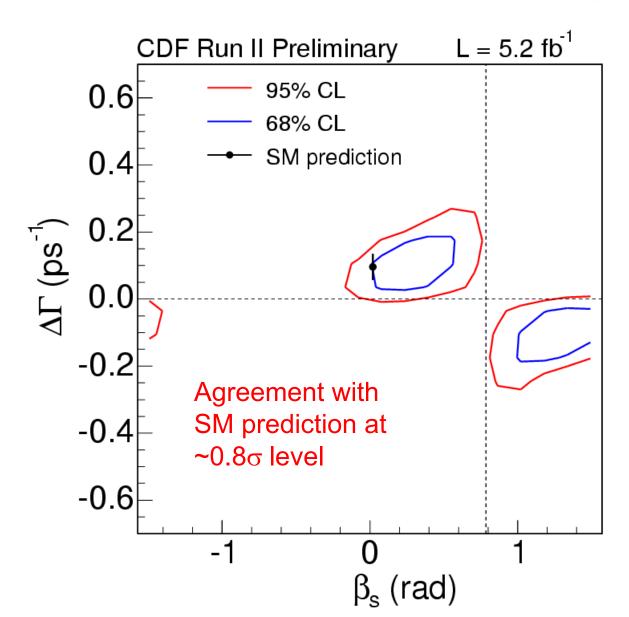
#### $\beta_s$ - $\Delta\Gamma$ Contours Without Coverage Adjustment

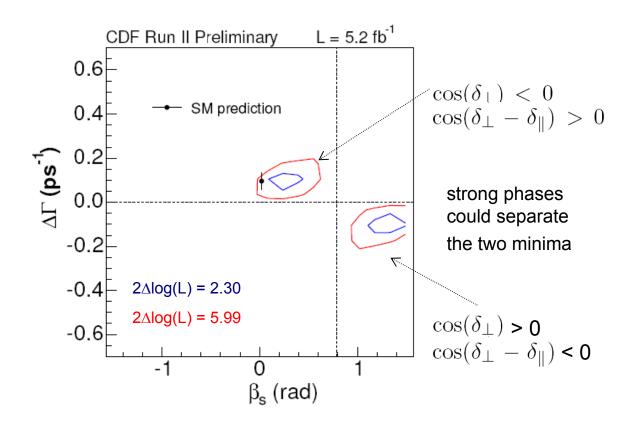


### $\beta_s$ - $\Delta\Gamma$ Contours With Coverage Adjustment



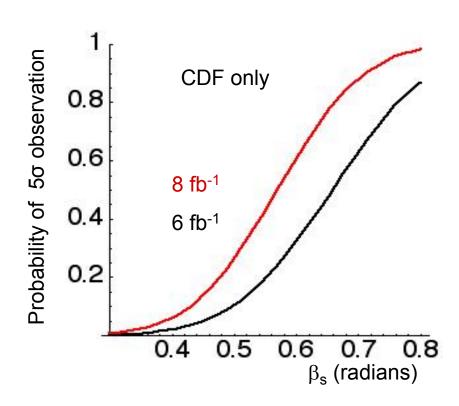
#### $\beta_s$ - $\Delta\Gamma$ Contours With Systematics on Coverage







#### Sensitivity



#### Introduction

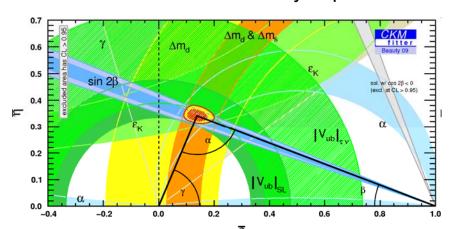
- *CP* violation means that the laws of nature are not invariant under the simultaneous transformation of Charge and Parity

- Charge conjugation transforms particles into anti-particles
- Parity transformation is a mirror reflection (space inversion)
- Parity conservation was first questioned by T.D. Lee and C.N. Yang in 1956 when they argued that there was no experimental evidence for parity conservation in weak interactions
- Same year, C.S. Wu showed that Parity is violated in beta decays of Cobalt nuclei
- The combined *CP* was soon adopted as the correct symmetry, just to be shown wrong by Cronin and Fitch in 1964 when they showed that *CP* is violated in neutral Kaon decays



#### Why Look for CPV in B<sub>s</sub> System?

- CP violation has been studied in various Kaon and B-meson decays
- CKM matrix is well constrained by experimental data



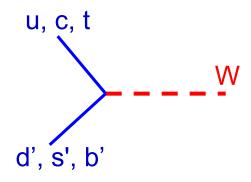


- Within the SM framework, ČP violation in the quark sector is too small to explain the matter antimatter asymmetry in the universe
- Could still find large *CP* violation within the SM in the lepton sector
  - initial asymmetry between leptons and anti-leptons may induce baryon asymmetry through baryon number violation processes (lepto-genesis)
  - long baseline neutrino experiments will investigate CP violation in neutrino sector
- Alternatively we look for sources of CP violation beyond the SM in the quark sector
- Promising place to look for non-SM *CP* violation is the neutral *B*<sub>s</sub> meson system



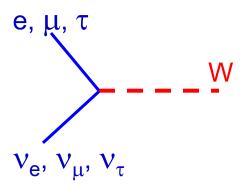
#### CP Violation in the Standard Model

- *CP* violation enters the Standard Model through complex phases in mixing matrices that connect up-type fermions with down-type fermions via W bosons:



$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

- Cabibbo-Kobayashi-Maskawa (CKM) quark mixing matrix transforms quark mass eigenstates into weak eigenstates and induces *CP* violation in the hadronic sector



Pontecorvo-Maki-Nakagawa-Sakata
 (PMNS) neutrino mixing matrix
 → induces neutrino oscillations and

$$\begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix}$$

possibly *CP* violation in lepton sector

#### **CKM Matrix**

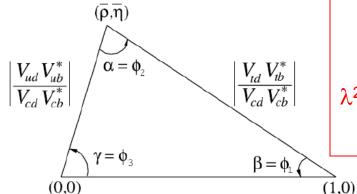
- Expand CKM matrix in  $\lambda = V_{us} = sin(\theta_{Cabibbo}) \approx 0.23$ 

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \approx \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + \frac{1}{2}A^2\lambda^5[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4(1 + 4A^2) & A\lambda^2 \\ A\lambda^3[1 - (1 - \frac{1}{2}\lambda^2)(\rho + i\eta)] & -A\lambda^2 + \frac{1}{2}A\lambda^4[1 - 2(\rho + i\eta)] & 1 - \frac{1}{2}A^2\lambda^4 \end{pmatrix}$$

- To conserve probability CKM matrix must be unitary
  - → Unitary relations can be represented as "unitarity triangles"

unitarity 
$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$
 relations:

unitarity triangles:



$$V_{td}V_{tb}^{*} = 0$$

$$V_{us}V_{ub}^{*} + V_{cs}V_{cb}^{*} + V_{ts}V_{tb}^{*} = 0$$

$$\frac{\left|\frac{V_{ts}V_{tb}^{*}}{V_{cs}V_{cb}^{*}}\right|}{\left|\frac{V_{ts}V_{tb}^{*}}{V_{cs}V_{cb}^{*}}\right|} \sim 1$$

$$\lambda^{2} \sim \left|\frac{V_{us}V_{ub}^{*}}{V_{cs}V_{cb}^{*}}\right| (0,0) = 1$$

$$\beta_{s} (1,0)$$

Small *CP* violation phase  $\beta_s$  accessible in  $B_s \rightarrow J/\psi \Phi$  decays



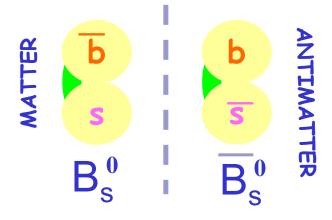
## Neutral B<sub>s</sub> System

- Time evolution of  $B_s$  flavor eigenstates described by Schrodinger equation:

$$i\frac{d}{dt} \begin{pmatrix} |B_s^0(t)\rangle \\ |\bar{B}_s^0(t)\rangle \end{pmatrix} = \left(\mathbf{M} - \frac{i}{2}\mathbf{\Gamma}\right) \begin{pmatrix} |B_s^0(t)\rangle \\ |\bar{B}_s^0(t)\rangle \end{pmatrix}$$

-Diagonalize mass (M) and decay ( $\Gamma$ ) matrices  $\rightarrow$  mass eigenstates :

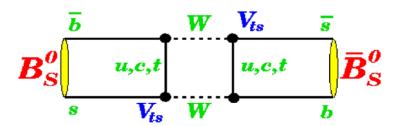
$$|B_s^H\rangle = p\,|B_s^0\rangle - q\,|\bar{B}_s^0\rangle \qquad |B_s^L\rangle = p\,|B_s^0\rangle + q\,|\bar{B}_s^0\rangle$$



- Flavor eigenstates differ from mass eigenstates and mass eigenvalues are also different:

$$\Delta m_s = m_H - m_L \approx 2|M_{12}|$$
 $\rightarrow B_s$  oscillates with frequency  $\Delta m_s$ 
precisely measured by

CDF  $\Delta m_s = 17.77 +/- 0.12 ps^{-1}$ 
DØ  $\Delta m_s = 18.56 +/- 0.87 ps^{-1}$ 

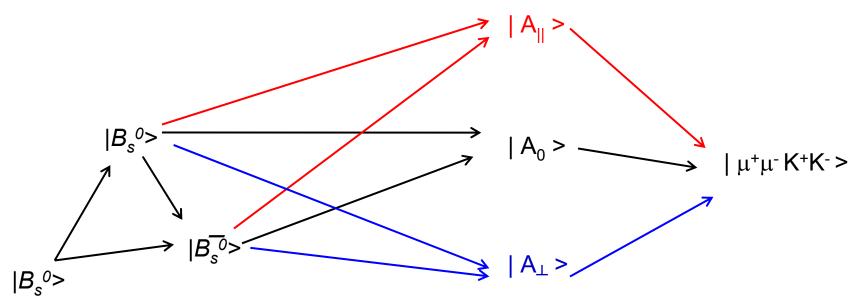


- Mass eigenstates have different decay widths

$$\Delta \Gamma = \Gamma_L - \Gamma_H \approx 2|\Gamma_{12}|\cos(\Phi_s)$$
 where  $\phi_s^{SM} = \arg\left(-\frac{M_{12}}{\Gamma_{12}}\right) \approx 4 \times 10^{-3}$ 

## Transversity Basis

- Use "transversity basis" in which the vector meson polarizations w.r.t. direction of motion are either (Phys. Lett. B 369, 144 (1996), 184 hep-ph/9511363 ):
  - transverse ( $\perp$  perpendicular to each other)  $\rightarrow$  *CP odd*
  - transverse ( parallel to each other) → CP even
  - longitudinal (0)  $\rightarrow$  *CP* even
- Corresponding decay amplitudes:  $A_{\it 0}$ ,  $A_{\parallel}$ ,  $A_{\perp}$



## Decay Rate

-  $B_s \rightarrow J/\Psi\Phi$  decay rate as function of time, decay angles and initial  $B_s$  flavor:

$$\frac{d^4P(t,\vec{\rho})}{dtd\vec{\rho}} \propto |A_0|^2 \mathcal{T}_+ f_1(\vec{\rho}) + |A_{||}|^2 \mathcal{T}_+ f_2(\vec{\rho}) \qquad \text{time dependence terms}$$
 
$$+ |A_\perp|^2 \mathcal{T}_- f_3(\vec{\rho}) + |A_{||}||A_\perp| \mathcal{U}_+ f_4(\vec{\rho}) \qquad \text{angular dependence terms}$$
 
$$+ |A_0||A_{||}|\cos(\delta_{||}) \mathcal{T}_+ f_5(\vec{\rho})$$
 
$$+ |A_0||A_\perp| \mathcal{V}_+ f_6(\vec{\rho}), \qquad \text{terms with } \beta_{\rm s} \text{ dependence}$$

$$\mathcal{T}_{\pm} = e^{-\Gamma t} \times \left[ \cosh(\Delta \Gamma t/2) \mp \cos(2\beta_s) \sinh(\Delta \Gamma t/2) + \sin(2\beta_s) \sin(\Delta m_s t) \right];$$

$$\mathcal{U}_{\pm} = \pm e^{-\Gamma t} \times \left[ \sin(\delta_{\perp} - \delta_{\parallel}) \cos(\Delta m_{s} t) \right] - \cos(\delta_{\perp} - \delta_{\parallel}) \cos(2\beta_{s}) \sin(\Delta m_{s} t)$$

$$\pm \cos(\delta_{\perp} - \delta_{\parallel}) \sin(2\beta_{s}) \sinh(\Delta \Gamma t / 2) \right]$$

$$\mathcal{V}_{\pm} = \pm e^{-\Gamma t} \times \left[ \sin(\delta_{\perp}) \cos(\Delta m_s t) - \cos(\delta_{\perp}) \cos(2\beta_s) \sin(\Delta m_s t) \right]$$

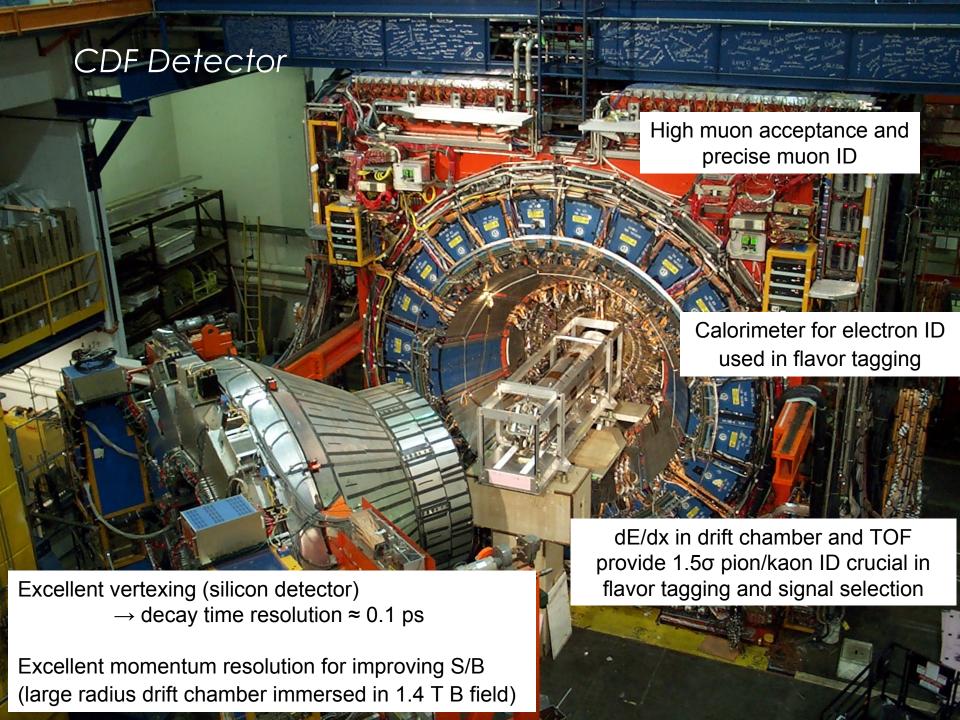
$$\pm \cos(\delta_{\perp}) \sin(2\beta_s) \sinh(\Delta \Gamma t/2) ].$$

terms with  $\Delta m_s$  dependence present if initial state of B meson (B vs anti-B) is determined (flavor tagged)

'strong' phases:

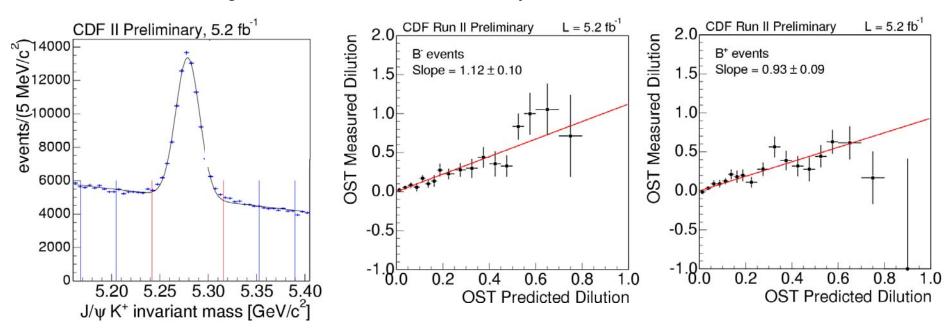
$$\delta_{\parallel} \equiv \operatorname{Arg}(A_{\parallel}(0)A_{0}^{*}(0))$$
  
$$\delta_{\perp} \equiv \operatorname{Arg}(A_{\perp}(0)A_{0}^{*}(0))$$

- Identification of B flavor at production (flavor tagging)  $\rightarrow$  better sensitivity to  $\beta_s$ 



# Opposite Side Tagging Calibration and Performance

- OST combines in a NN opposite side lepton and jet charge information
- Initially calibrated using a sample of inclusive semileptonic B decays
  - predicts tagging probability on event-by-event basis
- Re-calibrated using  $\approx$  52,000  $B^{+/-} \rightarrow J/\Psi K^{+/-}$  decays



- OST efficiency = 94.2 +/- 0.4%, OST dilution = 11.5 +/- 0.2 % (correct tag probability ~56%)
- Total tagging power = 1.2%

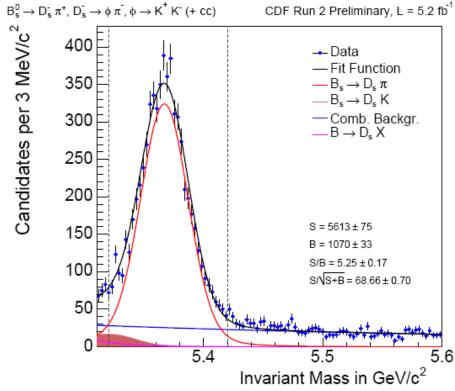
## Same Side Tagging Calibration

- Event-by-event predicted dilution based on simulation
- Calibrated with 5.2 fb<sup>-1</sup> of data
- Simultaneously measuring the  $B_s$  mixing frequency  $\Delta m_s$  and the dilution scale factor A

$$P_{Sig}(ct|\sigma_{ct}, \xi = \xi_D \cdot \xi_P, D) = \frac{1}{N} \cdot \left[ \frac{1}{\tau} e^{-\tilde{t}/\tau} \cdot (1 + \xi \mathcal{A}D \cdot \cos(\Delta m_s \tilde{t})) \right] \otimes \mathcal{G}(c\tilde{t}|\sigma_{ct}) \cdot \epsilon(ct|\sigma_{ct})$$

- *D* event by event predicted dilution
- $\frac{-\xi}{B_s}$  and un-tagged events
- Fully reconstructed  $B_s$  decays selected by displaced track trigger

Decay Channel	S
$B_s^0 \to D_s^- \pi^+, \ D_s^- \to \phi \pi^-$	$5613 \pm 75$
$B_s^0 \to D_s^- \pi^+, \ D_s^- \to K^* K^-$	$2761 \pm 53$
$B_s^0 \to D_s^- \pi^+, \ D_s^- \to (3\pi)^-$	$2652 \pm 52$
$B_s^0 \to D_s^-(3\pi)^+, \ D_s^- \to \phi\pi^-$	$1852 \pm 43$
Sum	$\boxed{12877 \pm 113}$





- $B_s$  oscillation frequency measured  $\Delta m_s = (17.79 \pm 0.07) \; \mathrm{ps}^{-1}$  (statistical error only)
- In good agreement with the published CDF measurement with 1 fb<sup>-1</sup> PRL 97, 242003 2006, PRL 97, 062003 2006

$$\Delta m_s = 17.77 \pm 0.10 \; (\mathrm{stat}) \pm 0.07 \; (\mathrm{syst}) \; \mathrm{ps}^{-1}$$
 used as external constraint in  $\beta_s$  measurement

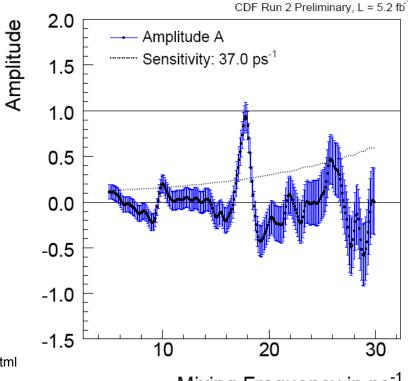
 Dilution scale factor (amplitude) in good agreement with 1:

$$A = 0.94 \pm 0.15 \ (stat.) \pm 0.13 \ (syst.)$$

- Largest systematic uncertainty from decay time resolution modeling
- Total SSKT tagging power:

$$\varepsilon \mathcal{A}^2 D^2 = (3.2 \pm 1.4) \%$$

http://www-cdf.fnal.gov/physics/new/bottom/100204.blessed-sskt-calibration/index.html CDF public note 10108



# ٧

#### S-Wave

- As noted in arxiv:0812.2832v3, the KK pair in  $B_s \to J/\Psi$  KK decays can be in an s-wave state with ~6% contribution in a +/-10 MeV window around the  $\Phi$  peak
- Systematic effects from neglecting such contribution were first investigated by Clarke et al in arxiv:0908.3627v1 where it is shown that:
  - 10% un-accounted s-wave contamination in the  $\phi$  region leads to
    - 10% bias in the measured  $2\beta_s$ , towards the SM prediction
    - 15% increase in statistical errors
- S-wave contribution can be either non-resonant or from the  $f^0(980)$  resonance
- To account for potential s-wave contribution, enhance the likelihood function to account for the s-wave amplitude  $A_s$  and interference between s-wave and p-wave
- Time dependence of the s-wave amplitude  $A_S$  is *CP-odd*, same as  $A_{\perp}$
- Mass and phase of s-wave component are assumed flat (good approximation in a narrow +/- 10 MeV around the  $\phi$  mass)

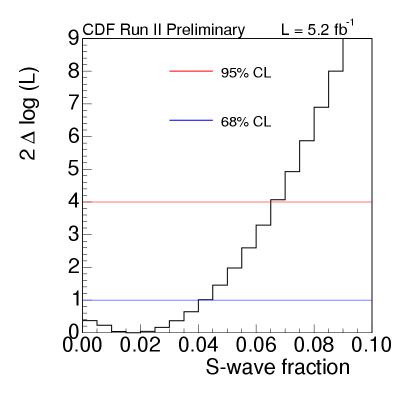


#### S-Wave Measurement

 The fitted s-wave fraction is found to be very small in the KK mass range used in this analysis: [1.009, 1.028] GeV

s-wave fraction < 6.7% at 95% C.L.

- To be compared with expectation from <a href="mailto:arxiv:0812.2832v3">arxiv:0812.2832v3</a> of 6.3% s-wave contribution in a range of +/- 10 MeV around the Φ peak

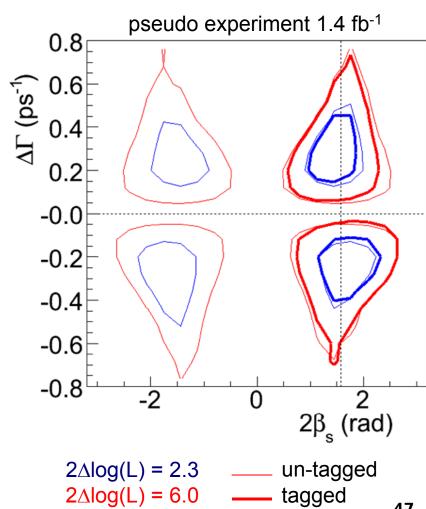


# CP Violation Phase $\beta_s$ in Tagged B<sub>s</sub> $\rightarrow$ J/ $\Psi\Phi$ Decays

- Without the s-wave the likelihood function is symmetric under the transformation

$$2\beta_s \to \pi - 2\beta_s$$
  $\Delta\Gamma \to -\Delta\Gamma$   
 $\delta_{\parallel} \to 2\pi - \delta_{\parallel}$   $\delta_{\perp} \to \pi - \delta_{\perp}$ 

- Study expected effect of tagging using pseudo-experiments
- Improvement of parameter resolution is small due to limited tagging power ( $\varepsilon D^2 \sim 4.5\%$ compared to B factories ~30%)
- However,  $\beta_s \rightarrow -\beta_s$  no longer a symmetry  $\rightarrow$  4-fold ambiguity reduced to 2-fold ambiguity
- Adding the s-wave "slightly" breaks the symmetry due to asymmetric Φ mass shape
- Symmetry still valid with good approximation...



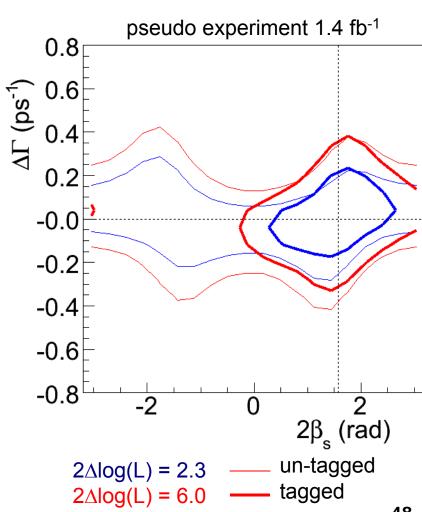
$$\frac{\log(L) = 2.3}{\log(L) = 6.0} \quad \text{tagged}$$

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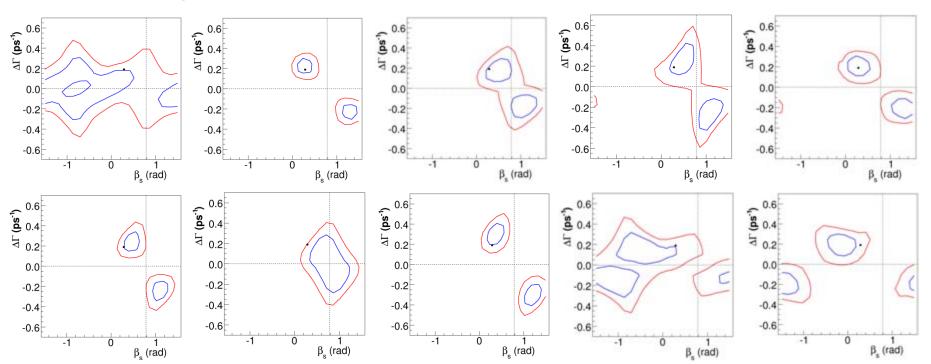
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## Cross Checks With Pseudo-Experiments

- Generate 10 pseudo-experiments with  $\beta_s$  = 0.3 and  $\Delta\Gamma$  = 0.2 corresponding to 1.4 fb<sup>-1</sup> same parameters, just different random seeds
- Large fluctuations expected in shape and size of confidence regions

$$--- 2\Delta \log(L) = 2.3$$
  $-- 2\Delta \log(L) = 6.0$ 





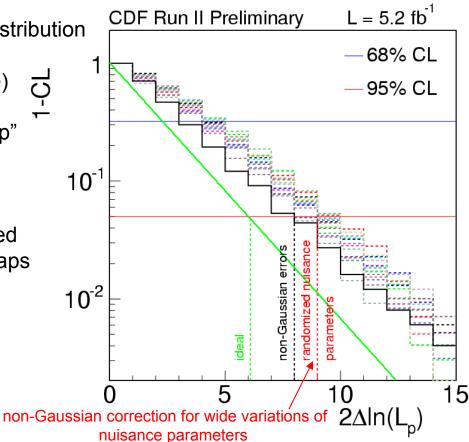
- Pseudo-experiments show that we are still not in perfect Gaussian regime
   → quote confidence regions instead of point estimates
- In ideal case (high statistics, Gaussian likelihood), to get the 2D 68% (95%) C.L. regions, take a slice through profiled likelihood at 2.3 (6.0) units up from minimum

- In this analysis integrated likelihood ratio distribution (black histogram) deviates from the ideal  $\chi^2$  distribution (green continuous curve)

 Using pseudo-experiments establish a "map" between Confidence Level and 2∆log(L)

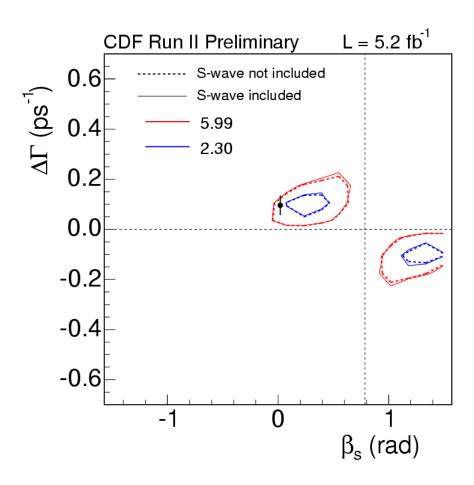
- All nuisance parameters are randomly varied within +/-  $5\sigma$  from their best fit values and maps of CL vs  $2\Delta log(L)$  re-derived

- To establish final confidence regions use most conservative case



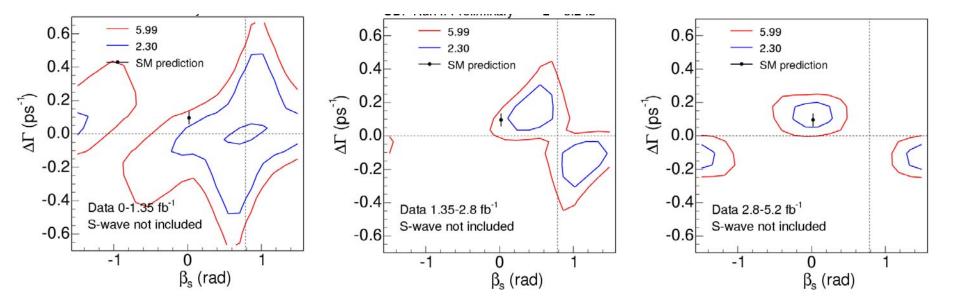
## $\beta_s$ - $\Delta\Gamma$ Contours with and without Including the S-Wave

- Compare likelihood contours with and without including the s-wave
- Very small effect on  $\beta_{\rm s}$  and  $\varDelta \Gamma$





- Divide 5.2 fb<sup>-1</sup> sample in three sub-samples corresponding to three public releases:
  - 0 1.4 fb<sup>-1</sup> (initial result released at the end of 2007, PRL 100, 161802 (2008), arXiv:0712.2397)
  - 1.4 2.8 fb<sup>-1</sup> (added for 2008 ICHEP update)
  - 2.8 5.2 fb<sup>-1</sup> (added for this update)
- Previous results reproduced with updated analysis
- Clearly, improved agreement with the SM expectation comes from the second half of data (2.8 5.2 fb<sup>-1</sup>)



## Comparison with Previous Results

- $-\beta_s$  and  $\Delta\Gamma_s$  allowed parameter space greatly reduced
- Agreement with SM expectation improves with higher statistics

