

NLO Event Generation with SHERPA

Marek Schönherr¹

IKTP TU Dresden

30/06/2009



¹for SHERPA: J. Archibald, T. Gleisberg, S. Höche, F. Krauss, MS, S. Schumann, F. Siegert, J. Winter

Contents

- 1 NLO Cross Section Calculations
- 2 The POWHEG Method
- 3 First Results

NLO Cross Section

- LO cross section

$$\sigma_{LO} = \int_N d\sigma_B$$

- NLO cross section

$$\sigma_{NLO} = \int_N d\sigma_B + \int_N d\sigma_V + \int_{N+1} d\sigma_R$$

- $d\sigma_V$ may be IR-divergent due to massless loop propagators
- $\int_1 d\sigma_R$ may be IR-divergent due to emission of massless line
→ sum is finite (KLN-theorem)

NLO Cross Section

Problem: $d\sigma_V$ and $d\sigma_R$ live in different phase spaces

- in $d\sigma_V$ the loop integral is IR divergent
 - every point in the N -dim. phase space carries the divergence structure as poles in $\frac{1}{\epsilon}$ and $\frac{1}{\epsilon^2}$ (in dim. renormalisation)
- for $d\sigma_R$ the phase space integral over the extra emission is IR divergent
 - a similar pole structure cannot be achieved by Monte Carlo integration (integration in integer number of dimensions only)

→ NLO cross section cannot be directly integrated using Monte Carlo methods

→ no obvious cancellation of divergences

Subtraction Schemes

$$\sigma_{NLO} = \int_N d\sigma_B + \int_N d\sigma_V + \int_{N+1} d\sigma_R$$

individual terms rendered integrable by subtraction scheme

→ needs to provide local counter terms $d\sigma_A$ which are integrable over extra emission phase space

$$\sigma_{NLO} = \int_N d\sigma_B + \int_N \underbrace{\left[d\sigma_V + \int_1 d\sigma_A \right]}_{\text{IR-finite loop integration}} + \int_{N+1} \underbrace{\left[d\sigma_R - d\sigma_A \right]}_{\text{IR-finite PS integration}}$$

The POWHEG Method

- generate N -jet event with

$$d\sigma_B + d\sigma_V + d\sigma_I + \int_1 [d\sigma_R - d\sigma_S]$$

- shower radiates real emission with Sudakov form factor

$$\Delta_t = \exp \left[- \int \theta(t_r - t) \frac{R}{B} d\Phi_r \right]$$

→ shower needs to cover the whole phase space

Available in SHERPA

- automated Catani-Seymour subtraction term generation (in AMEGIC++) [Gleisberg, Krauss: Eur.Phys.J.C53\(2008\)501-523](#)
- LO ME generators (AMEGIC++, COMIX)
[Krauss, Kuhn, Soff: JHEP02\(2002\)044](#), [Gleisberg, Hoeche: JHEP12\(2008\)039](#)
- few built-in virtual MEs + extensive tensor integral library
- interfaces to NLO codes (BLACKHAT, GOLEM, etc.)
[Bern et al: arXiv:0902.2760](#), [Binoth et al: arXiv:0810.0992](#)
- parton shower (CSSHOWER++) [Schumann, Krauss: JHEP03\(2008\)038](#)

→ rearrange terms to give N -jet event with NLO weight

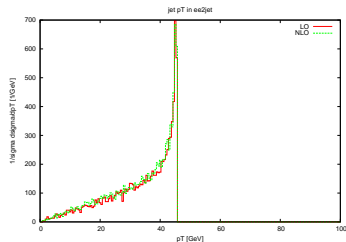
→ get correct real emission via POWHEG method

First Results on $ee \rightarrow 2jets$ at 91.25 GeV

fixed order result:

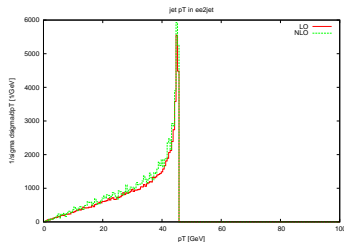
$$d\sigma_{NLO} = \left(1 + \frac{\alpha_S}{\pi}\right) d\sigma_{LO} \quad \text{with } (J = 1)$$

unweighted event generation



LO: $2e^-e^+ \rightarrow j\bar{j}$:
 NLO: $2e^-e^+ \rightarrow j\bar{j} + QCD(BVIRS)$:

weighted event generation



40189.9 pb +- (18.29 pb = 0.045509 %)
 41712.5 pb +- (20.79 pb = 0.049842 %)

Further Ideas

Resummation Effects:

- determine k_{\perp}^2 of most probable splitting for real emission
→ shower language
- consistently replace one power of α_S with $\alpha_S(k_T^2)$
→ capture resummed large logarithms from higher orders
→ enlarged possibility of softer emission

LO:	$2_{-2}e^{-}_{-}e^{+}_{-}j_{-}j_{-}$:	40189.9 pb +- (18.29 pb = 0.045509 %)
NLO:	$2_{-2}e^{-}_{-}e^{+}_{-}j_{-}j_{-}$ QCD(BVIRS) :	41712.5 pb +- (20.79 pb = 0.049842 %)
NLO+NLL:	$2_{-2}e^{-}_{-}e^{+}_{-}j_{-}j_{-}$ QCD(BVIRS) :	42757.4 pb +- (20.41 pb = 0.047748 %)

Conclusion

Done:

- consistent use of already built-in components in new approach
- integration of $R + S$ converges
- numbers seem to be right
- inclusion of resummed corrections real extra emission via α_S rescaling

Todo:

- determine $\left(\frac{R}{B}\right)_{max}$ for POWHEG matching
- imposing upper cut-off J on real emission to make cross sections exclusive for multijet merging