

# Extreme QCD in Heavy Ion Collisions

**Raju Venugopalan**

**BNL**

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# Quantum Chromodynamics (QCD)

- QCD - “nearly perfect” fundamental quantum theory of quark and gluon fields (F.Wilczek, hep-ph/9907340)
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- ii) Global “chiral” symmetry: exact for massless quarks
- iii) Baryon number and axial charge (m=0) are conserved
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  - v) Discrete C,P & T symmetries
- Chiral, Axial, Scale and (in principle) P & T broken by vacuum/quantum effects - “emergent” phenomena
  - What happens at finite temperature & density ?

# QCD: structure & consequences



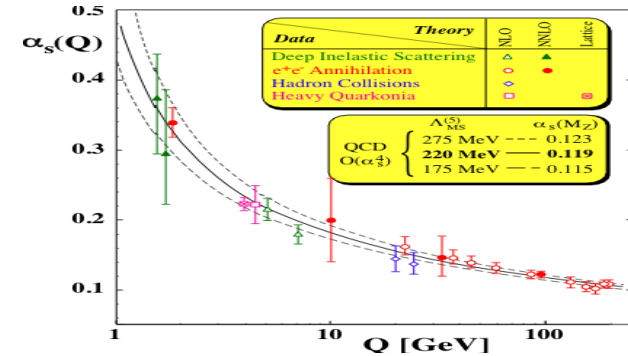
**Asymptotic freedom:**

**S:** Coupling strength of gluons weaker  
at short separation

**C:** Super-dense & super-hot QCD matter  
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Collins-Perry (1975); Cabibo-Parisi (1975)

Gross, Wilczek, Politzer (1973)



# QCD: structure & consequences



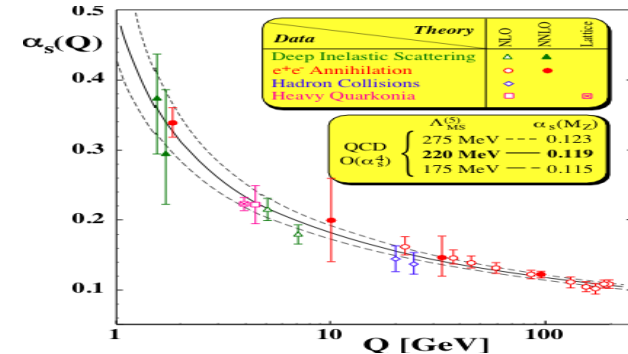
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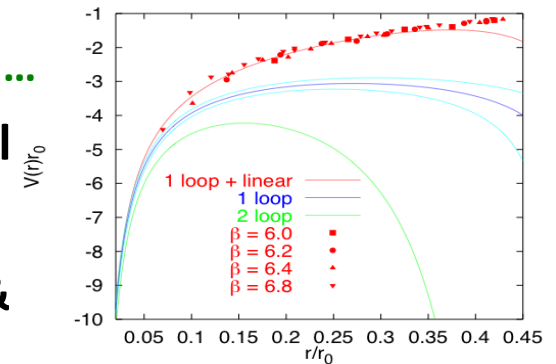


**Infrared slavery:**

Wilson (1974), Polyakov, ...

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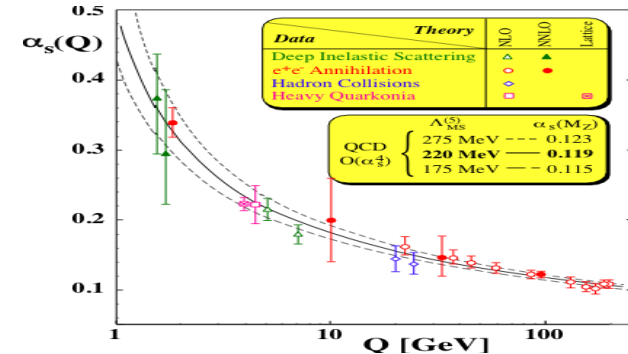


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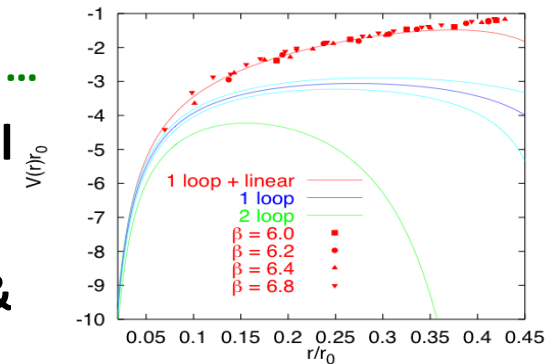
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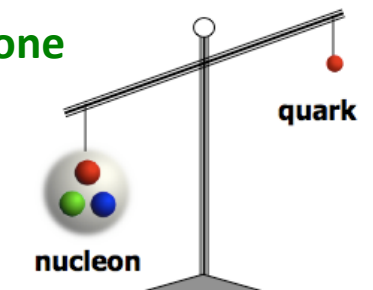


## (Broken) Chiral symmetry

- S: spontaneously generated Chiral condensate
- C: Chiral symmetry restored at finite T

Nambu-Goldstone

$$\langle \bar{q}_R q_L \rangle \neq 0$$





# QCD : brief pre-history

From E. Fermi: "Notes on Thermodynamics and Statistics" (1953)



Hagedorn (1965)



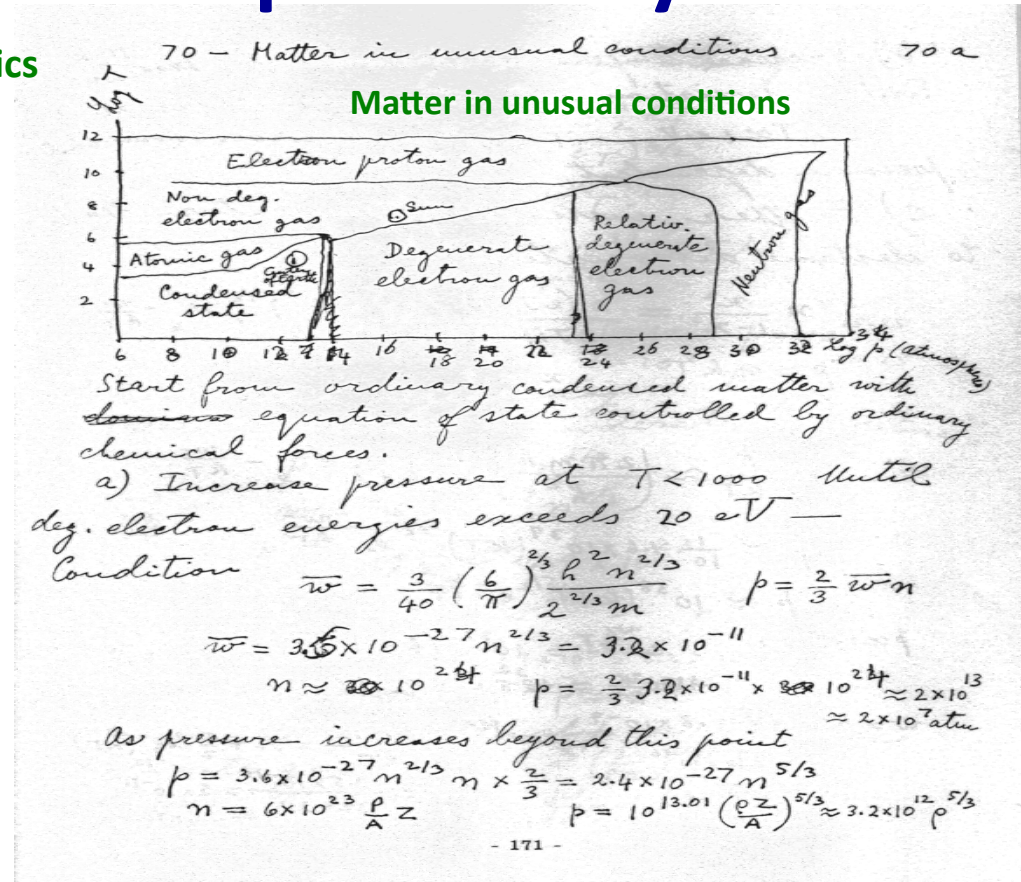
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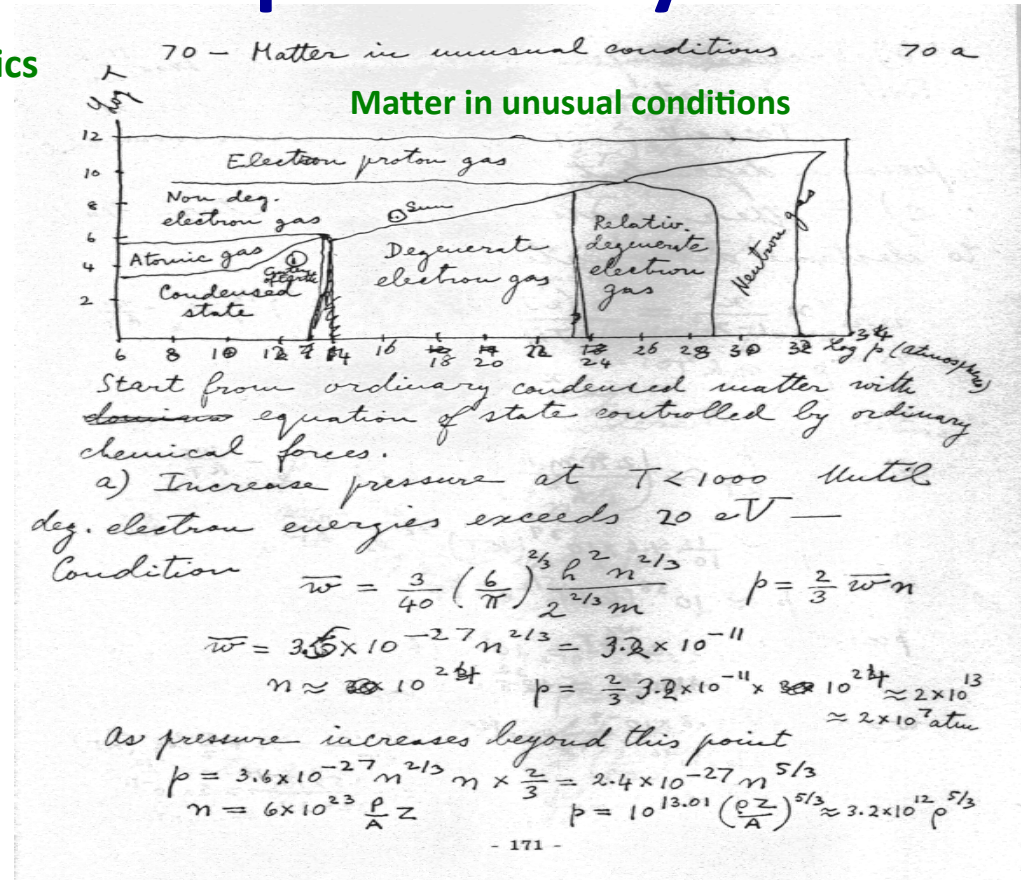
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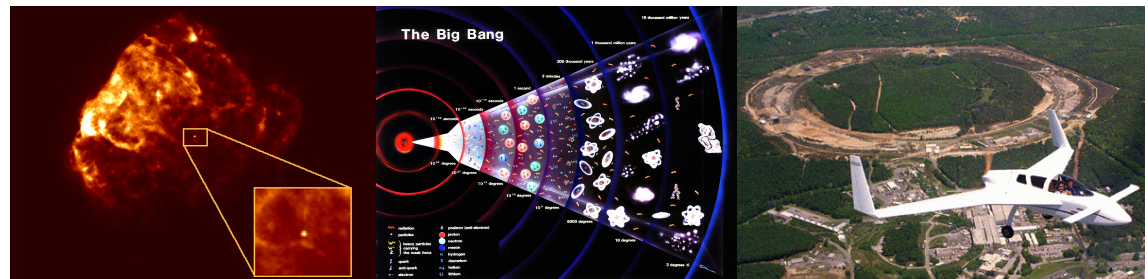
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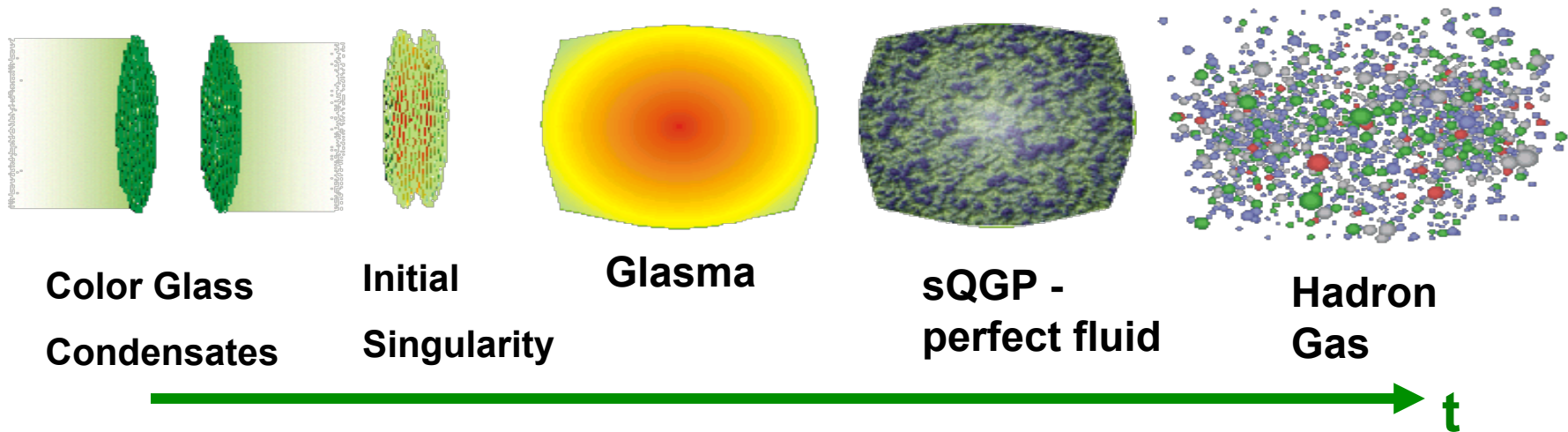
Neutron Stars

Big Bang

RHIC

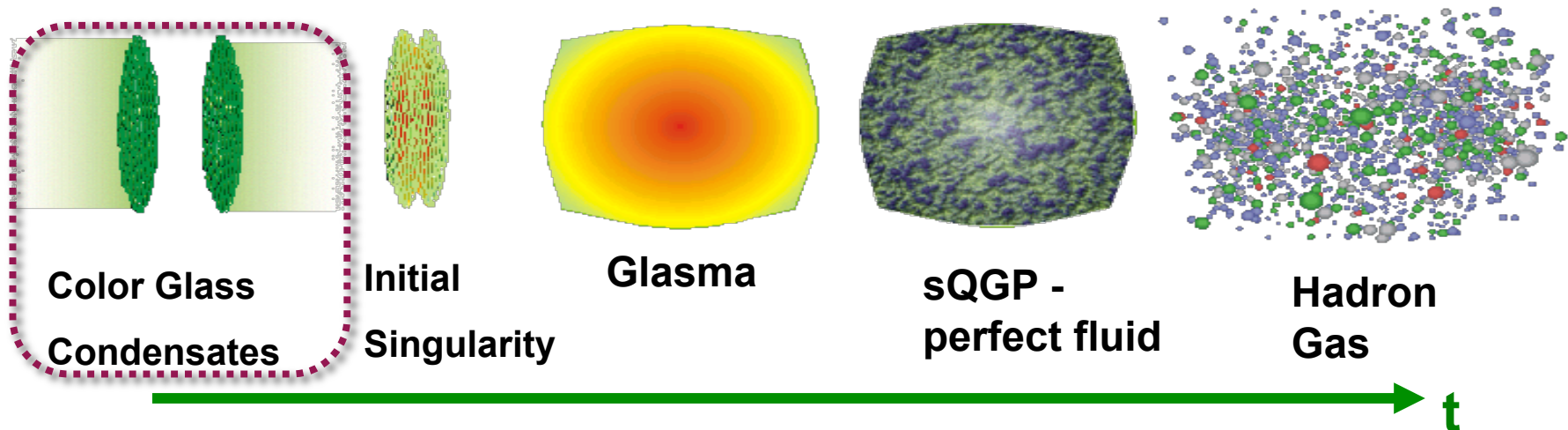


# Standard model of HI Collisions



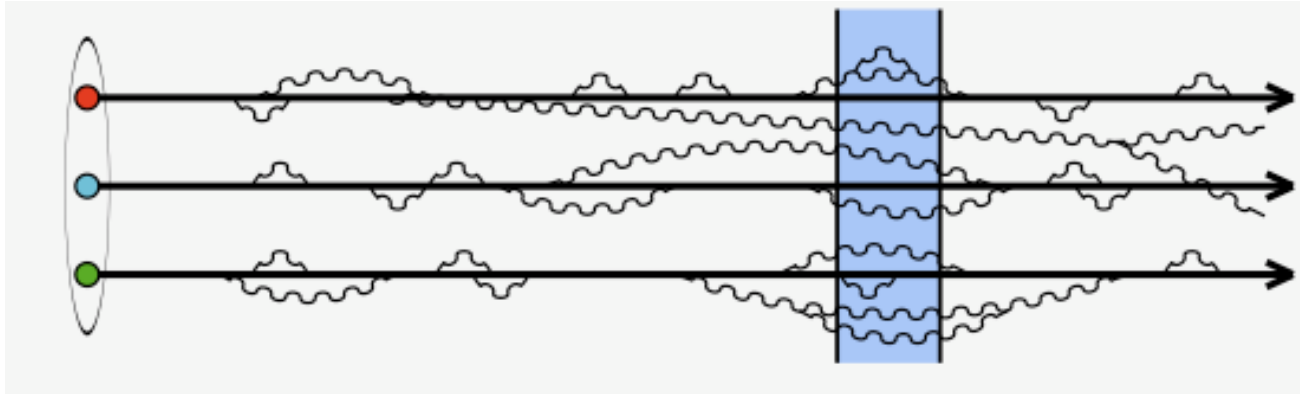
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- Solidify, Improve, Falsify ?

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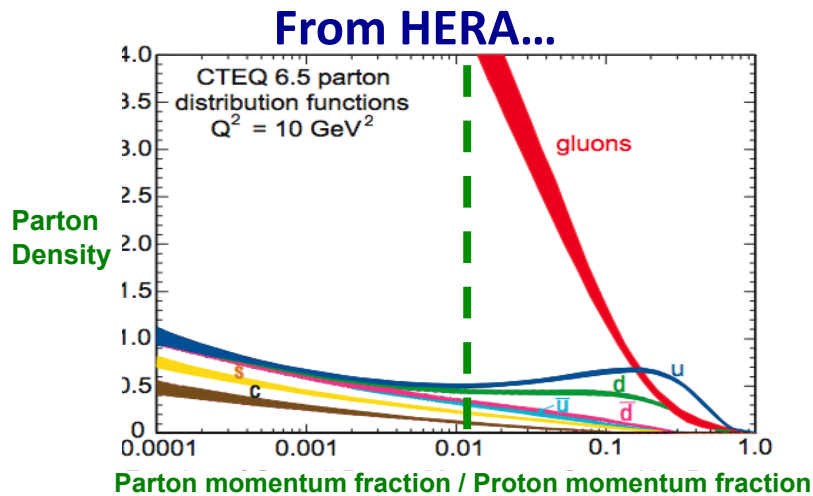


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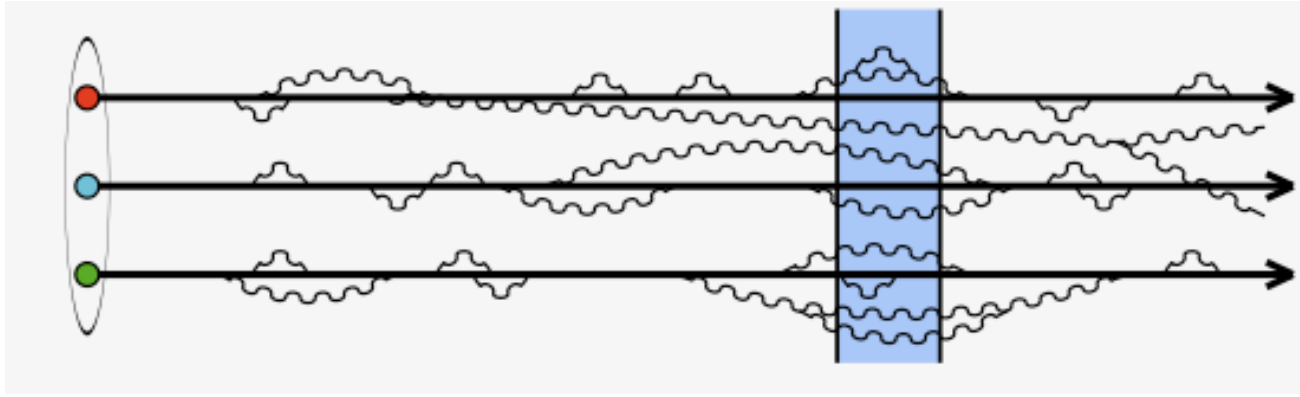
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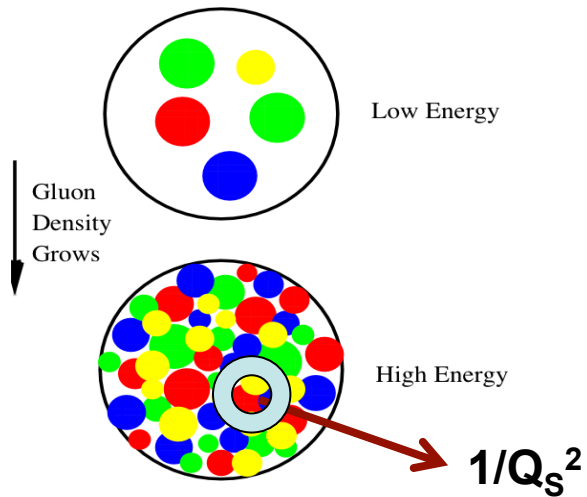
Wee parton fluctuations time dilated on strong interaction time scales



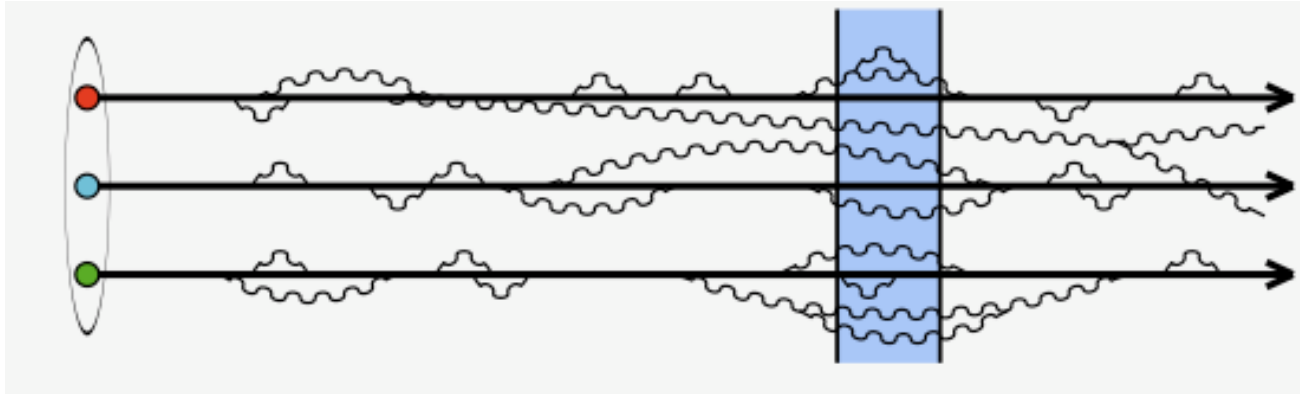
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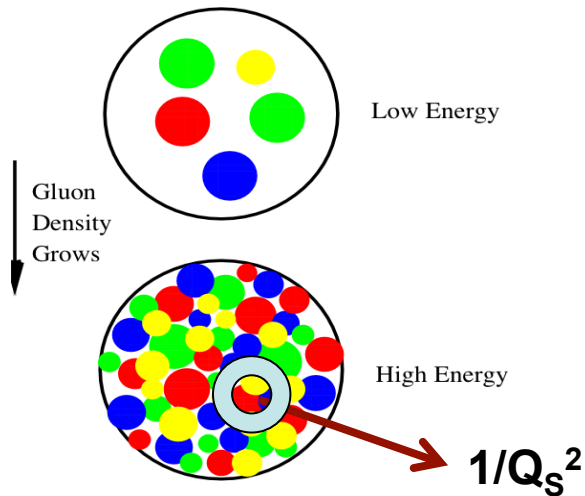
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# Gluon Saturation in a nucleus: classical coherence from quantum fluctuations



Wee parton fluctuations time dilated on strong interaction time scales



The gluon density saturates at a maximal value of  $\sim 1/\alpha_s$   
→ gluon saturation

Gribov, Levin, Ryskin; Mueller, Qiu

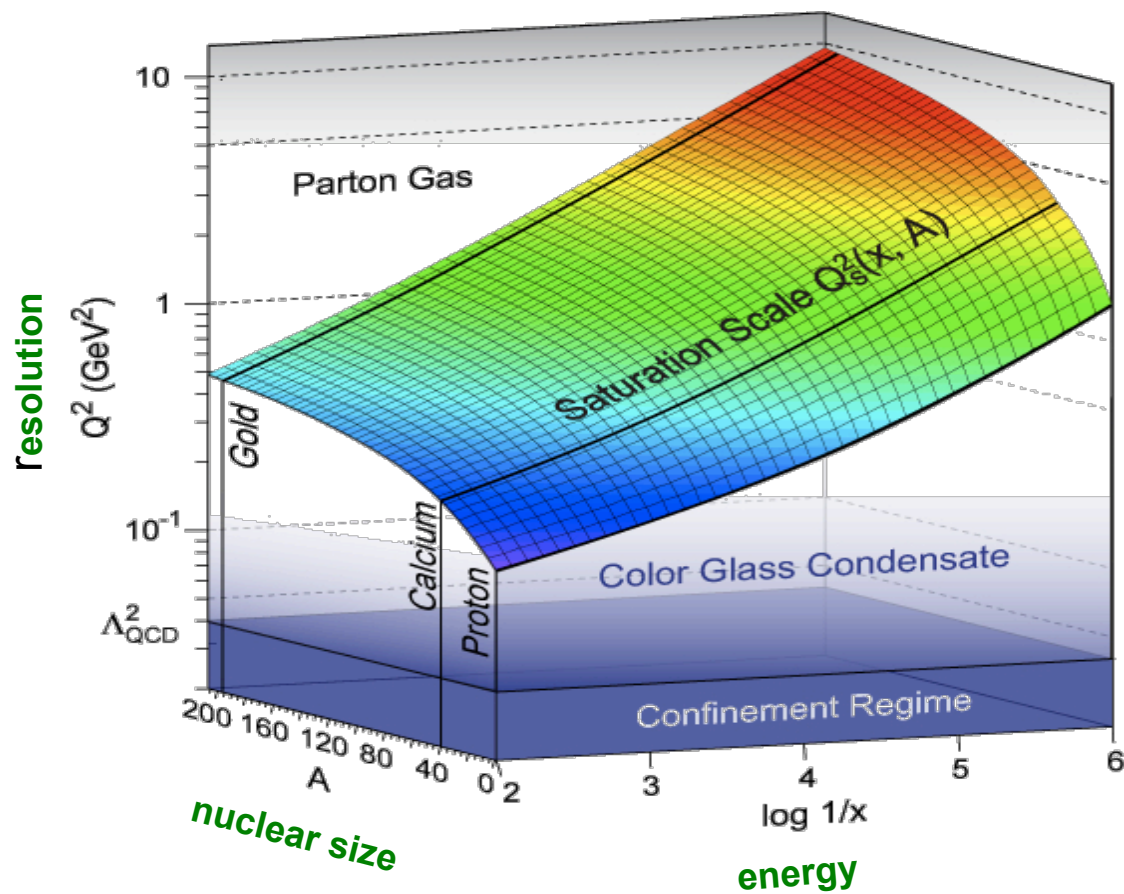
Large occupation # => classical color fields

McLerran, RV

# Many-body high energy QCD: The Color Glass Condensate

- QCD EFT framework to understand strongly correlated wee gluons in nuclei and their evolution with energy
- Dynamics characterized primarily by one universal semi-hard scale -- the saturation scale  $Q_s^A$  ( **geometrical scaling-Peschanski talk**)
- Physics highly non-perturbative, but can be described in weak coupling -- novel regime of many body QCD
- **Strong hints from HERA (esp. diffractive/exclusive data-M. Machado talk) AND RHIC data**

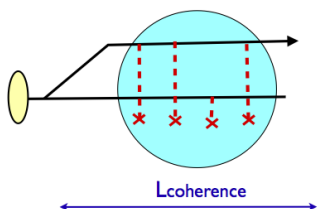
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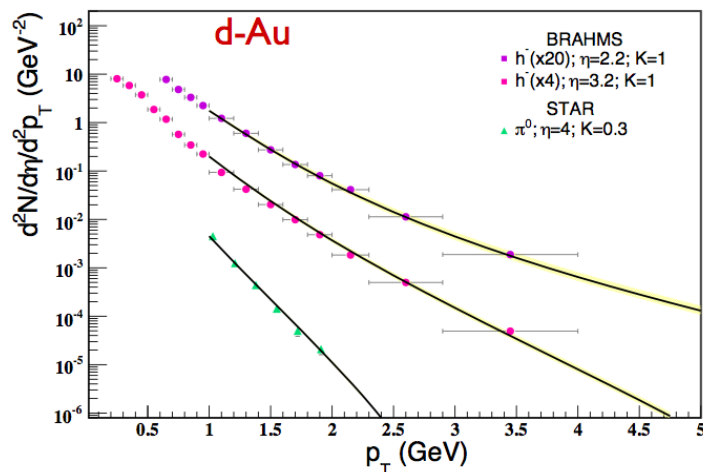


# Evidence for Saturation @ RHIC ?

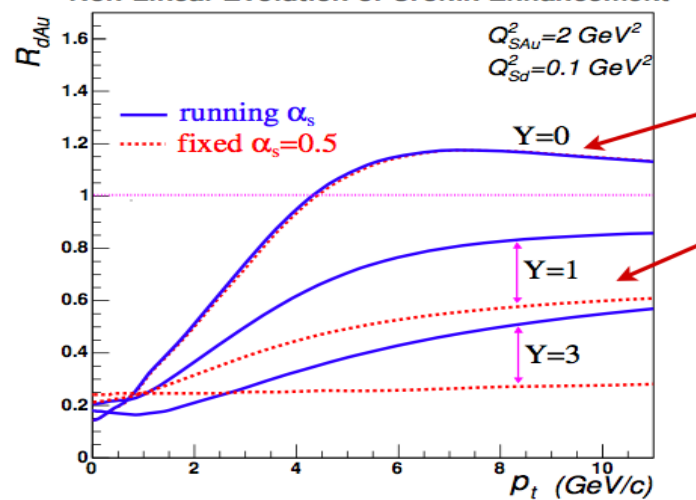
## Rich Deuteron+Gold physics program



Talks by J. Albacete  
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Caveats: M. Sumbara

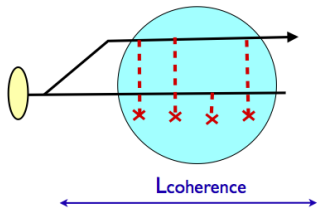


Non-Linear Evolution of Cronin Enhancement

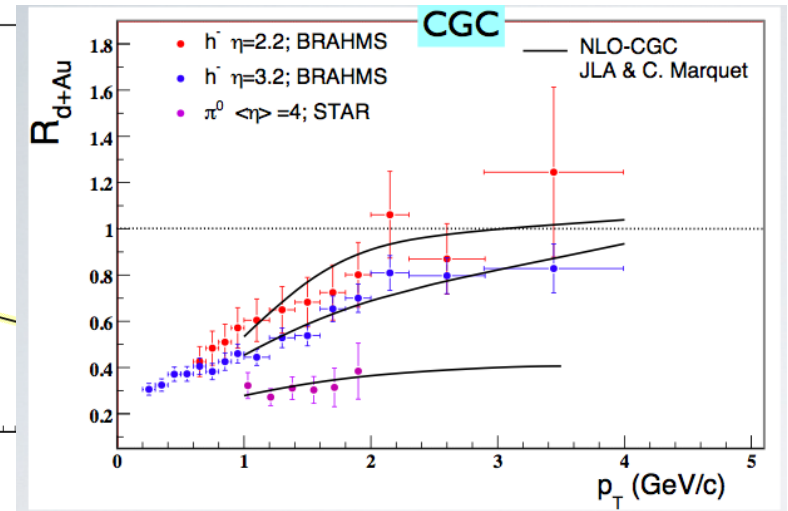
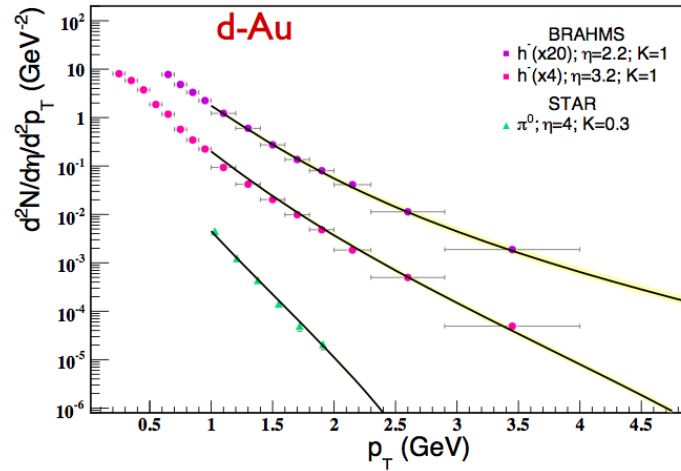


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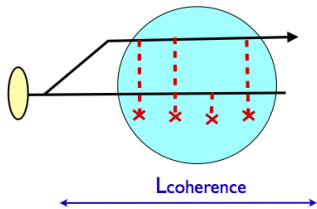


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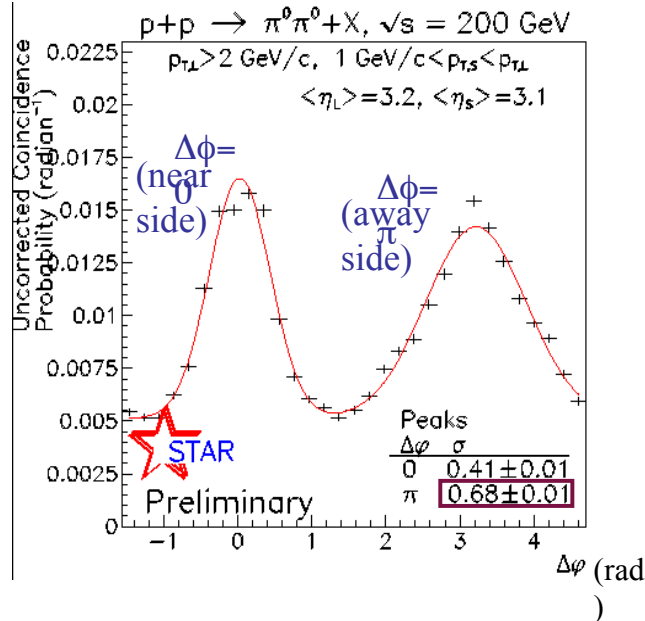
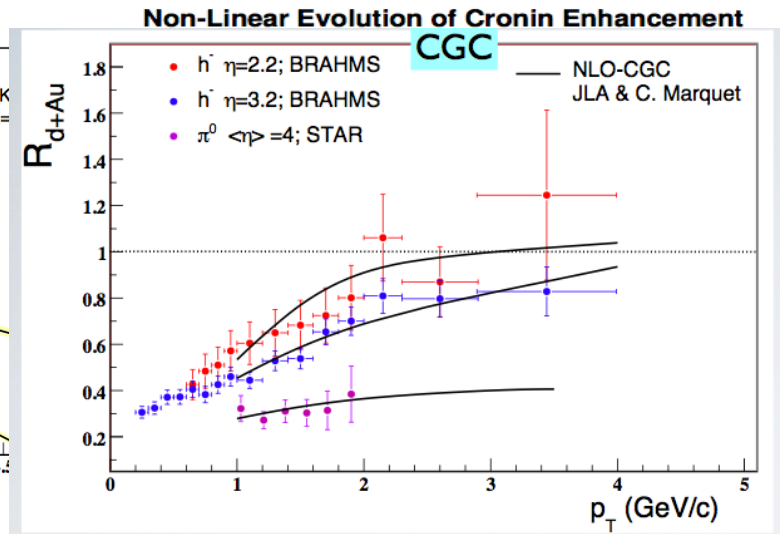
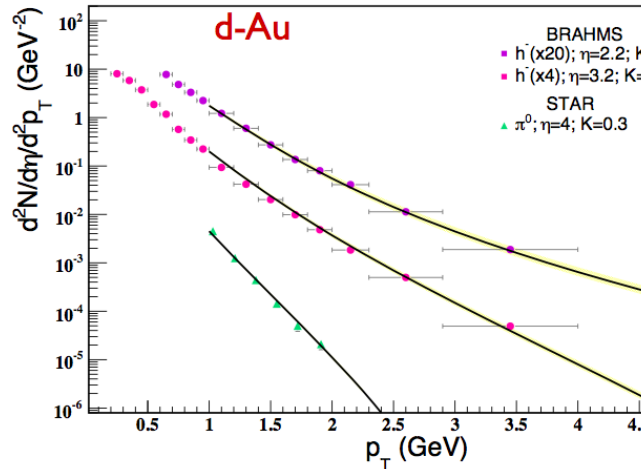
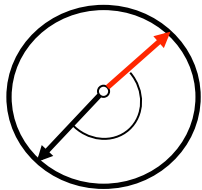


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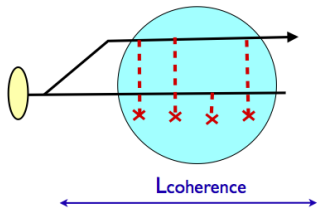


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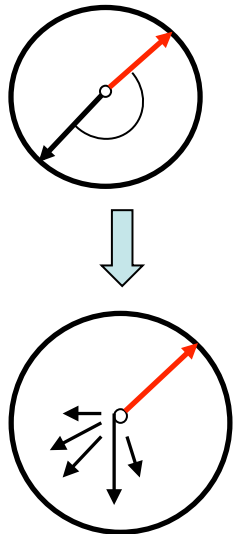


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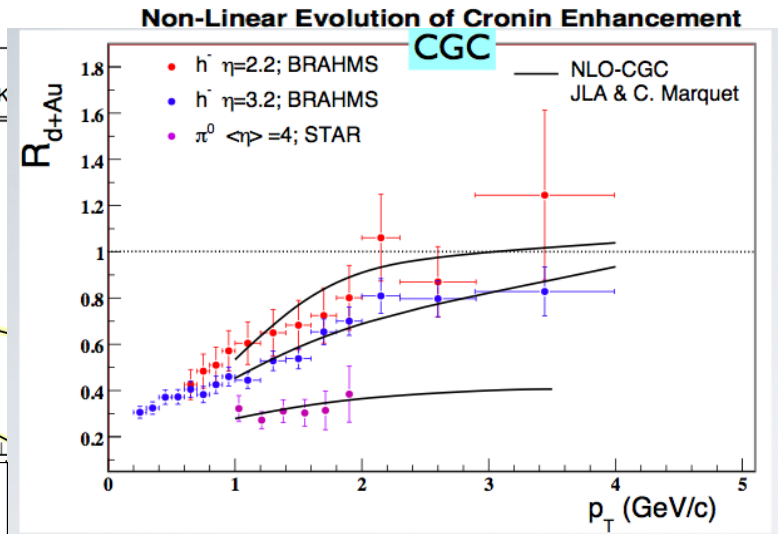
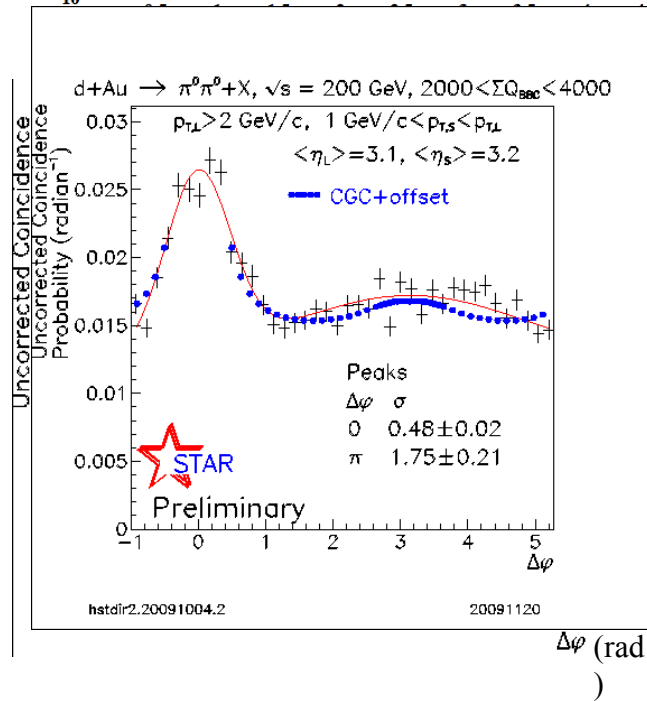
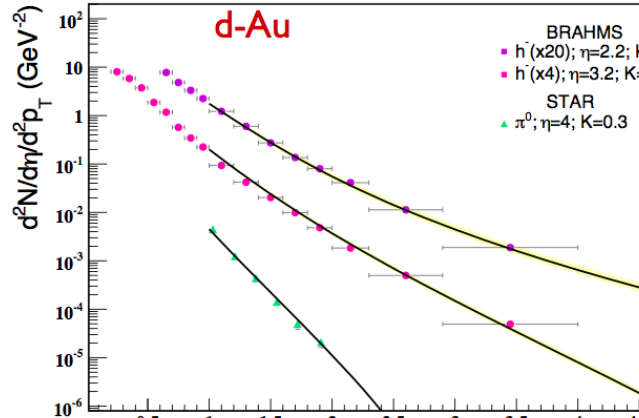
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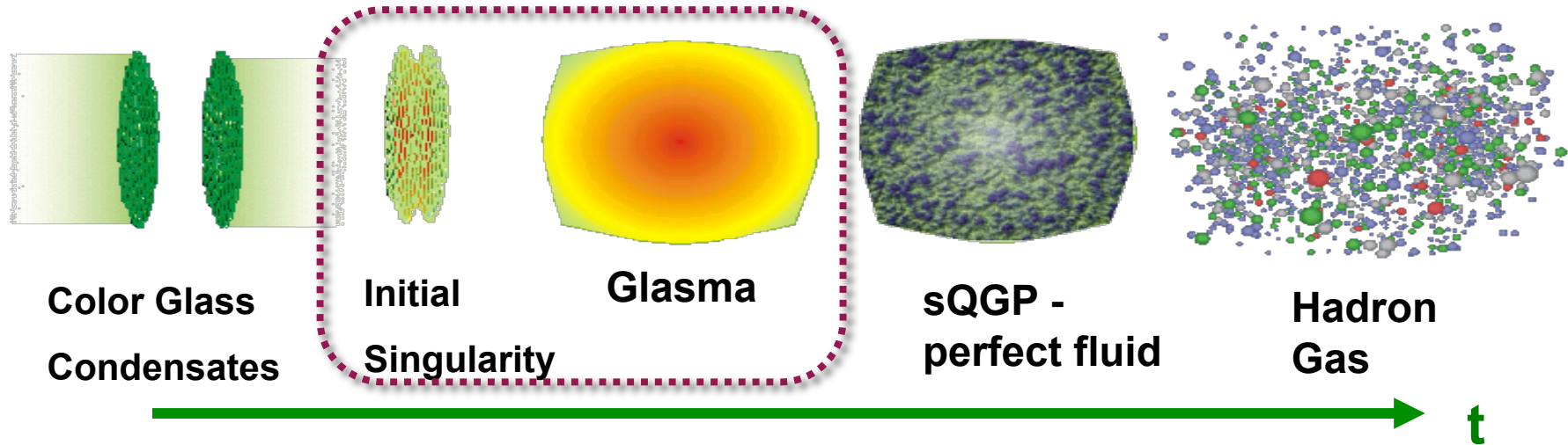
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Away side parton  
randomized  
by strong color field



# Standard model of HI Collisions



**Glasma** (\Glahs-maa\): *Noun*: non-equilibrium matter between Color Glass Condensate (CGC) & Quark Gluon Plasma (QGP)

# The big role of wee glue



# The big role of wee glue

## D. Nucleus-Nucleus Collisions at Fantastic Energies (Nucleus-Nucleus Collisions at Fantastic Energies)

Before leaving this subject it is fun to consider the collision of two nuclei at energies sufficiently high so that in addition to the fragmentation regions, a central plateau region can develop. Let us consider a central collision of a relatively small nucleus, say carbon, with a big one, say lead. Let us look at this collision in a center-of-mass frame for which the rapidities of both of the nucleus projectiles exceeds the critical rapidity. In such a frame they both possess the fur coat of wee-parton vacuum fluctuations. In such a central collision we see that the collision initially occurs between the fur of wee partons in each of the projectiles. Therefore the number of independent collisions will be of order of the area of overlap of the two projectiles; namely the cross-sectional area of the smaller nucleus.

At LHC, ~14 units in rapidity!



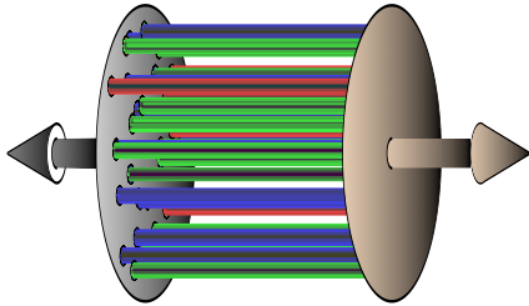
Bj, DESY lectures (1975)

# The big role of wee glue

- ❑ These humble wee gluons are responsible for producing the hottest matter on earth
- ❑ We now understand that they can be described as classical fields (**MV**) with (quantum) energy evolution given by “Wilsonian” renorm. group (**JIMWLK/BK**)

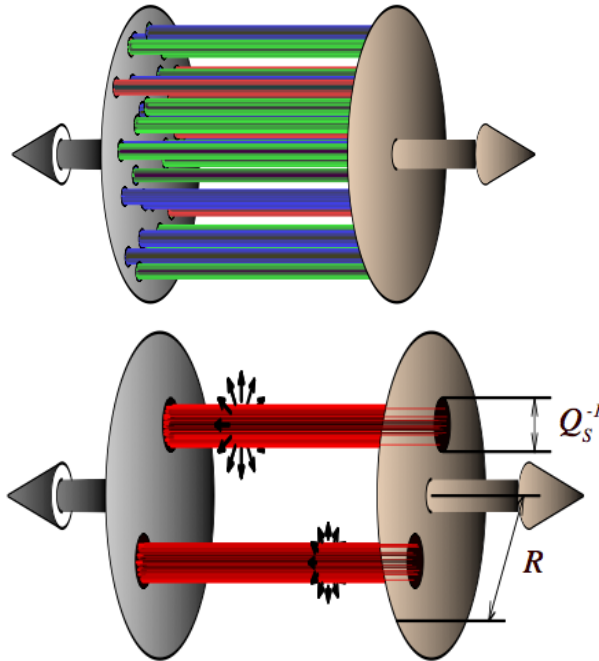


# Classical features of the Glasma



Solutions of Yang-Mills equations produce (nearly) boost invariant gluon field configurations: **“Glasma flux tubes”**

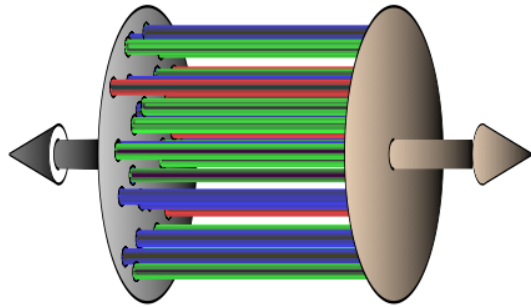
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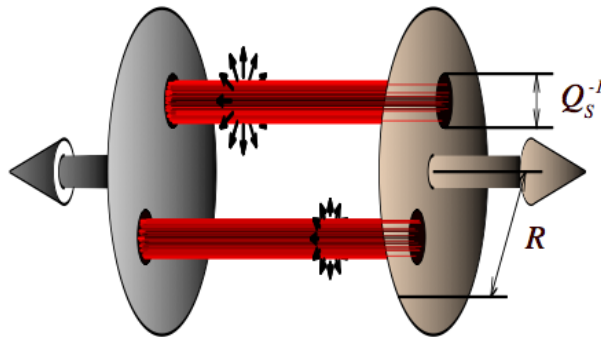
Solutions of Yang-Mills equations produce (nearly) boost invariant gluon field configurations: **“Glasma flux tubes”**

Lumpy gluon fields are **color screened** in transverse plane over distances  $\sim 1/Q_s$   
- Negative Binomial multiplicity distribution.

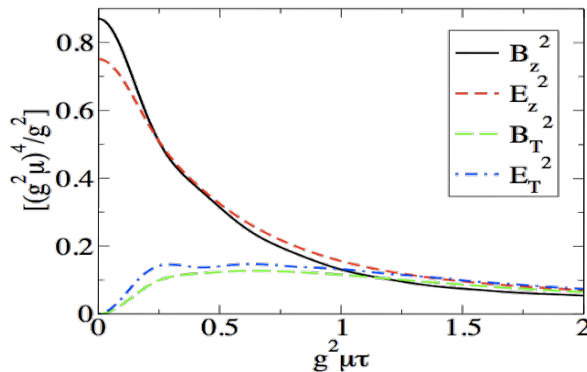
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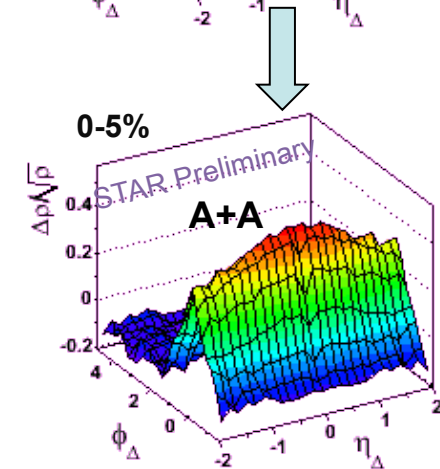
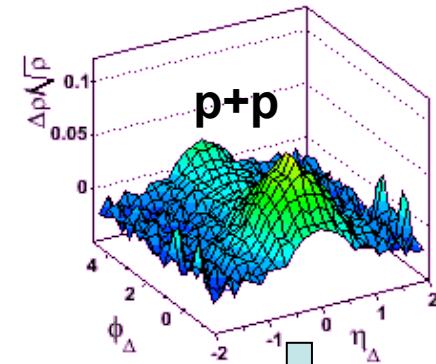
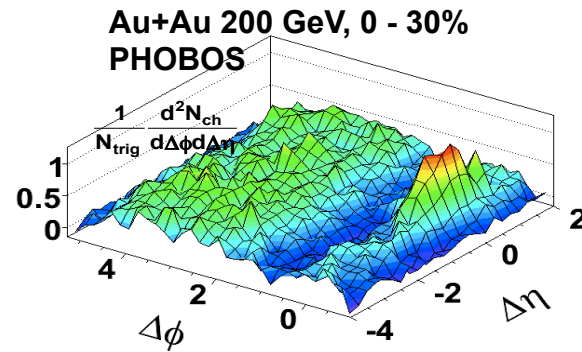
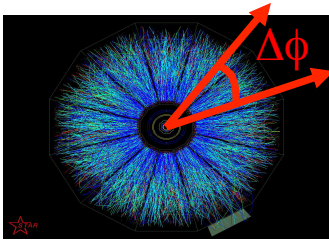


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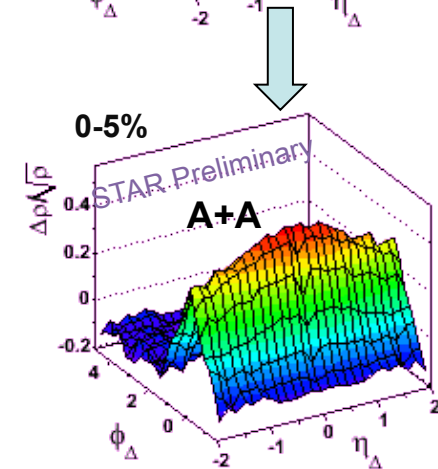
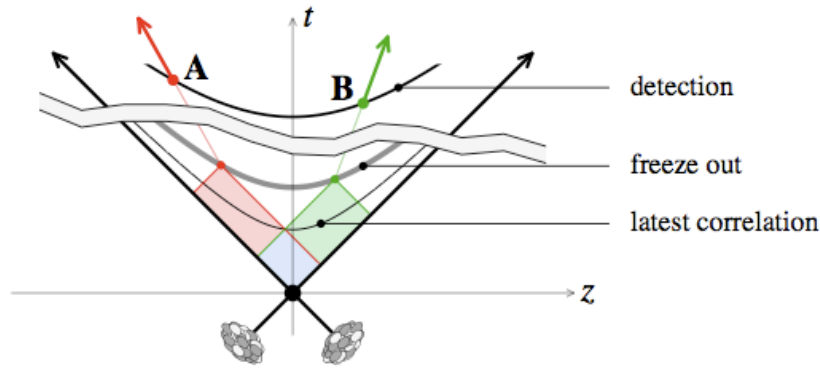
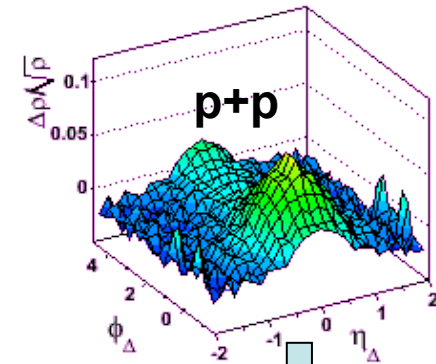
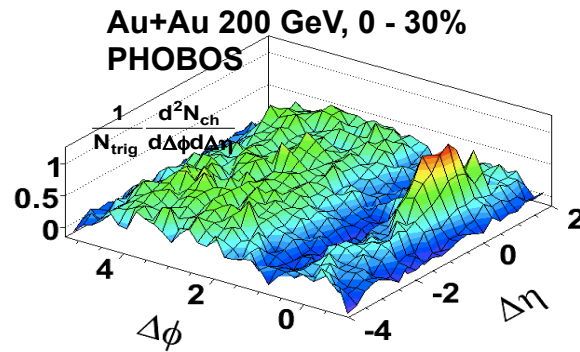
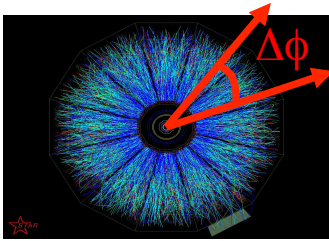


“Glasma flux tubes” have non-trivial longitudinal color E & B fields at early times  
 --generate **Chern-Simons** topological charge

# The Ridge: Glasma flux tubes+ Radial flow



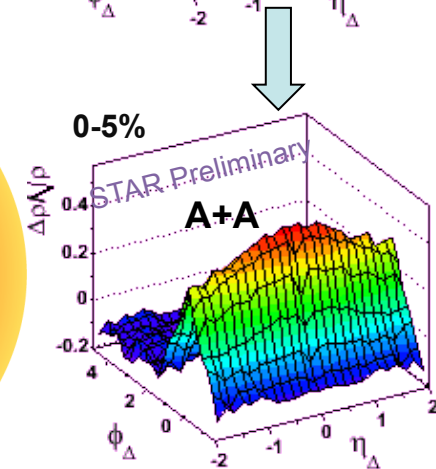
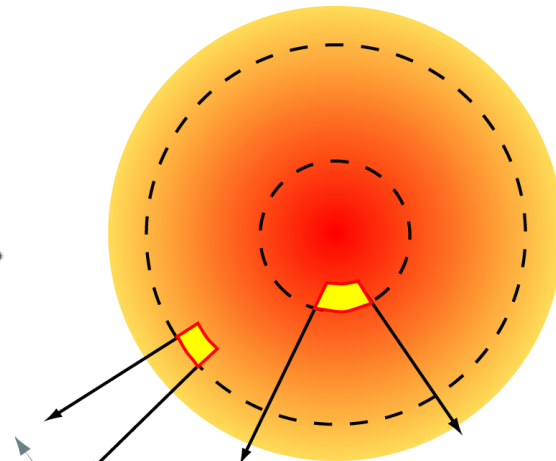
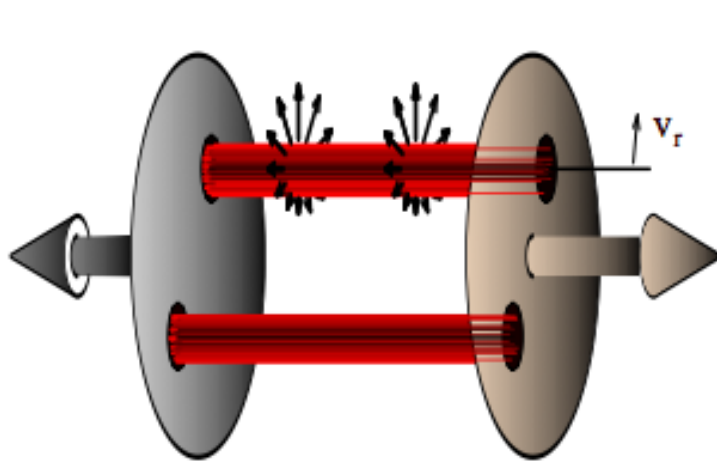
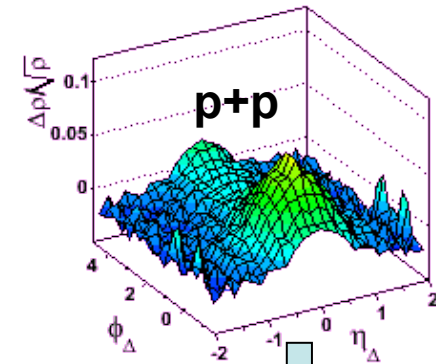
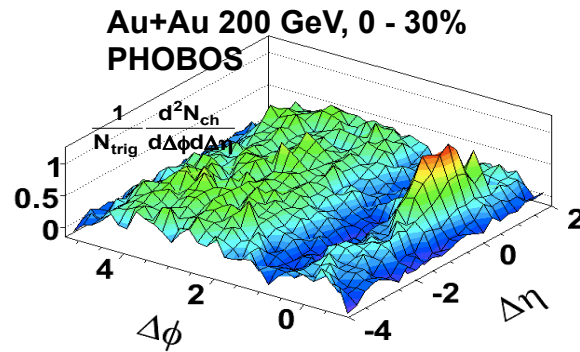
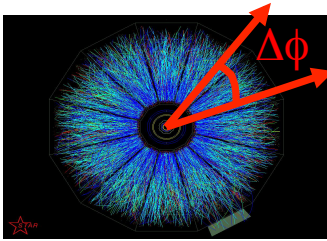
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Long range rapidity correlations are sensitive to Glasma dynamics at early times

$$\tau \leq \tau_{\text{freeze-out}} \exp\left(-\frac{1}{2}|y_A - y_B|\right)$$

# The Ridge: Glasma flux tubes+ Radial flow



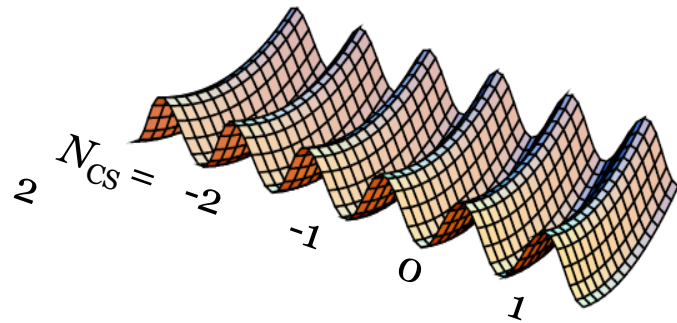
Glasma flux tubes provide the long range rapidity correlation

Dumitru, Gelis, McLerran, RV; Gavin, McLerran, Moschelli

Radial (“Hubble”) flow of the tubes provides the azimuthal collimation

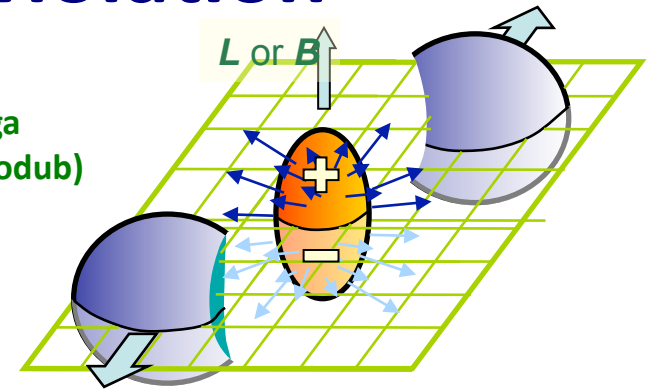
Voloshin; Shuryak

# Chiral Magnetic Effect: Local strong parity violation



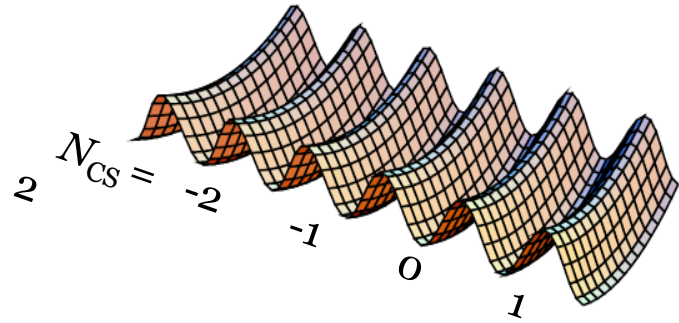
Topological fluctuations  
-sphaleron transitions in Glasma  
-as in EW baryogenesis

Kharzeev, McLerran, Warringa  
Kharzeev, Fukushima, Warringa  
(Talk by Fraga; Poster: Chernodub)



External (QED) magnetic field

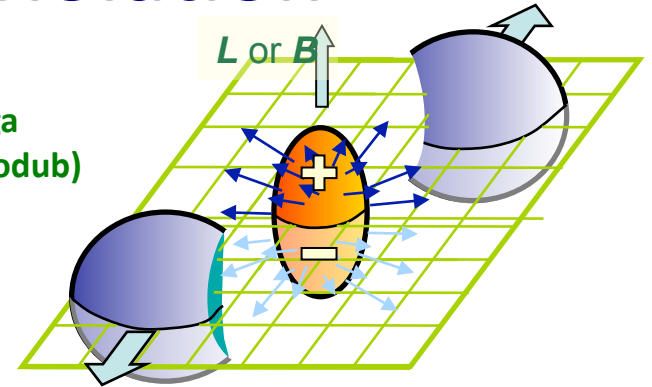
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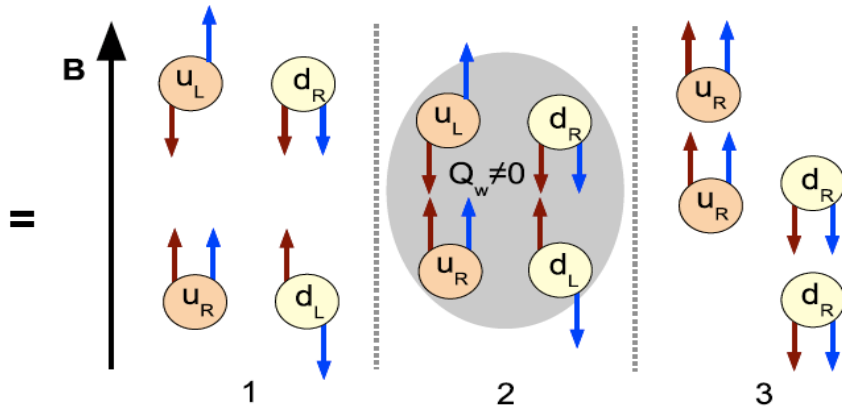
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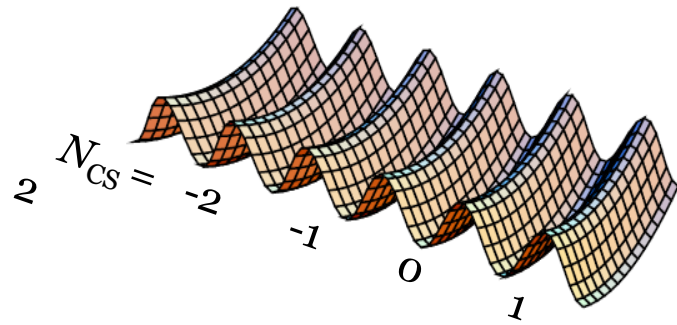
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Chiral magnetic effect

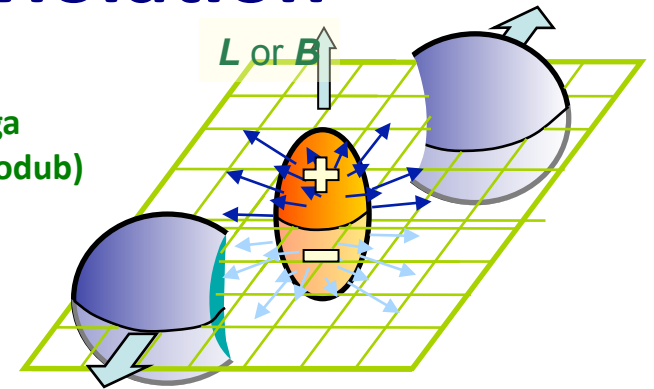


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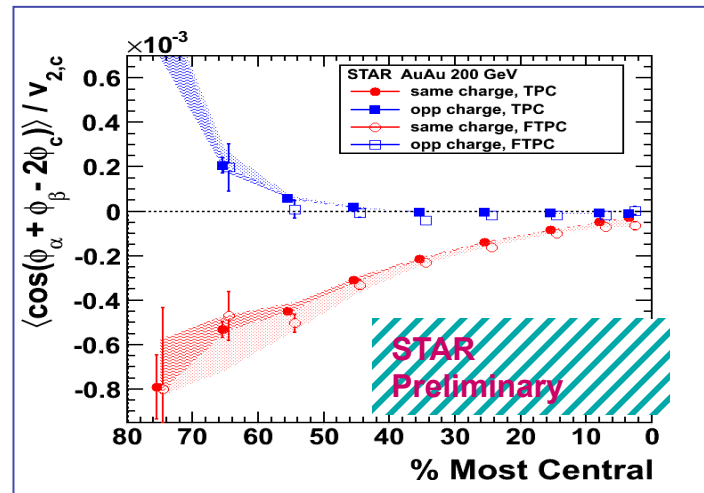


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-as in EW baryogenesis

Kharzeev, McLerran, Warringa  
Kharzeev, Fukushima, Warringa  
(Talk by Fraga; Poster: Chernodub)



External (QED) magnetic field

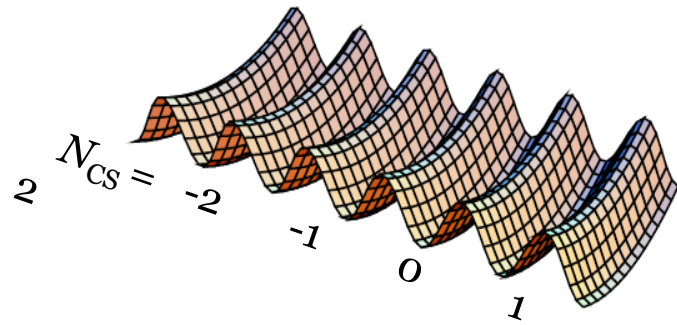


STAR,  
PRL 103,  
251601 (2010)

Possible experimental signal of charge separation

Important caveats (Wang: Bzdak, Koch, Liao; Pratt)

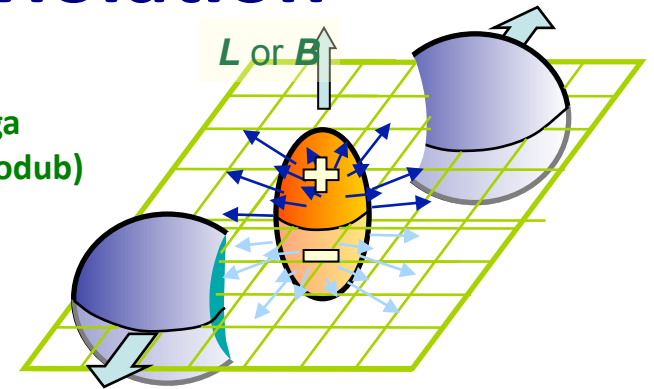
# Chiral Magnetic Effect: Local strong parity violation



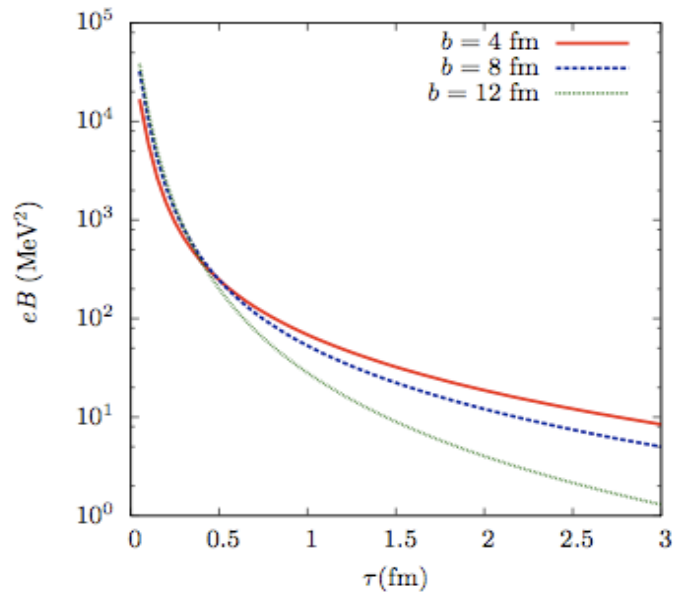
**Topological fluctuations**  
-sphaleron transitions in Glasma  
-as in EW baryogenesis

Khazzev, McLerran, Warringa  
Khazzev, Fukushima, Warringa  
(Talk by Fraga; Poster: Chernodub)

+

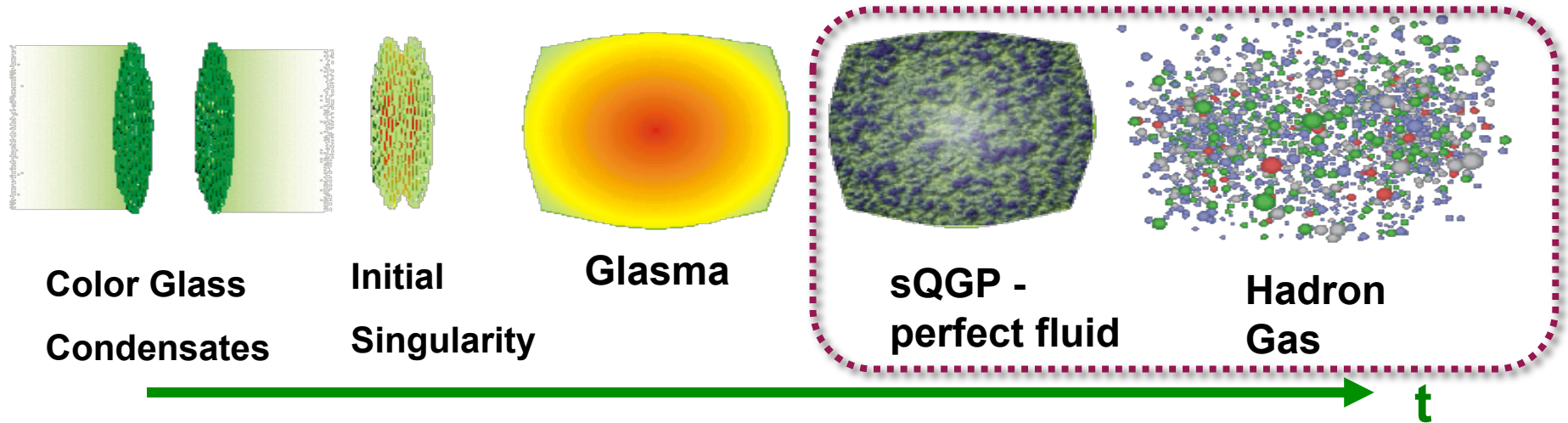


**External (QED) magnetic field**



**Effect most significant,  
for transitions at early times**

# Standard model of HI Collisions



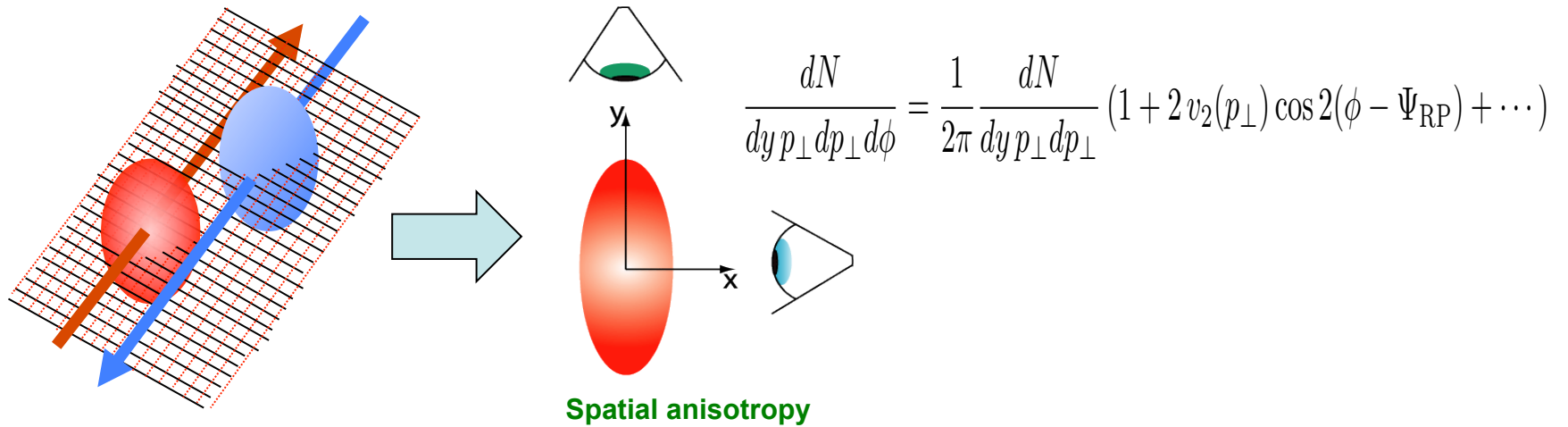
- Thermalization of Glasma still not understood
- Fast isotropization may occur due to instability induced field amplification  
- very analogous to **“preheating” phenomenon in inflationary cosmology**

Dusling, Gelis, Epelbaum, RV

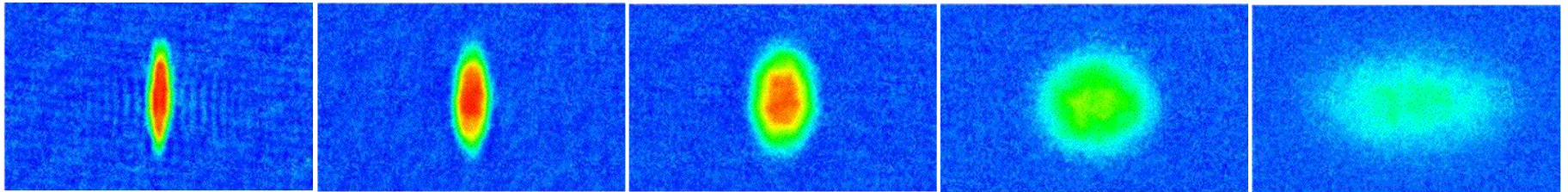
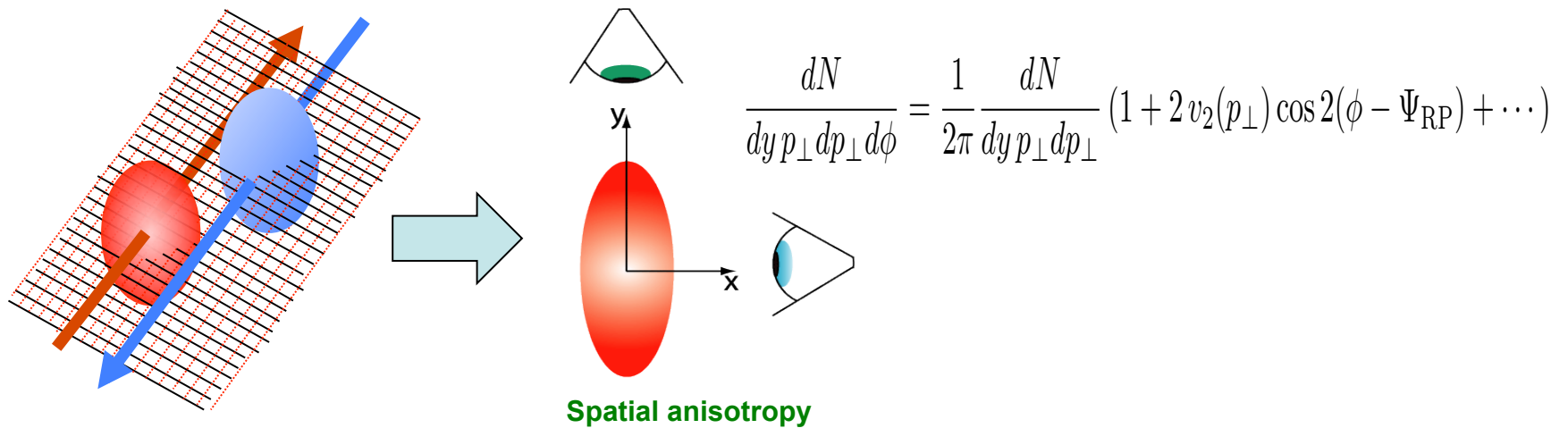
- Glasma provides initial conditions for hydrodynamic flow

Hirano et al.; Lappi, RV;  
Dumitru, Drescher, Nara

# Strong flow = (nearly) ideal hydrodynamics



# Strong flow = (nearly) ideal hydrodynamics



$v_2$  measures how efficiently hot matter converts spatial anisotropies to momentum anisotropy – most efficient way is hydrodynamics

$v_2$  at RHIC is large (B.Cole talk)  
– hydro flow must set in early



# A perfect fluid at RHIC

“Bjorken Hydrodynamics”

$$\frac{d\varepsilon}{d\tau} = - \frac{(\varepsilon + P - \frac{4}{3} \frac{\eta}{\tau})}{\tau}$$

Viscous term smaller than ideal term for

$$\frac{\eta}{\varepsilon + P} \frac{1}{\tau} \equiv \frac{\eta}{s} \frac{1}{\tau T} \ll 1$$

From kinetic theory

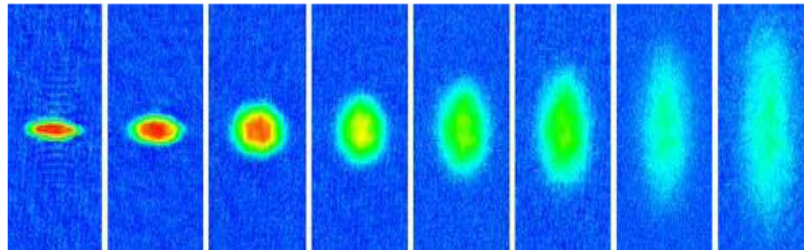
$$\frac{\eta}{s} \sim \frac{\hbar}{k_B} \frac{\tau_{\text{relax.}}}{\tau_{\text{quant.}}}$$



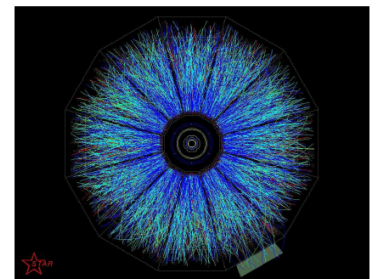
H<sub>2</sub>O



<sup>4</sup>He



<sup>6</sup>Li



sQGP

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Schafer,Teaney

Fluid	$T$ [K]	$\eta$ [Pa · s]	$\eta/n$ [ $\hbar$ ]	$\eta/s$ [ $\hbar/k_B$ ]
H <sub>2</sub> O	370	$2.9 \times 10^{-4}$	85	8.2
<sup>4</sup> He	2	$1.2 \times 10^{-6}$	0.5	1.9
<sup>6</sup> Li ( $ a_s  \simeq \infty$ )	$23 \times 10^{-6}$	$\leq 1.7 \times 10^{-15}$	$\leq 1$	$\leq 0.5$
QGP	$2 \times 10^{12}$	$\leq 5 \times 10^{11}$	-	$\leq 0.4$

# AdS/CFT and the KSS bound

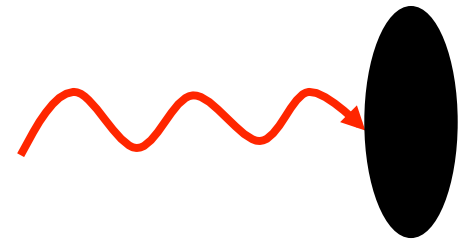
AdS/CFT conjecture: Maldacena

**Duality** between strongly coupled N=4 supersymmetric Yang-Mills theory at large coupling and Nc & classical 10 dimensional gravity in the background of D3 branes

KSS bound: Kovtun, Son, Starinets

Classical absorption cross-section of a graviton with energy  $\omega$  on a black brane

$$\sigma(\omega) = \frac{8\pi G}{\omega} \int dt d\mathbf{x} e^{i\omega t} \langle [T_{xy}(t, \mathbf{x}), T_{xy}(0, 0)] \rangle$$





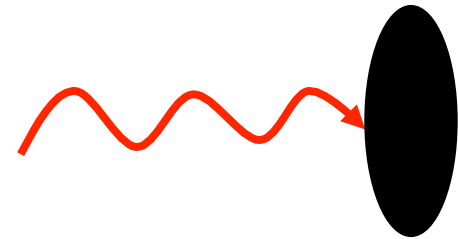
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From Kubo:  $\eta = \frac{\sigma(0)}{16\pi G}$

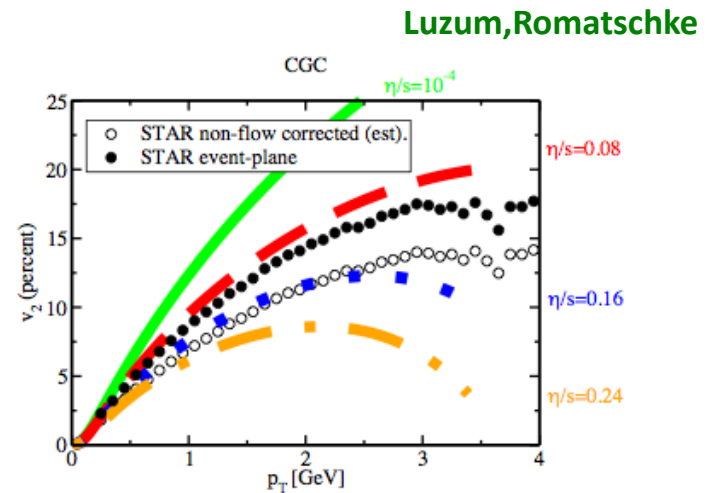
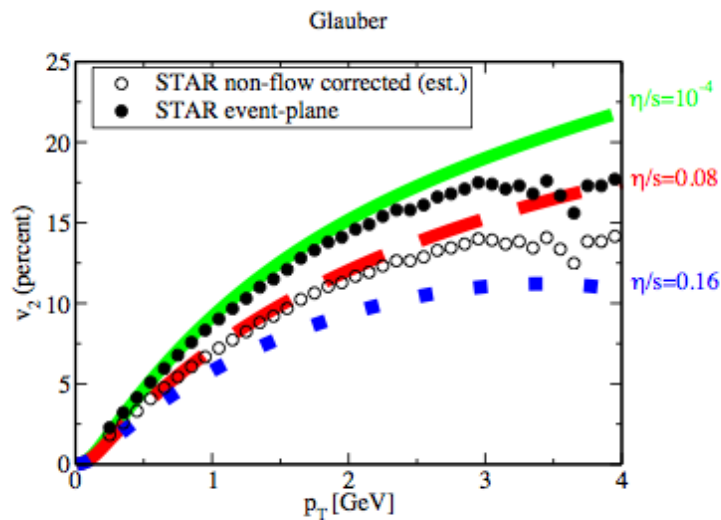
From Beckenstein:  $s = \frac{a}{4G}$

Theorem:  $\sigma(0) = a$

where  $a =$  area of black brane

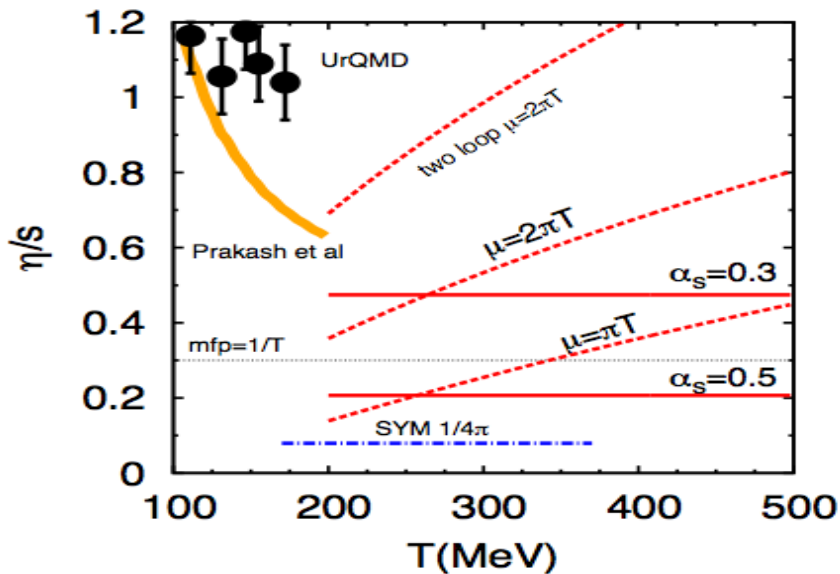
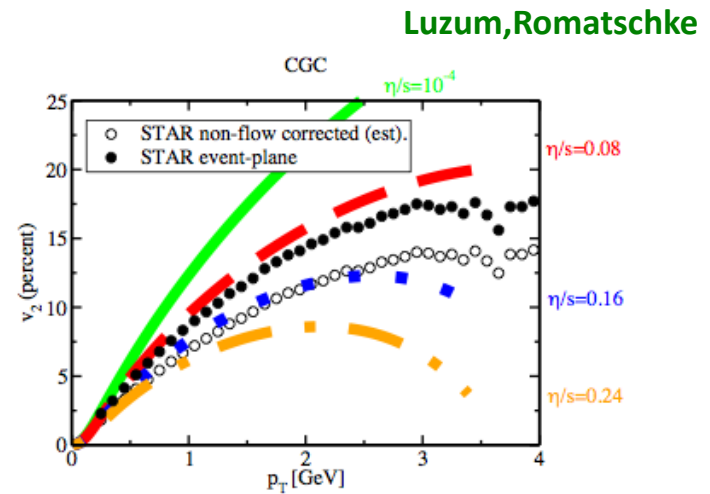
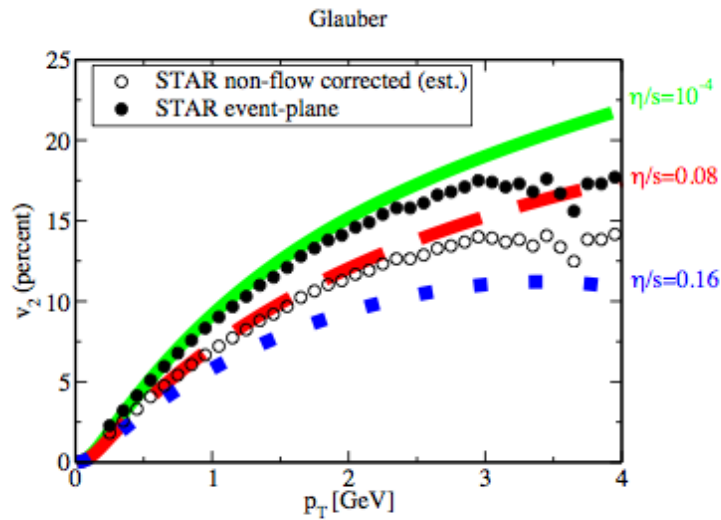
Putting it all together:  $\frac{\eta}{s} \geq \frac{\hbar}{k_B} \frac{1}{4\pi}$

# $\eta/s$ and viscous hydro @ RHIC



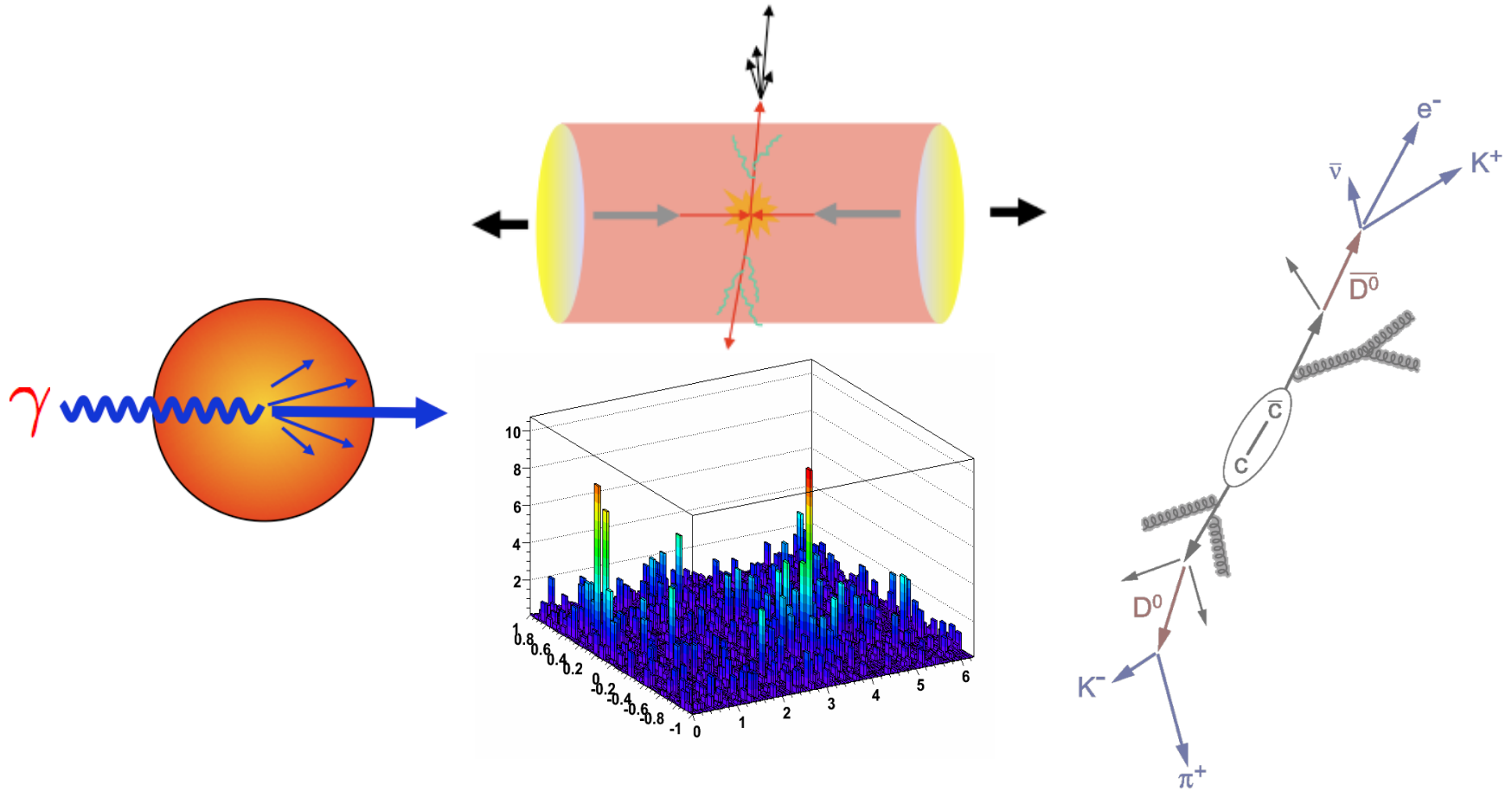
- $\eta/s$  from comparisons of viscous hydro codes to RHIC data suggest it nearly saturates the KSS bound – perfect fluid
- Quantitative value quite sensitive to initial conditions – Glasma!
- Probably safe to say its lower than any fluid with possible exception of cold  ${}^6\text{Li}$  atoms in the unitarity limit (Feshbach resonance)

# $\eta/s$ and viscous hydro @ RHIC



Viscosity from perturbative estimates is significantly larger

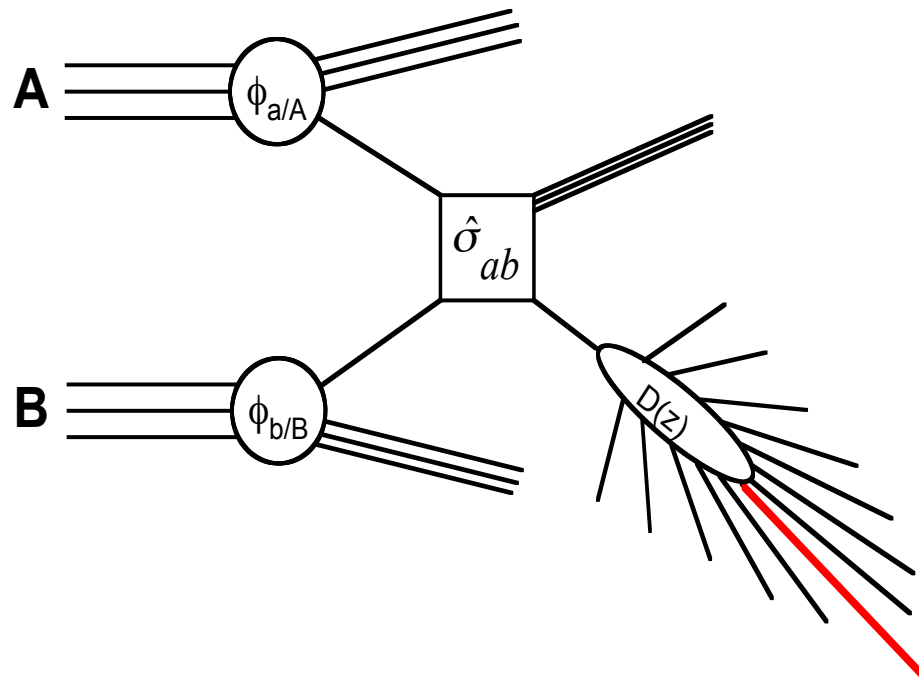
# Hard Probes of the sQGP



High  $p_T$  ( $\gg T$ ) hadrons, heavy mesons  
(charm and beauty),  $J/\psi$  and Upsilon, hard photons and JETS...

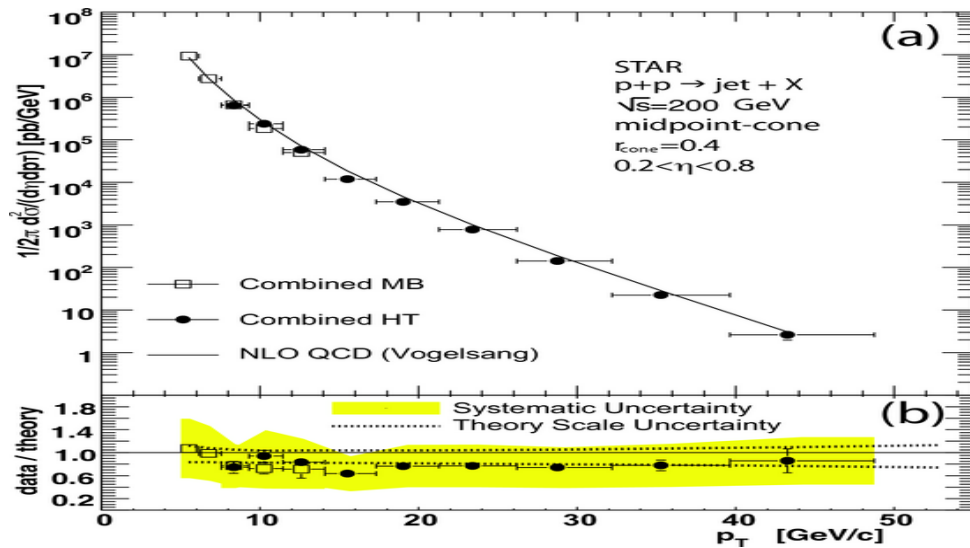
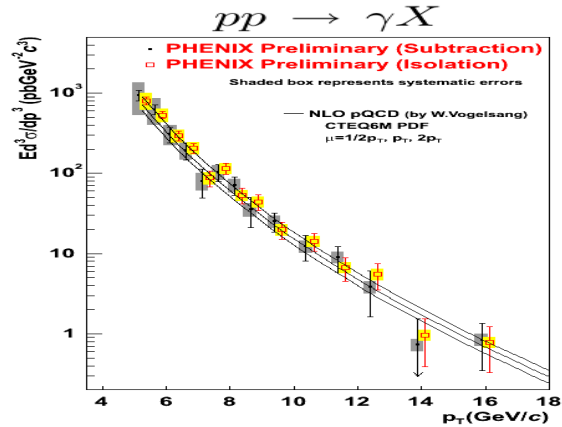
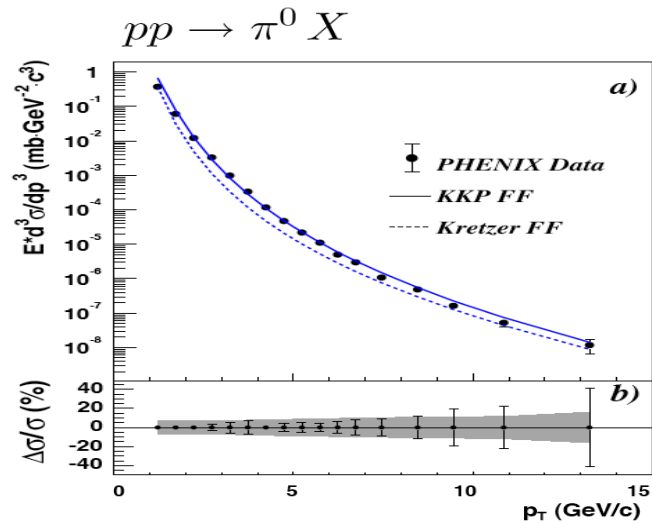
(Jet reconstruction techniques in HI collisions – M. Cacciari talk  
MCs for soft background- F. Krauss talk)

# The pp Benchmark



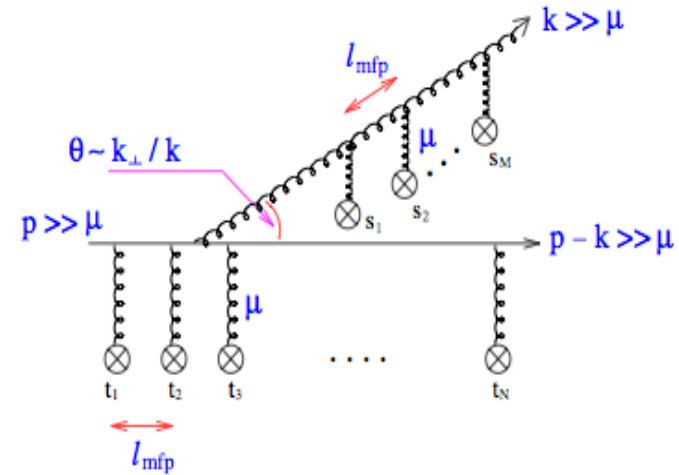
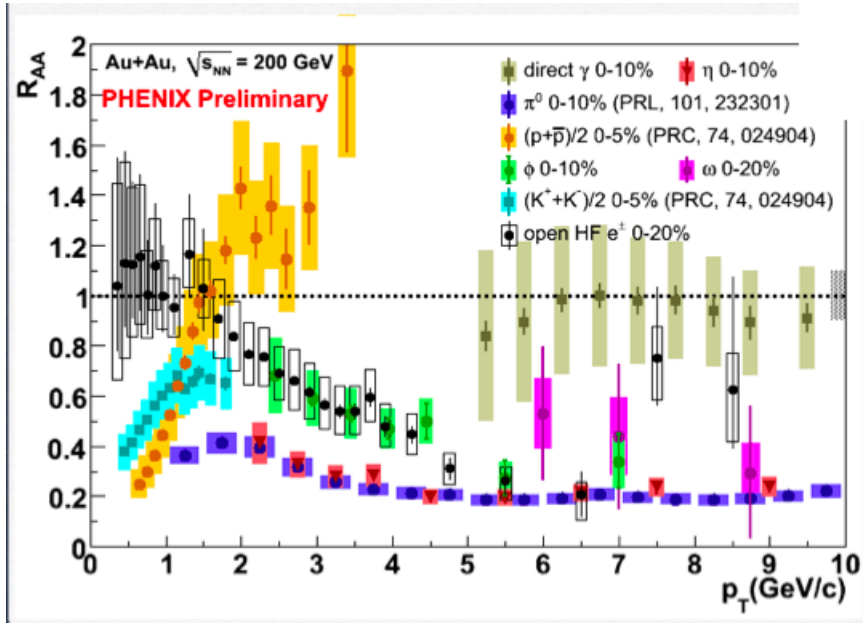
$$E \frac{d^3\sigma}{dp^3} = \sum_{abc} \int dx_a dx_b \phi_{a/A}(x_a, Q^2, \mu) \phi_{b/B}(x_b, Q^2, \mu) \frac{D_{\pi^0/c}(z, Q^2, \mu)}{z\pi} \frac{d\hat{\sigma}}{dt}$$

# The pp Benchmark



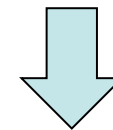
pQCD works exceptionally well @ RHIC to fairly low  $p_T$

# Parton energy loss



$$\frac{dE}{dz} = -\frac{\alpha_S N_c}{4} p_{\perp}^2$$

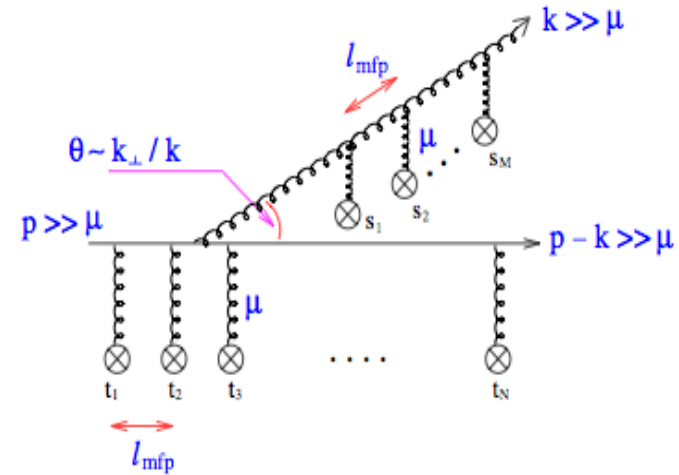
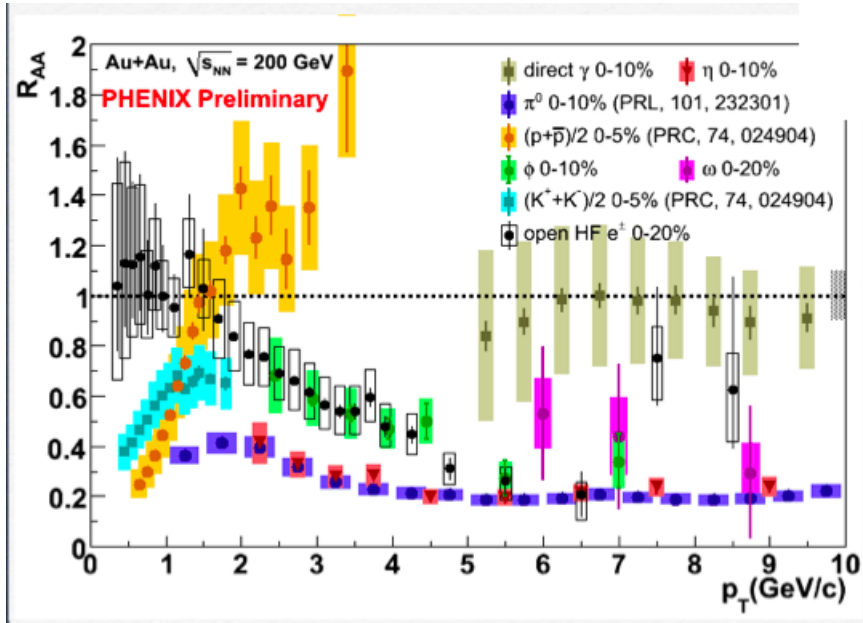
$$p_{\perp}^2 = \hat{q} L$$



Characterizes properties  
of hot medium

Bjorken; Gyulassy, Wang;  
Baier, Dokshitzer, Mueller, Peigne, Schiff

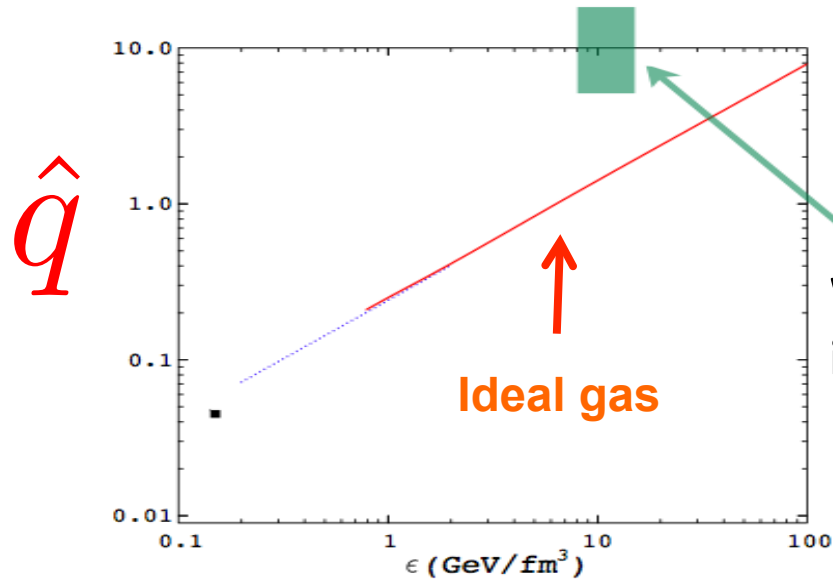
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# Weak versus Strong coupling in the sQGP

Weak coupling approaches fit the jet quenching data  
-- however, computations are not under control  
& large heavy quark energy loss is difficult to reconcile

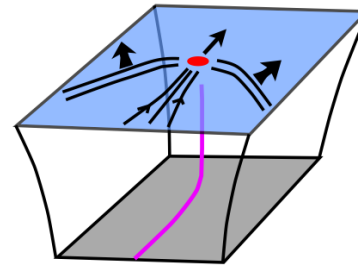
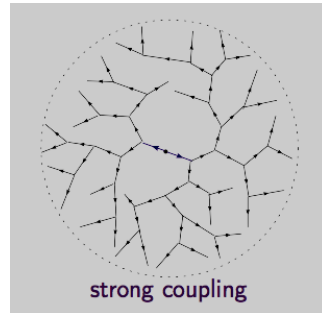
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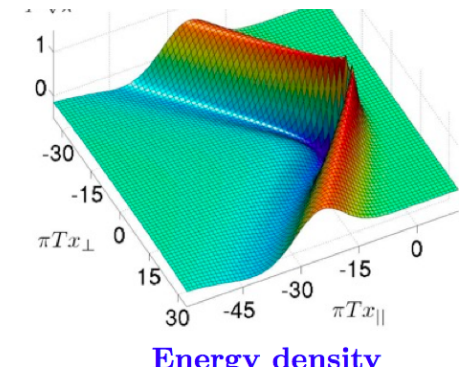
Large energy loss for light and heavy quarks “natural” in  
AdS/CFT approach

But no jets in AdS !

Hatta,Iancu,Mueller  
Hofman,Maldacena



Liu et al; Gubser et al;  
Casalderruy-Solana,Teaney;  
Chesler et al



# Weak versus Strong coupling in the sQGP

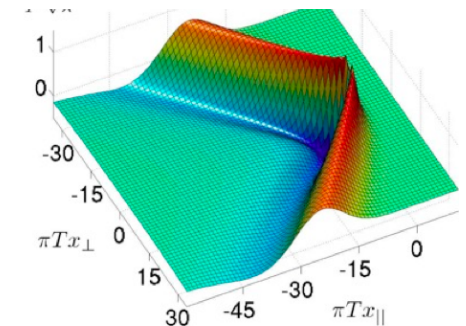
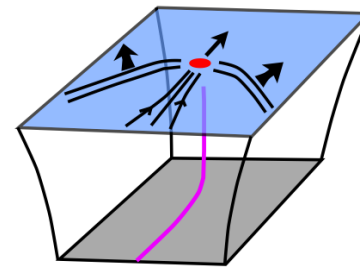
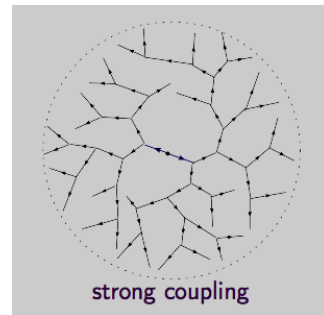
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Large energy loss for light and heavy quarks “natural” in AdS/CFT approach

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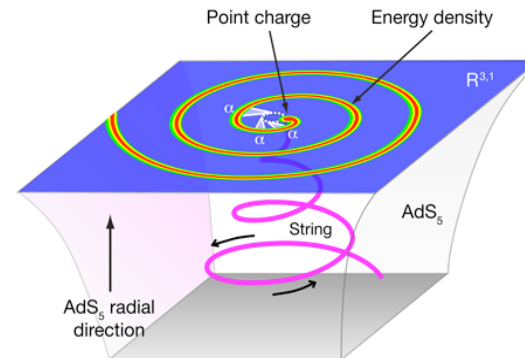
Hatta, Iancu, Mueller  
 Hofman, Maldacena



Energy density

Synchrotron radiation in AdS  
 may be more like in QCD-collimated jets

Chesler et al.

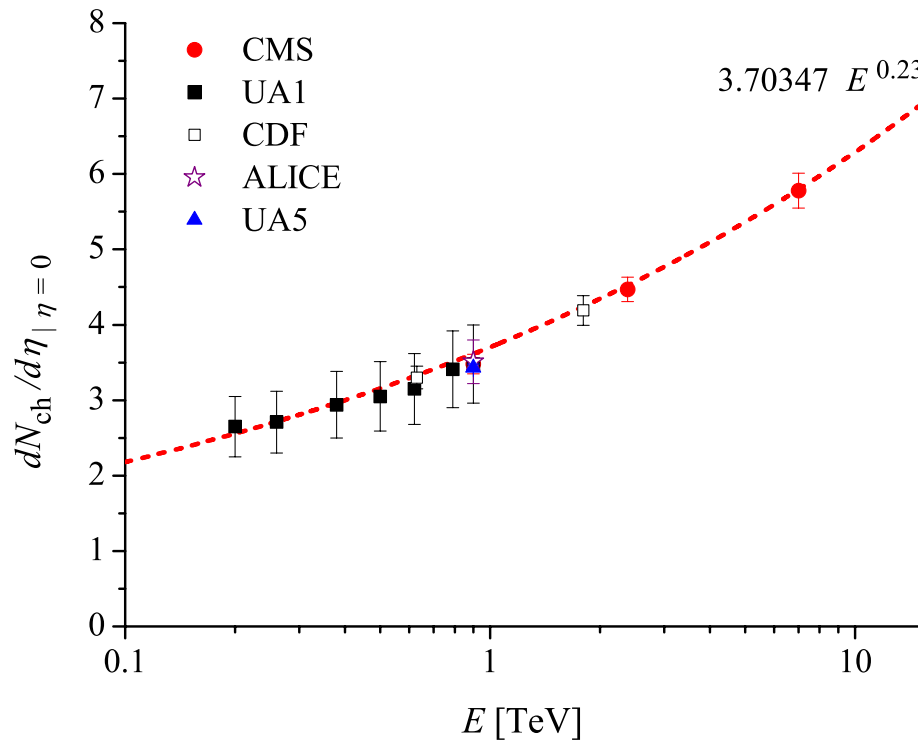


Class Window

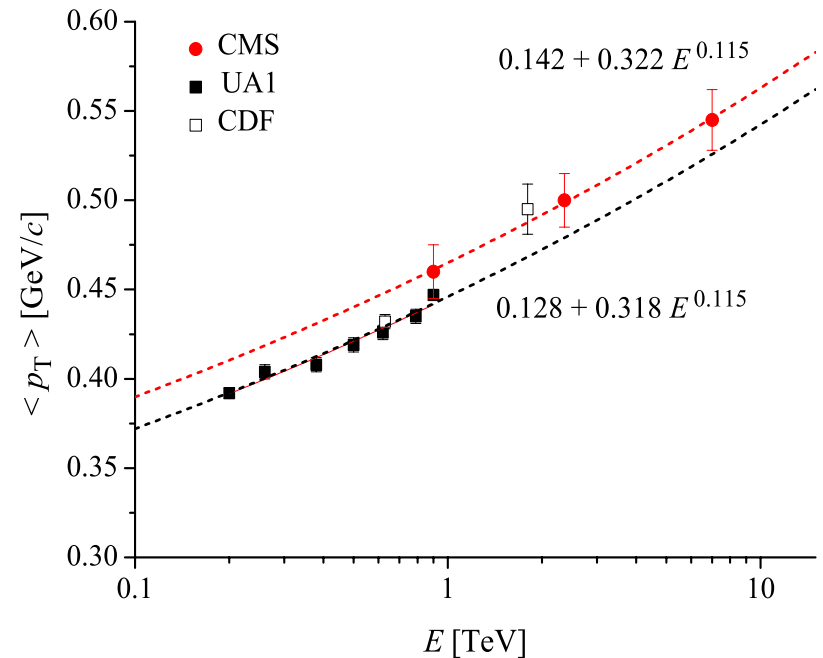
Lots of exciting ideas but no clear picture of transport in sQGP

# Looking forward: Heavy Ions at the LHC

Initial CGC based comparisons to LHC pp data are promising



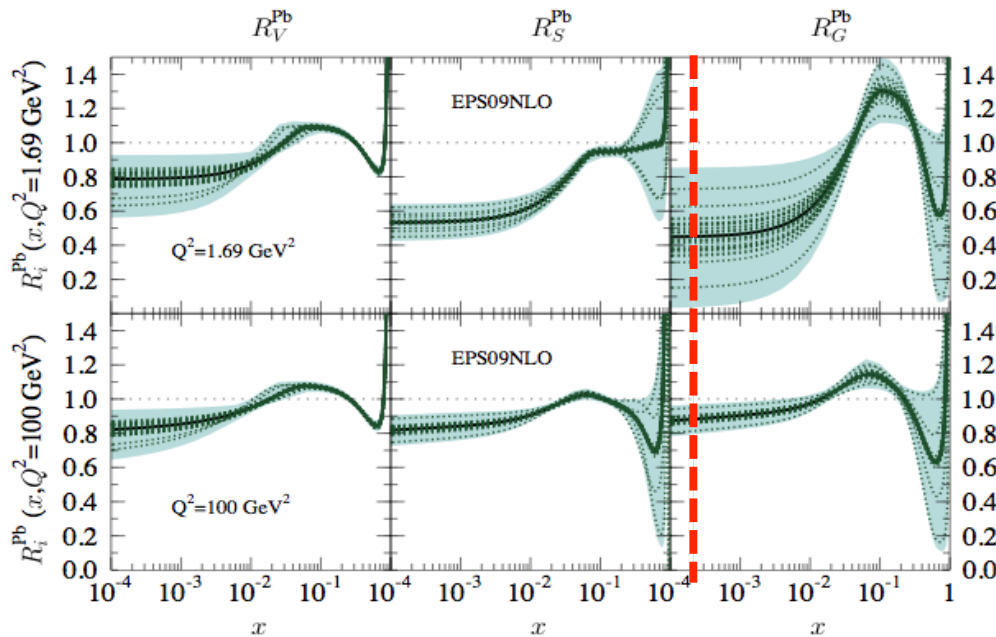
$p_T \propto Q_s(\sqrt{s})$   
 $dN/d\eta \propto Q_s^2(\sqrt{s})$



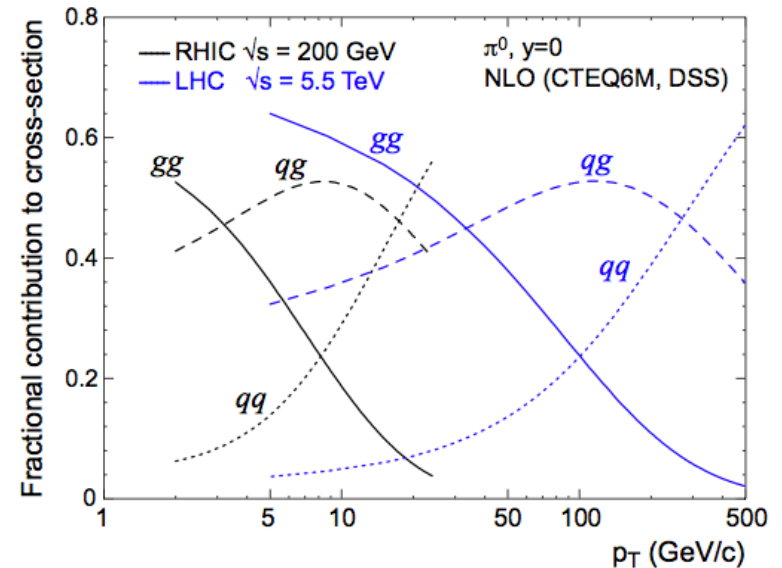
L. McLerran, M. Praszalowicz, arXiv:1006.4293  
 E. Levin, A. Rezaian, arXiv:1005.0631

# Looking forward: Heavy Ions at the LHC

$x \sim 5 \cdot 10^{-4}$  for semi-hard physics at LHC – large gluon shadowing



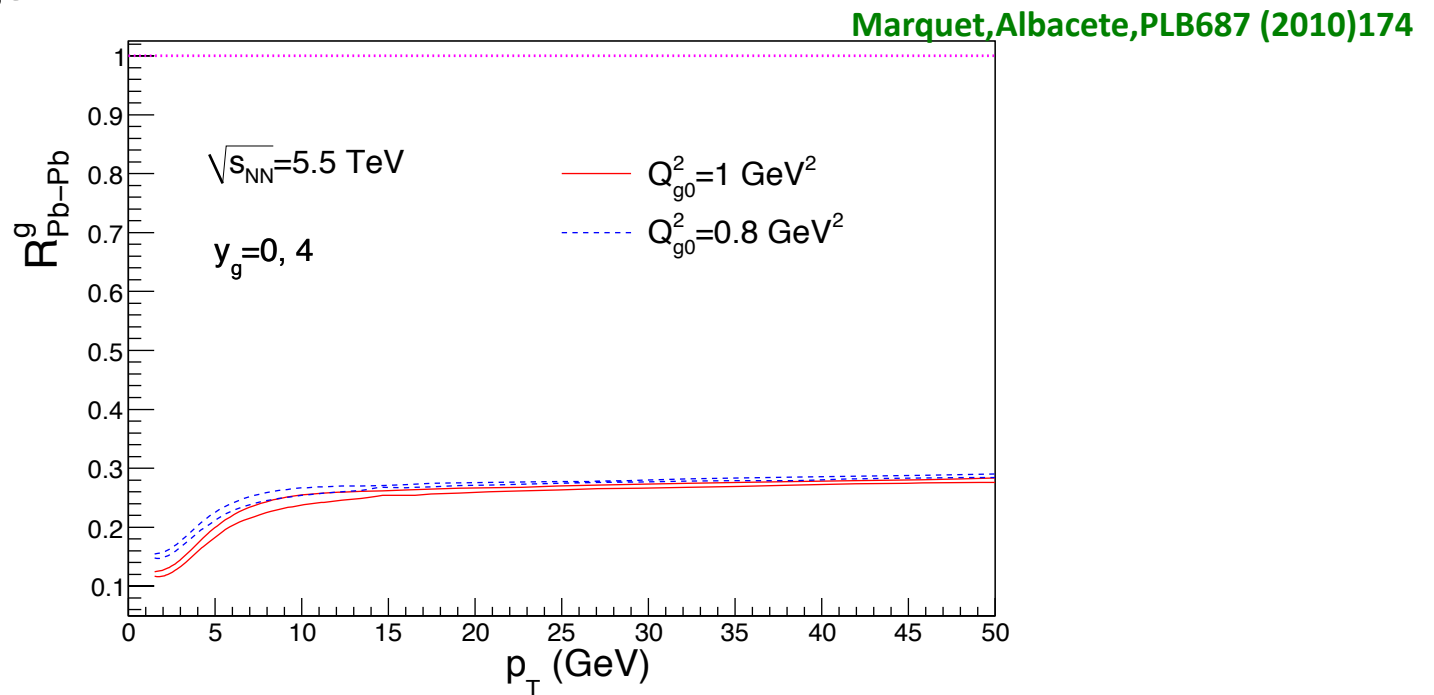
## H. Paukkunen's talk



d+A RHIC data and (possibly future) p+A LHC will help quantify results

# Looking forward: Heavy Ions at the LHC

Quenching of hadron spectrum (rel. to pp) from initial state effects alone is large

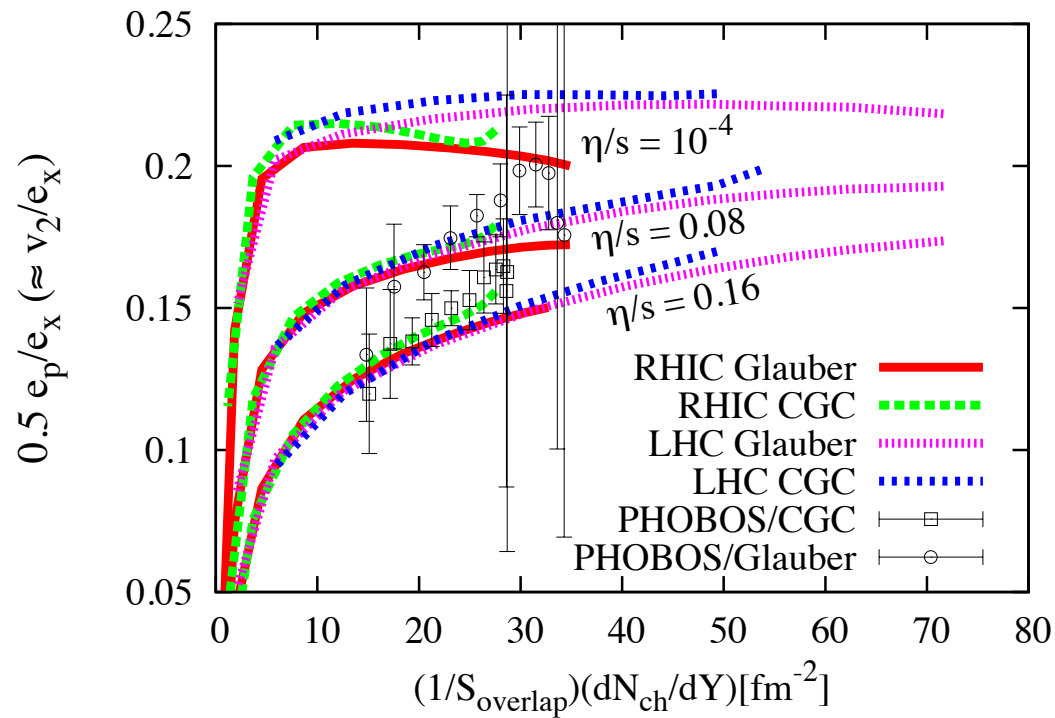


Challenging to extract sQGP properties –will need full array of impressive final states feasible

# Looking forward: Heavy Ions at the LHC

Early flow analyses – confirm if RHIC has saturated the “hydro bound”

Luzum, Romatschke, 0901.4588



# Outlook

- We now have a standard model of HI collisions with unanticipated features like the **CGC**, the **Glasma** & **sQGP**
- The sQGP produced at RHIC may be the most **perfect fluid** in nature
- Many approaches to understand and quantify further the properties of the sQGP – but no consensus yet
- **LHC Pb+Pb data and future RHIC upgrades will significantly enrich (alter?) the picture outlined here**



**THANK YOU FOR YOUR ATTENTION!**

# Strong flow = (nearly) ideal hydrodynamics

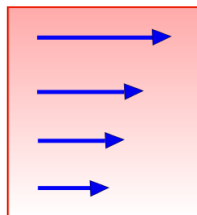
- QCD Hydrodynamics = low energy **effective theory of QCD**  
valid at large distances ( $L \gg \lambda_{\text{mfp}}$ )

- Relativistic hydrodynamics:  
Conservation laws  $\partial_\mu T^{\mu\nu} = 0$        $\partial_\mu j^\mu = 0$

- Constitutive equations:  $T^{\mu\nu} = (\varepsilon + P)U^\mu U^\nu + P g^{\mu\nu} + \tau^{\mu\nu}$   
 $j^\mu = n U^\mu + \nu^\mu$

- Dissipative terms:  $\tau^{ij} = -\eta \left( \partial^i U^j + \partial^j U^i - \frac{2}{3} \delta^{ij} \vec{\nabla} \cdot \vec{U} \right) - \zeta \delta^{ij} \vec{\nabla} \cdot \vec{U}$

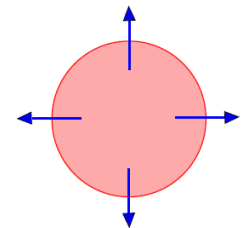
Shear viscosity



$$\nu^i = -\sigma T \partial^i \left( \frac{\mu}{T} \right)$$

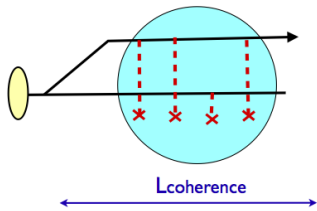
Conductivity (diffusion)

Bulk viscosity

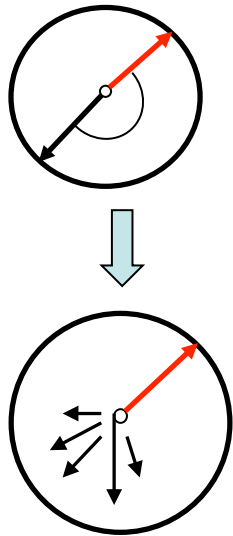


# Evidence for Saturation @ RHIC ?

## Rich Deuteron+Gold physics program



Talks by J. Albacete  
& C. Marquet  
Caveats: M. Sumbara



Away side parton  
randomized  
by strong color field

