# CP Violation and the Determination of the CKM Matrix Frank Porter (Caltech, BABAR)

- Cabibbo-Kobayashi-Maskawa (CKM) matrix "V"
  - Fundamental in Standard Model (SM)
  - Four parameters  $(\theta_{12}, \theta_{13}, \theta_{23}, \phi \leftrightarrow A, \lambda, \rho, \eta)$
  - Source of CP violation in SM
  - Testing the SM V is unitary  $3 \times 3$  matrix in SM
    - Additional generations can make non-unitary
    - Can test unitarity relations with measurements of magnitudes and/or phases
- New physics can show up in loops, often at same order as SM graphs
  - Look for differences among quantities that should be the same in SM, or for deviations from SM predictions

**Scope**, with apologies for the many topics left out

- Heavy flavors  $(s,c,b,t,\tau)$
- Nothing on EDM
- For neutrino sector (PMNS matrix), see talks by Lisi, Bellerive, Nakaya, and Piquemal
- For  $\beta_s$ , like sign di-muon asymmetry, see Borissov's talk [Also Belle (Wicht, 1204)]
- Not much discussion beyond the SM (but an underlying theme)
- For theory, see talk by Isidori (Lattice Kuramashi)
- Omit CPT tests (see Lusiani, 1173; Kundu, 270)
- Omit future prospects
- **CKM** magnitudes of elements
- CKM CP violation

The Cabibbo-Kobayashi-Maskawa mixing matrix

Relates quark mass eigenstates to weak eigenstates.

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + O(\lambda^4)$$

(Wolfenstein parameterization)

Often define  $\bar{\rho} \equiv \rho(1 - \lambda^2/2), \ \bar{\eta} \equiv \eta(1 - \lambda^2/2)$ 

Magnitudes

Phases (i.e., "angles of unitarity triangles")

Determinations assume standard model, but not using unitarity. Inconsistencies could be signs of new physics.



The magnitudes:  $|V_{ud}|$ 



2008 RPP

 $|V_{ud}| = 0.97418 \pm 0.00027$ 

Best determinations in superallowed  $0^+ \rightarrow 0^+$  nuclear  $\beta$  decays. Recent analysis from Hardy and Towner PRC **79** (2009) 055502 yields:

 $|V_{ud}| = 0.97425 \pm 0.00022$ 

# The magnitudes: $|V_{us}|$ $\begin{pmatrix} V_{us} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$

2008 RPP

 $|V_{us}| = 0.2255 \pm 0.0019$  $|V_{us}|$  from kaon decays

New averages from FlaviaNet Kaon Working Group, arXiv:1005.2323 [hep-ph] (2010), see also KLOE (Archilli, 1085)



- $K_{\ell 3}$ :  $|V_{us}|f_{+}(0) = 0.2163(5)$  or  $|V_{us}| = 0.2254 \pm 0.0013$  with  $f_{+}(0) = 0.959(5)$  (lattice, Boyle et al., arXiv1004:0886 (2010))
- $K_{\ell 2}: \frac{|V_{us}|}{|V_{ud}|} \frac{f_K}{f_{\pi}} = 0.2758(5) \text{ or } \frac{|V_{us}|}{|V_{ud}|} = 0.2312 \pm 0.0013$ with  $f_K/f_{\pi} = 1.193(6)$  (lattice average)
- Combining, obtain  $|V_{us}|(K) = 0.02253 \pm 0.0009$

#### $|V_{us}|$ from tau decays

BABAR (Lusiani, 1173) Measure in exclusive  $\tau$  decays with 467 fb<sup>-1</sup>

$$R_{K/\pi} \equiv \frac{\mathcal{B}(\tau^- \to K^- \nu_\tau)}{\mathcal{B}(\tau^- \to \pi^- \nu_\tau)}$$
  
= 0.06531 ± 0.00056 ± 0.00093  
$$= \frac{f_K^2 |V_{us}|^2 \left(1 - \frac{m_K^2}{m_\tau^2}\right)^2}{f_\pi^2 |V_{ud}|^2 \left(1 - \frac{m_\pi^2}{m_\tau^2}\right)^2} (1 - \delta_{LD})$$
  
$$M_{T \to \pi} \mathcal{V} \xrightarrow{P_1 \text{-prong}} \mathcal{V}_{T \to \pi} \mathcal{V}_{\tau}$$
  
$$= \frac{f_K^2 |V_{us}|^2 \left(1 - \frac{m_K^2}{m_\tau^2}\right)^2}{(1 - \frac{m_\pi^2}{m_\tau^2})^2} (1 - \delta_{LD})$$
  
$$M_{T \to \pi} \mathcal{V} \xrightarrow{P_1 \text{-prong}} \mathcal{V}_{T \to \pi} \mathcal{V}_{\tau}$$
  
$$= \frac{f_K^2 |V_{us}|^2 \left(1 - \frac{m_\pi^2}{m_\tau^2}\right)^2}{(1 - \frac{m_\pi^2}{m_\tau^2})^2} (1 - \delta_{LD})$$
  
$$M_{T \to \pi} \mathcal{V} \xrightarrow{P_1 \text{-prong}} \mathcal{V}_{T \to \pi} \mathcal{V}_{\tau}$$
  
$$= \frac{f_K^2 |V_{us}|^2 \left(1 - \frac{m_\pi^2}{m_\tau^2}\right)^2}{(1 - \delta_{LD})}$$
  
$$M_{T \to \pi} \mathcal{V} \xrightarrow{P_1 \text{-prong}} \mathcal{V}_{\tau}$$
  
$$= \frac{f_K^2 |V_{us}|^2 \left(1 - \frac{m_\pi^2}{m_\tau^2}\right)^2}{(1 - \delta_{LD})}$$

- Approach avoids absolute strange decay constant  $(f_K^2)$ , replacing with ratio to pion. Use  $f_K/f_{\pi} = 1.189 \pm 0.007$  and  $\delta_{LD} = 0.0003 \pm 0.0044$ - Result is:  $|V_{us}| = 0.2255 \pm 0.0024$
- $\tau \to s$  inclusive
- At ICHEP08,  $3.2\sigma$  discrepancy:  $|V_{us}| = 0.2159 \pm 0.0030$
- 2010 preliminary evaluation (Lusiani, 1173)  $|V_{us}| = 0.2165 \pm 0.0023$
- Discrepancy =  $3.6\sigma$

[see also BABAR,  $\Lambda_c$  decays (Hartmann, 557)]

# $|V_{us}|$ summary

## Lusiani 1173 (HFAG- $\tau$ ) compilation, Preliminary



My average  $|V_{us}| = 0.2253 \pm 0.0008$ , does not include  $\tau \rightarrow s$  inclusive

# The magnitudes: $|V_{ub}|$

 $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$ 

2008 RPP  $|V_{ub}| = 0.00393 \pm 0.00036$ , combined exclusive and inclusive (dominant)

- Inclusive semileptonic decays  $B \to X \ell \nu$  where  $X = X_u$ 
  - Select B decays by reconstructing recoil B, either fully or partially
  - Huge background from  $b \to c$  transitions  $(X = X_c)$
  - Can restrict kinematic region, e.g., to  $m_X < m_D$
  - Can use  $MM^2$  to preferentially select single missing  $\nu$  (and low multiplicity)
  - Use theory to extrapolate from restricted kinematic region to full phase space
    - BLNPPRD 72 (2005) 073006DGEarXiv:0806.4524 [hep-ph]GGOUJHEP 0710 (2007) 058ADFREur Phys J C 59 (2009) 831(and references therein)

– Belle inclusive (PRL 104 (2010) 021801) on full sample 657M  $B\bar{B}$ :



TABLE II. Values for  $|V_{ub}|$  with relative errors (in %).

Theory	$ V_{ub}  \times 10^3$	Stat	Syst	$m_b$	Th.
BLNP [5]	4.37	4.3	4.0	+3.1 -2.7	+4.3 -4.0
DGE [6]	4.46	4.3	4.0	$+3.2 \\ -3.3$	$^{+1.0}_{-1.5}$
GGOU [7]	4.41	4.3	4.0	1.9	$^{+2.1}_{-4.5}$

My average Belle inclusive:  $0.00441 \pm 0.00026(\text{expt}) \pm 0.00024(\text{thy})$ 

# Inclusive $|V_{ub}|$ (continued)



– BABAR inclusive (Sigamani, 732):

Measure Partial Branching Fractions for  $B \to X_u \ell \bar{\nu}$ 

*B* tag is via exclusive reconstruction of recoil *B* in  $B \to \overline{D}^{(*)}h$ , where  $h = \pi$  or h = KFor  $p_{\ell}^* > 1.0$  GeV, with a 2-D fit to  $(M_X, q^2)$ , and averaging (consistent) results according to (BLNP, DGE, GGOU, ADFR), obtain

 $|V_{ub}| = 0.00431 \pm 0.00035$  (preliminary)

Background-subtracted lepton momentum distribution in  $B \to X_u \ell \bar{\nu}$  decays



# $|V_{ub}|$ in exclusive semileptonic decays

Exclusive semileptonic decays to light quark states

- Constraints reduce background, but also lower statistics
- Theory for form factors

E.g., for  $B \to \pi \ell \nu$  with  $\ell = e$  or  $\mu$ , to good approximation a single form factor contributes:

$$\frac{d\Gamma(B^0 \to \pi^- \ell^+ \nu)}{dq^2 d\cos\theta_{W\ell}} = |V_{ub}|^2 \frac{G_F^2 p_\pi^3}{32\pi^3} \sin^2\theta_{W\ell} |f_+(q^2)|^2.$$

- Belle (Ha, 944) Exclusive  $B^0 \to \pi^- \ell^+ \nu$ , untagged 605 fb<sup>-1</sup>  $\mathcal{B}(B^0 \to \pi^- \ell \nu) =$ (1.49±0.04(stat)±0.07(syst))×10<sup>-4</sup>  $|V_{ub}f_+(0)| = (9.24\pm0.18(stat)\pm0.20(syst)\pm0.07(\tau_B))$ ×10<sup>-4</sup>  $|V_{ub}| = (0.00343\pm0.00033)$  (using FNAL-MILC PRD **79** (2009) 054507)

- **BABAR** (Wulsin, 1180) 
$$B \to \pi \ell \nu \ (\rho \ell \nu)$$

TABLE XIII:  $|V_{ub}|$  derived from  $B \to \pi \ell \nu$  and  $B \to \rho \ell \nu$ decays for various  $q^2$  regions and form-factor calculations. Quoted errors are experimental uncertainties and theoretical uncertainties of the form-factor integral  $\Delta \zeta$ . (Uncertainties for the  $B \to \rho \ell \nu$  form-factor integrals are not available.)

$\mathcal{B}(B^0 \to \pi^- \ell^+ \nu) = (1.41 \pm 0.05 \pm 0.07) \times 10^{-4}$	
$\mathcal{B}(B^0 \to \rho^- \ell^+ \nu) = (1.75 \pm 0.15 \pm 0.27) \times 10^{-4}$	

	$q^2$ Range (GeV <sup>2</sup> )	$\frac{\Delta \zeta}{(\mathrm{ps}^{-1})}$	$ V_{ub} $ (10 <sup>-3</sup> )
$B \to \pi \ell \nu$			
LCSR [15]	0 - 16	$5.44{\pm}1.43$	$3.63 \pm 0.12^{+0.59}_{-0.40}$
HPQCD [22]	16 - 26.4	$2.02{\pm}0.55$	$3.21\pm0.17^{+0.55}_{-0.36}$
$B \to \rho \ell \nu$			
LCSR [16]	0 - 16.0	13.79	$2.75\pm0.24$
ISGW2 [14]	0 - 20.3	14.20	$2.83\pm0.24$

For  $B \to \pi \ell \nu$  and simult. fit to FNAL/MILC lattice,  $|V_{ub}| = 0.00295 \pm 0.00031$ 

My average for BABAR  $\pi \ell \nu$ , including error for spread:  $|V_{ub}| = 0.00326 \pm 0.00054$ 



# $V_{ub}$ in leptonic *B* decays

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

For  $Q_q^+ = \pi^+$ ,  $K^+$ ,  $D^+$ ,  $D_s^+$ ,  $B^+$ , with  $V_{(Qq)} = V_{Qq}$  or  $V_{qQ}$  as appropriate:

$$\Gamma(Q_q^+ \to \ell^+ \nu_\ell) = \frac{G_F^2}{8\pi} m_{Q_q}^3 \left(\frac{m_\ell}{m_{Q_q}}\right)^2 \left(1 - \frac{m_\ell^2}{m_{Q_q}^2}\right)^2 |V_{(Qq)}|^2 f_{Q_q}^2,$$

Belle 711 fb<sup>-1</sup> (Stypuła, 1097)  $B \to \tau \nu$  (and  $B \to D^* \tau \nu$ ); exclusive semileptonic tag measure  $f_B |V_{ub}| = (9.3^{+1.2}_{-1.1} \pm 0.9) \times 10^{-4} \text{ GeV}$ , from  $\mathcal{B}(B^- \to \tau^- \bar{\nu}_{\tau}) = (1.54^{+0.38}_{-0.37}(\text{stat})^{+0.29}_{-0.31}(\text{syst})) \times$   $10^{-4}$  (significance 3.6 $\sigma$ ) Gives  $|V_{ub}| = 0.00489 \pm$  0.00079 for  $f_B = 0.19$  GeV  $E_{\text{ECL}} = \text{residual energy in calorimeter}$ 



BABAR (De Nardo, 581)  $\mathcal{B}(B^- \to \tau^- \bar{\nu}_{\tau}) = (1.80^{+0.57}_{-0.54}(\text{stat}) \pm 0.26(\text{syst})) \times 10^{-4},$ significance  $3.6\sigma$ 

Combine with semileptonic tags:  $\mathcal{B}(B^- \to \tau^- \bar{\nu}_{\tau}) = (1.76 \pm 0.49) \times 10^{-4}$ 



Measurement	Experiment	$V_{ub}$
Inclusive	Belle	$0.00441 \pm 0.00024$
Inclusive	BABAR	$0.00431 \pm 0.00035$
Exclusive $\pi \ell \nu$	Belle	$0.00343 \pm 0.00033$
Exclusive $\pi \ell \nu$	BABAR	$0.00326 \pm 0.00054$
$B \to \tau \nu$	Belle	$0.00484 \pm 0.00079$
$B \to \tau \nu$	BABAR	$0.0057 \pm 0.0019$

#### Recent measurements

Longstanding inclusive/exclusive discrepancy remains. For example, comparing Belle inclusive with Belle exclusive the difference is  $2.3\sigma$ 

CKMfitter average  $|V_{ub}| = 0.00392 \pm 0.00009 \pm 0.00045$  (based on HFAG end of 2009 preliminary)

## First row unitarity

In SM (V is 3 × 3 unitary), must have:  

$$1 = |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2$$

$$= 0.99995 \pm 0.00057$$

Limit (Bayesian) on possible 4th generation:

$$|V_{u4}| = \sqrt{1 - |V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2}$$
  
< 0.031 (90% CL, flat prior in  $|V_{u4}|^2$ )  
< 0.061 (90% CL, flat prior in  $|V_{u4}|$ )

In spite of "four-nines" sum, numbers from first two generations not sufficiently precise to require the third generation

# The magnitudes: $|V_{cd}|$

 $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$ 

2008 RPP remains up-to-date  $|V_{cd}| = 0.230 \pm 0.011$ 

- From neutrino charm production (di-muons/single muons, CDHS, CCFR, CHARM II + CHORUS)
- Prospects for leptonic and semileptonic D (and  $D_s$  for  $|V_{cs}|$ ) to contribute, once theoretical uncertainties in decay constants and form factors are reduced further. (see also Melikhov 254)

The magnitudes: 
$$|V_{cs}|$$
  $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$ 

RPP 2008: Leptonic  $D_s$  decays; semileptonic D decays  $|V_{cs}| = 1.04 \pm 0.06$ 

# The magnitudes: $|V_{cb}|$ $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$

2008 RPP  $|V_{cb}| = 0.0412 \pm 0.0011$  (combined exclusive and inclusive)

#### New results in exclusive $B \to \text{charm}$

– Belle (Dungel, 943) New result for  $B^0 \to D^{*-} \ell^+ \nu$ , signal side reconstructed, 711 fb<sup>-1</sup>

$$\mathcal{F}(1)|V_{cb}| = 0.0345 \pm 0.0002 \pm 0.0010$$

 $\mathcal{F}(1)$  is the hadronic form factor at zero recoil ( $w = v_B \cdot v_D^* = 1$ ) Use HQET (Caprini, Lellouch, Neubert NPB **530** (1998) 153) for w-dependence of form factor. Lattice QCD (Bernard et al., PRD **79** (2009) 014506):  $\mathcal{F}(1) = 0.921 \pm 0.013 \pm 0.020$ 

 $|V_{cb}| = 0.0375 \pm 0.0015$ 

− **BABAR** (Petrella, 1179) [PRL **104** (2010) 011802]  $B \rightarrow D\ell\nu$ , fully reconstructed tags (average of charged and neutral D modes)

 $G(1)|V_{cb}| = 0.0423 \pm 0.0019 \pm 0.0014$ 

G(1) is the hadronic form factor at zero recoil  $(w = v_B \cdot v_D = 1)$ 

 $V_{cb} = 0.0392 \pm 0.0018 \pm 0.0013 \pm 0.0009$  (lattice)

Lattice form factor: Okamoto et al., NucPhysB 140 (2005) 461

# The magnitudes: $|V_{cb}|$ $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$

- New results in inclusive  $B \rightarrow \text{charm}$ 
  - BABAR (Petrella, 1179) [PRD 81 (2010) 032003] Measurement and Interpretation of Moments in Inclusive Decays B → X<sub>c</sub>ℓν Rates and Moments analysis of inclusive B → X<sub>c</sub>ℓν, based on (OPE) Benson, Bigi, Mannel, Uraltsev, NP B665 (2003) 367

#### $|V_{cb}| = 0.04205 \pm 0.0045 \pm 0.0070$

As with  $|V_{ub}|$  the inclusive results tend to be higher than the exclusive results CKMfitter average  $|V_{cb}| = 0.04089 \pm 0.00038 \pm 0.00059$  (based on HFAG end of 2009 preliminary)

# The magnitudes: $|V_{td}|$



#### $2008 \ \mathrm{RPP}$

 $|V_{td}| = 0.0081 \pm 0.0006$ 

#### $|V_{td}|$ from B mixing

- Uncertainty dominated by lattice QCD uncertainties.
- Some uncertainty cancels in ratio  $|V_{td}/V_{ts}|$ , measured using B and  $B_s$  mixing:

 $|V_{td}/V_{ts}| = 0.209 \pm 0.001 \pm 0.006 \ (2008 \ \text{RPP})$ 

- Using this, and  $|V_{ts}|$  obtain slightly more precise result:  $V_{td} = 0.0081 \pm 0.0005$
- **BABAR** (Bard, 1177) Another approach: Measure  $|V_{td}/V_{ts}|$  in "inclusive" ratio of radiative *B* decays related by  $d \leftrightarrow s$  with 471M  $B\bar{B}$ 
  - Penguin decays, so possible NP in loop, hence tests SM in comparison with other determination
  - For example, compare  $B^0 \to \pi^+ \pi^- \gamma$  with  $B^0 \to K^+ \pi^- \gamma$ . Analysis uses 7 such pairs of modes.
  - Result is

$$\frac{\mathcal{B}(b \to d\gamma)}{\mathcal{B}(b \to s\gamma)} = 0.033 \pm 0.009 \pm 0.003$$

from which we obtain (using (NLO) Ali, Asatrian Greub PLB **429** (1998) 87):

 $|V_{td}/V_{ts}| = 0.199 \pm 0.022(\text{stat}) \pm 0.024(\text{syst}) \pm 0.002(\text{thy})$ 

The magnitudes:  $|V_{ts}|$ 

 $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$ 

 $|V_{ts}|$  from  $B_s$  mixing 2008 RPP

 $|V_{ts}| = 0.0387 \pm 0.0023$ 

Dominant uncertainties from lattice QCD

The magnitudes:  $|V_{tb}|$   $\begin{pmatrix} v_{ud} & v_{us} & v_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$ 

2008 RPP  $|V_{tb}| > 0.74 \ 90\% \ \text{CL}, \ \sigma(p\bar{p} \to tX)$ 

 $|V_{tb}| = 0.77^{+0.18}_{-0.24}$  EW fit, top loops in  $Z \rightarrow b\bar{b}$ 

Can be measured in single top production, without assuming 3 generation unitarity (but assuming  $|V_{tb}| \gg |V_{td}|, |V_{ts}|$ )



 $- \frac{\text{CDF}/\text{D0}}{\text{fb}^{-1}} (\text{Quinn}, 1132) \text{ arXiv:} /0908.2171 \text{ [hep-ex] Combined CDF}(3.2 \text{ fb}^{-1}) \& \text{D0}(2.3 \text{ fb}^{-1}) |V_{tb}| = 0.88 \pm 0.07$ 

# CP violation, the unitarity triangles $\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{tc} & V_{tc} \end{pmatrix}$



All CP violation from CKM in SM Manifests as "unitarity triangle" relations with area  $\neq 0$  $VV^{\dagger} = V^{\dagger}V = 1$ 

Yields six distinct relations from the off-diagonal components. Two of these are under active investigation:

$$0 = V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = O(\lambda^3) + O(\lambda^3) + O(\lambda^3)$$



$$0 = V_{us}V_{ub}^* + V_{cs}V_{cb}^* + V_{ts}V_{tb}^* = O(\lambda^4) + O(\lambda^2) + O(\lambda^2)$$

$$V_{us}V_{ub}^{*}\frac{\alpha_{s}}{\gamma_{s}} \quad \frac{V_{ts}V_{tb}^{*}}{V_{cs}V_{cb}^{*}}\beta_{s}$$

# The angles: $\beta/\phi_1$

RPP 2008  $\sin 2\beta = 0.681 \pm 0.025, b \rightarrow c\bar{c}s$  decays to CP eigenstates

- Belle (Higuchi, 1094) Analysis of  $\sin 2\phi_1$  in  $B \to c\bar{c}K^0$  [ie, the "golden modes"] on final data sample of 772M  $B\bar{B}$ , in progress; expected error  $\delta(\sin 2\phi_1) \approx 0.024$ .
- **BABAR** (Latham, 559) BaBar Dalitz-plot analysis of  $B^0 \to \overline{D}{}^0\pi^+\pi^-$  Understanding time-dependent DP for  $B^0 \to D_{CP}\pi^+\pi^-$  towards measurement of sin 2 $\beta$  and cos 2 $\beta$ . Preliminary BFs presented.

Belle (Higuchi, 1094) Time-dependent Dalitz plot analysis of  $B^0 \to K^+ K^- K_S^0 \ (b \to ss\bar{s} \text{ penguin})$ 

– Find four solutions; preferred solution yields

$$\phi_1^{\text{eff}}(\phi(1020)K_S^0) = (32.2 \pm 9.0 \pm 2.6 \pm 1.4(\text{DP model}))^\circ$$
  
 $\phi_1^{\text{eff}}(f_0(980)K_S^0) = (31.3 \pm 9.0 \pm 3.4 \pm 4.0(\text{DP model}))^\circ$ 



#### Compare with Penguin modes

 $\sin 2\beta$  from the  $b \to c\bar{c}s$  "golden" modes



#### NP in loop can give rise to deviations from $\beta/\phi_1$

The angles: Measuring  $\alpha/\phi_2$ 

RPP 2008  $\alpha = (88^{+6}_{-5})^{\circ}$  from  $B \to \pi\pi, \rho\rho, \rho\pi$ Measure in  $b \to u \bar{u} d$ 

- E.g.,  $B \to \pi^+\pi^-, \ \rho^+\rho^-, \ \pi^+\pi^-\pi^0, \ a_1^{\pm}\pi^{\mp}$ 



b

- Penguin contributions (involving different CKM phase) complicate analysis. Isospin analysis permits isolation of tree amplitude [Gronau and London, PRL 65 (1990) 3381]



## The angles: Measuring $\gamma/\phi_3$

$$\gamma \equiv \arg\left(-rac{V_{ud}V_{ub}^{*}}{V_{cd}V_{cb}^{*}}
ight)$$

Accessible in interference between  $b \to c\bar{u}s$   $(O(\lambda^3))$  and  $b \to u\bar{c}s$   $(O(\lambda^3))$ , colorsuppressed) amplitudes. A suitable pair of channels is  $B^- \to D^{(*)0}K^-$  and  $B^- \to \bar{D}^{(*)0}K^-$ , where interference may occur when the D and  $\bar{D}$  decay to common final states.



Compare  $B^-$  and  $B^+$ 

Various approaches  $(D^0 \overline{D}^0 \text{ mixing is neglected})$ :

# The angles: Measuring $\gamma/\phi_3$ (GLW)

GLW (Gronau, London, Wyler): Uses  $D, \overline{D}$  decays to CP eigenstates, eg,  $K^+K^-$  or  $K_S\pi^0$ . In this case, both D and  $\overline{D}$  decays are Cabibbo suppressed.



**BABAR** (Martinez-Vidal, 1175) Preliminary  $B^{\pm} \to D_{CP}K^{\pm}$ , with  $D_{CP+} \to \pi^{-}\pi^{+}, K^{-}K^{+}$ and  $D_{CP-} \to K^{0}_{S}\pi^{0}, K^{0}_{S}\phi, K^{0}_{S}\omega$ :



# The angles: Measuring $\gamma/\phi_3$ (ADS)

ADS (Atwood, Dunietz, Soni): Use  $D^0 \to K^+\pi^-$  (doubly Cabibbo suppressed);  $\bar{D}^0 \to K^+\pi^-$  (Cabibbo favored), giving interfering amplitudes of similar order, although branching fractions are small.

**BABAR** (Martinez-Vidal, 1175)  $B^- \to D^{(*)}K^- r_B = (9.5^{+5.1}_{-4.1})\%, r_B^* = (9.6^{+3.5}_{-5.1})\%.$ 



# The angles: Measuring $\gamma/\phi_3$ (GGSZ)

GGSZ (Giri, Grossman, Soffer, Zupan): Look at the Dalitz plot for three-body D decays, eg,  $D \to K_S \pi^+ \pi^-$ . This mode is Cabibbo favored for both  $D^0$  and  $\bar{D}^0$ .

Belle (Joshi, 1096) PRD **81** (2010) 112002

Dalitz Plot analysis  $B \to D^{(*)}K, D \to K_S \pi^+ \pi^-$ (Cabibbo allowed; large strong phases; need Dalitz plot analysis)  $B \to DK \to K_S \pi^+ \pi^-$ 657M  $B\bar{B}$   $m_{\pm} = m(K_S \pi^{\pm})$ 



 $\phi_3 \pmod{180} = [78.4^{+10.8}_{-11.6} \pm 3.6 (\text{syst}) \pm 8.9 (\text{model})]^{\circ}$ 

*P*-value for *CP* conservation is  $5 \times 10^{-4}$  (combined  $B^{\pm} \to D^{(*)} K^{\pm}$ )

BABAR (Martinez-Vidal, 1175)  $B^{\mp} \rightarrow D^{(*)}K^{(*)\mp}$ exclude  $\gamma = 0$  at  $3.5\sigma$  $\gamma \pmod{180} = [68 \pm 14 \pm 4(\text{syst}) \pm 3(\text{model})]^{\circ}$ 

 $r_B = 0.096 \pm 0.029$ 



Frank Porter, ICHEP2010, Paris, 27 July 2010

## Understanding D decays

We have seen that measuring  $\gamma/\phi_3$  is intimately connected with D decays; motivated to understand D decays to reduce model dependence. CLEOc (Wilkinson, 702) use quantum correlations at  $\psi(3770) \rightarrow D^0 \bar{D}^0$  to measure strong phase differences between  $D^0 \to K^0_S \pi^+ \pi^-$  and  $\overline{D}^0 \to K^0_S \pi^+ \pi^ (818 \text{ pb}^{-1})$ . Updated analysis; new analysis of  $K_S K^+ K^-$ . Idea is can tag D eigenstate (either flavor or CP), eg, with tag D going to CP eigenstate such as  $K^+K^-$  (*CP*-even), hence signal  $D \rightarrow$  $K_S \pi^+ \pi^-$  is *CP*-odd *D* state.

 $c_i$  and  $s_i$  are cosines and sines of strong  $D - \overline{D}$ decay phase differences, averaged over bin i



# Searches for new physics in CP violation

- CP violation in B decays
  - Belle (Higuchi, 1094) Direct CP in  $B^+ \to J/\psi K^+$
  - Belle (Sahoo, 969) New result for time-dependent CP analysis of  $B^0 \to \phi K_S \gamma$
- CP violation in D mixing and decay
  - **BABAR** (Bellis, 1172)  $D^0 \to K^0_S \pi^+ \pi^-$  and  $D^0 \to K^0_S K^+ K^-$  Dalitz plot analysis
  - Belle (Ko, 1092) *CP* violation in  $D \to K_S(\pi, K, \eta, \eta')$  and  $D_{(s)} \to \phi \pi$
  - CDF (Mattson, 1082) CP violation in  $D^0 \rightarrow h^+ h^-$
  - CP violation in kaons
    - KEK E391a (Watanabe, 734) Final results on  $K_L \to \pi^0 \nu \bar{\nu}$
    - NA48 (Winhart, 1080) CP measurements in  $K^{\pm} \to \pi \ell^+ \ell^-$  and  $K_S \pi \pi ee$  decays
    - CP violation in  $\tau$  decays
    - Belle (Shapkin, 1093) CP violation in  $\tau^{\pm} \to K_S \pi^{\pm} \nu_{\tau}$

# The global fits



Both UTfit (Tarantino, 1081) and CKMfitter (T'Jampens, 190) identify  $\sin 2\beta$  (2.6 $\sigma$ /2.6 $\sigma$ ) and  $\mathcal{B}(B \to \tau \nu)$  (3.2 $\sigma$ /2.8 $\sigma$ ) as areas of discrepancy.

- UTfit in addition mentions  $\epsilon_K$  as discrepant by  $1.7\sigma$ .
- Global consistency from CKM fitter at  $2\sigma$

# Characterizing the discrepancy

Two-dimensional value of  $(\sin 2\beta, \mathcal{B}(B \to \tau \nu))$  in conflict with  $B_{B_d}, \alpha, \gamma$  constraints.



- Lattice error
- New physics

See also (Soni, 908)

# The $B_s$ sector

UTfit with new D0 results (awaiting CDF likelihood),  $3.1\sigma$  from SM in  $\phi_{B_s}$  (but new CDF result should pull it closer to SM).



# Conclusions

 $|V| = \begin{pmatrix} 0.97418 \pm 0.00027 & 0.2253 \pm 0.0008 & 0.00392 \pm 0.00046\\ 0.230 \pm 0.011 & 1.04 \pm 0.06 & 0.0409 \pm 0.0007\\ 0.0081 \pm 0.0005 & 0.0387 \pm 0.0023 & 0.88 \pm 0.07 \end{pmatrix}$ 

Still plenty of room for a fourth generation.

 $\begin{array}{c|c} \text{ICHEP 2010 averages (assuming 3 \times 3 unitarity, SM)} \\ & \text{CKMfitter, ICHEP10} & \text{UTfit, ICHEP10} \\ \hline A & 0.812^{+0.013}_{-0.027} \\ \hline \lambda & 0.22543 {\pm} 0.00077 \\ \hline \bar{\rho} & 0.144 {\pm} 0.025 & 0.132 {\pm} 0.020 \\ \hline \bar{\eta} & 0.342^{+0.016}_{-0.015} & 0.358 {\pm} 0.012 \\ \hline \alpha(^{\circ}) & 91.0 {\pm} 3.9 \\ \sin 2\beta & 0.689^{+0.023}_{-0.021} \\ \hline \gamma(^{\circ}) & 67.2 {\pm} 3.9 \end{array}$ 

Warning: errors may not scale as normal errors; see references.



Some " $2\sigma$ " hints

 $\tau$  to s inclusive puzzle

Exclusive vs inclusive differences for  $|V_{ub}|$  and  $|V_{cb}|$ 

 $\sin 2\beta$  and  $B \to \tau \nu$  discrepancy with SM

Like sign dimuon discrepancy with SM

Heavy flavors will continue to offer insights/constraints on possible new physics [LHC, high intensity kaons, super B factories, tau/charm threshold]