

Do we need HOM dampers on superconducting cavities in p linacs?

“Yes, we ~~do~~ can”: a **personal** view(*)

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11-12 Dec 08

(*) where I disagree with many people but agree in essence with retired ‘heavy-weights’ Ernst Haebel and Ron Sundelin
... however, this presentation is my personal justification for it ...

‘Prelude’

”The question is already settled, have a look at ..”

... a thesis work (author unfortunately not present, but he knows)

It examines only the **power** coupled out from the beam.

NO DIRECT BEAM STABILITY CONSIDERATIONS !!!

Conclusion: HOM coupler are not necessary for (p-linac)

The treatment of the power-aspect is all OK but there is a

Logical Problem:

One cannot ignore the **main benefit** of an object – **impedance reduction** – and conclude from the study of secondary effects – power outflow – that one does not need this object !

(Maybe one really doesn’t need it, but everything has to be included in a valid analysis !)

—> **work does NOT answer the question**

Sneak in the power study as beam instability study:

“To drive the beam unstable, one needs sufficiently high fields.” (**agreed**)

“The latter can only be produced when there is an HOM sitting on a machine line (.., what else G. Clooney ?)”

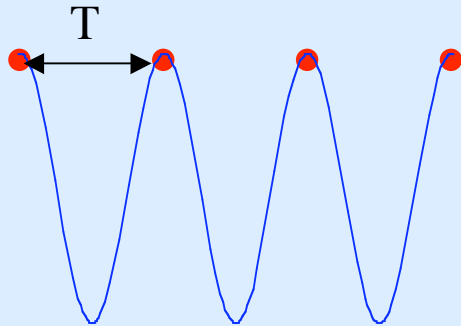
NO, we shall see why!

(“The probability for a high- Q_{ext} mode hitting such a machine line (within a BW...) is very, very low, hence”).

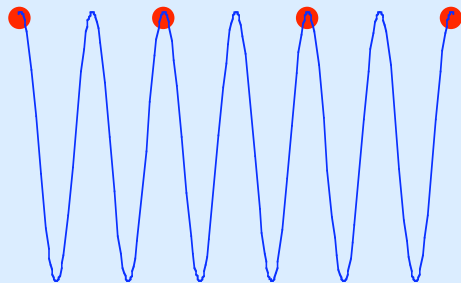
(“In SNS no power output seen” **Bad argument: damping was good enough to keep it so low !!**

(**Agreed fact: the more one damps the less comes out!**)

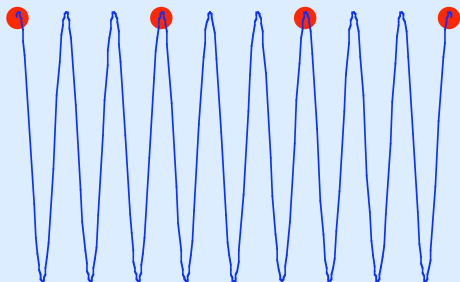
A **continuous** sine-wave has **only one** resonant frequency: its own.
But bunches arrive **stroboscopically**: 1 bunch each time **T** that
'tick' the observing oscillator: spectrum analyzer, cavity mode, ...



1st machine line: 1 oscillation
between 2 consecutive bunches
Resonant on $f_{ML} = 1/T$



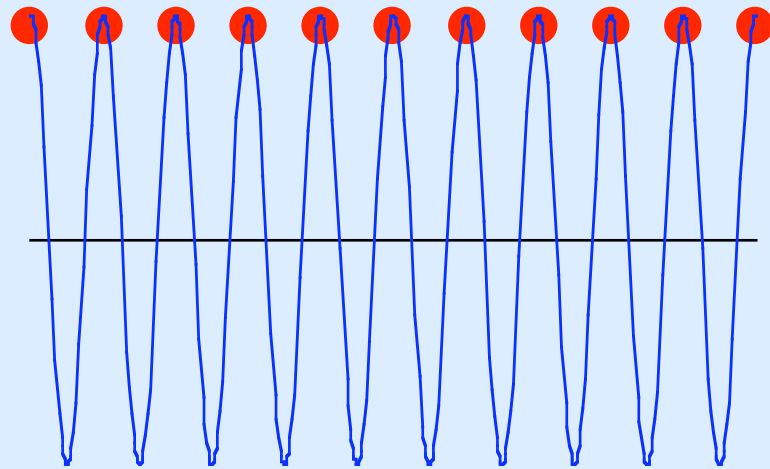
2nd machine line: 2 oscillation
between 2 consecutive bunches
Resonant on $f_{ML} = 2/T$



3rd machine line: 3 oscillation
between 2 consecutive bunches
Resonant on $f_{ML} = 3/T$

Bunches have transverse displacement: Express any pattern by
Fourier components: **harmonic position modulation**

- Example 1: Constant (transverse) displacement



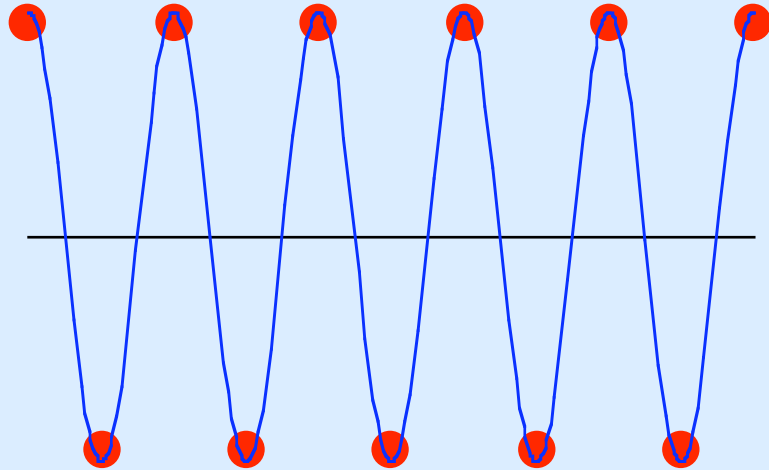
Bunch n : **$t_n = n \cdot T$**

Displacement **$x_n = 1 = \cos(0 \cdot 2\pi n)$**

Resonant oscillation: **$\cos(\omega_0 \cdot t)$**

$f = f_0 = 1/T$; $\omega_0 = 2\pi \cdot f_0$

- Example 2: Up-down (right-left) modulation (' $\xi = 1/2$ ')



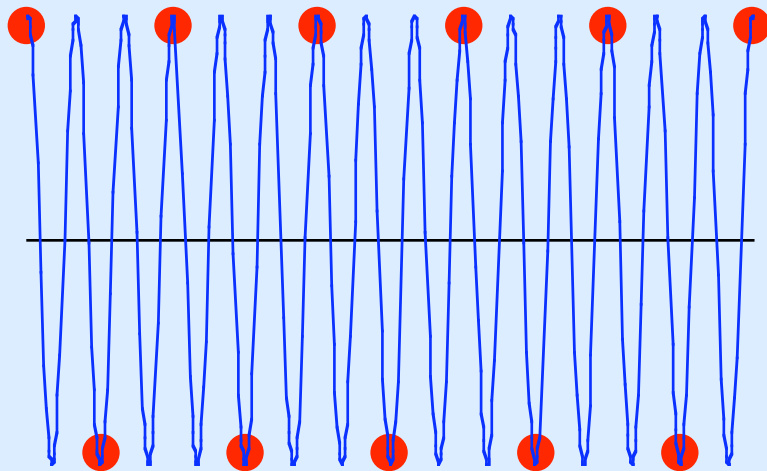
Bunch n: $t_n = n \cdot T$

Displacement $\cos(0.5 \cdot 2\pi \cdot n)$

Res. Osc. $\cos((1-0.5) \cdot \omega_0 \cdot t)$

$$f = f_0 \cdot (1-0.5) = 0.5 \cdot f_0$$

Lower sideband: Not on ML



Bunch n: $t_n = n \cdot T$

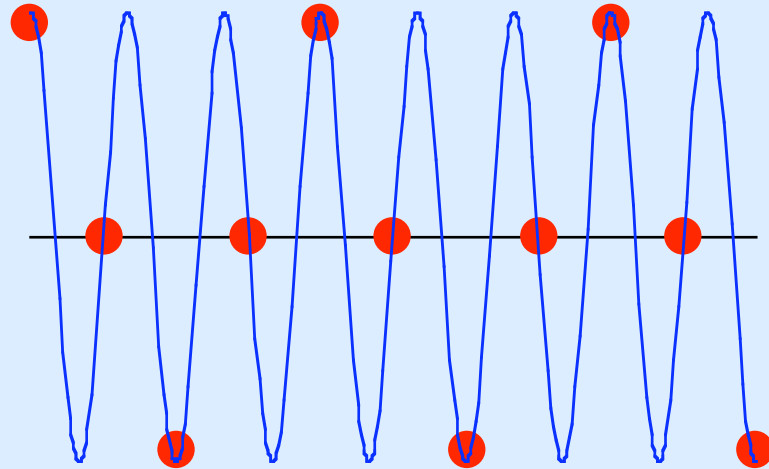
Displacement $\cos(0.5 \cdot 2\pi \cdot n)$

Res. Osc. $\cos((1+0.5) \cdot \omega_0 \cdot t)$

$$f = f_0 \cdot (1+0.5) = 1.5 \cdot f_0$$

Upper sideband: Not on ML

- Example 3: Up-zero-down-zero modulation (' $\xi = 1/4$ ')



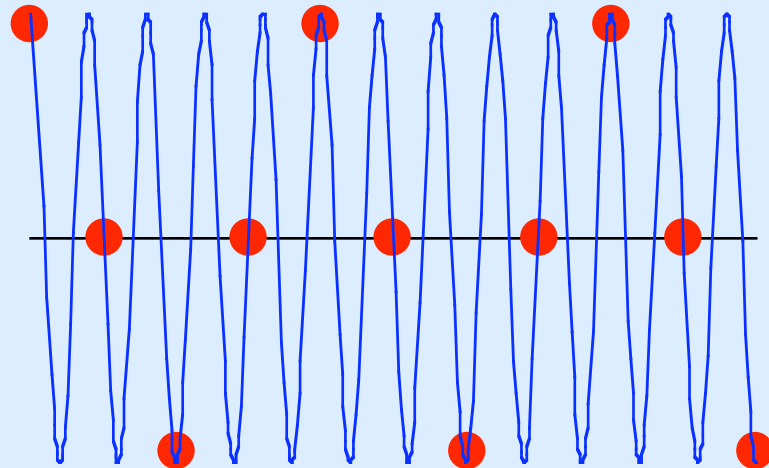
Bunch n: $t_n = n \cdot T$

Displacement $\cos(0.25 \cdot 2\pi \cdot n)$

Res. Osc. $\cos((1-0.25) \cdot \omega_0 \cdot t)$

$f = f_0 \cdot (1-0.25) = 0.75 \cdot f_0$

Lower sideband : Not on ML



Bunch n: $t_n = n \cdot T$

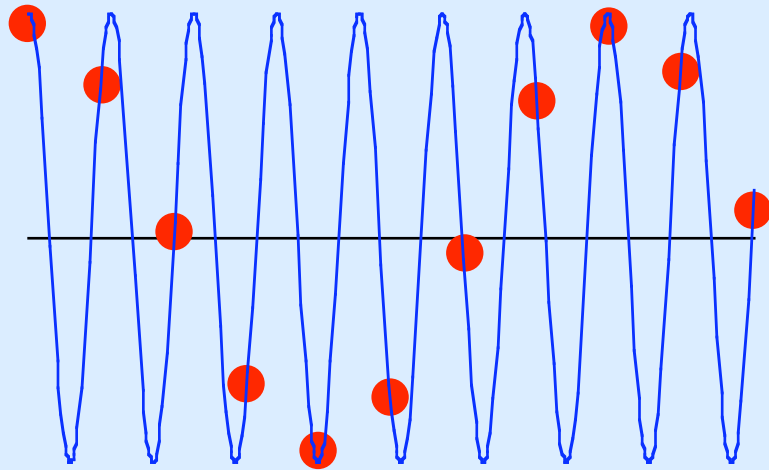
Displacement $\cos(0.25 \cdot 2\pi \cdot n)$

Res. Osc. $\cos((1+0.25) \cdot \omega_0 \cdot t)$

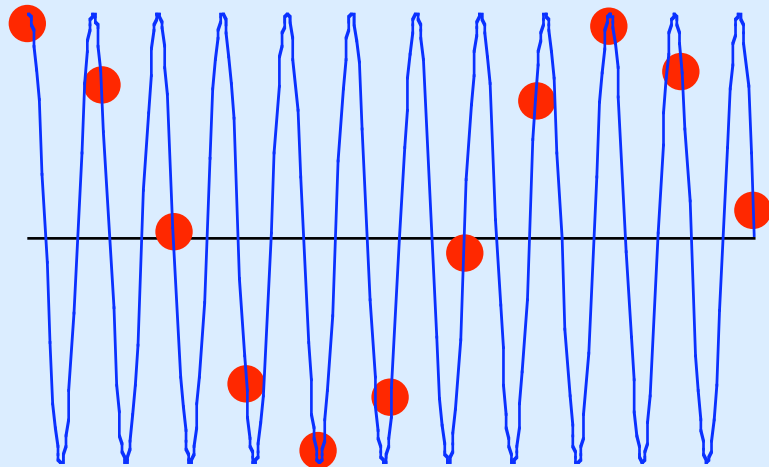
$f = f_0 \cdot (1+0.25) = 1.25 \cdot f_0$

Upper sideband : Not on ML

- Example 4: **Anything** (even irrational number) (' $\xi = 0.123$ ')

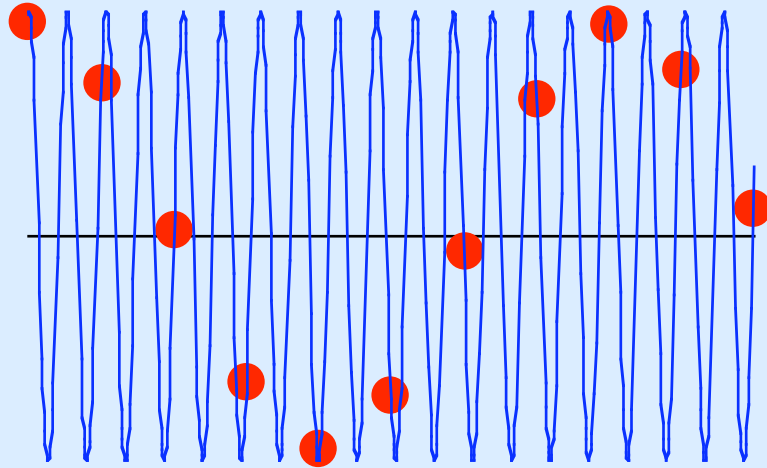


Bunch n: $t_n = n \cdot T$
 Displacement $\cos(0.123 \cdot 2\pi \cdot n)$
 Res. Osc. $\cos((1-0.123) \cdot \omega_0 \cdot t)$
 $f = f_0 \cdot (1-0.123) = 0.877 \cdot f_0$
 Lower sideband : Not on ML



Bunch n: $t_n = n \cdot T$
 Displacement $\cos(0.123 \cdot 2\pi \cdot n)$
 Res. Osc. $\cos((1+0.123) \cdot \omega_0 \cdot t)$
 $f = f_0 \cdot (1+0.123) = 1.123 \cdot f_0$
 Upper sideband : Not on ML

- Example 5: An additional integer number of oscillations
(between other machine lines)



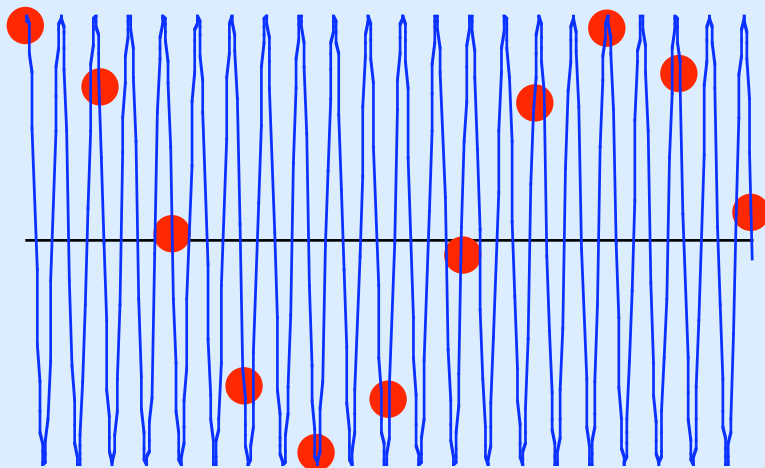
Bunch n: $t_n = n \cdot T$

Displacement $\cos(0.123 \cdot 2\pi \cdot n)$

Res. Osc. $\cos((2-0.123) \cdot \omega_0 \cdot t)$

$f = f_0 \cdot (2-0.123) = 1.877 \cdot f_0$

Lower sideband : Not on ML



Bunch n: $t_n = n \cdot T$

Displacement $\cos(0.123 \cdot 2\pi \cdot n)$

Res. Osc. $\cos((2+0.123) \cdot \omega_0 \cdot t)$

$f = f_0 \cdot (2+0.123) = 2.123 \cdot f_0$

Upper sideband : Not on ML

Numerical example:

Given: Bunch repetition rate 350 MHz, $T=1/350\text{MHz}$
Assume $f_{\text{HOM}} = 1234.567890 \text{ MHz}$ (any number)

$$1234.567.. / 350 = 3.527... = 3+0.527 (= 4 - 0.473)$$

f_{HOM} above 3rd, below 4th machine line (ML)

$$1234.567... - 3*350 = 184.567... \text{ [MHz] above 3rd ML}$$
$$(1234.567... - 4*350 = -165.433... \text{ [MHz] below 4th ML})$$

Bunch pattern has $\xi=0.527$ oscillations per T (527 on 1000)

For ANY mode frequency exists matching pattern !!!

Possible excitation of an Instability (circular machines):

- Any HOM-f has a matching pattern: hence any HOM-f can be resonantly excited provided this pattern is present (*)
- If an HOM is excited, it ‘momentum modulates’ (kicks) the passing beam, **exciting this matching pattern**
- ‘Momentum modulation’ transforms (over ‘drift’, focussing) into ‘position modulation’
- In **circular machines** the ‘position modulated’ beam passes the cavity(s) again (and again), possibly increasing the HOM field.
- An increased HOM field may increase ‘momentum modulation’
 - > causes **well known Coupled Bunch Instability** (even normal conducting cavities – low Q – do the job !!)
 - > **HOM damping (sc. cav) in circular machines accepted fact**

(*) neglect shift by betatron frequency, not an essential point for the present discussion

Claim: It is improbable that HOM-f and a given pattern match

Wrong argument:

- A once existing HOM excites THE matching pattern
 - This pattern excites HOM on the initially exciting frequency
- > not ‘independent events’, **probability is not applicable**

In circ. machine HOM amplitude A increases by factor g per turn:

$$A_n = A_0 \cdot g^n \quad \dots \text{ and pattern amplitude with it}$$

Where is first A_0 coming from ? **Noise on the (injected bunches) !**

HOM filters its own frequency out of the noise: initial step A_0

Example: Beam of 2.5 mA at $T=1/350$ MHz : **10^8 protons per bunch**

Average position of 10^8 particles has relative scatter of 10^{-4} :

Bunch centre position jitters by 10^{-4} bunch-lengths

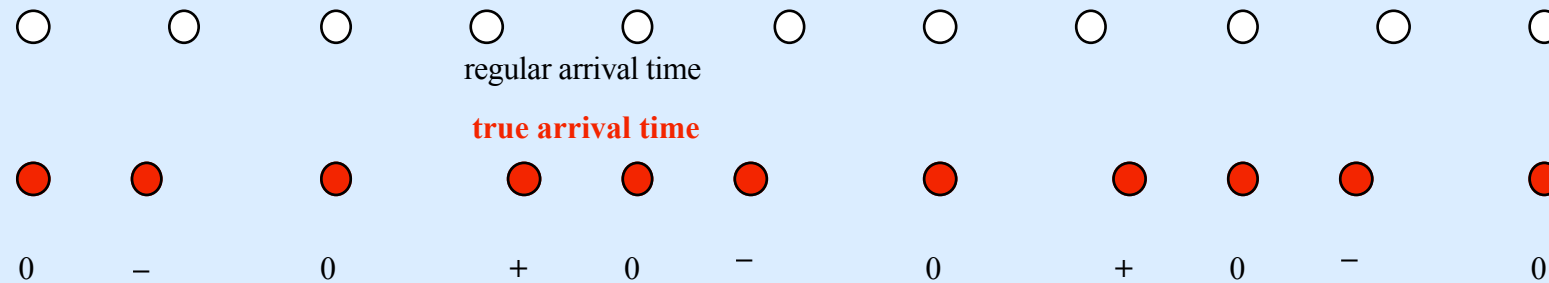
(cum grano salis)

There is much more ... RF noise on main field, injector jitter, ...

no hope for $A_0 \equiv 0$

CBI exists also in longitudinal plane (easier to draw transverse):
 Position modulation by time advance/delay of bunch arrival:

Longitudinal position modulation



$$t_n = T \cdot (n + a \cdot \cos(2\pi \cdot n \cdot \xi)) \quad \text{arrival time modulation}$$

at k^{th} machine line (assume $a \ll 1$)

$$\begin{aligned} f_{\text{ML}} &= k \cdot f_0 && \text{amplitude 1} \\ f_{\text{SB}} &= (k \pm \xi) \cdot f_0 && \text{amplitude } a/(2\pi) \end{aligned}$$

Is a similar mechanism possible in a linac ?

Seems not: each bunch encounters each cavity only once, but

- HOM gets excited by the (tiny) noise on the arriving bunches (the complex voltage vector ‘random walks’ in the complex plane, the ‘noise power’ drives $\langle V^2 \rangle$ linearly up in time till compensated by *damping*: not reached for high Q_{ext} modes)
- Excited cavities ‘kick’ later bunches imposing the matching ‘momentum pattern’ sent to more downstream cavities
- Over drift/focussing transforms into ‘displacement pattern’
- More downstream cavities are coherently excited (Coherent: time of flight from cavity M to cavity N constant)
- even more downstream cavities get even more excited

Much simplified model to understand principle:

- Each cavity is driven only by the beam deviation due to nearest upstream cavity with very **weak** coupling factor κ

(neglect more upstream, transported through focusing system (*))

- HOM field changes by complex factor d during T , $|d| \approx 1$

d complex: HOM not on ML!!

(Consider noise only on first cavity; study first 3 cavities: U, V, W)

$$U^{(n+1)} = d \cdot (U^{(n)} + \delta U^{(n)}); \quad \delta U^{(n)} = \text{random noise signal}$$

$$V^{(n+2)} = d \cdot (V^{(n+1)} + \kappa * \text{Re}(U^{(n+1)}));$$

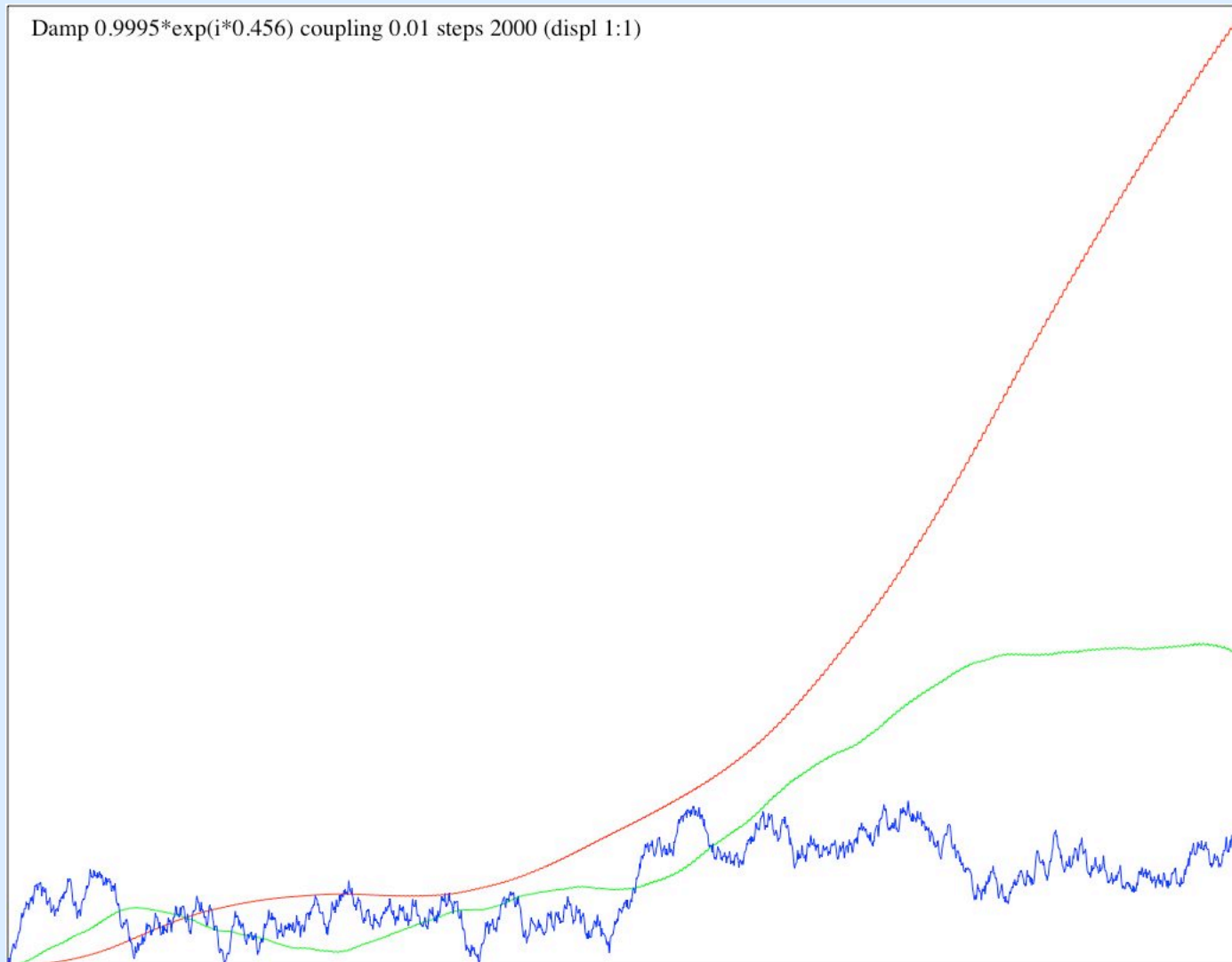
$$W^{(n+3)} = d \cdot (W^{(n+2)} + \kappa * \text{Re}(V^{(n+2)})); \quad \text{can be analyzed mathemat.}$$

(1 ms beam pulse at 352 MHz: 350,000 bunches)

(*) All deviations coming from more upstream are mutually coherent: All particles have

1) fixed time of flight 2) see same focusing system

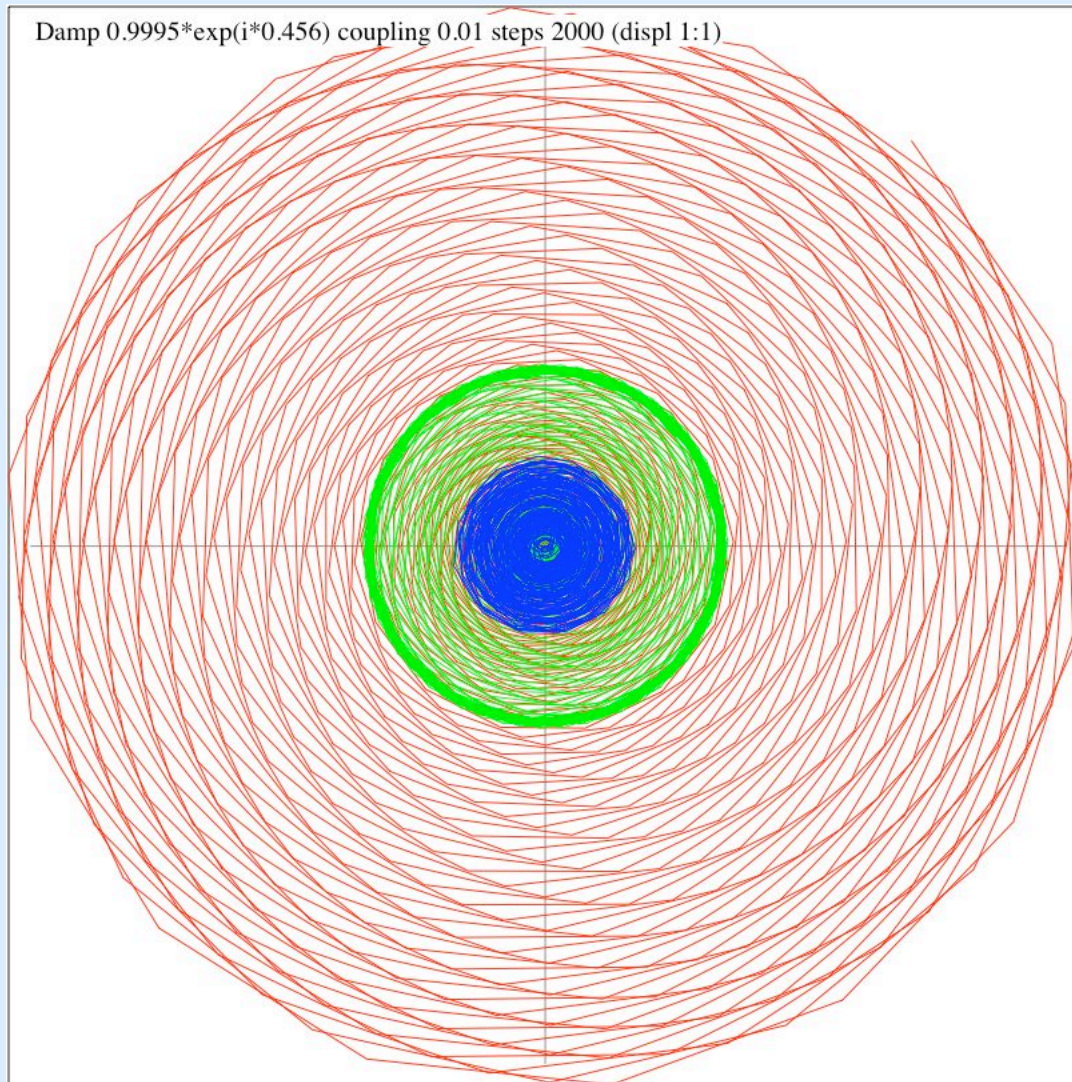
|V| as function of time/bunch number, 2000 bunches



U:Cavity 1 **V:Cavity 2** **W:Cavity 3**

Time-wise 'increase' and also increase cavity-to-cavity

Polar display of complex cavity voltage: reference main Main Line

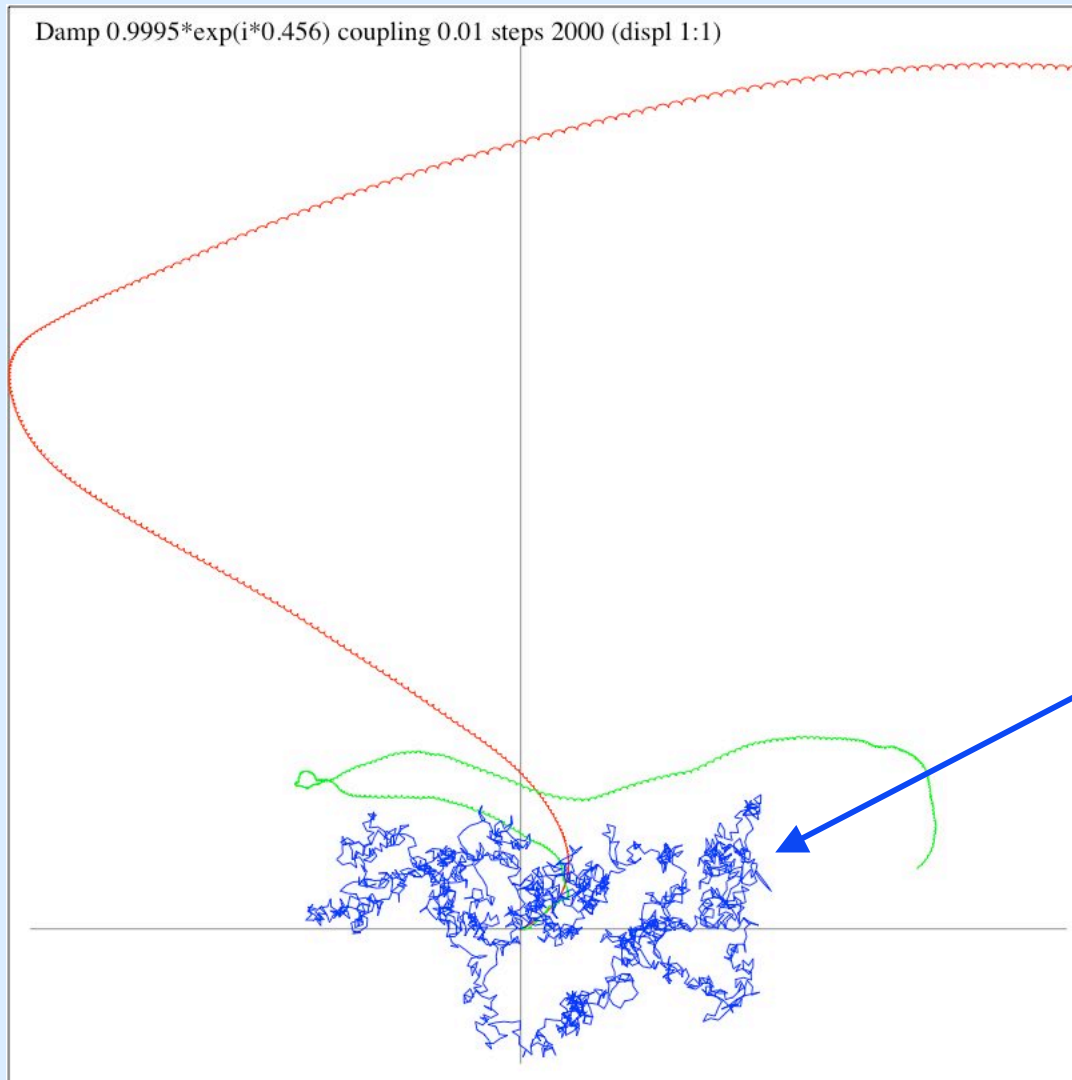


↕
Each vector is
turned by a further
complex d per step

2000 steps

U:Cavity 1 **V:Cavity 2** **W:Cavity 3**

Polar display of complex cavity voltage: reference HOM

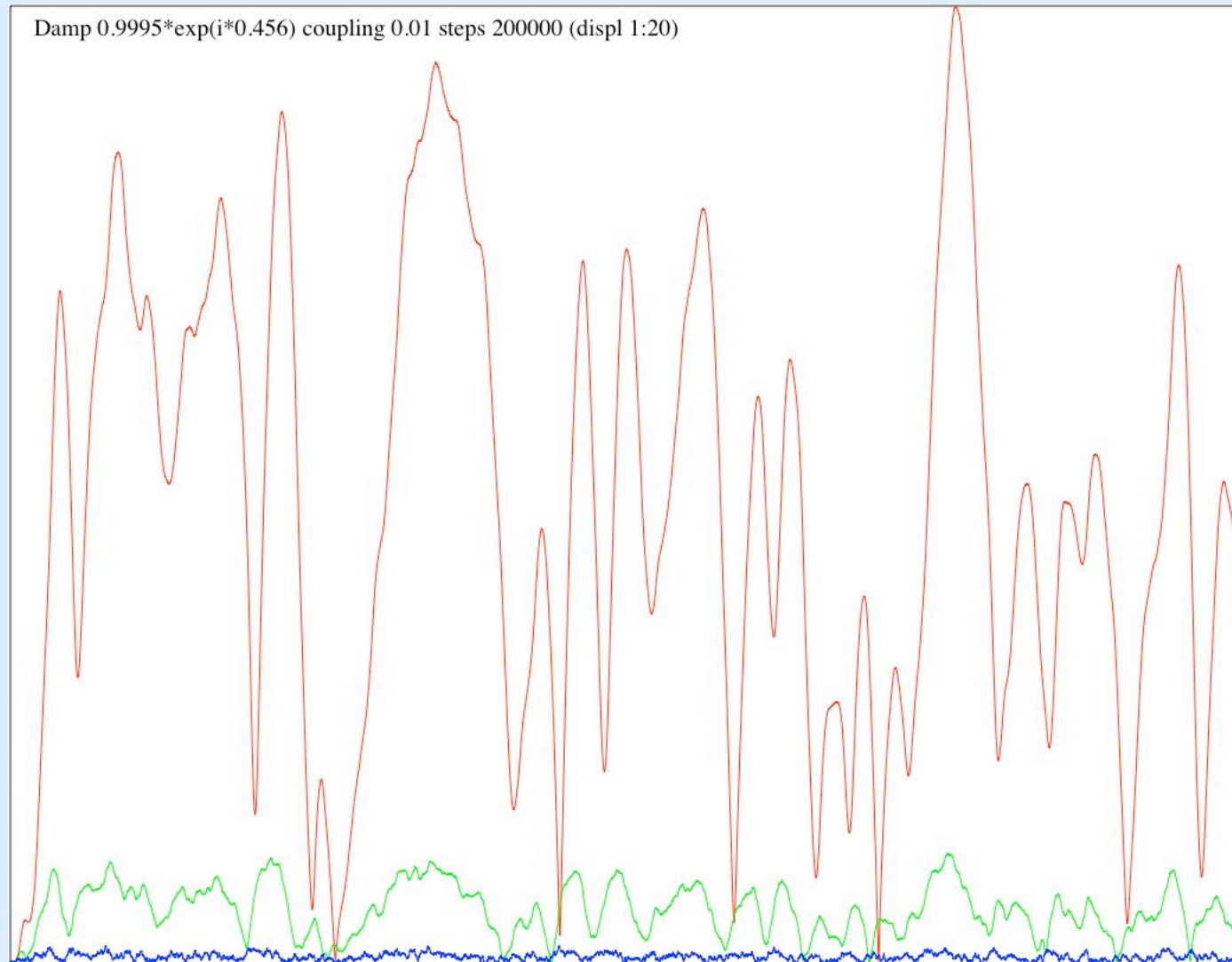


2000 steps

Random walk of
noise driven voltage
in cavity 1

U:Cavity 1 **V:Cavity 2** **W:Cavity 3**

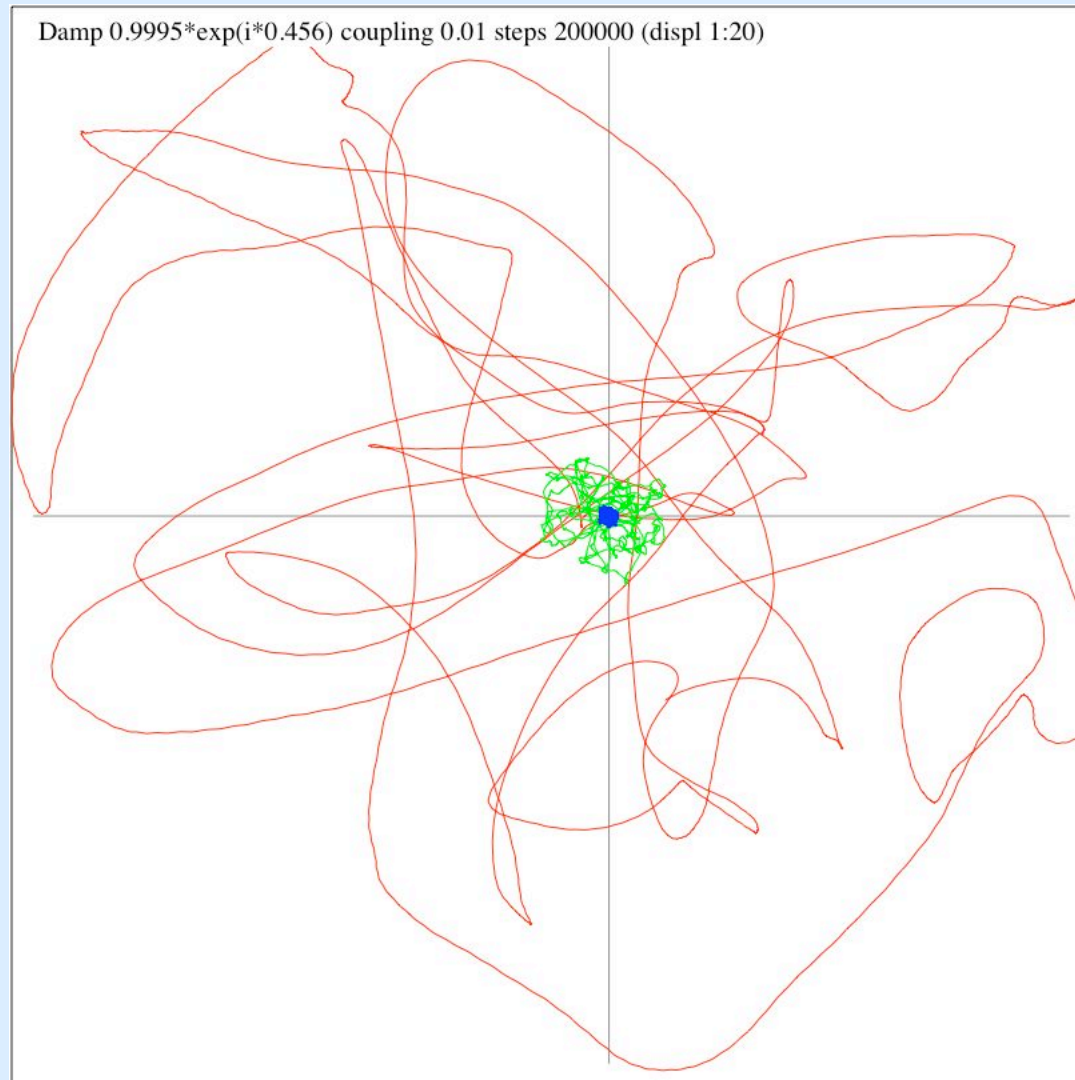
$|V|$ as function of time/bunch number, 200,000 bunches ($\approx 1/2$ ms)



U:Cavity 1 **V:Cavity 2** **W:Cavity 3**: in saturation (depends Q_{ext})

Time-wise 'constant' but increase cavity-to-cavity

Polar display of complex cavity voltage: reference HOM



200,000 steps

U:Cavity 1 **V:Cavity 2** **W:Cavity 3**

In a linac noise driven beam blow-up
might be possible, to be checked !!!

Detailed analysis with true machine conditions

First simulations with true physics (for a generic p-linac similar
SNS and SPL) for **only one** HOM (at a **random frequency**)
indicate that

Beam-loss due to such effects

cannot be excluded

(beam driven out of acceptance for some pulses)

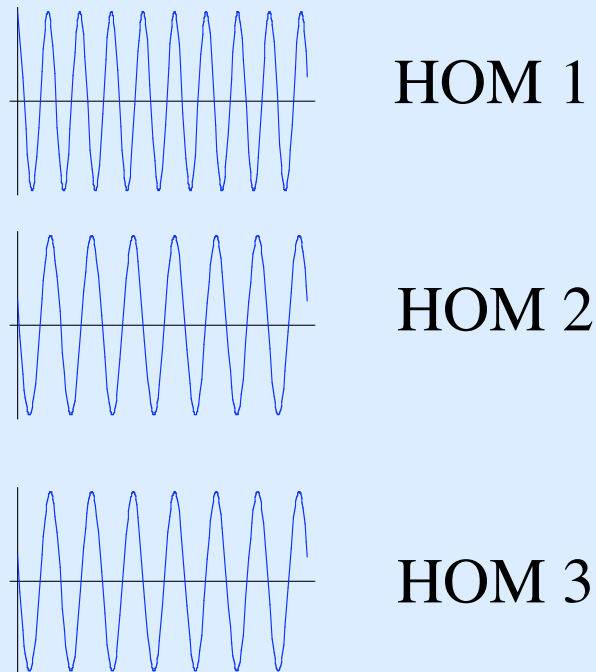
Best means to fight it:

HOM frequency scatter from cavity to cavity

Comparison circular-linear:

- In a circular machine bunch pattern passes again and again, in a linac only once:
less efficient excitation mechanism
- In a circular machine with N bunches in the turn, there are (periodicity condition) only the ‘rational’ pattern with $\xi=k/N$ oscillations per T possible (k integer $0 \leq k < N$,) Hence HOM-f can be matched only ‘about’, i.e. not necessarily exactly on the HOM peak while in a linac (no periodicity condition) any mode can be perfectly matched with a bunch displacement pattern: each mode participates ‘at best of its possibilities’

- In a circular machine generally the ‘worst’ mode makes the race, others have no time to show before beam is lost
- In Linac all modes, excited on HOM peak, may collaborate to kick out the beam:



Sum of action of **incoherent**
HOM 1, 2 and 3

