## Lifetime measurements in ${ }^{53} \mathbf{C a}$

S. Chen<br>on behalf of the HiCARI collaboration and the RIBF-170 collaboration

University of York

Liverpool, $10^{\text {th }}$ April 2024

## Shell structures and nuclear forces



Maria Goeppert-Mayer, Phys. Rev. 75, 1969 (1949) O. Haxel, Phys. Rev. 75, 1766 (1949)

## Shell structures and nuclear forces



- Shell structure and magic numbers were the cornerstones of the shell model for many decades
- Experiments on exotic nuclei found magic numbers are not immutable throughout the nuclear chart
- Shell structure is now recognized as local concept
J. Dobaczewski et al., Prog. Part. Nucl. Phys. 59 (2007) 432-445


## Shell structures and nuclear forces


T. Otsuka et al., Phys. Rev. Lett. 87, 082502 (2001)

T. Otsuka et al., Phys. Rev. Lett. 95, 232502 (2005)

Tensor force: attractive when the intrinsic spins of neutron and proton are anti-parallel and repulsive when they are parallel

${ }^{30} \mathrm{Si} \rightarrow{ }^{24} \mathrm{O}$ : absence of strong $\pi 0 d_{5 / 2}-v 0 d_{3 / 2}$ attraction $\Rightarrow N=16$ new magic number in oxygen

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## $N=32,34$ shell closure in Ca isotopes

D. Steppenbeck et al., Nature 502,207-10 (2013)

${ }^{60} \mathrm{Fe}(Z=26)$

b

${ }^{58} \mathrm{Cr}(Z=24)$


- Absence of strong $\pi f_{7 / 2}-v f_{5 / 2}$ attraction make $N=32,34$ new magic number in calcium

(sub-)shell closure at $N=16,20,28,32,34$



## $N=32$ shell closure in ${ }^{52} \mathrm{Ca}$

D. Steppenbeck et al., Nature 502,207-10 (2013)

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${ }^{58} \mathrm{Cr}(Z=24)$


- Absence of strong $\pi f_{7 / 2}-v f_{5 / 2}$ attraction make $N=32,34$ new magic number in calcium
- Experimental evidences for $\boldsymbol{N}=\mathbf{3 2}$ magic
- Large $\mathrm{E}\left(2^{+} \frac{1}{1}\right)$
- Mass measurement: large empirical two-neutron shell gap
A. Huck et al., Phys. Rev. C 31,2226-2237 (1985)

F. Wienholtz et al., Nature 498, 346-349 (2013)



## $N=34$ shell closure in ${ }^{54} \mathbf{C a}$

D. Steppenbeck et al., Nature 502,207-10 (2013)

${ }^{60} \mathrm{Fe}(Z=26)$
b

${ }^{58} \mathrm{Cr}(Z=24)$


- Absence of strong $\pi f_{7 / 2}-v f_{5 / 2}$ attraction make $N=32,34$ new magic number in calcium
- Experimental evidences for $\boldsymbol{N}=34$ magic
- $E\left(2^{+}\right)$; first evidence of magicity
- Mass measurement: $N=34$ shell gap similar size with $N=32$ shell gap
D. Steppenbeck et al., Nature 502,207-10 (2013)

S. Michimasa et al., Phys. Rev. Lett. 121, 022506 (2018)




## Excitation of ${ }^{53} \mathbf{C a}$

D. Steppenbeck et al., Nature 502,207-10 (2013)

g.s. (1/2-)

| $\pi p_{1 / 2}$ | 34 |
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${ }^{53} \mathrm{Ca}(\mathrm{Z}=20)$

2220 keV (3/2-)

${ }^{53} \mathrm{Ca}(\mathrm{Z}=20)$
$E_{x}=2220 \mathrm{keV}$ via $\beta$-decay,
F. Perrot et al. PRC 74,014313 (2006)
S. Chen et al., PRL 123,142501 (2019)

| State | Energy (keV) | GXPF1Bs | NNLOsat | NN+3N (Inl) |
| :---: | :---: | :---: | :---: | :---: |
| $3 / 2-$ | $2220(13)$ | 2061 | 2635 | 2611 |
| $5 / 2-$ | $1738(17)$ | 1934 | 1950 | 2590 |

## Lifetime measurements:

- E2 transition probability
- Benchmarks to test different theoretical descriptions beyond excitation energies

1753 keV (5/2-)

## Lifetime measurement method

- Detect prompt gamma rays from fast moving ( $\sim 0.5 c$ ) particles
- Doppler correction for prompt gamma

$$
E_{\gamma \theta}=E_{\gamma} \frac{1-\beta \cos \theta}{\sqrt{1-\beta^{2}}}
$$

- Finite excitation state lifetime lead to $\theta$ and $\beta$ distribution different from zero lifetime

$$
=>\text { asymmetric peak shape after }
$$

Doppler-correction

- Peak shape analysis to extract excitation state lifetime


Figure: Schematic diagram of lifetime measurement method (Thesis, S. Heil, 2019)

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Doppler-correction

- Peak shape analysis to extract excitation state lifetime
$==>$ complicated lifetime responses, need compare with detailed simulations


Figure: Doppler broadening as a function of gamma-emission angle (P. Doornenbal, PTEP 2012, 03C004)

## Experiment Setup at RIBF


arget


## Particle Identification



## Experiment Setup - HiCARI

- High-resolution Cluster Array at the RIBF (RIKEN Accel. Prog., (2021), K. Wimmer, et. al.)
- Hybrid HPGe array:
- $6 \times$ Miniball cluster ( 6 segments)
- 4 x SuperClover cluster (4 segments)
- 1 x Gretina Quad cluster (position sensitive)
- $1 \times$ Gretina P3 cluster (position sensitive)


Figure: HiCARI array

## Benchmark - ${ }^{36} \mathrm{Ar}$ from ${ }^{37} \mathrm{~K}-1 p$

${ }^{36} \mathbf{A r}$ low-lying excitation states from NNDC



- Spectra of CH2 target have worse energy resolutions, due to large $\beta$ uncertainty
- ${ }^{36} \mathrm{Ar}$ low-lying excitation states: known energy and lifetime
$==>$ Benchmark Geant4 simulations


Figure: Doppler-corrected gamma-ray energy spectra

## Benchmark - ${ }^{36} \mathrm{Ar}$ from ${ }^{37} \mathrm{~K}-1 p$

${ }^{36}$ Ar low-lying excitation states from NNDC


- Fitting function:

Geant4 simulations + double-expo background




## ${ }^{53}$ Ca from ${ }^{55} \mathrm{Sc}-1 p 1 n$



Previous measurement:
Gamma-ray spectra measured with $\mathrm{NaI}(\mathrm{TI})$ detectors
D. Steppenbeck et al., Nature 502,207-10 (2013)



Figure: Doppler-corrected gamma-ray energy spectra

## ${ }^{53}$ Ca from ${ }^{55}$ Sc - $1 p 1$ n





## ${ }^{53} \mathrm{Ca}$ from ${ }^{55} \mathrm{Sc}-1 p 1 n$



Preliminary results:

|  | Energy $/ \mathrm{keV}$ | Halflife /ps |
| :--- | :---: | :---: |
| Miniball | $1752(8)$ | $15(8)$ |
| Clover | $1730(6)$ | $8(8)$ |
| tracking | $1744(8)$ | $11(8)$ |
| weighted | $1740(4)$ | $11(5)$ |

- only statistic uncertainties are considered
- weighted by $1 / \sigma^{2}$
- deduced $\mathrm{B}\left(\mathrm{E} 2,5 / 2^{-} \rightarrow 1 / 2^{-}\right)=3.2_{-1.0}{ }^{+2.8} \mathrm{e}^{2} \mathrm{fm}^{4}$

Theoretical calculations:

|  | Energy / keV | B(E2) / $\mathrm{e}^{2} \mathrm{fm}^{4}$ | Halflife_E $\mathrm{E}_{\text {exp }} / \mathrm{ps}$ |
| :--- | :---: | :---: | :---: |
| UFP-CA | 1767 | 5.11 | 7.0 |
| VS-IMSRG | 2116 | 0.785 | 45 |

- Shell Model calculations using UFP-CA interaction (by Alex Brown)
- Valence-space in-medium similarity renormalization group (VS-IMSRG) using chiral effective field theory (EFT) interaction (by Jason Holt, 2019)


## Summary

- Performed in-beam gamma-ray spectroscopy measurement in neutron-rich Ca isotopes with a Hybrid HPGe array at RIBF
- Benchmarked the simulation with ${ }^{36} \mathrm{Ar}$ spectra, the obtained gamma-ray response functions well reproduced the peak shapes
- Analysed the ${ }^{53} \mathrm{Ca}$ spectra with detailed simulations, the lifetime of the $5 / 2$ state is extracted to be $11(5) \mathrm{ps}$, leading to a $\mathrm{B}\left(\mathrm{E} 2,5 / 2^{-} \rightarrow 1 / 2^{-}\right)=3.2_{-1.0}{ }^{+2.8} \mathrm{e}^{2} \mathrm{fm}^{4}$
- Experimental results are compared with Shell-Model calculations using UFP-CA interaction, and VS-IMSRG approach using EFT interactions with 2 N and 3 N forces


## Collaborations

university University of York：S．Chen，R．Crane，W．Marshall，R．Taniuchi，M．Petri， S．Paschalis，M．Bentley，L．Tetley

RIKEN：P．Doornenbal，H．Baba，F．Browne，B．Mauss，B．Moon，H．Sakurai， D．Suzuki
RCNP
RCNP：N．Aoi，E．Ideguchi，S．Iwazaki，A．Kohda，Y．Yamamoto
universität zu KÖLN
$\hat{*}$｜A berkeley lab
CSIC
KOREA
UNIVERSITY
東京大学
universityof
SURREY
University of Surrey：T．Parry
KULEUVEN KU Leuven：H．de Witte
®TRIUMF TRIUMF：J．Holt
Thank you for your attention
$\int \frac{\text { MICHIGAN STATE }}{u N \text { IVER S I T Y }}$

MSU：B．A．Brown

Backup

## BigRIPS Separator

Big RIKEN Projectile Fragment Separator

- Bp- $\Delta E-B \rho$ separation
- Large acceptance


$$
\begin{array}{ll}
0 & \Delta \theta= \pm 40 \mathrm{mrad} \\
\text { - } & \Delta \varphi= \pm 50 \mathrm{mrad} \\
\text { - } & \Delta \mathrm{p} / \mathrm{p}= \pm 3 \%
\end{array}
$$

- Event-by-event $B \rho-T O F-\Delta E$ particle identification


## Shell structures and nuclear forces



- spin-orbital splitting:
$E\left(j_{>}=l+s\right)<E\left(j_{c}=I-s\right)$
- conventional magic numbers:
$2,8,20,28,50, \ldots$

(b)


Tensor force

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## Shell structures and nuclear forces

| Woods-Saxon potential | with Spin-Orbital (L-S) force |
| :---: | :---: |
| $4 \hbar \omega-0 g$ |  |
|  | (10) $[50]$ (2) $[40]$ (6) $[38]$ (4) $[32]$ (8) $[28]$ |
|  |  |
| $1 \hbar \omega-0 p-=:-0 p_{1 / 2}$ | $\varlimsup^{\text {(2) }} \text { (4) }[8]$ |
| $0 \hbar \omega$ - $0 s$ - $-\cdots-0 s_{1 / 2}$ | - (2) [2] |

- spin-orbital splitting:
$E(j=1+s)<E\left(j_{<}=l-s\right)$
- conventional magic numbers:
$2,8,20,28,50, \ldots$

(sub-)shell closure at $N=16,20,28,32,34$



## Particle Identification



## Doppler-correction

P. Doornenbal, PTEP 2012, 03C004

- Doppler correction for prompt gamma

$$
\frac{E_{\gamma}}{E_{\gamma 0}}=\frac{\sqrt{1-\beta^{2}}}{1-\beta \cos \vartheta_{\gamma}}
$$

- Doppler-corrected gamma energy resolution:

$$
\begin{aligned}
\left(\frac{\Delta E_{\gamma 0}}{E_{\gamma 0}}\right)^{2}= & \left(\frac{\beta \sin \vartheta_{\gamma}}{1-\beta \cos \vartheta_{\gamma}}\right)^{2} \times\left(\Delta \vartheta_{\gamma}\right)^{2}+\left(\frac{\beta-\cos \vartheta_{\gamma}}{\left(1-\beta^{2}\right)\left(1-\beta \cos \vartheta_{\gamma}\right)}\right)^{2} \times(\Delta \beta)^{2} \\
& +\left(\frac{\Delta E_{\text {int }}}{E_{\gamma}}\right)^{2} .
\end{aligned}
$$



Fig. 1. Doppler broadening due to $\Delta \beta, \Delta \vartheta_{\gamma}$, and $\Delta E_{\text {intr }}$ as a function of the $\gamma$-ray emission (detector) angle $\vartheta_{\gamma}$. Three different velocities were assumed. The upper panel displays only the velocity uncertainty effect for $\Delta \beta / \beta=0.1$, while the middle panel displays the broadening due to a detector opening angle of $\Delta \vartheta_{\gamma}=122 \mathrm{mrad}$. In the bottom panel, the sum effect including an intrinsic energy resolution of $6 \%$ at 1.33 MeV is displayed. The calculations were performed for a $1 \mathrm{MeV} \gamma$-ray energy assuming a square-root dependence of the intrinsic energy resolution.

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${ }^{36}$ Ar low-lying excitation states from NNDC


- Fix $4^{+} \rightarrow 2^{+}$and $3 \rightarrow 2^{+}$energies and lifetimes
- Study $2^{+} \rightarrow 0^{+}$energy and tracking detectors position resolution
- Red markers: $1-\sigma$ region $\left(\right.$ min-chi $\left.{ }^{2}+1\right)$
- MinChi2 at

$$
\begin{aligned}
& 1967 \mathrm{keV}, 3.3 \mathrm{~mm} \text { (P3) } \\
& 1969 \mathrm{keV}, 3.3 \mathrm{~mm} \text { (Quad) }
\end{aligned}
$$

chi2 P3



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$\chi^{2}$ surface \{Clover\}

halflife (ps)

