

Physics of Glue

30 years
of the Lund String

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Lund
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DEPARTMENT OF THEORETICAL PHYSICS



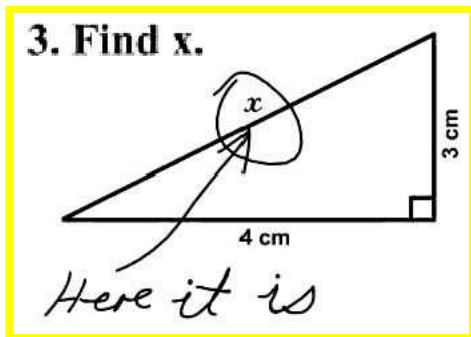
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Gösta Gustafson is a happy man whose quiver is packed with such arrows

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Understanding the interface — *metamorphosis* of coloured quarks into “white” hadrons — remains the main, most difficult, quest and headache.



What do we know (if anything) about this “*metamorphosis*” ?

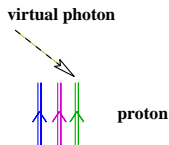
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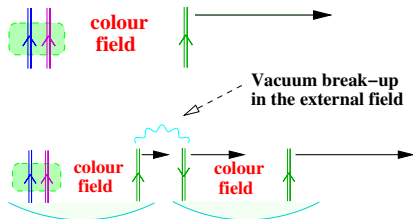
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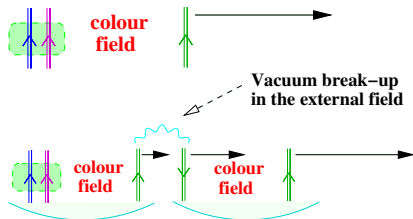
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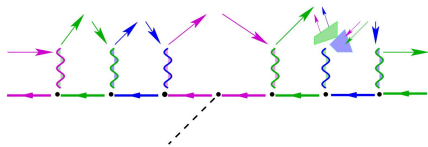
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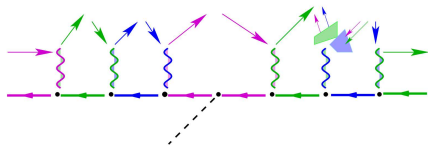
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Feynman Hadron Plateau: “one” hadron per unit $\Delta\omega/\omega$

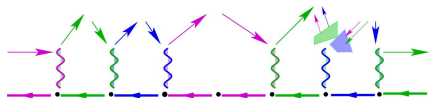


⇒ a “String” of hadrons



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The core concept of the
Lund Model

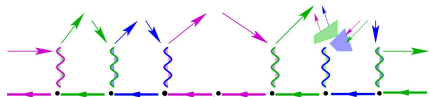


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The key features of the Lund (string) hadronization picture:

- ▶ Uniformity in *rapidity*: $dN_h = \text{const} \times d\omega_h/\omega_h$
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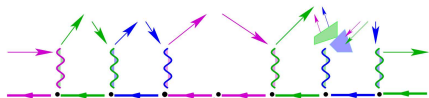
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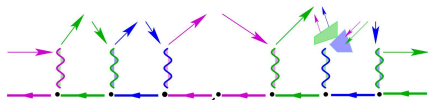
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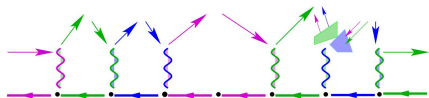
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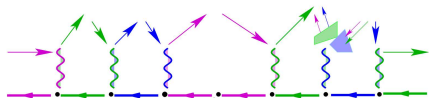
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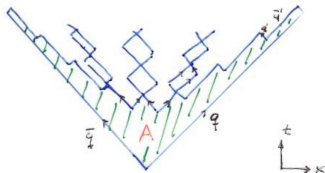
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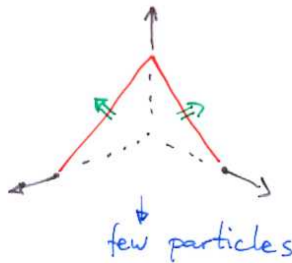
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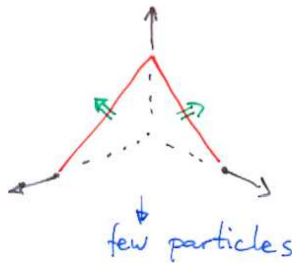
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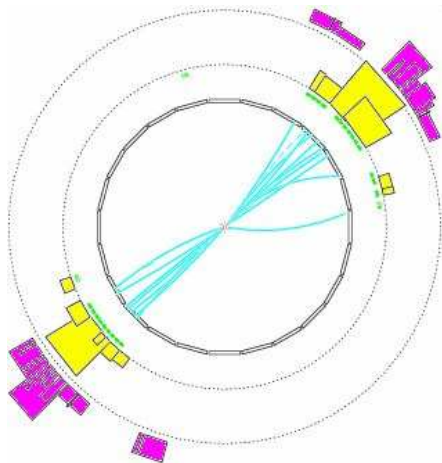
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- ▶ Breakup and Hadrons (*Yo-yo mesons*)
- ▶ Fluctuations (*Gluon as a kink*)



The crucial step:

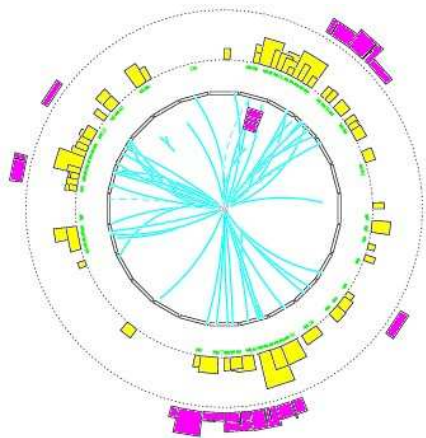
Stressing the rôle of *colour topology* in multiple hadroproduction

Hadrons between Jets

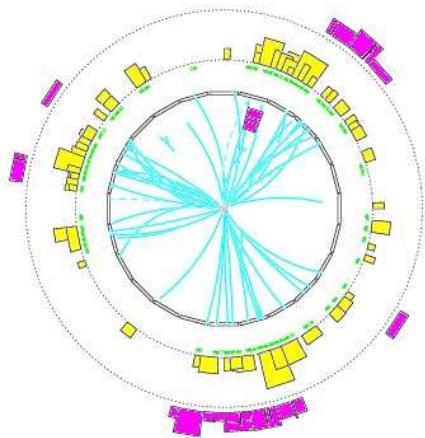


Near 'perfect' 2-jet event

2 well collimated jets of particles.



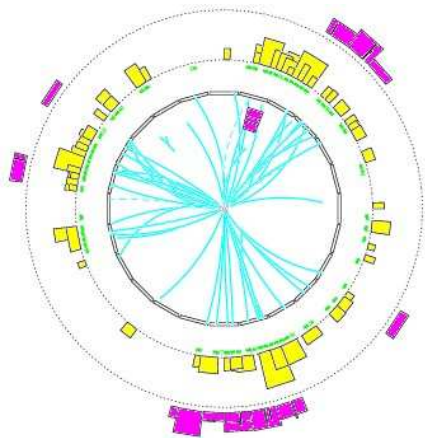
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The first QCD analysis was done by J.Ellis, M.Gaillard & G.Ross (1976)

- ▶ Planar events with large k_{\perp} ;
- ▶ How to measure gluon spin ;
- ▶ Gluon jet – softer, more populated.

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What sort of **final hadronic state** will it produce?

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Gluon \simeq quark-antiquark pair:

$$3 \otimes \bar{3} = N_c^2 = 9 \simeq 8 = N_c^2 - 1.$$

Relative mismatch : $\mathcal{O}(1/N_c^2) \ll 1$ (the large- N_c limit)

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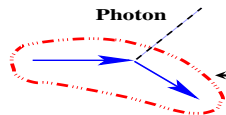
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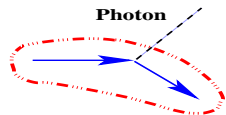
Lund model interpretation of a *gluon* —

Gluon – a “*kink*” on the “string” (colour tube)
that connects the quark with the antiquark

Look at hadrons produced in a $q\bar{q} + \text{photon}$
 e^+e^- annihilation event (recall Tornbjörn's)



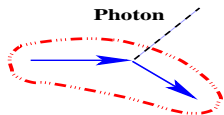
Look at hadrons produced in a $q\bar{q} + \text{photon}$
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The hot-dog of hadrons that was “*cylindric*” in
the cms, is now *lopsided* [boosted string]



Look at hadrons produced in a $q\bar{q} + \text{photon}$
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Now substitute a **gluon** for the photon in the same kinematics.

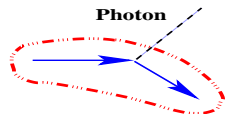




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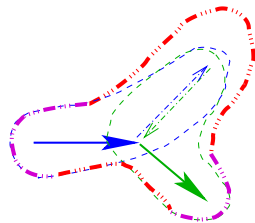
The gluon carries “double” colour charge;
quark pair is *repainted* into octet colour state.



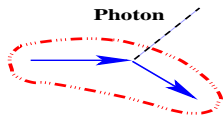
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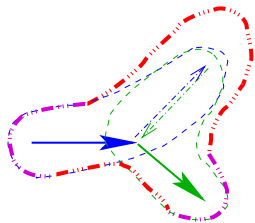
Lund: hadrons = the sum of two independent
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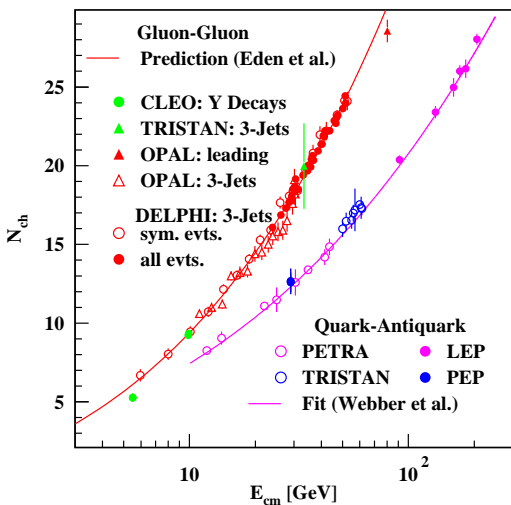
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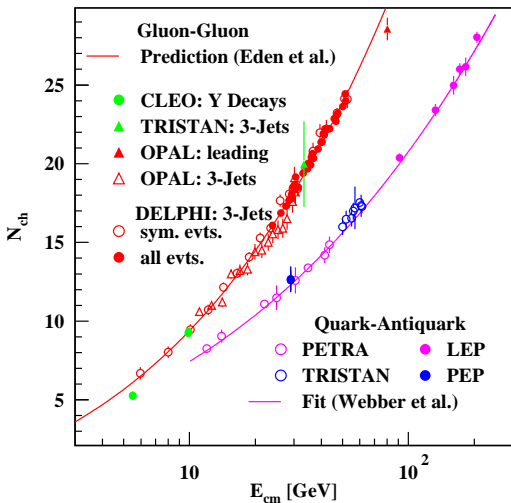
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The first immediate consequence :

Double Multiplicity of hadrons in fragmentation of the *gluon*



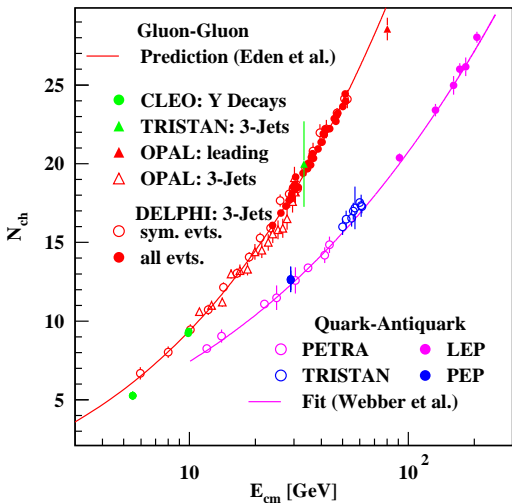
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Lessons :

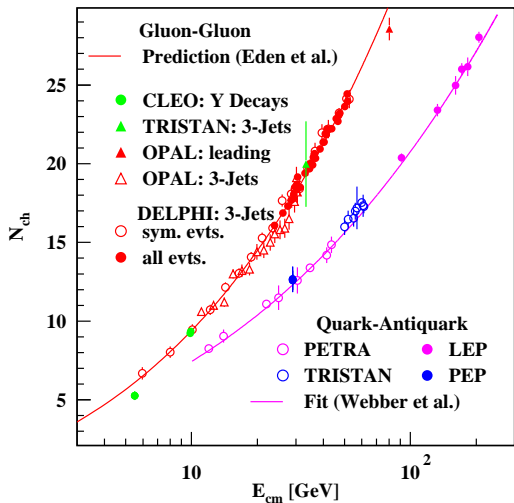
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Lessons :

- ▶ N increases *faster* than $\ln E$
 (\implies Feynman was wrong)
- ▶ $N_g/N_q < 2$

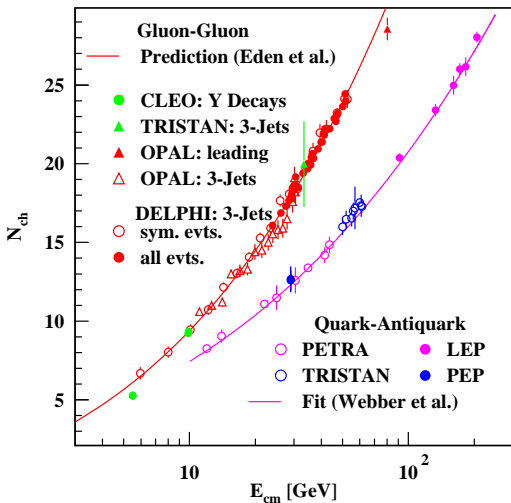


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Lessons :

- ▶ N increases *faster* than $\ln E$
 (\implies Feynman was wrong)
- ▶ $N_g/N_q < 2$ however
- ▶ $\frac{dN_g}{dN_q} = \frac{N_c}{C_F} = \frac{2N_c^2}{N_c^2-1} = \frac{9}{4} \simeq 2$
 (\implies bremsstrahlung gluons add to the hadron yield; QCD respecting parton cascades)

Comparing hadron multiplicities

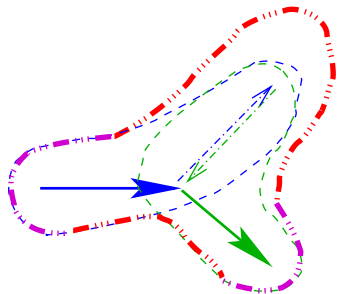


Look at experimental findings

Lessons :

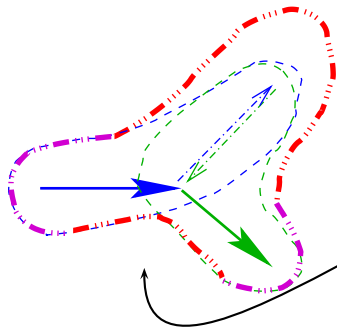
- ▶ N increases *faster* than $\ln E$
 (⇒ Feynman was wrong)
- ▶ $N_g/N_q < 2$ however
- ▶ $\frac{dN_g}{dN_q} = \frac{N_c}{C_F} = \frac{2N_c^2}{N_c^2-1} = \frac{9}{4} \simeq 2$
 (⇒ bremsstrahlung gluons add to the hadron yield; QCD respecting parton cascades)

Now let's look at a more subtle consequence of **Lund wisdom**



Lund: final hadrons are given by the sum of **two independent substrings** made of

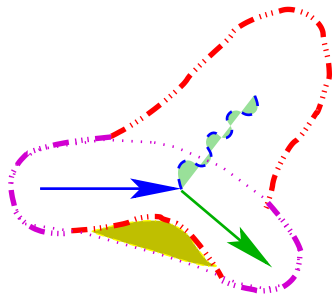
$$q + \frac{1}{2}g \quad \text{and} \quad \bar{q} + \frac{1}{2}g .$$



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Let's look into the *inter-quark valley* and compare the hadron yield with that in the $q\bar{q}\gamma$ event.

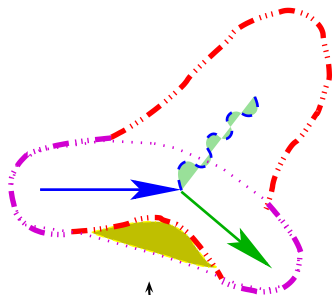


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The **overlay** results in a magnificent "*String effect*" — **depletion of particle production** in the $q\bar{q}$ valley !



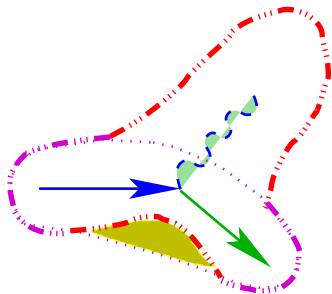
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Destructive interference
from the QCD point of view



Lund: final hadrons are given by the sum of two independent substrings made of $q + \frac{1}{2}g$ and $\bar{q} + \frac{1}{2}g$.

Let's look into the *inter-quark valley* and compare the hadron yield with that in the $q\bar{q}\gamma$ event.

The overlay results in a magnificent "*String effect*" — depletion of particle production in the $q\bar{q}$ valley !

Destructive interference from the QCD point of view

Ratios of hadron flows between jets in various multi-jet processes — example of non-trivial CIS (collinear-and-infrared-safe) QCD observable

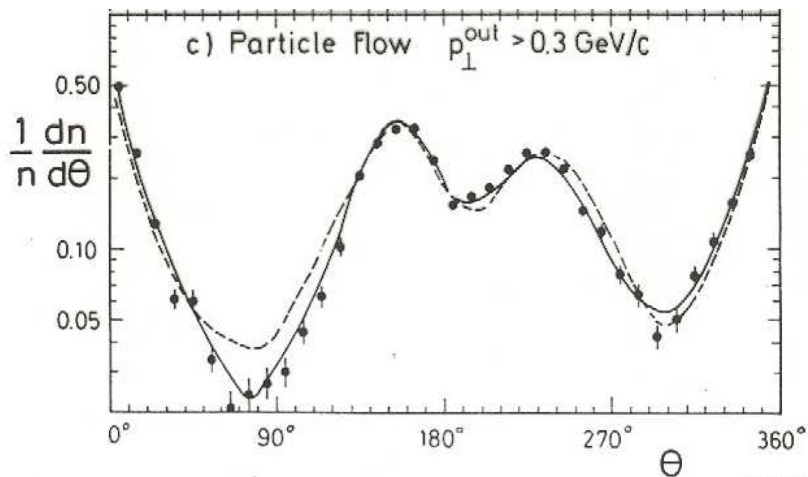
QCD prediction :

$$\frac{dN_{q\bar{q}\gamma}}{dN_{q\bar{q}g}} \simeq \frac{2(N_c^2 - 1)}{N_c^2 - 2} = \frac{16}{7}$$

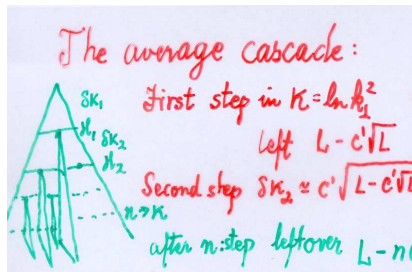
(experiment: 2.3 ± 0.2)

Prediction: 1978 (Lund)

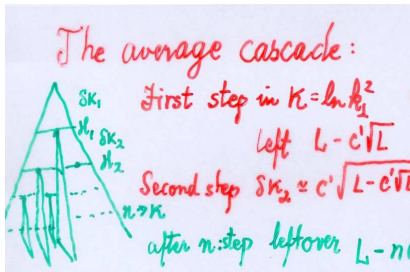
Measurement: 1981 (JADE)



Gösta's Origami

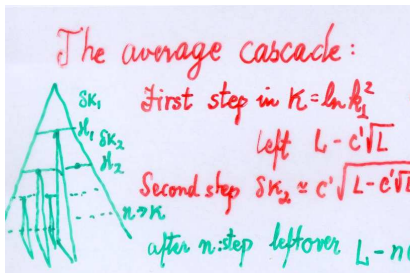


Gösta's Origami



- ▶ *Fractal structure* of parton cascades
- ▶ Multiplicity *anomalous dimension*
- ▶ Fragmentation functions

Gösta's Origami

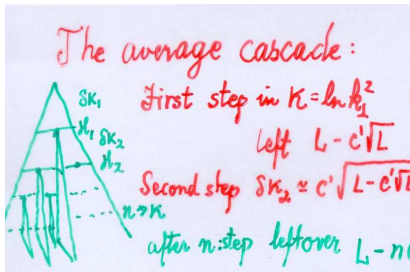


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A dual description:

radiation of a gluon \equiv dipole \rightarrow two dipoles

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The base for the *Ariadne* Monte Carlo generator

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The case of $2 \rightarrow 2$ hard parton scattering is more involved (4 emitters), especially so for *gluon-gluon* scattering.

The difficult quest of sorting out large angle gluon radiation in all orders in $(\alpha_s \log Q)^n$ was set up and solved by George Sterman and collaborators.

Here one encounters 6 (5 for $SU(3)$) colour channels that mix with each other under soft gluon radiation, and the classical picture of gluon (or dipole) multiplication is likely to fail.

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Soft anomalous dimension ,

$$\frac{\partial}{\partial \ln Q} M \propto \left\{ -N_c \ln \left(\frac{t u}{s^2} \right) \cdot \hat{\Gamma} \right\} \cdot M, \quad \hat{\Gamma} V_i = E_i V_i.$$

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$$\left[E_i - \frac{4}{3} \right]^3 - \frac{(1 + 3b^2)(1 + 3x^2)}{3} \left[E_i - \frac{4}{3} \right] - \frac{2(1 - 9b^2)(1 - 9x^2)}{27} = 0,$$

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Mark the *mysterious symmetry* w.r.t. to $x \rightarrow b$: interchanging internal (group rank) and external (scattering angle) variables of the problem ...

Some news concerning
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Have a look at the *simplest* element of the parton multiplication
Hamiltonian (non-singlet anomalous dimension) in three loops, α_s^3

$$\begin{aligned}
 P_{\text{ns}}^{(2)+}(x) = & 16 C_A C_F n_f \left(\frac{1}{6} p_{\text{qq}}(x) \left[\frac{10}{3} \zeta_2 - \frac{209}{36} - 9 \zeta_3 - \frac{167}{18} H_0 + 2 H_0 \zeta_2 - 7 H_0 \right. \right. \\
 & \left. \left. + 3 H_{1,0,0} - H_3 \right] + \frac{1}{3} p_{\text{qq}}(-x) \left[\frac{3}{2} \zeta_3 - \frac{5}{3} \zeta_2 - H_{-2,0} - 2 H_{-1} \zeta_2 - \frac{10}{3} H_{-1,0} - H_{-1,0,0} \right. \right. \\
 & \left. \left. + 2 H_{-1,2} + \frac{1}{2} H_0 \zeta_2 + \frac{5}{3} H_{0,0} + H_{0,0,0} - H_3 \right] + (1-x) \left[\frac{1}{6} \zeta_2 - \frac{257}{54} - \frac{43}{18} H_0 - \frac{1}{6} H_{0,0} \right. \right. \\
 & \left. \left. - (1+x) \left[\frac{2}{3} H_{-1,0} + \frac{1}{2} H_2 \right] + \frac{1}{3} \zeta_2 + H_0 + \frac{1}{6} H_{0,0} + \delta(1-x) \left[\frac{5}{4} - \frac{167}{54} \zeta_2 + \frac{1}{20} \zeta_2 \right. \right. \right. \\
 & \left. \left. + 16 C_A C_F^2 \left(p_{\text{qq}}(x) \left[\frac{5}{6} \zeta_3 - \frac{69}{20} \zeta_2^2 - H_{-3,0} - 3 H_{-2} \zeta_2 - 14 H_{-2,-1,0} + 3 H_{-2,0,0} \right. \right. \right. \right. \\
 & \left. \left. - 4 H_{-2,2} - \frac{151}{48} H_0 + \frac{41}{12} H_0 \zeta_2 - \frac{17}{2} H_0 \zeta_3 - \frac{13}{4} H_{0,0} - 4 H_{0,0} \zeta_2 - \frac{23}{12} H_{0,0,0} + 5 H_{0,0,0,0} \right. \right. \\
 & \left. \left. - 24 H_1 \zeta_3 - 16 H_{1,-2,0} + \frac{67}{9} H_{1,0} - 2 H_{1,0} \zeta_2 + \frac{31}{3} H_{1,0,0} + 11 H_{1,0,0,0} + 8 H_{1,1,0,0} \right. \right.
 \end{aligned}$$

$$\begin{aligned}
 & + \frac{67}{9}H_2 - 2H_2\zeta_2 + \frac{11}{3}H_{2,0} + 5H_{2,0,0} + H_{3,0} \Big] + p_{\text{qq}}(-x) \left[\frac{1}{4}\zeta_2^2 - \frac{67}{9}\zeta_2 + \frac{31}{4}\zeta_3 \right. \\
 & - 32H_{-2}\zeta_2 - 4H_{-2,-1,0} - \frac{31}{6}H_{-2,0} + 21H_{-2,0,0} + 30H_{-2,2} - \frac{31}{3}H_{-1}\zeta_2 - 42H_{-1,0} \\
 & - 4H_{-1,-2,0} + 56H_{-1,-1}\zeta_2 - 36H_{-1,-1,0,0} - 56H_{-1,-1,2} - \frac{134}{9}H_{-1,0} - 42H_{-1,1} \\
 & + 32H_{-1,3} - \frac{31}{6}H_{-1,0,0} + 17H_{-1,0,0,0} + \frac{31}{3}H_{-1,2} + 2H_{-1,2,0} + \frac{13}{12}H_0\zeta_2 + \frac{29}{2}H_0 \\
 & + 13H_{0,0}\zeta_2 + \frac{89}{12}H_{0,0,0} - 5H_{0,0,0,0} - 7H_2\zeta_2 - \frac{31}{6}H_3 - 10H_4 \Big] + (1-x) \left[\frac{133}{36} + \right. \\
 & - \frac{167}{4}\zeta_3 - 2H_0\zeta_3 - 2H_{-3,0} + H_{-2}\zeta_2 + 2H_{-2,-1,0} - 3H_{-2,0,0} + \frac{77}{4}H_{0,0,0} - \frac{20}{6} \\
 & + 4H_{1,0,0} + \frac{14}{3}H_{1,0} \Big] + (1+x) \left[\frac{43}{2}\zeta_2 - 3\zeta_2^2 + \frac{25}{2}H_{-2,0} - 31H_{-1}\zeta_2 - 14H_{-1,0} \right. \\
 & + 24H_{-1,2} + 23H_{-1,0,0} + \frac{55}{2}H_0\zeta_2 + 5H_{0,0}\zeta_2 + \frac{1457}{48}H_0 - \frac{1025}{36}H_{0,0} - \frac{155}{6}H_2
 \end{aligned}$$

$$\begin{aligned}
 & +2H_{2,0,0} - 3H_4 \Big] - 5\zeta_2 - \frac{1}{2}\zeta_2^2 + 50\zeta_3 - 2H_{-3,0} - 7H_{-2,0} - H_0\zeta_3 - \frac{37}{2}H_0\zeta_2 \\
 & - 2H_{0,0}\zeta_2 + \frac{185}{6}H_{0,0} - 22H_{0,0,0} - 4H_{0,0,0,0} + \frac{28}{3}H_2 + 6H_3 + \delta(1-x) \left[\frac{151}{64} + \right. \\
 & \left. - \frac{247}{60}\zeta_2^2 + \frac{211}{12}\zeta_3 + \frac{15}{2}\zeta_5 \right] + 16 C_A^2 C_F \left(p_{\text{qq}}(x) \left[\frac{245}{48} - \frac{67}{18}\zeta_2 + \frac{12}{5}\zeta_2^2 + \frac{1}{2}\zeta_3 \right] \right. \\
 & \left. + H_{-3,0} + 4H_{-2,-1,0} - \frac{3}{2}H_{-2,0} - H_{-2,0,0} + 2H_{-2,2} - \frac{31}{12}H_0\zeta_2 + 4H_0\zeta_3 + \frac{389}{72} \right. \\
 & \left. - H_{0,0,0,0} + 9H_1\zeta_3 + 6H_{1,-2,0} - H_{1,0}\zeta_2 - \frac{11}{4}H_{1,0,0} - 3H_{1,0,0,0} - 4H_{1,1,0,0} + 4H_{1,2,0,0} \right. \\
 & \left. + \frac{11}{12}H_3 + H_4 \right] + p_{\text{qq}}(-x) \left[\frac{67}{18}\zeta_2 - \zeta_2^2 - \frac{11}{4}\zeta_3 - H_{-3,0} + 8H_{-2}\zeta_2 + \frac{11}{6}H_{-2,0} \right. \\
 & \left. - 3H_{-1,0,0,0} + \frac{11}{3}H_{-1}\zeta_2 + 12H_{-1}\zeta_3 - 16H_{-1,-1}\zeta_2 + 8H_{-1,-1,0,0} + 16H_{-1,-1,2,0} \right. \\
 & \left. - 8H_{-2,2} + 11H_{-1,0}\zeta_2 + \frac{11}{6}H_{-1,0,0} - \frac{11}{3}H_{-1,2} - 8H_{-1,3} - \frac{3}{4}H_0 - \frac{1}{6}H_0\zeta_2 - 4 \right.
 \end{aligned}$$

$$\begin{aligned}
 & -3H_{0,0}\zeta_2 - \frac{31}{12}H_{0,0,0} + H_{0,0,0,0} + 2H_2\zeta_2 + \frac{11}{6}H_3 + 2H_4 \Big] + (1-x) \left[\frac{1883}{108} - \frac{1}{2} \right. \\
 & -H_{-2,-1,0} + \frac{1}{2}H_{-3,0} - \frac{1}{2}H_{-2}\zeta_2 + \frac{1}{2}H_{-2,0,0} + \frac{523}{36}H_0 + H_0\zeta_3 - \frac{13}{3}H_{0,0} - \frac{5}{2}H_{0,0,0} \\
 & \left. -2H_{1,0,0} \right] + (1+x) \left[8H_{-1}\zeta_2 + 4H_{-1,-1,0} + \frac{8}{3}H_{-1,0} - 5H_{-1,0,0} - 6H_{-1,2} - \frac{13}{3} \right. \\
 & -\frac{43}{4}\zeta_3 - \frac{5}{2}H_{-2,0} - \frac{11}{2}H_0\zeta_2 - \frac{1}{2}H_2\zeta_2 - \frac{5}{4}H_{0,0}\zeta_2 + 7H_2 - \frac{1}{4}H_{2,0,0} + 3H_3 + \frac{3}{4} \\
 & \left. + \frac{1}{4}\zeta_2^2 - \frac{8}{3}\zeta_2 + \frac{17}{2}\zeta_3 + H_{-2,0} - \frac{19}{2}H_0 + \frac{5}{2}H_0\zeta_2 - H_0\zeta_3 + \frac{13}{3}H_{0,0} + \frac{5}{2}H_{0,0,0} \right. \\
 & \left. -\delta(1-x) \left[\frac{1657}{576} - \frac{281}{27}\zeta_2 + \frac{1}{8}\zeta_2^2 + \frac{97}{9}\zeta_3 - \frac{5}{2}\zeta_5 \right] \right) + 16 C_F n_f^2 \left(\frac{1}{18} p_{\text{qq}}(x) \left[H_{0,0} \right. \right. \\
 & \left. \left. + (1-x) \left[\frac{13}{54} + \frac{1}{9}H_0 \right] - \delta(1-x) \left[\frac{17}{144} - \frac{5}{27}\zeta_2 + \frac{1}{9}\zeta_3 \right] \right) + 16 C_F^2 n_f \left(\frac{1}{3} p_{\text{qq}}(x) \left[\right. \right.
 \end{aligned}$$

$$\begin{aligned}
 & -\frac{55}{16} + \frac{5}{8}H_0 + H_0\zeta_2 + \frac{3}{2}H_{0,0} - H_{0,0,0} - \frac{10}{3}H_{1,0} - \frac{10}{3}H_2 - 2H_{2,0} - 2H_3 \Big] + \frac{2}{3} \\
 & -\frac{3}{2}\zeta_3 + H_{-2,0} + 2H_{-1}\zeta_2 + \frac{10}{3}H_{-1,0} + H_{-1,0,0} - 2H_{-1,2} - \frac{1}{2}H_0\zeta_2 - \frac{5}{3}H_{0,0} - \\
 & -(1-x) \left[\frac{10}{9} + \frac{19}{18}H_{0,0} - \frac{4}{3}H_1 + \frac{2}{3}H_{1,0} + \frac{4}{3}H_2 \right] + (1+x) \left[\frac{4}{3}H_{-1,0} - \frac{25}{24}H_0 + \right. \\
 & \left. + \frac{7}{9}H_{0,0} + \frac{4}{3}H_2 - \delta(1-x) \left[\frac{23}{16} - \frac{5}{12}\zeta_2 - \frac{29}{30}\zeta_2^2 + \frac{17}{6}\zeta_3 \right] \right) + 16 C_F^3 \left(p_{\text{qq}}(x) \left[\right. \right. \\
 & \left. \left. + 6H_{-2}\zeta_2 + 12H_{-2,-1,0} - 6H_{-2,0,0} - \frac{3}{16}H_0 - \frac{3}{2}H_0\zeta_2 + H_0\zeta_3 + \frac{13}{8}H_{0,0} - 2H_0 \right. \right. \\
 & \left. \left. + 12H_1\zeta_3 + 8H_{1,-2,0} - 6H_{1,0,0} - 4H_{1,0,0,0} + 4H_{1,2,0} - 3H_{2,0} + 2H_{2,0,0} + 4H_{2,1,0} \right. \right. \\
 & \left. \left. + 4H_{3,0} + 4H_{3,1} + 2H_4 \right] + p_{\text{qq}}(-x) \left[\frac{7}{2}\zeta_2^2 - \frac{9}{2}\zeta_3 - 6H_{-3,0} + 32H_{-2}\zeta_2 + 8H_{-2,0} \right. \right. \\
 & \left. \left. - 26H_{-2,0,0} - 28H_{-2,2} + 6H_{-1}\zeta_2 + 36H_{-1}\zeta_3 + 8H_{-1,-2,0} - 48H_{-1,-1}\zeta_2 + 40H_{-1,-1,0} \right] \right)
 \end{aligned}$$

$$\begin{aligned}
 &+48H_{-1,-1,2} + 40H_{-1,0}\zeta_2 + 3H_{-1,0,0} - 22H_{-1,0,0,0} - 6H_{-1,2} - 4H_{-1,2,0} - 32 \\
 &- \frac{3}{2}H_0\zeta_2 - 13H_0\zeta_3 - 14H_{0,0}\zeta_2 - \frac{9}{2}H_{0,0,0} + 6H_{0,0,0,0} + 6H_2\zeta_2 + 3H_3 + 2H_{3,0} - \\
 &+(1-x) \left[2H_{-3,0} - \frac{31}{8} + 4H_{-2,0,0} + H_{0,0}\zeta_2 - 3H_{0,0,0,0} + 35H_1 + 6H_1\zeta_2 - H_1, \right. \\
 &+(1+x) \left[\frac{37}{10}\zeta_2^2 - \frac{93}{4}\zeta_2 - \frac{81}{2}\zeta_3 - 15H_{-2,0} + 30H_{-1}\zeta_2 + 12H_{-1,-1,0} - 2H_{-1,0} \right. \\
 &- 24H_{-1,2} - \frac{539}{16}H_0 - 28H_0\zeta_2 + \frac{191}{8}H_{0,0} + 20H_{0,0,0} + \frac{85}{4}H_2 - 3H_{2,0,0} - 2H_3 \\
 &- H_4 \left. \right] + 4\zeta_2 + 33\zeta_3 + 4H_{-3,0} + 10H_{-2,0} + \frac{67}{2}H_0 + 6H_0\zeta_3 + 19H_0\zeta_2 - 25H_{0,0} \\
 &- 2H_2 - H_{2,0} - 4H_3 + \delta(1-x) \left[\frac{29}{32} - 2\zeta_2\zeta_3 + \frac{9}{8}\zeta_2 + \frac{18}{5}\zeta_2^2 + \frac{17}{4}\zeta_3 - 15\zeta_5 \right] \left. \right)
 \end{aligned}$$

2×2 anomalous dimension matrix occupies

1 st loop: 1/10 page

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Moch, Vermaseren and Vogt

[waterfall of results launched
March 2004, and counting]

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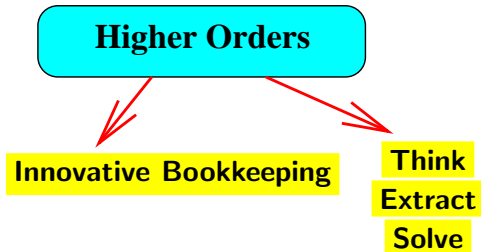
not too encouraging a trend ...



How to reduce complexity ?

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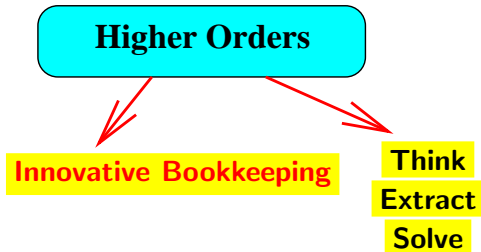
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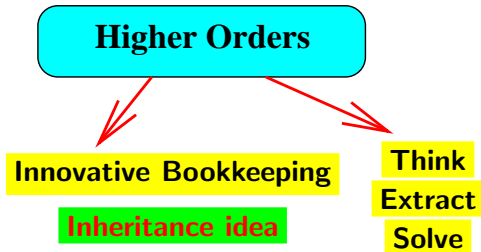
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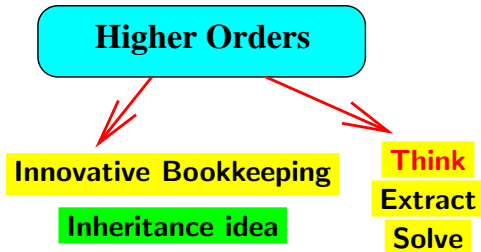
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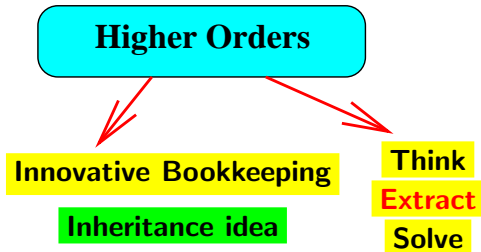
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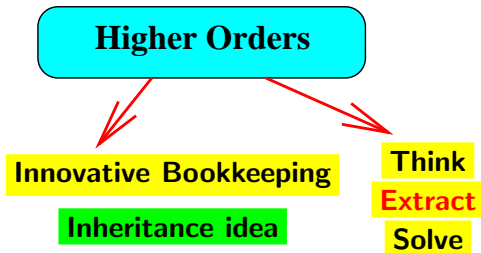
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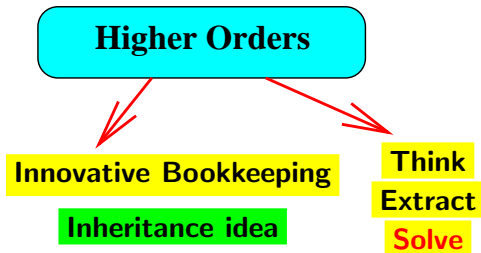
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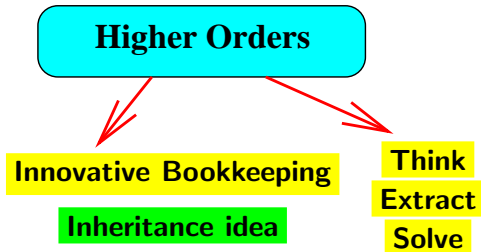
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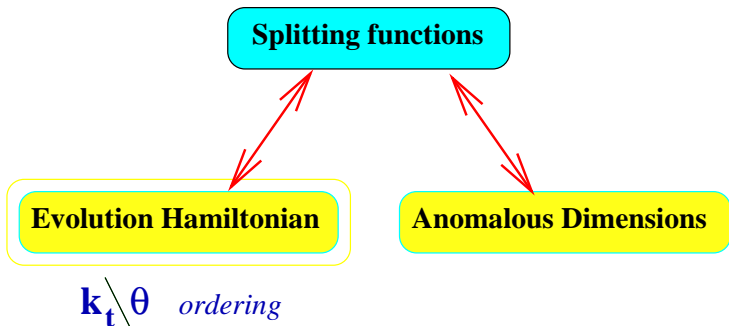
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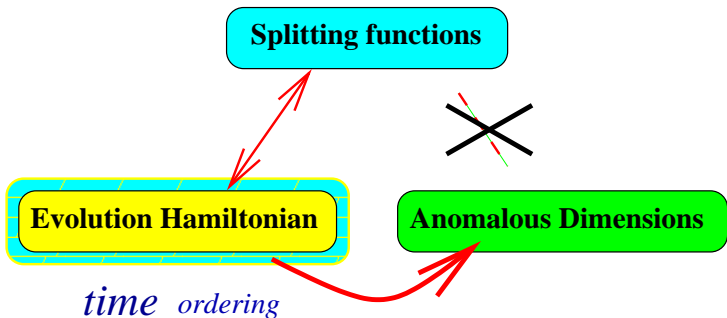
➡ A playing ground for theoretical theory: SUSY, AdS/CFT, ...

In the standard approach,



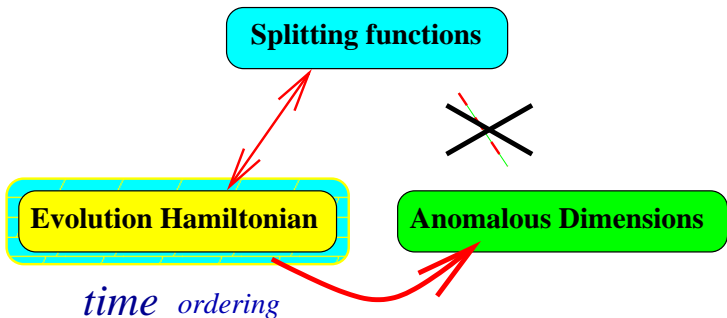
- ▶ parton splitting functions are equated with anomalous dimensions;
- ▶ they are different for DIS and e^+e^- evolution;
- ▶ “clever evolution variables” are different too

In the new approach,



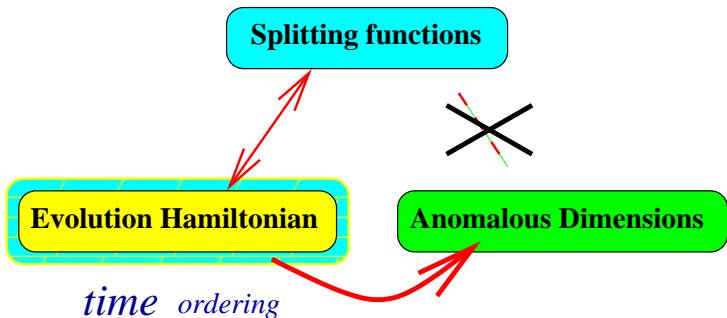
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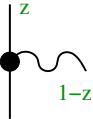
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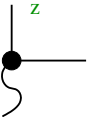
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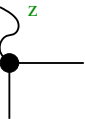
Recall an old hint from QCD ...



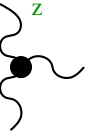
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$$= N_c \cdot \frac{1+z^4+(1-z)^4}{z(1-z)}$$

Four “parton splitting functions”

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► Exchange the **decay products** : $z \rightarrow 1-z$

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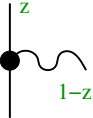
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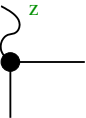
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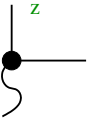
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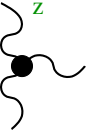
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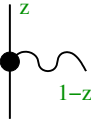
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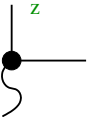
Three (QED) “kernels” are inter-related; gluon self-interaction stays put :

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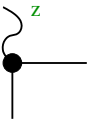
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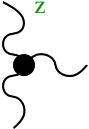
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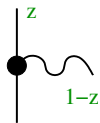


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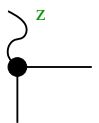
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All four are related !

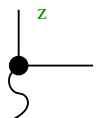
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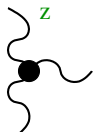
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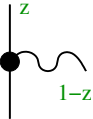


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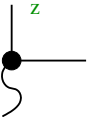
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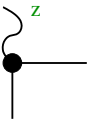
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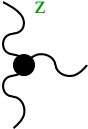
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WHY and WHAT FOR ?

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And here we arrive at the second — **Divide and Conquer** — issue

Recall the diagonal first loop anomalous dimensions:

$$\tilde{\gamma}_{q \rightarrow q(x)+g} = \frac{C_F \alpha_s}{\pi} \left[\frac{x}{1-x} + (1-x) \cdot \frac{1}{2} \right],$$

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Let us look at the rôles these animals play on the QCD stage

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- ✗ Classical Field
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- ✓ define the physical coupling
- ✓ responsible for
 - ➔ DL radiative effects,
 - ➔ reggeization,
 - ➔ QCD/Lund string (gluons)
- ✓ play the major rôle in evolution

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In addition,

- ✗ Tree multi-clagon (Parke–Taylor) amplitudes are *known exactly*
- ✗ It is clagons which dominate in all the *integrability cases*

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► $\beta(\alpha) \equiv 0$ in all orders ! \implies $\gamma \Rightarrow \frac{x}{1-x}$ + **no quagons !**

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$$\frac{\text{clever 2nd loop}}{\text{clever 1st loop}} < 2\% \quad \left(\begin{array}{l} \text{Heavy quark fragmentation} \\ \text{D-r, Khoze \& Troyan , PRD 1996} \end{array} \right)$$

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Employ $\mathcal{N}=4$ SYM to simplify the major part of the QCD dynamics !

- ▶ A steady progress in high order perturbative QCD **calculations** is worth accompanying by **reflections** upon the origin and the structure of higher loop correction effects
- ▶ Reformulation of parton cascades in terms of Gribov–Lipatov reciprocity respecting evolution equations (RREE)
 - ▶ reduces complexity by (at least) an order of magnitude
 - ▶ improves perturbative series (less singular, better “convergent”)
 - ▶ links interesting phenomena in the DIS and e^+e^- annihilation channels
- ▶ The Low theorem should be part of theor.phys. curriculum, worldwide
- ▶ Complete solution of the $\mathcal{N}=4$ SYM QFT should provide us with a *one-line-all-orders* description of the major part of QCD dynamics
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