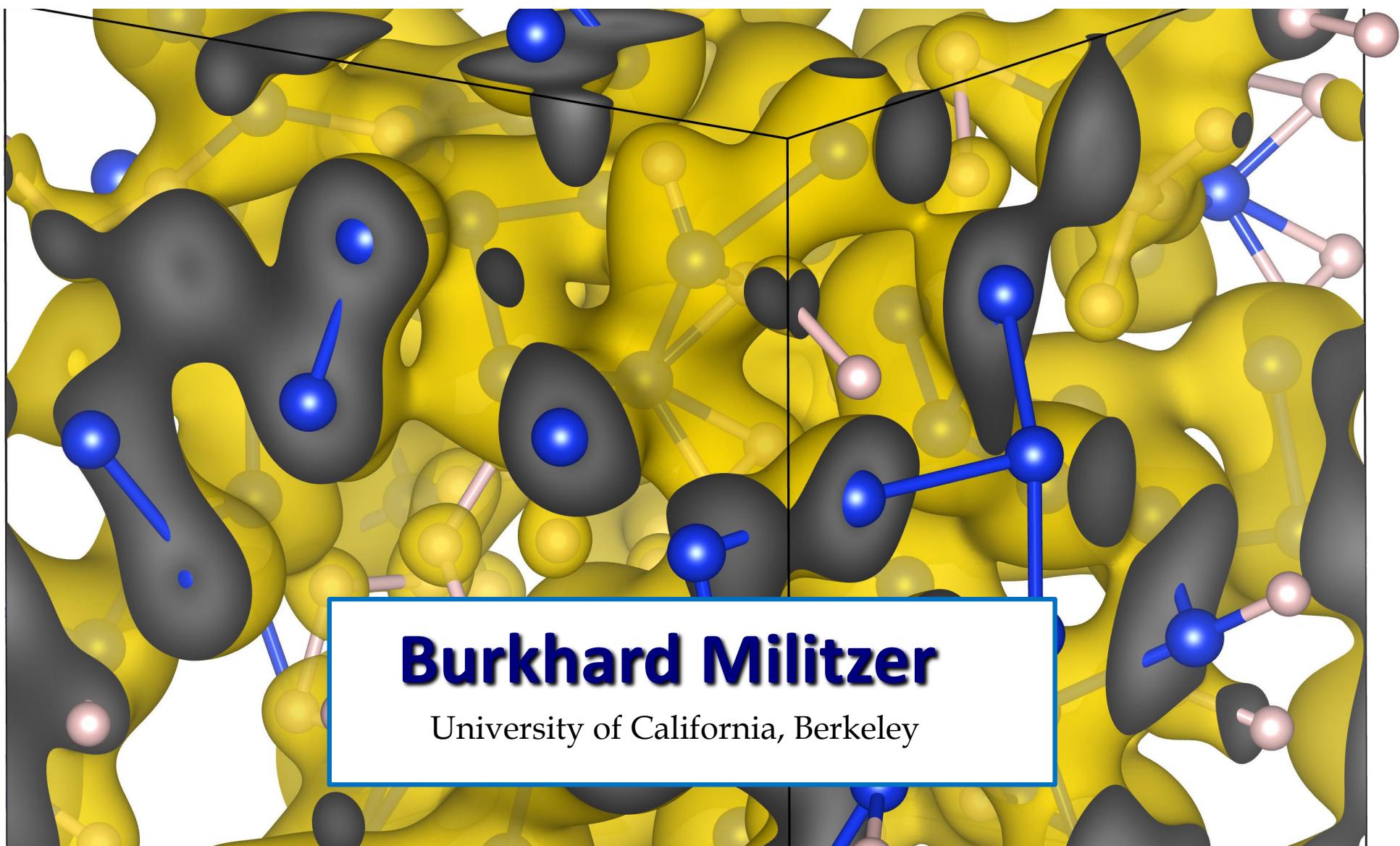


# **Path Integral Monte Carlo Simulations of Warm Dense Matter**



# **Plan for L5, L9, L12 and T1 & T3**

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**L5: Path integral Monte Carlo (PIMC)  
simulations and First Principles Equation of  
state (FPEOS) database**

**L9: NASA mission Juno to Jupiter, dilute core**

**T1: “Build that Planet” with SPH method**

**L12: NASA mission Cassini to Saturn. How did  
that planet become the Lord of the Rings?**

**T3: FPEOS tutorial**

# **Software needed for T1 & T3**

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**T1: “Built that Planet” with SPH method**

**Python + Jupyter notebooks (installation on laptop required, for example with Anaconda)**

**Alternative: use Google Colab (no installation, no animation)**

**T3: FPEOS tutorial**

**Requires a C++ compiler for all calculations, uses Python for all graphics.**

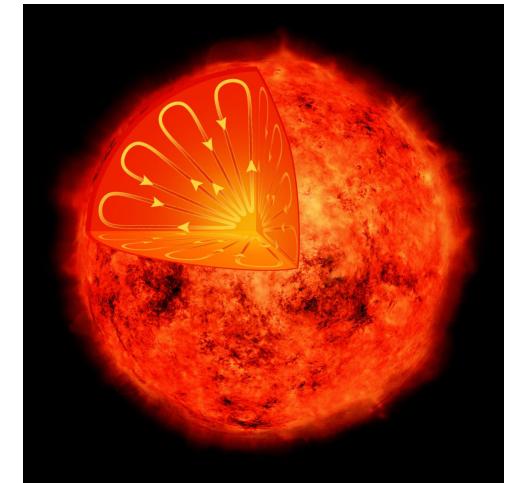
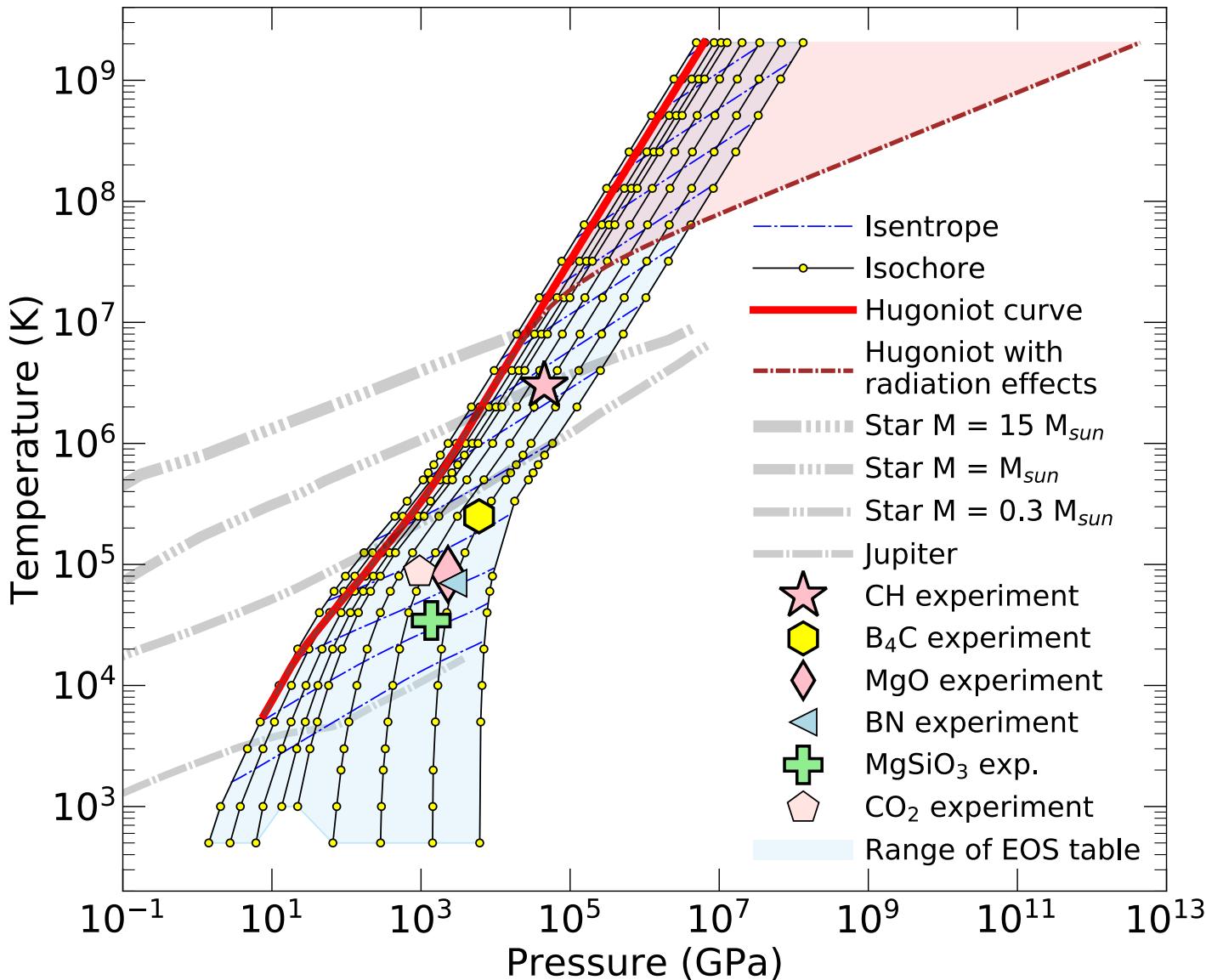
# Outline of lecture 1

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- Path integral Monte Carlo (PIMC) method
- Comparison with different experiments
- First Principles Equation of state (FPEOS) database

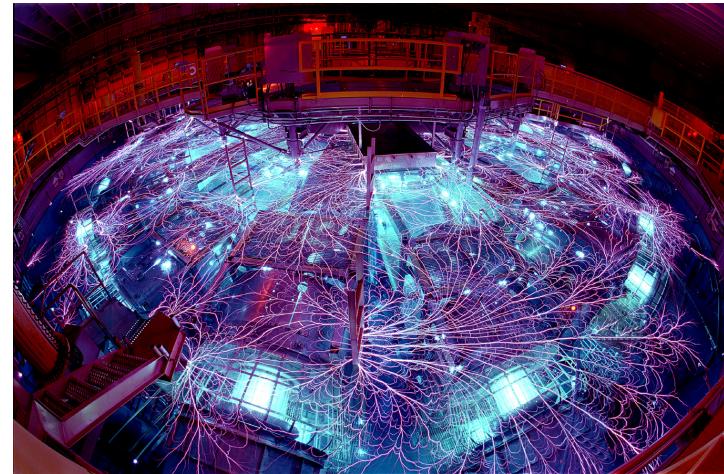
# Center for Matter under Extreme Conditions (CMEC)



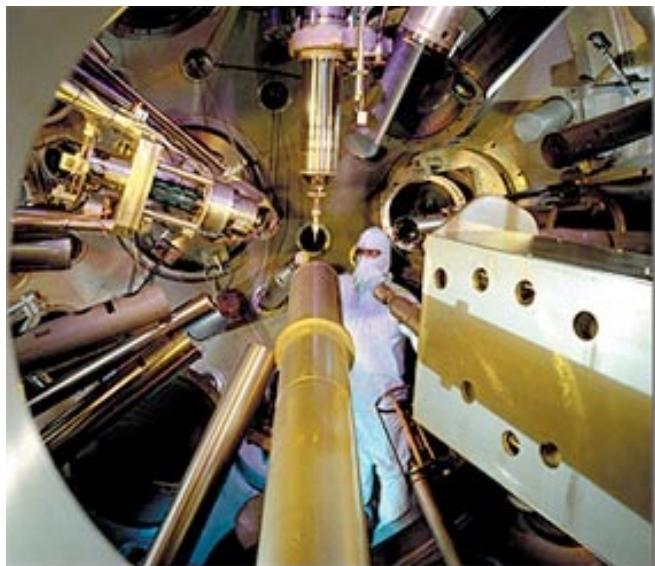
# Study planetary interiors in the laboratory: shock wave experiments



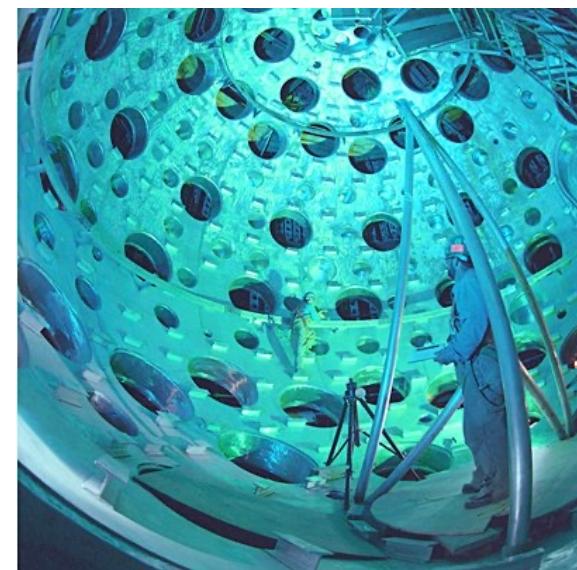
Two-stage gas gun (Livermore) 0.2 Mbar



Z capacitor bank (Sandia) 2 Mbar

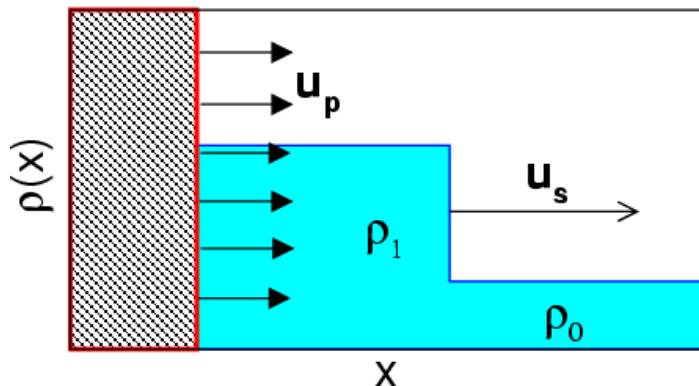


Nova laser (Livermore) 3.4 Mbar



National Ignition Facility 700 Mbar

# Shock wave measurements determine the Equation of State on the Hugoniot curve

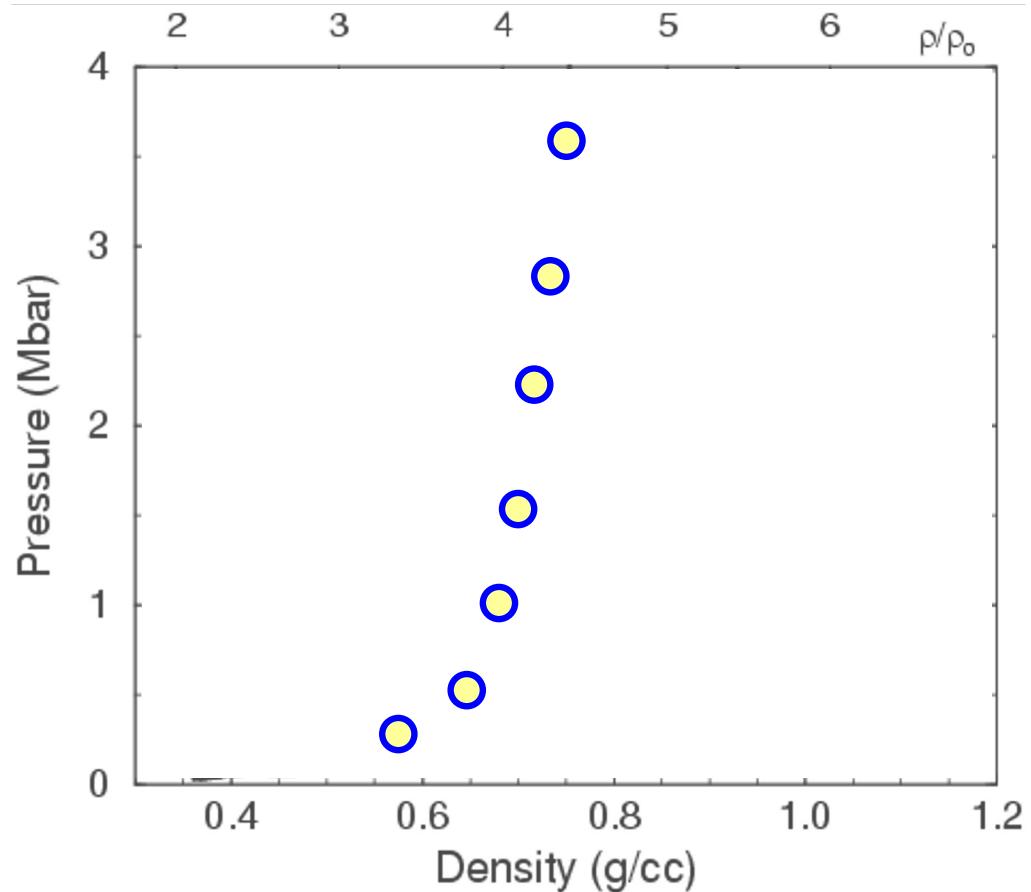


Conservation of mass, momentum and energy yields:

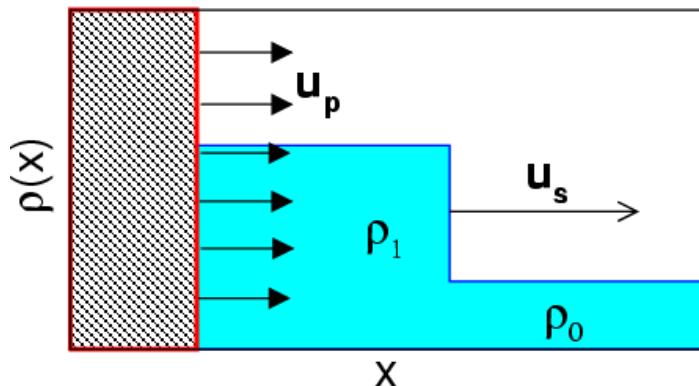
$$\rho = \rho_0 \frac{u_s}{u_s - u_p}$$

$$P = P_0 + \rho_0 u_s u_p$$

$$E = E_0 + \frac{1}{2} (V_0 - V) (P + P_0)$$



# Shock wave measurements determine the Equation of State on the Hugoniot curve

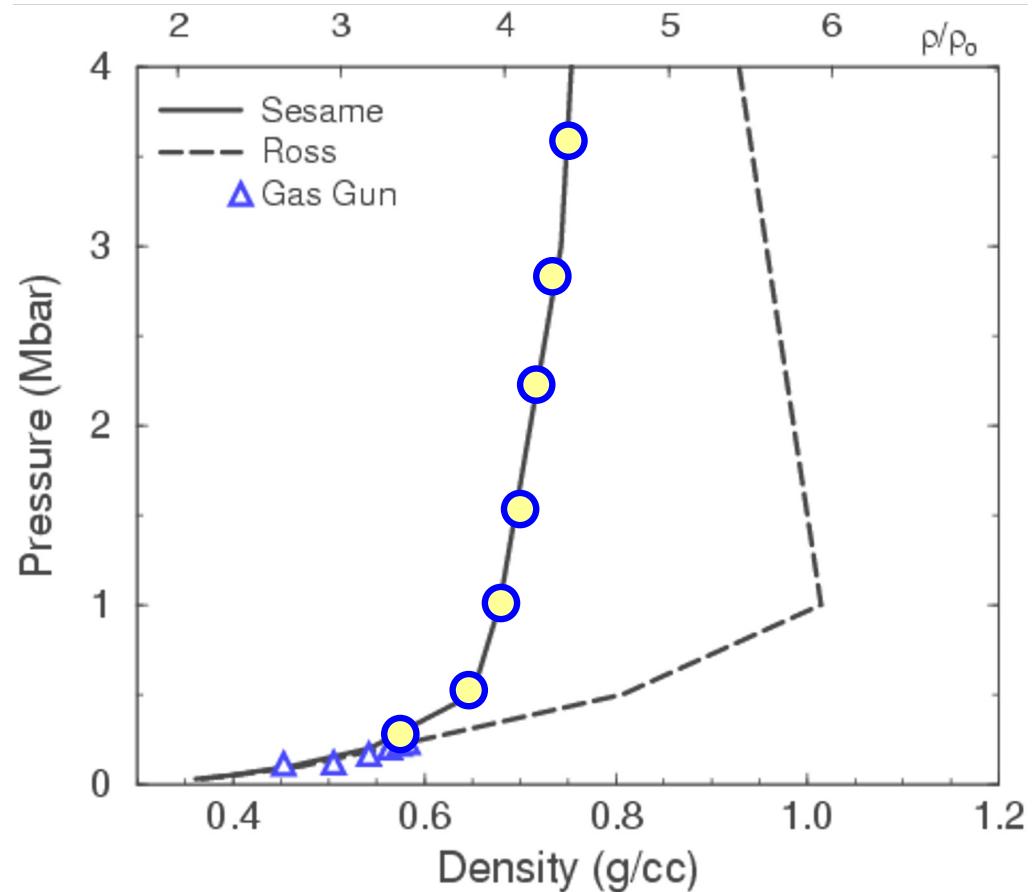


Conservation of mass, momentum and energy yields:

$$\rho = \rho_0 \frac{u_s}{u_s - u_p}$$

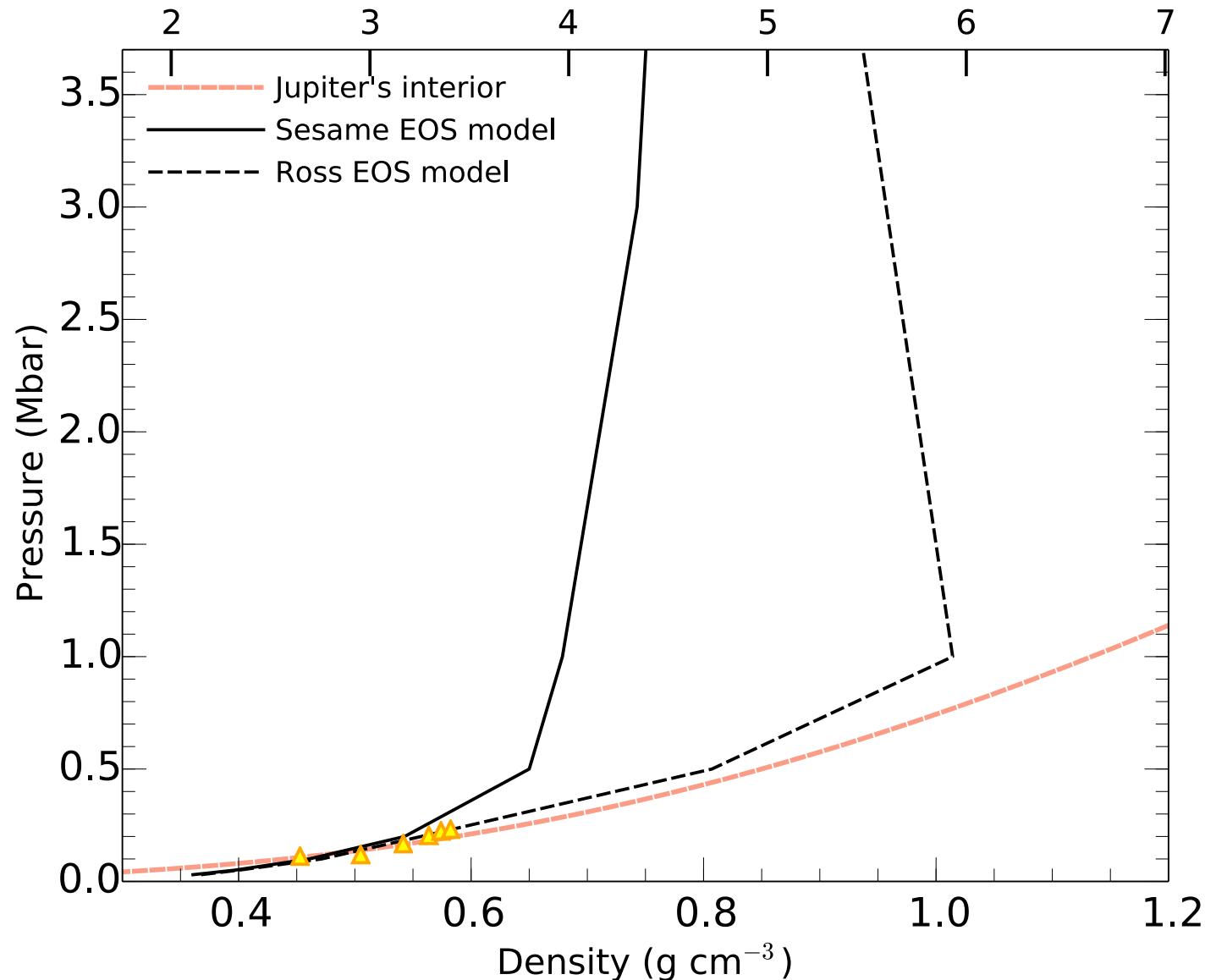
$$P = P_0 + \rho_0 u_s u_p$$

$$E = E_0 + \frac{1}{2} (V_0 - V) (P + P_0)$$

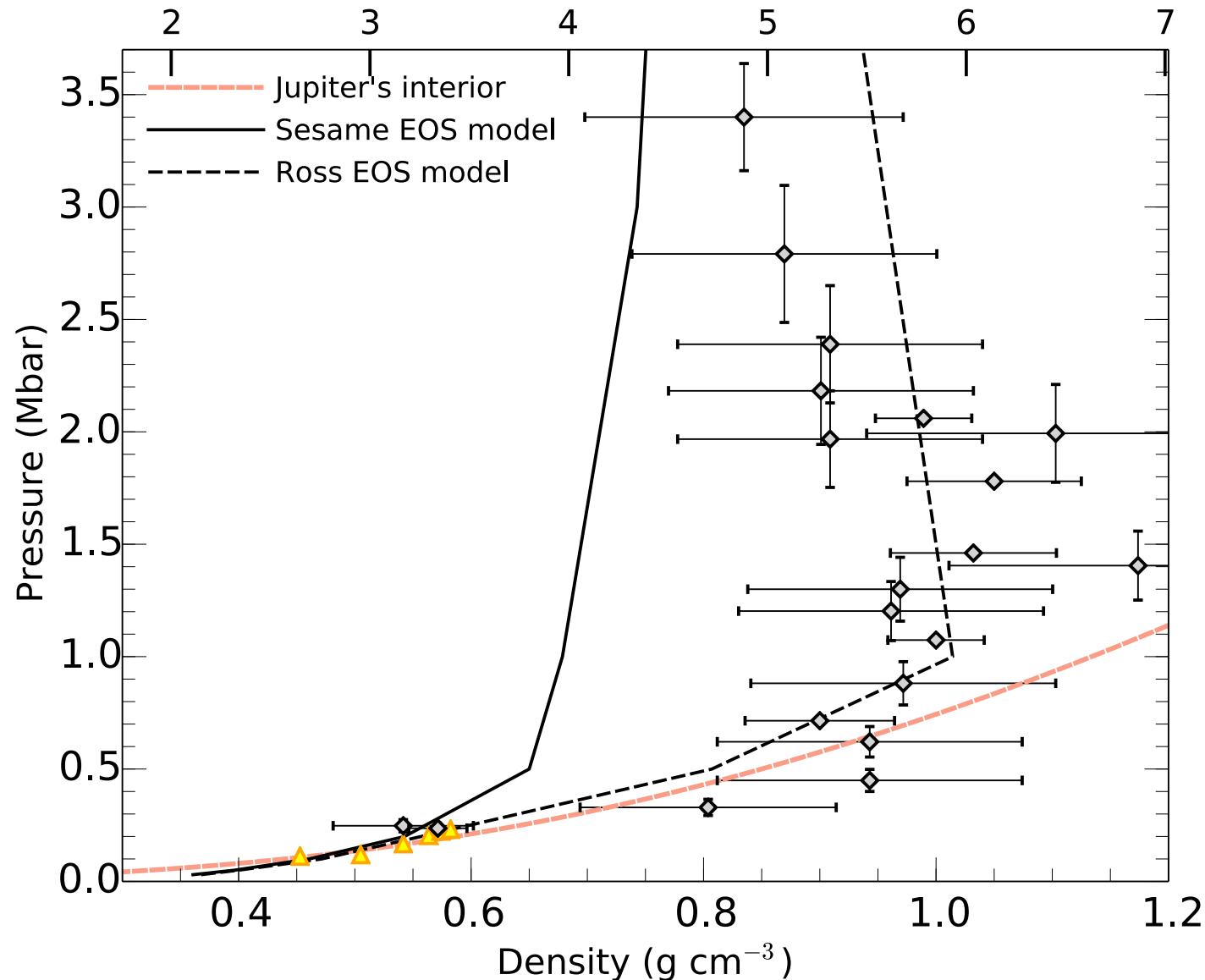


# Comparison of Simulation Results and Shock Wave experiments of Deuterium

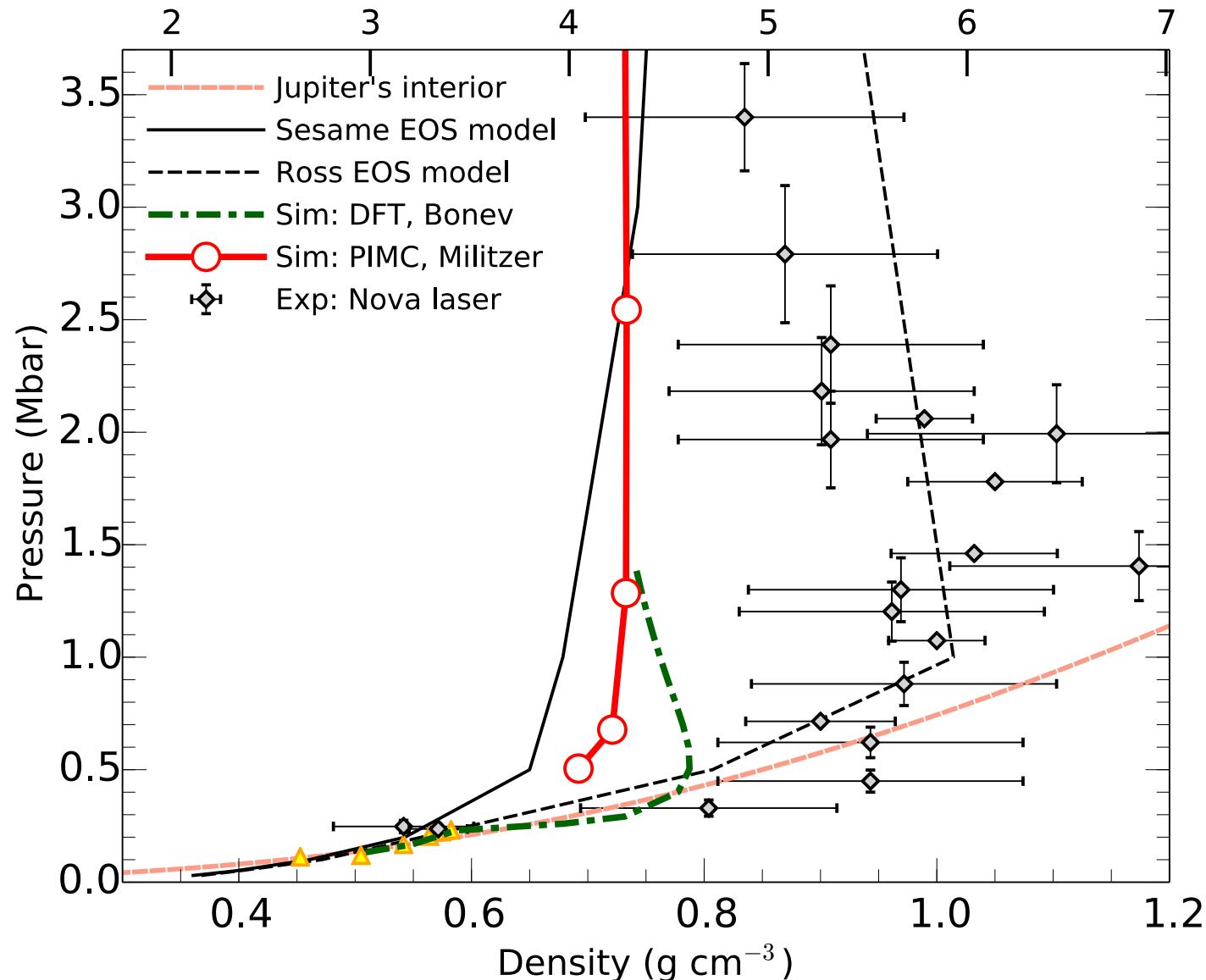
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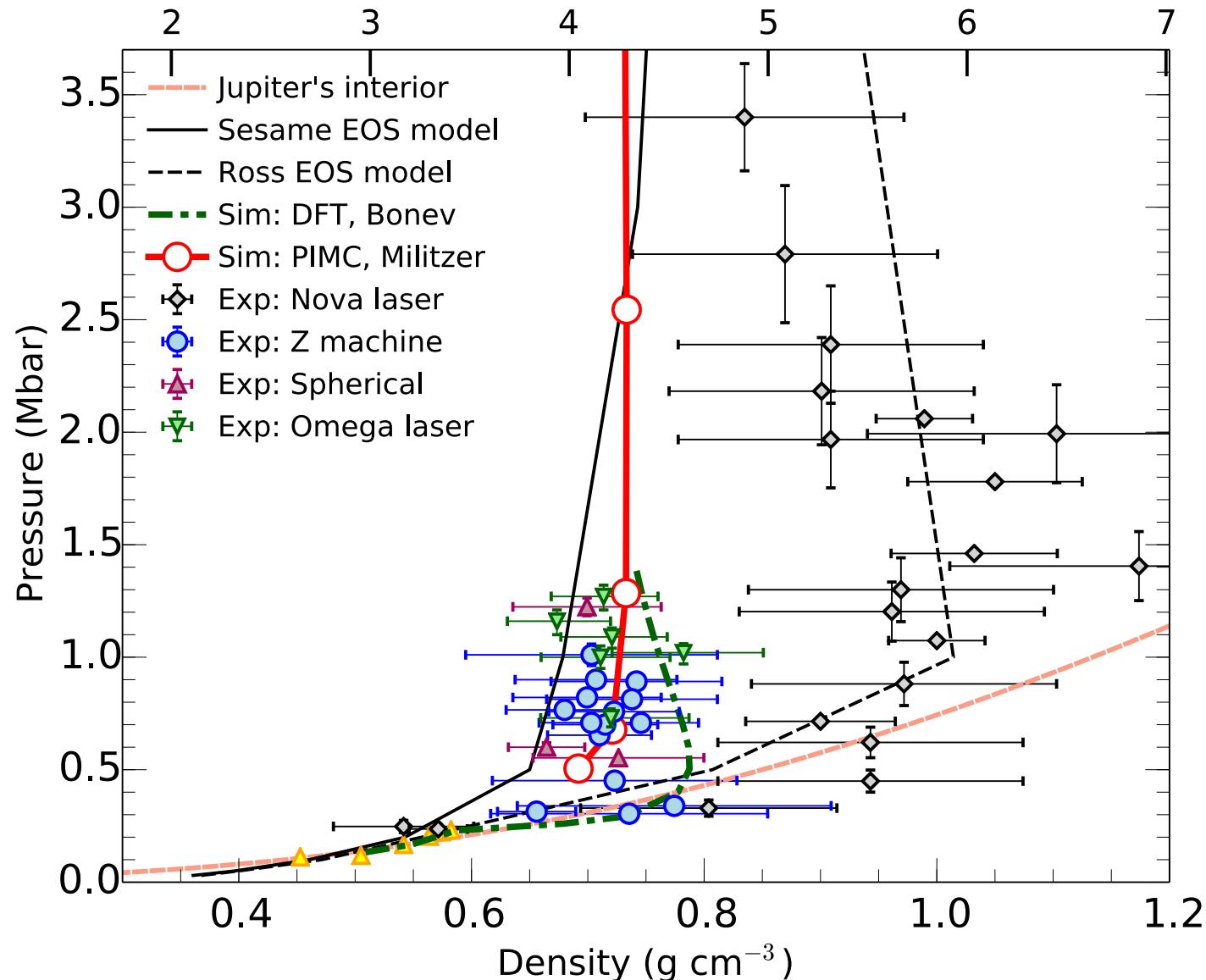
# Comparison of Simulation Results and Shock Wave experiments of Deuterium



# Comparison of Simulation Results and Shock Wave experiments of Deuterium



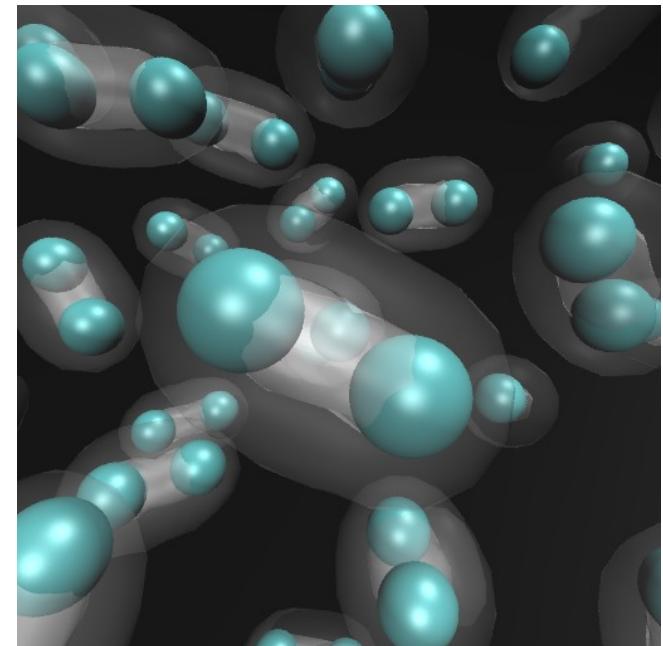
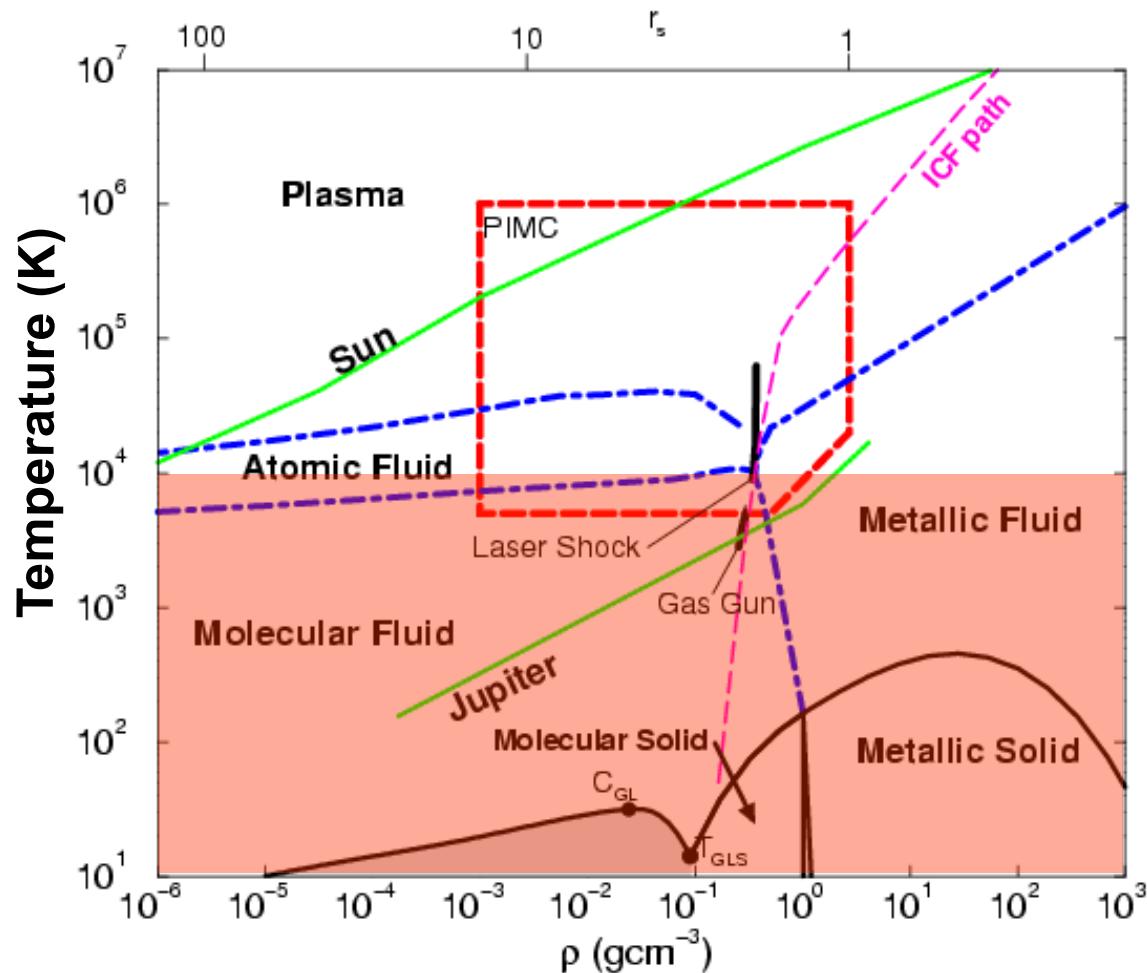
# Comparison of Simulation Results and Shock Wave experiments of Deuterium



I.

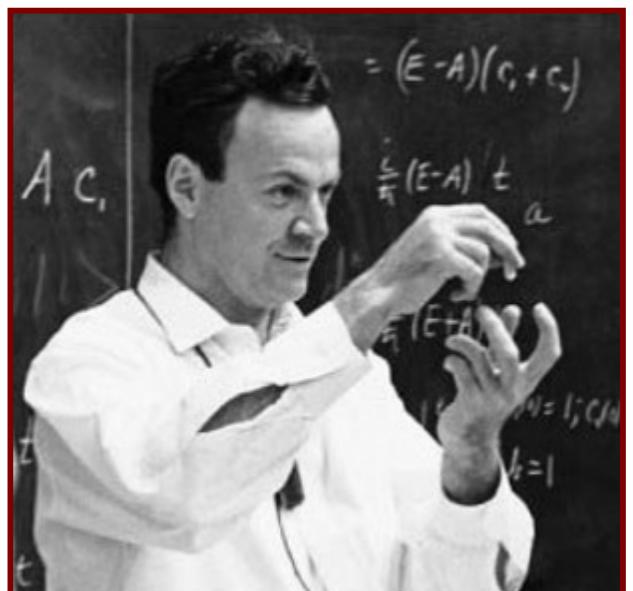
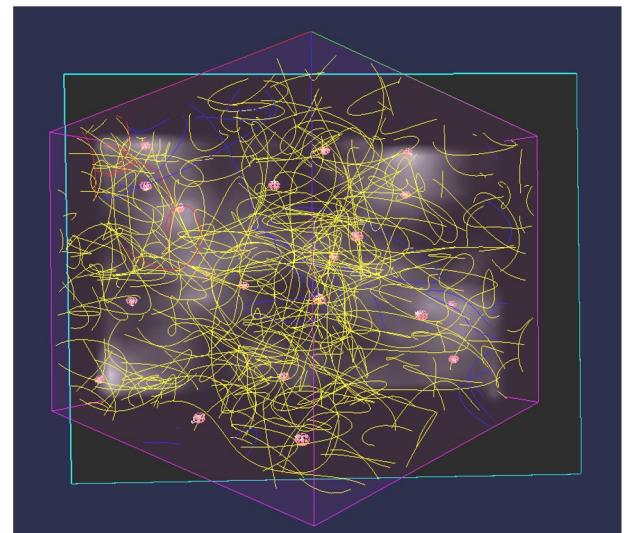
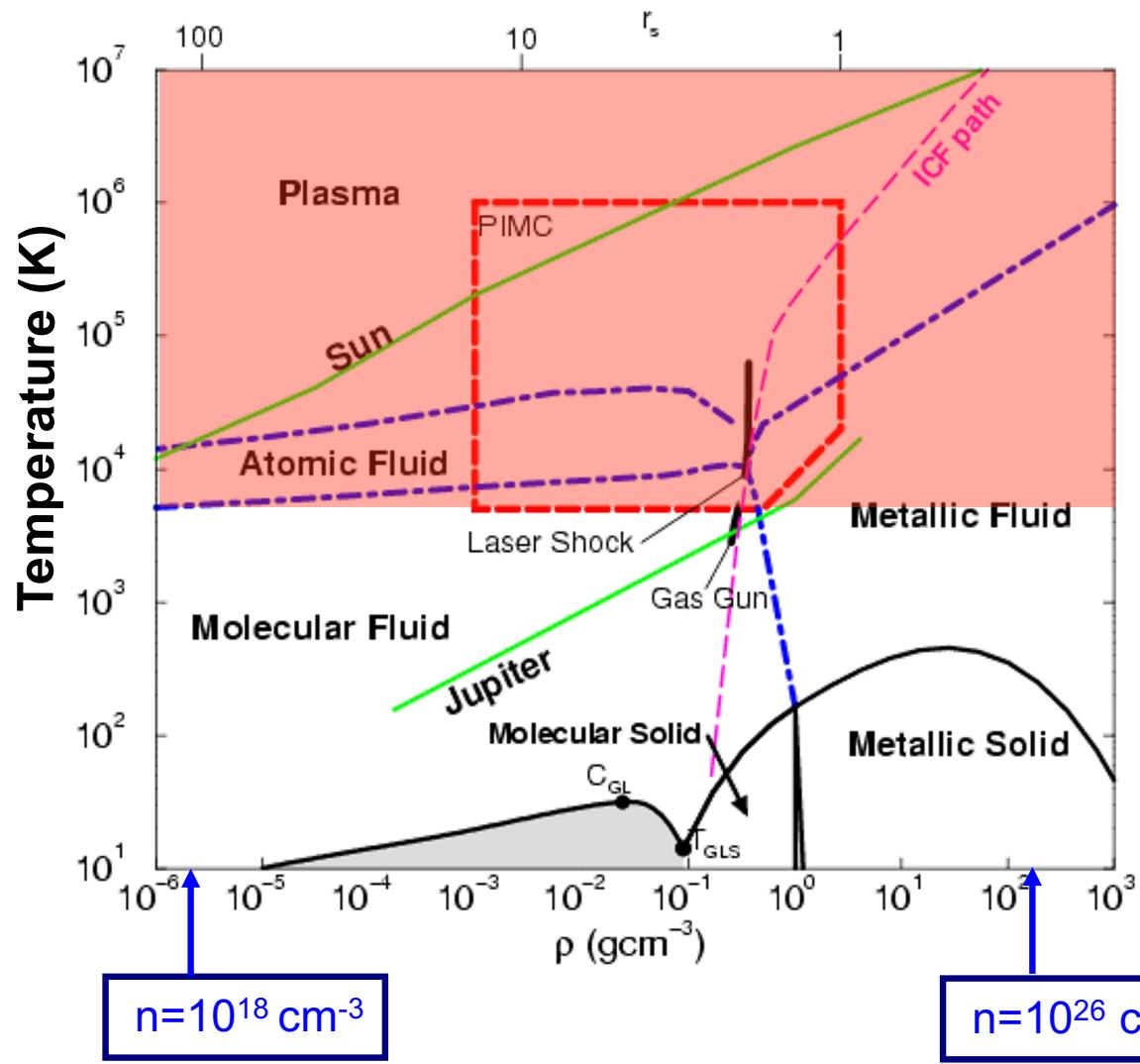
# Path Integral Monte Carlo

# Density functional molecular dynamics at lower T



Born-Oppenheimer approx.  
MD with classical nuclei:  
 $F = m a$   
Forces derived DFT with  
electrons in the instantaneous  
ground state.

# Path integral Monte Carlo at high $T > 10^4 \dots 10^6$ K



# Starting from Restricted PIMC Simulations of Hydrogen

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PHYSICAL REVIEW  
LETTERS

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VOLUME 73

17 OCTOBER 1994

NUMBER 16

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**Equation of State of the Hydrogen Plasma by Path Integral Monte Carlo Simulation**

C. Pierleoni,<sup>1,2,\*</sup> D. M. Ceperley,<sup>3</sup> B. Bernu,<sup>1</sup> and W. R. Magro<sup>3</sup>

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VOLUME 76, NUMBER 8

PHYSICAL REVIEW LETTERS

19 FEBRUARY 1996

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**Molecular Dissociation in Hot, Dense Hydrogen**

W. R. Magro,<sup>1</sup> D. M. Ceperley,<sup>2</sup> C. Pierleoni,<sup>3</sup> and B. Bernu<sup>4</sup>

# Canonical Ensembles: Classical

---

**Boltzmann factor**

$$e^{-E / k_B T}$$

**Thermodynamic averages:**

$$Z_{Cl} = \sum_S e^{-\beta E_S}$$

# Canonical Ensembles: Classical                          Quantum

Boltzmann factor

$$e^{-E / k_B T}$$

Density matrix

$$\hat{\rho} = e^{-\beta \hat{H}}$$

$$\rho(R, R', \beta) = \langle R | e^{-\beta \hat{H}} | R' \rangle$$

$$\rho(R, R', \beta) = \sum_S e^{-\beta E_S} \Psi_S^*(R) \Psi_S(R')$$

Thermodynamic averages:

$$Z_{Cl} = \sum_S e^{-\beta E_S}$$

$$Z_Q = Tr[\hat{\rho}] = \int dR \langle R | e^{-\beta \hat{H}} | R \rangle$$
$$\langle \hat{O} \rangle = \frac{Tr[\hat{O} \hat{\rho}]}{Tr[\hat{\rho}]}$$

# Step 1 towards the path integral

## Matrix squaring property of the density matrix

---

Matrix squaring in operator notation:

$$\hat{\rho} = e^{-\beta \hat{H}} = \left( e^{-(\beta/2)\hat{H}} \right) \left( e^{-(\beta/2)\hat{H}} \right), \quad \beta = \frac{1}{k_B T}$$

Matrix squaring in real-space notation:

$$\langle R | \hat{\rho} | R' \rangle = \int dR_1 \langle R | e^{-(\beta/2)\hat{H}} | R_1 \rangle \langle R_1 | e^{-(\beta/2)\hat{H}} | R' \rangle$$

Matrix squaring in matrix notation:

$$\begin{bmatrix} \dots & R' & \dots \\ R & \ddots & \vdots \\ \dots & \dots & \dots \end{bmatrix} = \begin{bmatrix} \dots & R_1 & \dots \\ R & \ddots & \vdots \\ \dots & \dots & \dots \end{bmatrix} * \begin{bmatrix} \dots & R' & \dots \\ R_1 & \ddots & \vdots \\ \dots & \dots & \dots \end{bmatrix}$$

# Repeat the matrix squaring step

---

Matrix squaring in operator notation:

$$\hat{\rho} = e^{-\beta \hat{H}} = \left( e^{-(\beta/4)\hat{H}} \right)^4, \quad \beta = \frac{1}{k_B T}$$

Matrix squaring in real-space notation:

$$\langle R | \hat{\rho} | R' \rangle = \int dR_1 \int dR_2 \int dR_3 \langle R | e^{-(\beta/4)\hat{H}} | R_1 \rangle \langle R_1 | e^{-(\beta/4)\hat{H}} | R_2 \rangle \langle R_2 | e^{-(\beta/4)\hat{H}} | R_3 \rangle \langle R_3 | e^{-(\beta/4)\hat{H}} | R' \rangle$$

# Path Integrals in Imaginary Time

Every particle is represented by a path, a ring polymer.

**Density matrix:**

$$\hat{\rho} = e^{-\beta \hat{H}} = \left( e^{-\tau \hat{H}} \right)^M, \quad \beta = \frac{1}{k_B T}, \quad \tau = \frac{\beta}{M}$$

$$\langle \hat{O} \rangle = \frac{\text{Tr}[\hat{O}\hat{\rho}]}{\text{Tr}[\hat{\rho}]}$$

**Trotter break-up:**

$$\langle R | \hat{\rho} | R' \rangle = \langle R | (e^{-\tau \hat{H}})^M | R' \rangle = \int dR_1 \dots \int dR_{M-1} \langle R | e^{-\tau \hat{H}} | R_1 \rangle \langle R_1 | e^{-\tau \hat{H}} | R_2 \rangle \dots \langle R_{M-1} | e^{-\tau \hat{H}} | R' \rangle$$

# Path Integrals in Imaginary Time

Simplest form for the paths' action: primitive approx.

**Density matrix:**

$$\hat{\rho} = e^{-\beta \hat{H}} = \left( e^{-\tau \hat{H}} \right)^M, \quad \beta = \frac{1}{k_B T}, \quad \tau = \frac{\beta}{M}$$

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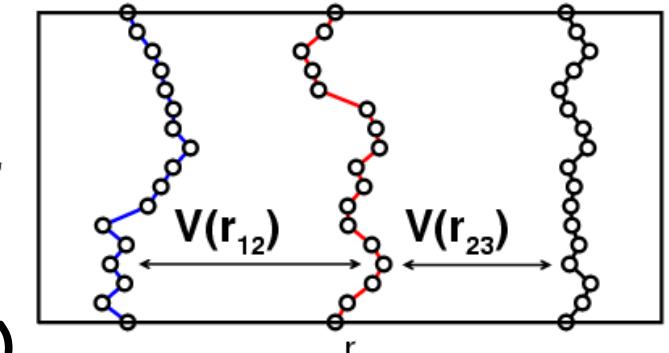
**Trotter formula:**

$$e^{-\beta(\hat{T} + \hat{V})} = \lim_{M \rightarrow \infty} \left[ e^{-\tau \hat{T}} e^{-\tau \hat{V}} \right]^M$$

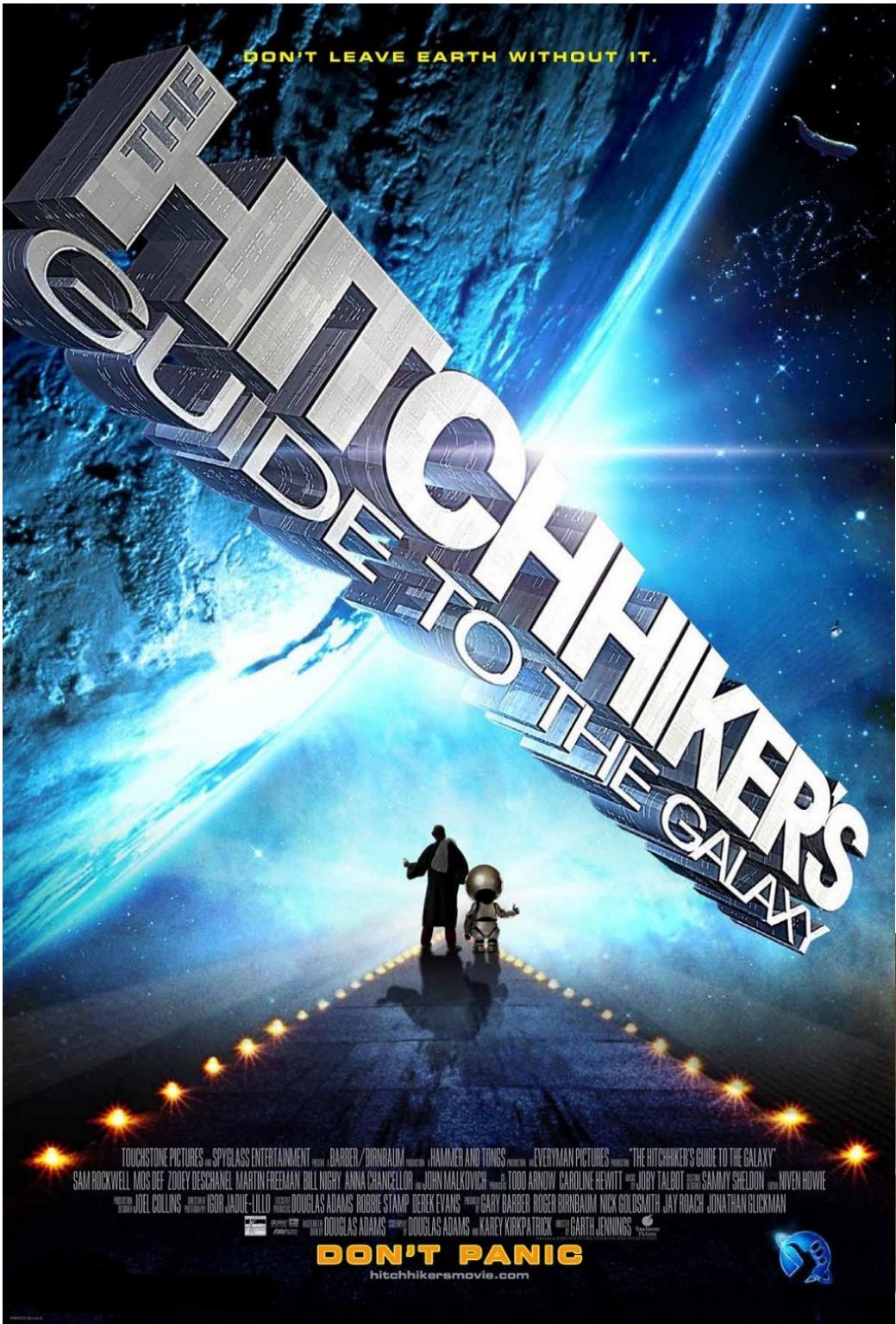
**Path integral and primitive action  $S$ :**

$$\langle R | \hat{\rho} | R' \rangle = \oint_{R \rightarrow R'} dR_t e^{-S[R_t]}$$

$$S[R_t] = \sum_{i=1}^M \frac{(R_{i+1} - R_i)^2}{4\lambda\tau} + \frac{\tau}{2} [V(R_i) + V(R_{i+1})]$$



Pair action: Miltzner, Comp. Phys. Comm. (2016)



DON'T  
PANIC

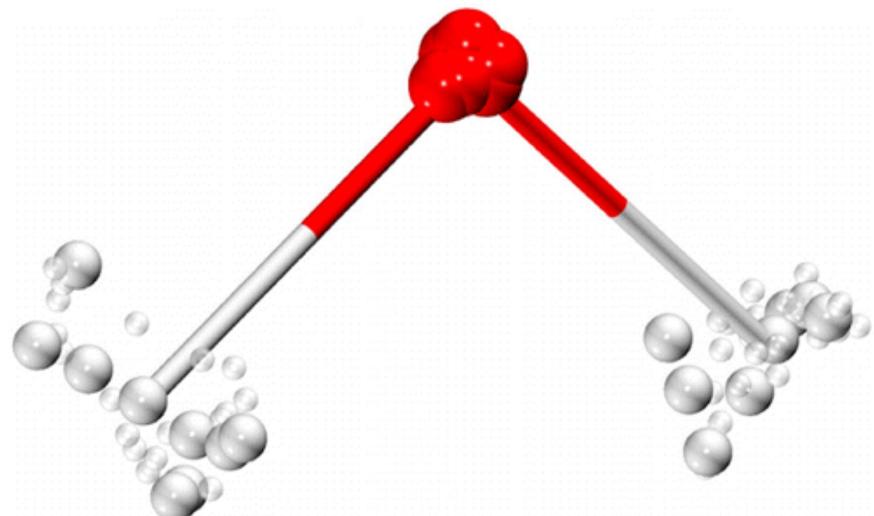
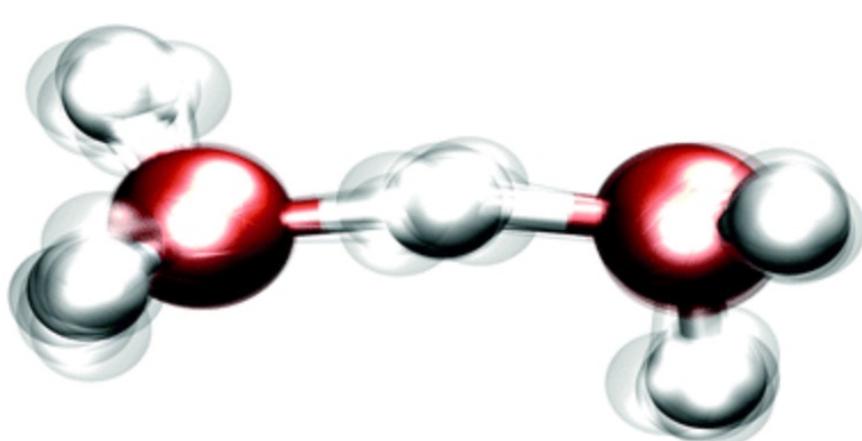
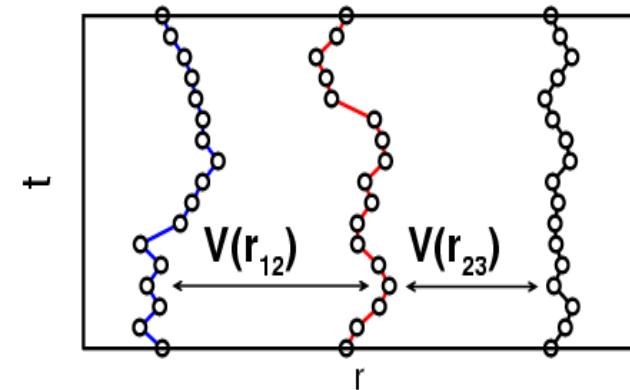
Douglas Adams:  
*"Infinite Improbability Drive"*  
of spaceship  
*"Heart of Gold"*



# Path Integrals in Imaginary Time include Zero-Point Motion and some Tunnelling Effects

$$\langle R | \hat{\rho} | R' \rangle = \oint_{R \rightarrow R'} dR_t e^{-S[R_t]}$$

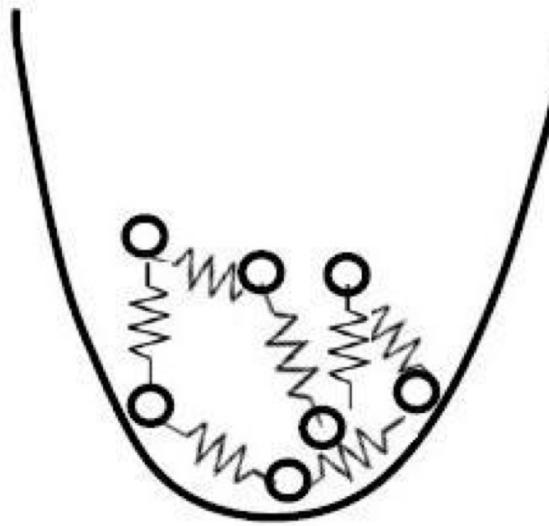
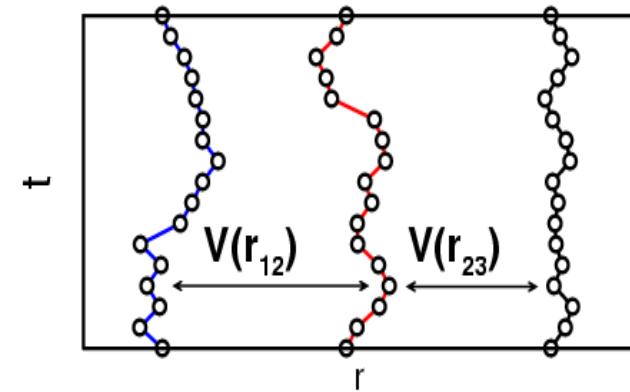
$$S[R_t] = \sum_{i=1}^M \frac{(R_{i+1} - R_i)^2}{4\lambda\tau} + \frac{\tau}{2} [V(R_i) + V(R_{i+1})]$$



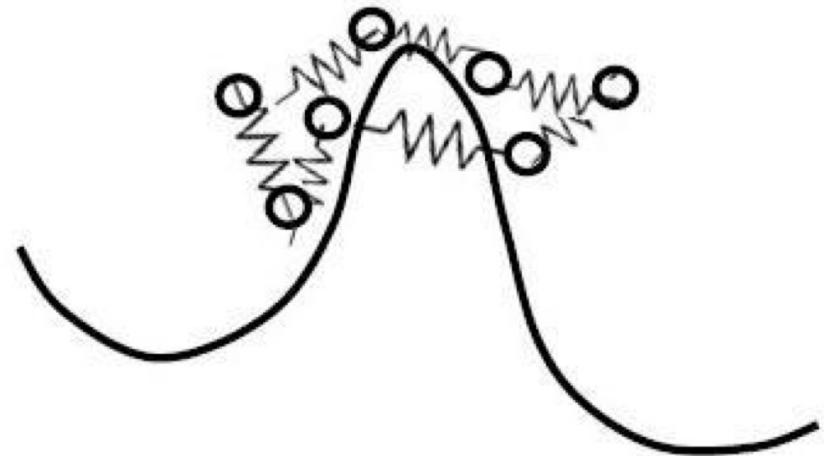
# Path Integrals in Imaginary Time include Zero-Point Motion and some Tunnelling Effects

$$\langle R | \hat{\rho} | R' \rangle = \oint_{R \rightarrow R'} dR_t e^{-S[R_t]}$$

$$S[R_t] = \sum_{i=1}^M \frac{(R_{i+1} - R_i)^2}{4\lambda\tau} + \frac{\tau}{2} [V(R_i) + V(R_{i+1})]$$



zero-point energy

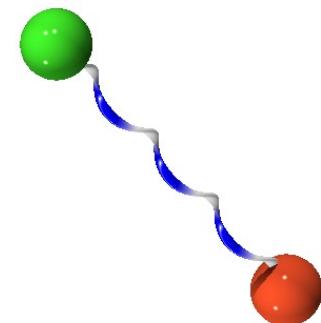
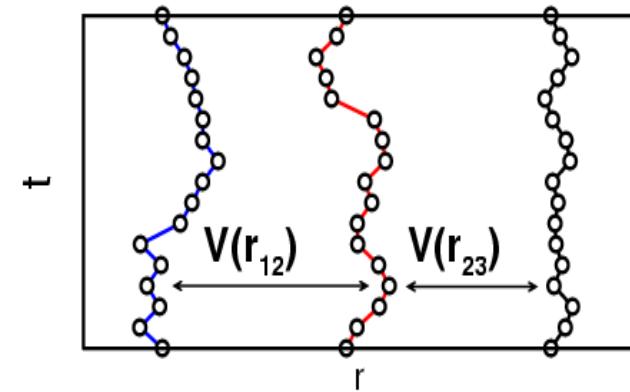


tunneling

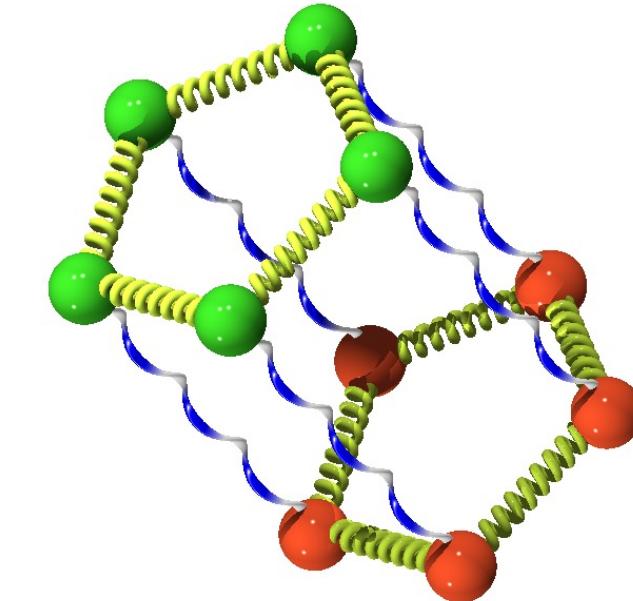
# Path Integrals in Imaginary Time include Zero-Point Motion and some Tunnelling Effects

$$\langle R | \hat{\rho} | R' \rangle = \oint_{R \rightarrow R'} dR_t e^{-S[R_t]}$$

$$S[R_t] = \sum_{i=1}^M \frac{(R_{i+1} - R_i)^2}{4\lambda\tau} + \frac{\tau}{2} [V(R_i) + V(R_{i+1})]$$



Classical limit ( $P=1$ )



Path integral (here with  $P=5$ )

# Bosonic and Fermionic Density Matrices

---

**Bosonic density matrix:**

Sum over all symmetric eigenstates.

$$\rho_B(R,R',\beta) = \sum_i e^{-\beta E_i} \Psi_S^{[i]*}(R) \Psi_S^{[i]}(R')$$

**Fermionic density matrix:**

Sum over all antisymmetric eigenstates.

$$\rho_F(R,R',\beta) = \sum_i e^{-\beta E_i} \Psi_{AS}^{[i]*}(R) \Psi_{AS}^{[i]}(R')$$

# Bosonic and Fermionic Path Integrals

**Bosonic density matrix:**

Sum over all symmetric eigenstates.

$$\rho_B(R, R', \beta) = \sum_i e^{-\beta E_i} \Psi_S^{[i]*}(R) \Psi_S^{[i]}(R')$$

Project out the symmetric states:

$$\rho_B(R, R', \beta) = \sum_P (+1)^P \rho_D(R, PR', \beta)$$

**Fermionic density matrix:**

Sum over all antisymmetric eigenstates.

$$\rho_F(R, R', \beta) = \sum_i e^{-\beta E_i} \Psi_{AS}^{[i]*}(R) \Psi_{AS}^{[i]}(R')$$

Project out the antisymmetric states:

$$\rho_F(R, R', \beta) = \sum_P (-1)^P \rho_D(R, PR', \beta)$$

$$\langle R | \hat{\rho}_{F/B} | R' \rangle = \sum_P (\pm 1)^P \int dR_1 \dots \int dR_{M-1} \langle R | e^{-\tau \hat{H}} | R_1 \rangle \dots \langle R_{M-1} | e^{-\tau \hat{H}} | PR' \rangle$$

# Bosonic and Fermionic Path Integrals

**Bosonic density matrix:**

Sum over all symmetric eigenstates.

$$\rho_B(R, R', \beta) = \sum_i e^{-\beta E_i} \Psi_S^{[i]*}(R) \Psi_S^{[i]}(R')$$

Project out the symmetric states:

$$\rho_B(R, R', \beta) = \sum_P (+1)^P \rho_D(R, PR', \beta)$$

**Fermionic density matrix:**

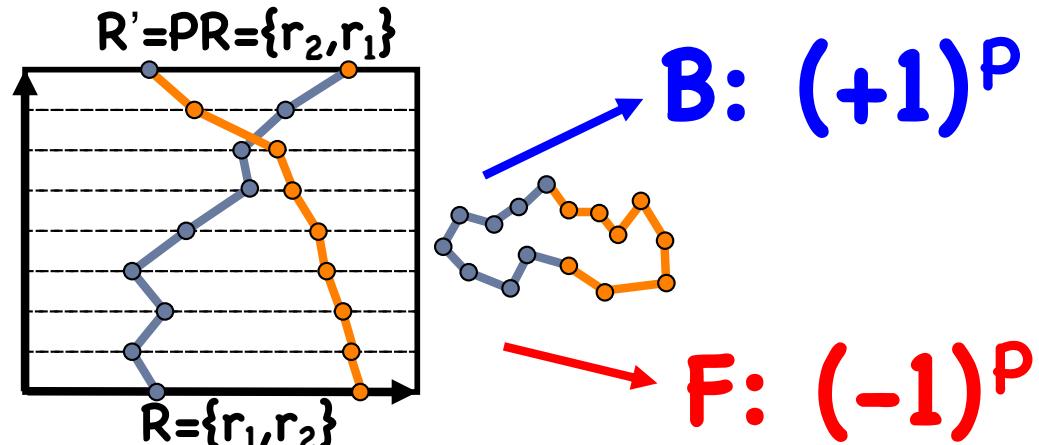
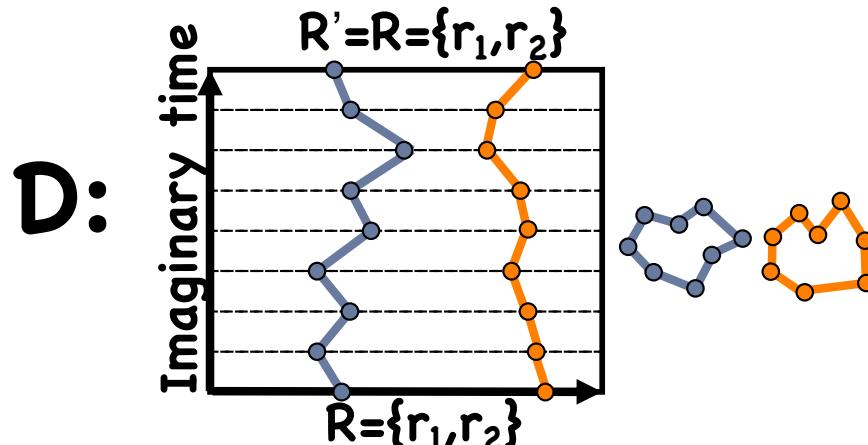
Sum over all antisymmetric eigenstates.

$$\rho_F(R, R', \beta) = \sum_i e^{-\beta E_i} \Psi_{AS}^{[i]*}(R) \Psi_{AS}^{[i]}(R')$$

Project out the antisymmetric states:

$$\rho_F(R, R', \beta) = \sum_P (-1)^P \rho_D(R, PR', \beta)$$

$$\langle R | \hat{\rho}_{F/B} | R' \rangle = \sum_P (\pm 1)^P \int dR_1 \dots \int dR_{M-1} \langle R | e^{-\tau \hat{H}} | R_1 \rangle \dots \langle R_{M-1} | e^{-\tau \hat{H}} | PR' \rangle$$



# Bosonic and Fermionic Path Integrals

Bosonic density matrix:

Sum over all symmetric eigenstates.

$$\rho_B(R, R', \beta) = \sum_i e^{-\beta E_i} |\Psi_{BS}^{[i]}(R)\rangle \langle \Psi_{BS}^{[i]}(R')|$$

Project onto symmetric states:

$$|\rho_B\rangle = P |\rho\rangle$$

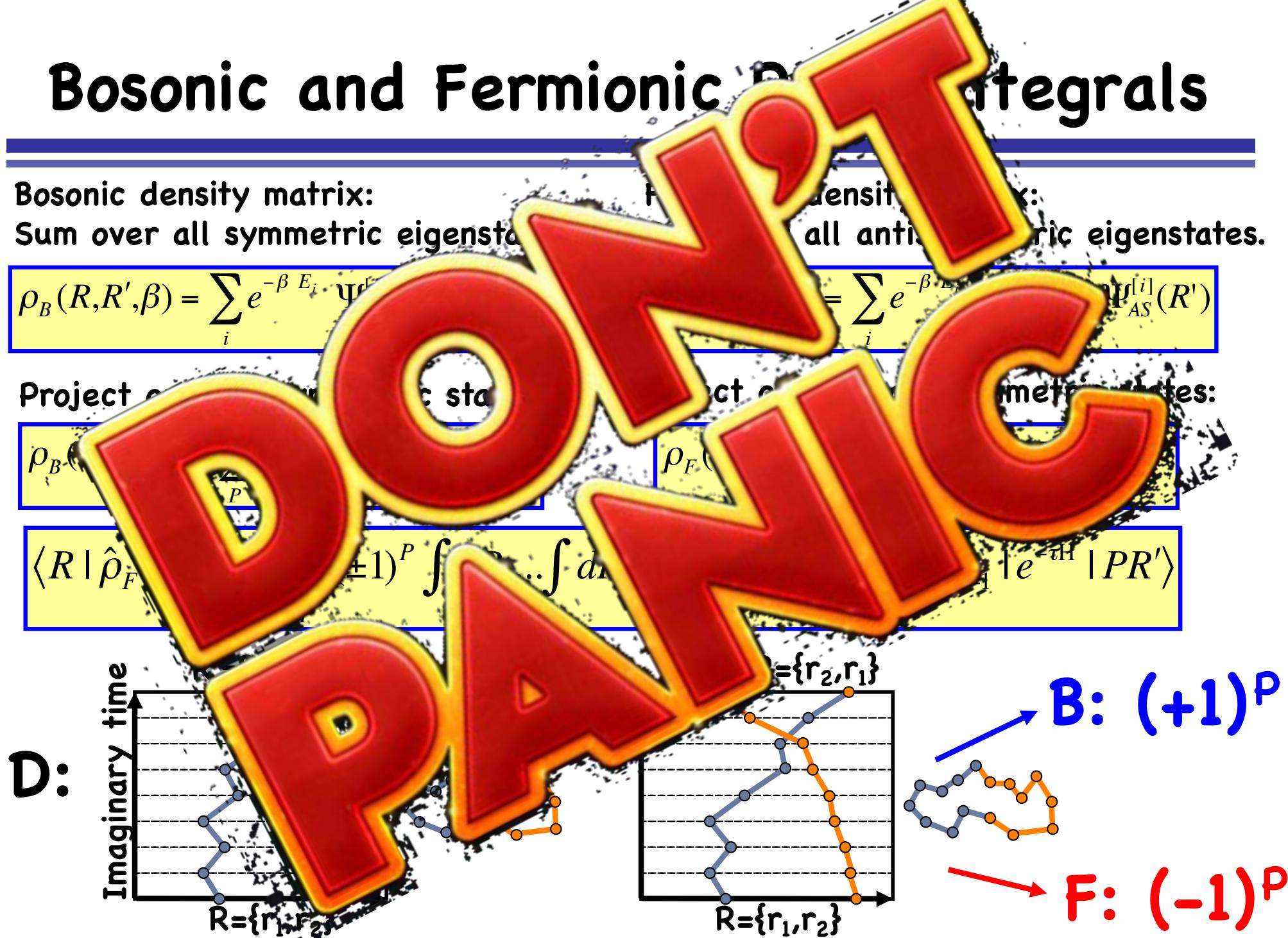
$$\langle R | \hat{\rho}_F | R' \rangle = (-1)^P \int dx_1 \dots \int dx_N e^{-\beta H} |e^{-\beta H} |PR'\rangle$$

Fermionic density matrix:

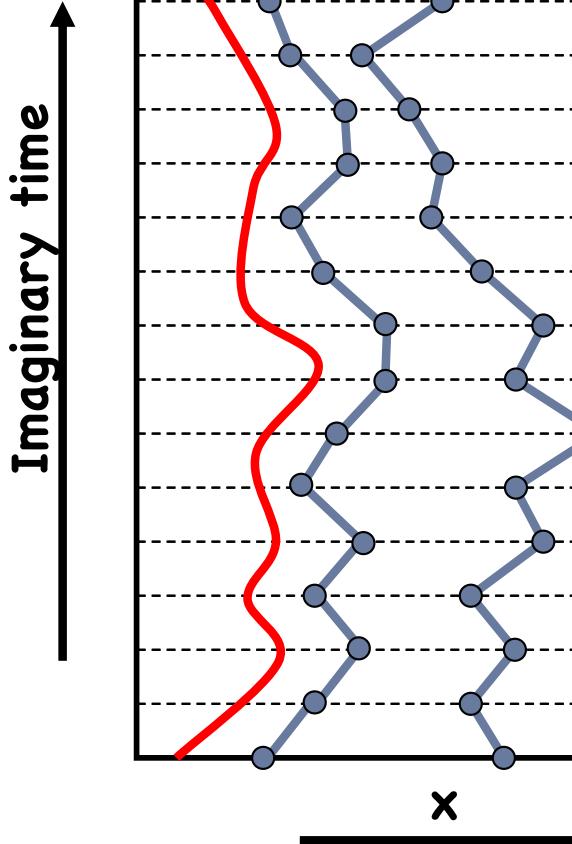
$$\rho_F(R, R', \beta) = \sum_i e^{-\beta E_i} |\Psi_{AF}^{[i]}(R)\rangle \langle \Psi_{AF}^{[i]}(R')|$$

Sum over all antisymmetric eigenstates.

$$\langle R | \hat{\rho}_F | R' \rangle = (-1)^P \int dx_1 \dots \int dx_N e^{-\beta H} |e^{-\beta H} |PR'\rangle$$



# Restricted PIMC for fermions: How is the restriction applied?



Free-particle nodes:

Construct a **fermionic trial density matrix** in form of a Slater determinant of single-particle density matrices:

$$\rho_T(R, R', \beta) = \begin{vmatrix} \rho(r_1, r'_1, \beta) & \cdots & \rho(r_1, r'_N, \beta) \\ \vdots & \ddots & \vdots \\ \rho(r_N, r'_1, \beta) & \cdots & \rho(r_N, r'_N, \beta) \end{vmatrix}$$

Enforce the following nodal condition for all time slices along the paths:

$$\rho_T[R(t), R(0), t] > 0$$

This 3N-dimensional conditions eliminates all negative and some positive contribution to the path → Solves the fermion sign problem approx.

$$\rho_0^{[1]}(r, r'; \beta) = \sum_k e^{-\beta E_k} \Psi_k(r) \Psi_k^*(r')$$

# Fermionic Path Integrals

Example: **Closed paths** of 2 free particles

Distinguishable particles:

Consider path types: **NA** + **NX**

Bosons:

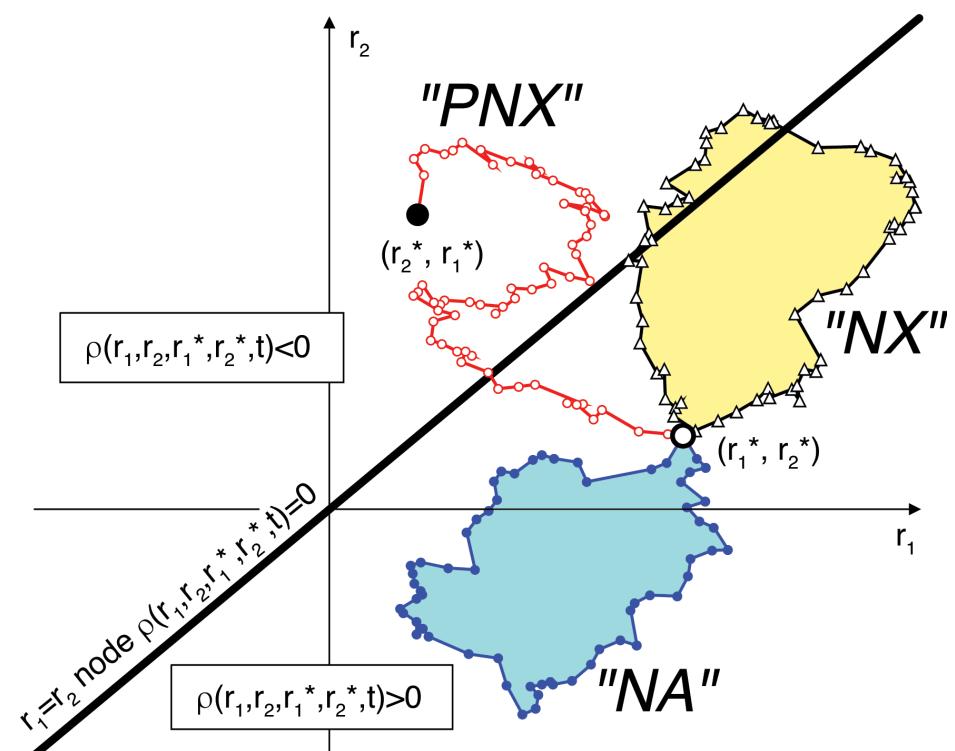
Consider path types: **NA** + **NX** + **PNX**

Direct fermions:

Consider path types: **NA** + **NX** - **PNX**

Restricted fermions:

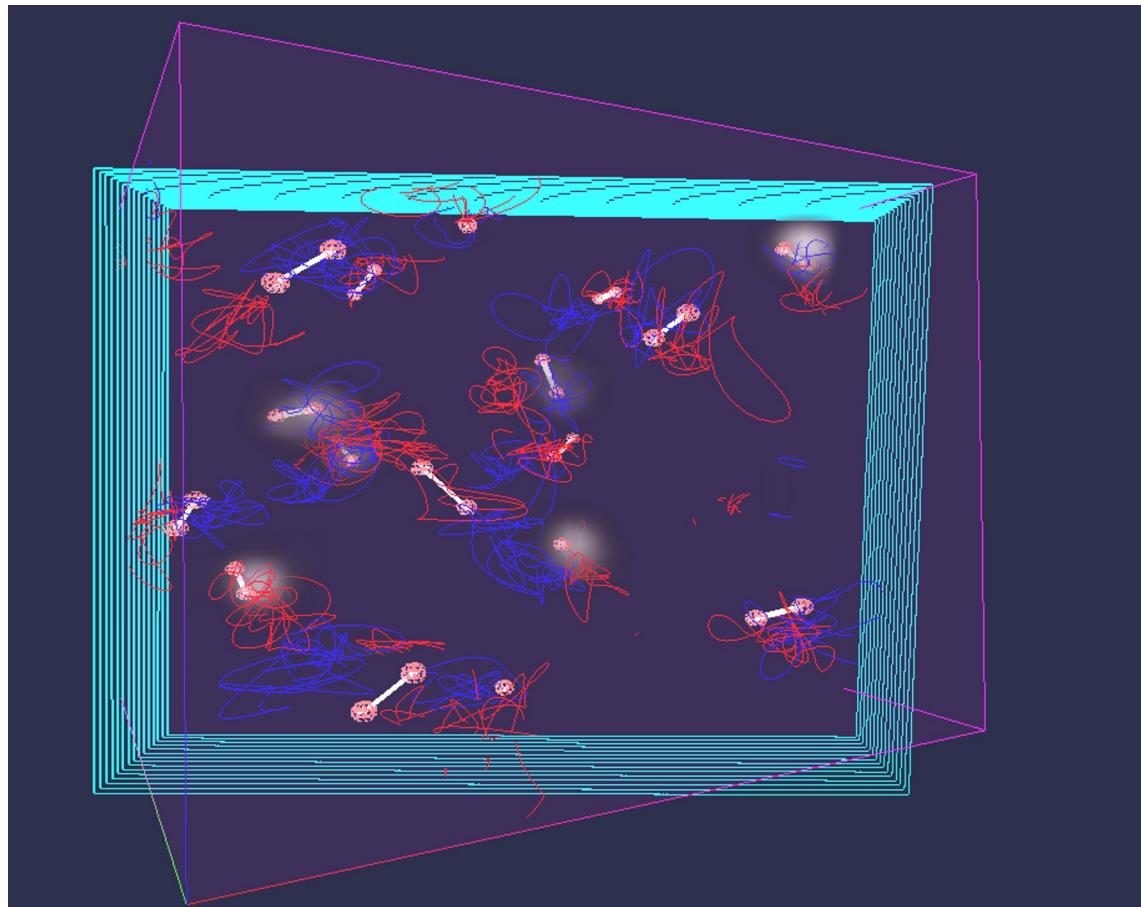
Consider only path type: **NA**



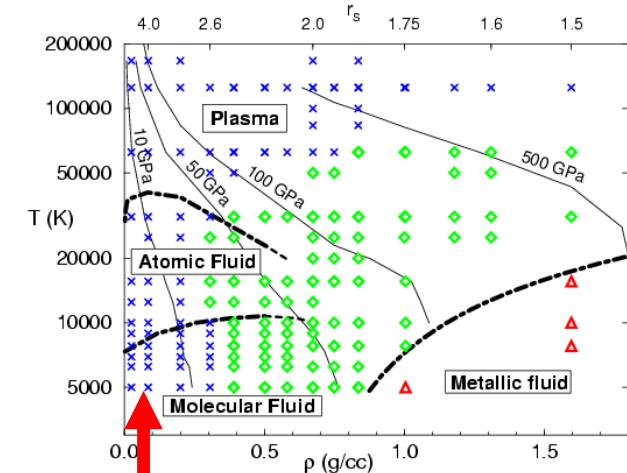
# *I. Hydrogen*

# Molecular Hydrogen

Snapshot from a PIMC simulation with 32 protons and electrons



2 protons (pink spheres) and spin-up and one spin-down electron form one  $\text{H}_2$  molecule.

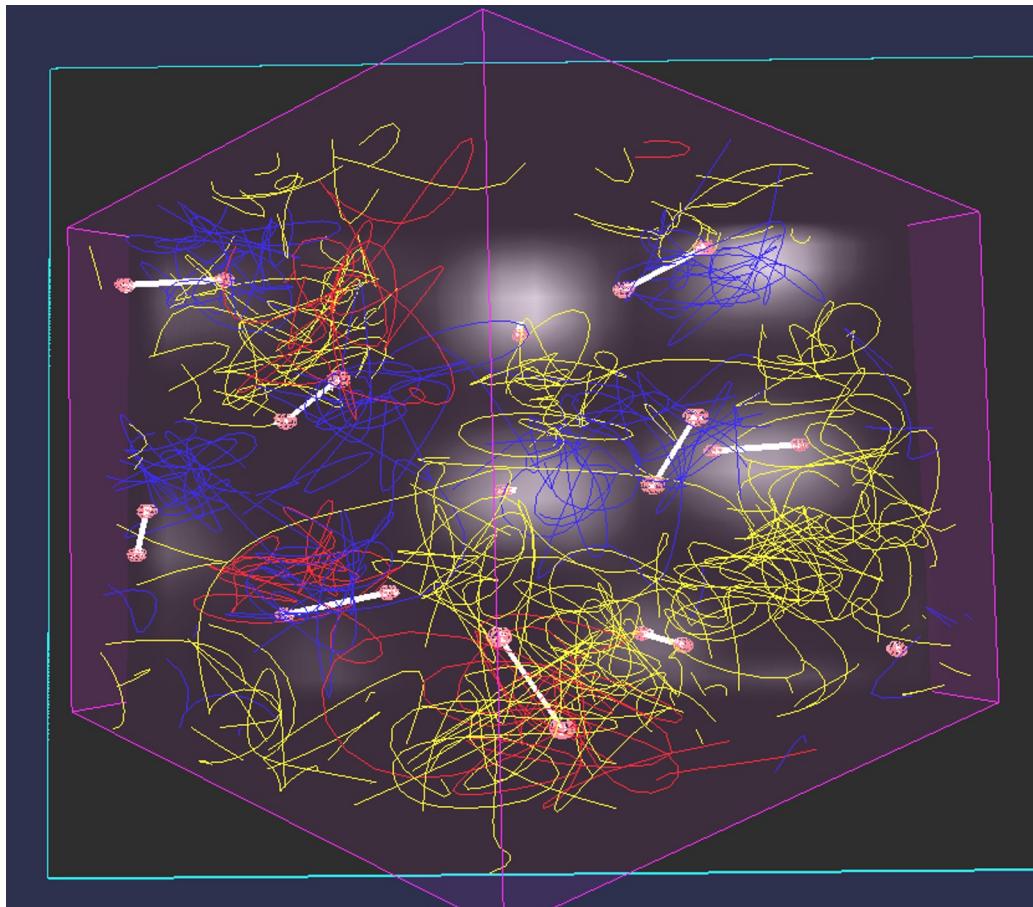


T=5000K,  $r_s=4$

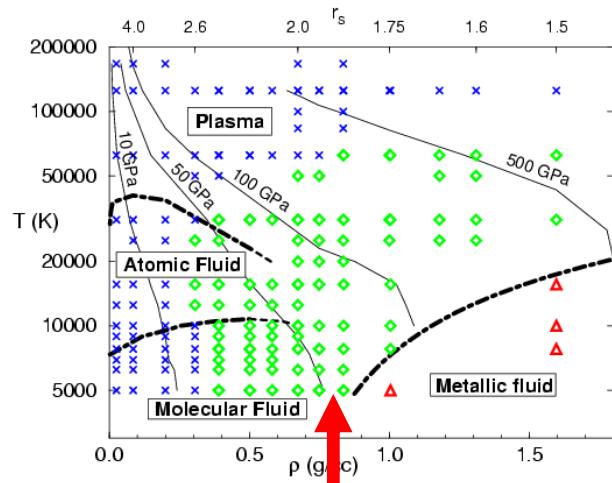
100% molecules,  
weakly interacting

# Molecular Hydrogen

Snapshot from a PIMC simulation with 32 protons and electrons



2 protons (pink spheres) and spin-up and one spin-down electron form one  $H_2$  molecule.

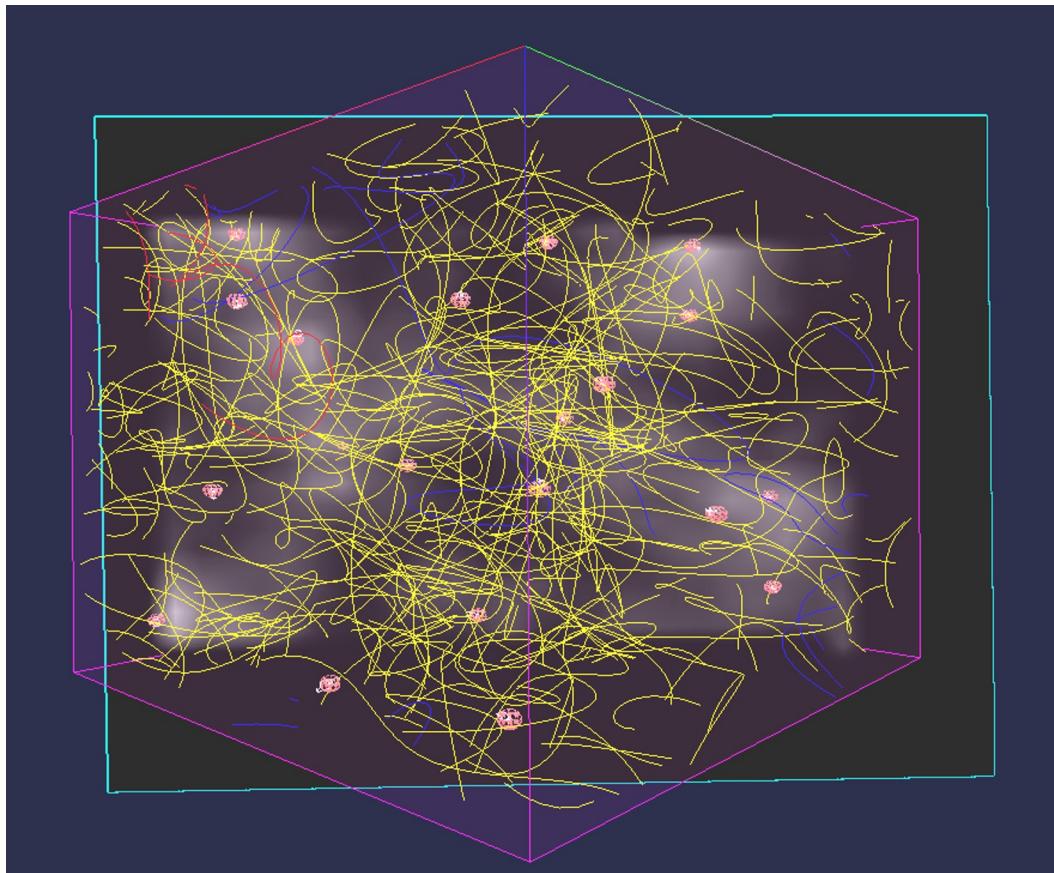


T=5000K,  $r_s=1.86$

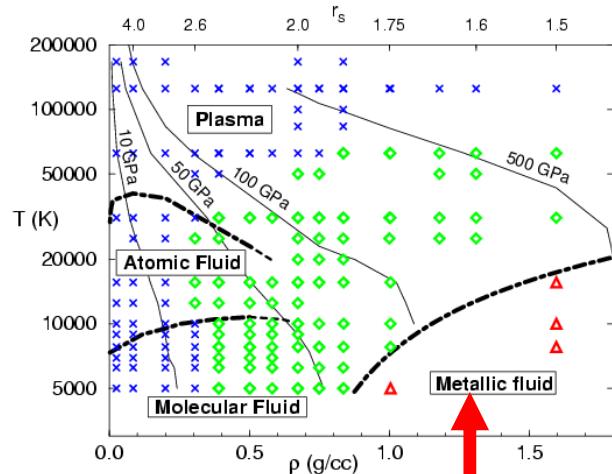
- strongly interacting molecules, close to pressure dissociation
- Electrons are degenerate, partially delocalized
- Electron paths are permuting

# Metallic Hydrogen

Snapshot from a PIMC simulation with 32 protons and electrons



Free protons (**pink spheres**) and delocalized electrons.



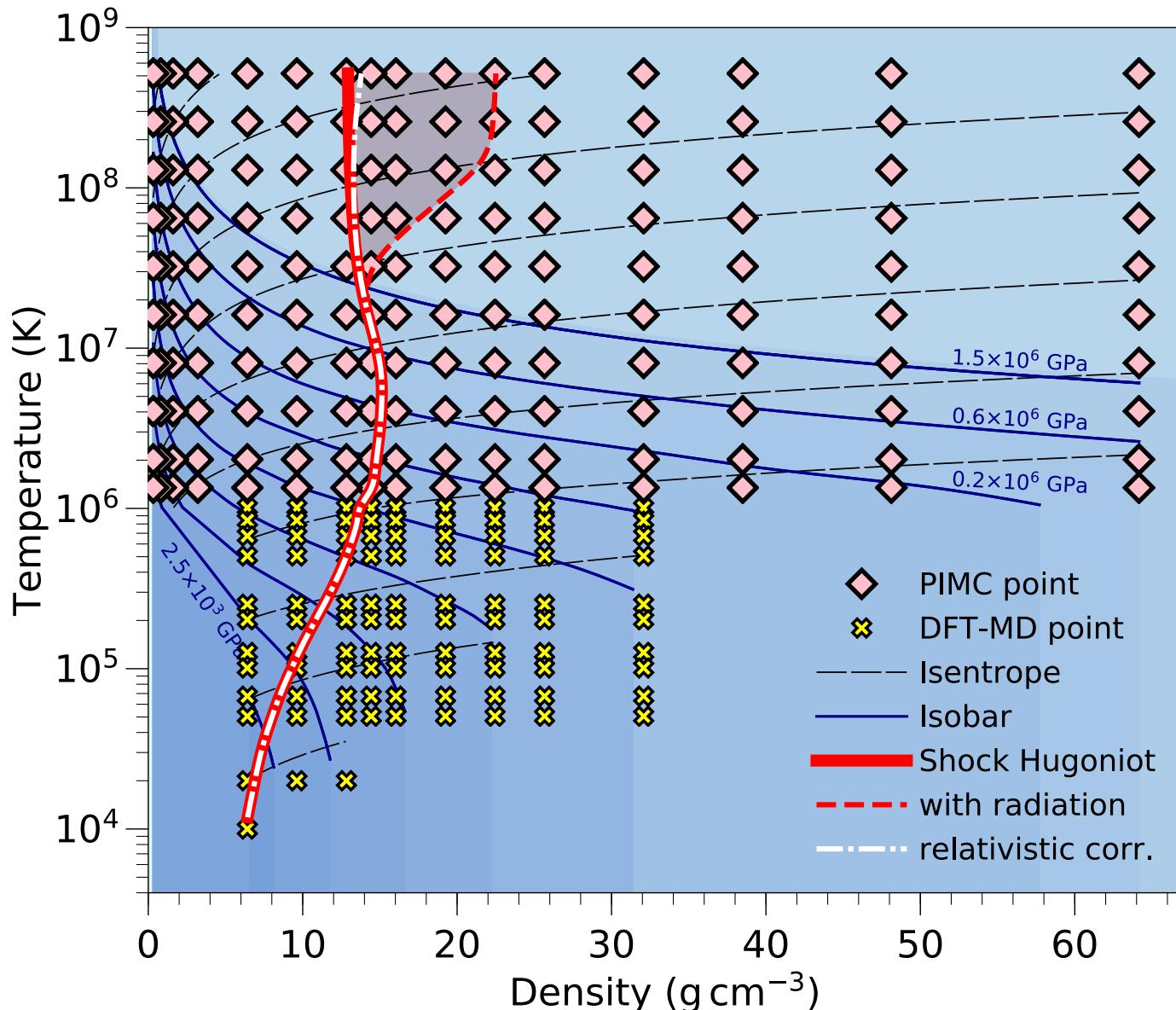
$T=5000\text{K}, r_s=1.6$

- Pressure dissociation, free protons
- Degenerate electron gas
- High number of permutations

# **Silicates:**

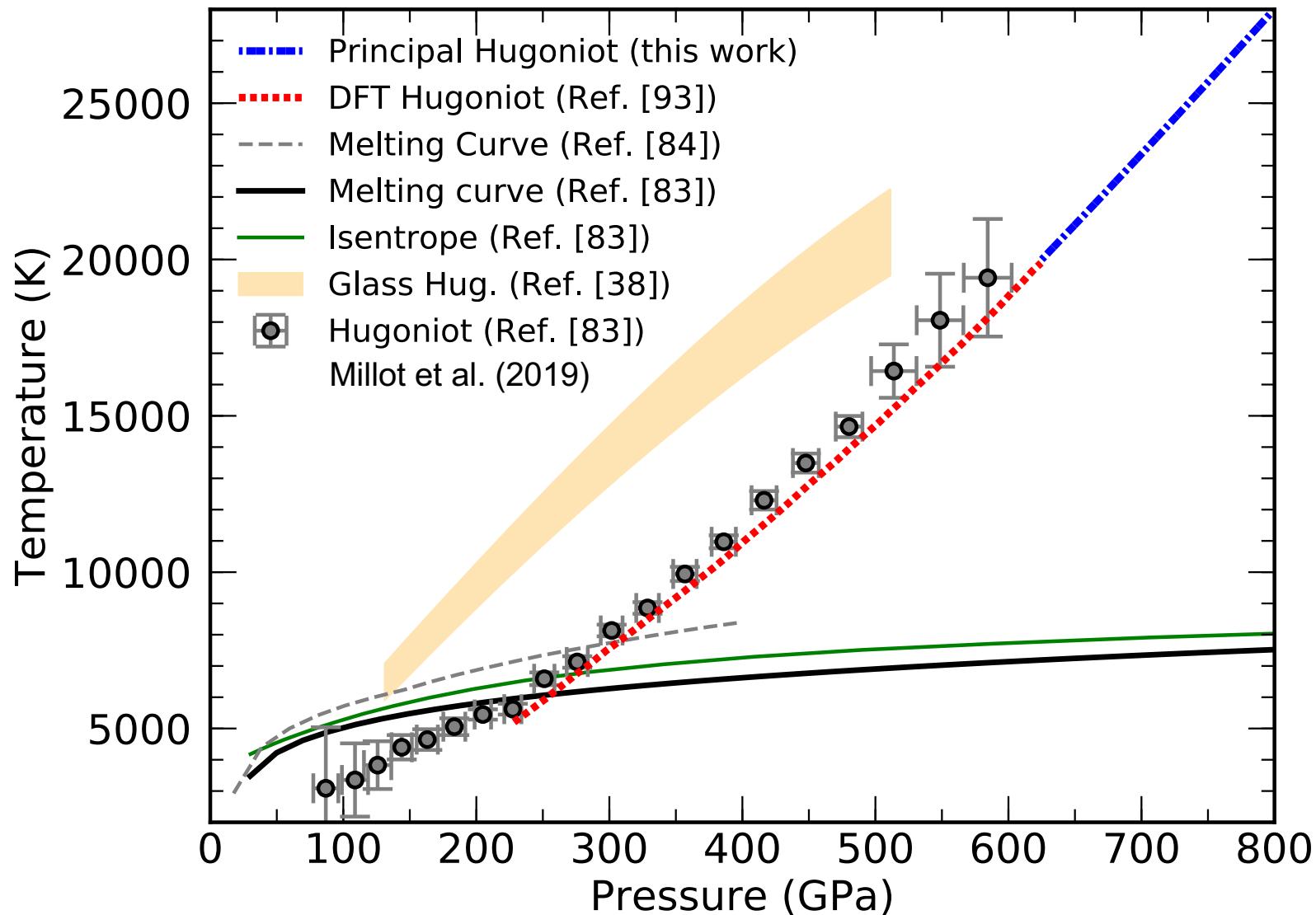
# **MgSiO<sub>3</sub>**

# MgSiO<sub>3</sub> : Principal Hugoniot Curve

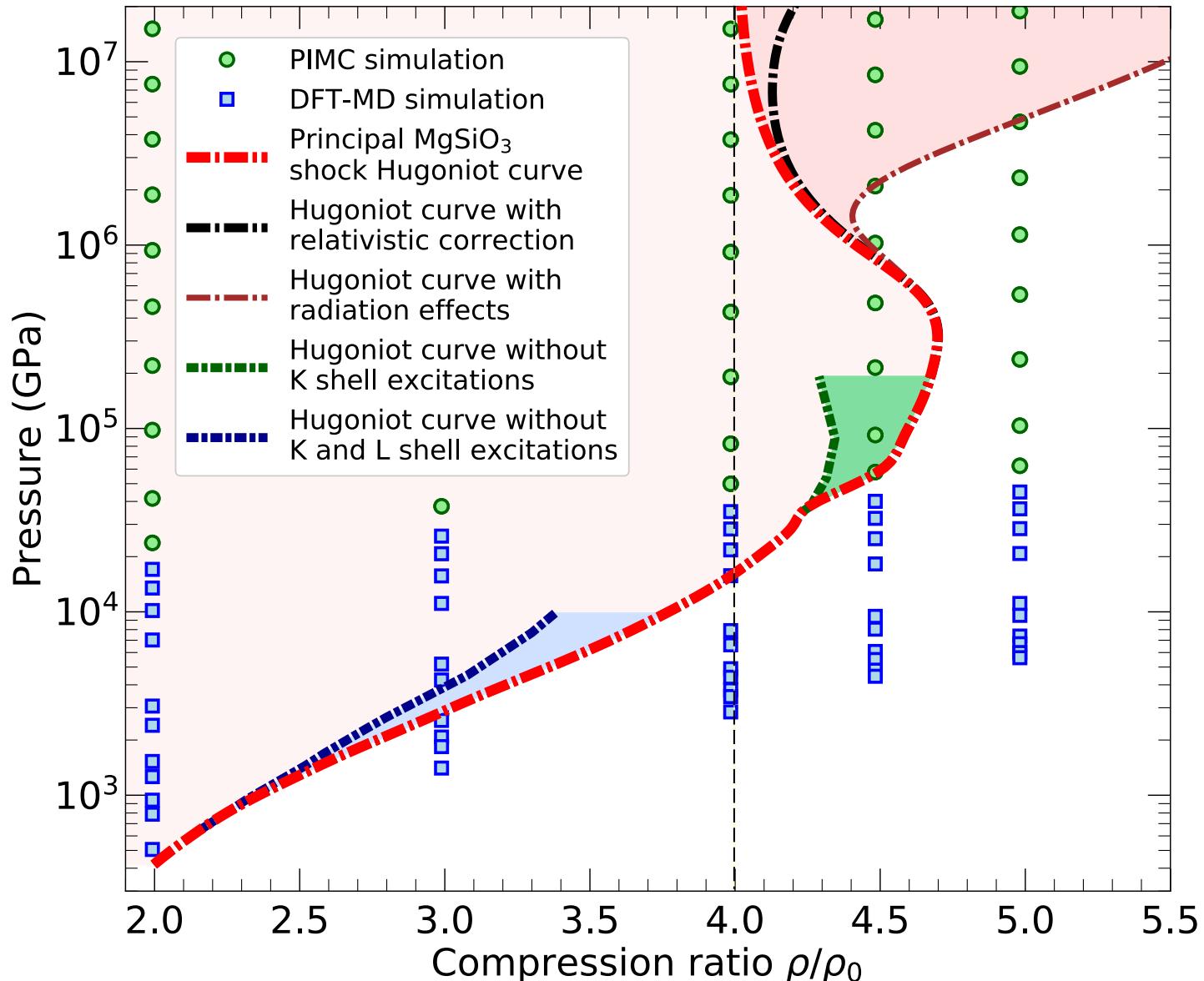


Gonzalez,  
Soubiran,  
Peterson,  
Militzer,  
*Phys. Rev. B*  
101 (2020)  
024107

# MgSiO<sub>3</sub> : Principal Hugoniot Curve



# MgSiO<sub>3</sub> : Principal Hugoniot Curve

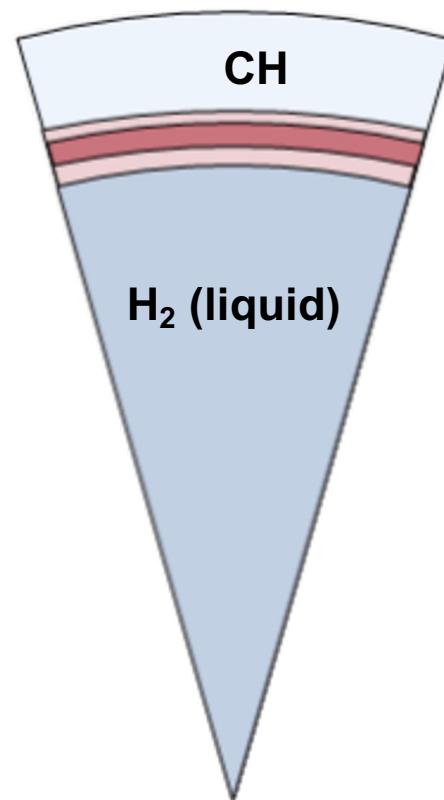
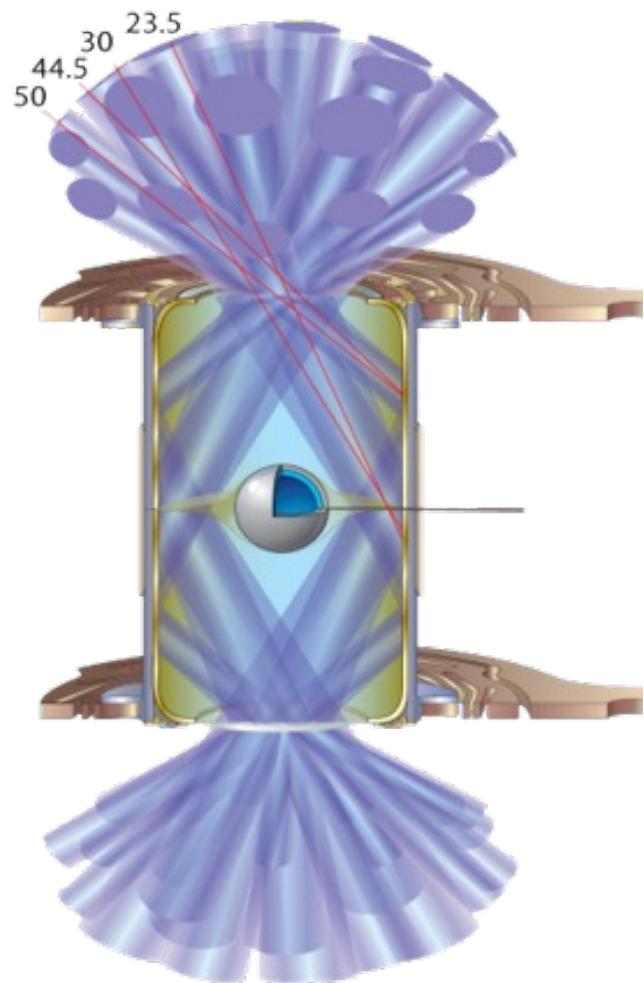


Gonzalez,  
Soubiran,  
Peterson,  
Militzer,  
*Phys. Rev. B*  
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*CH plastics*

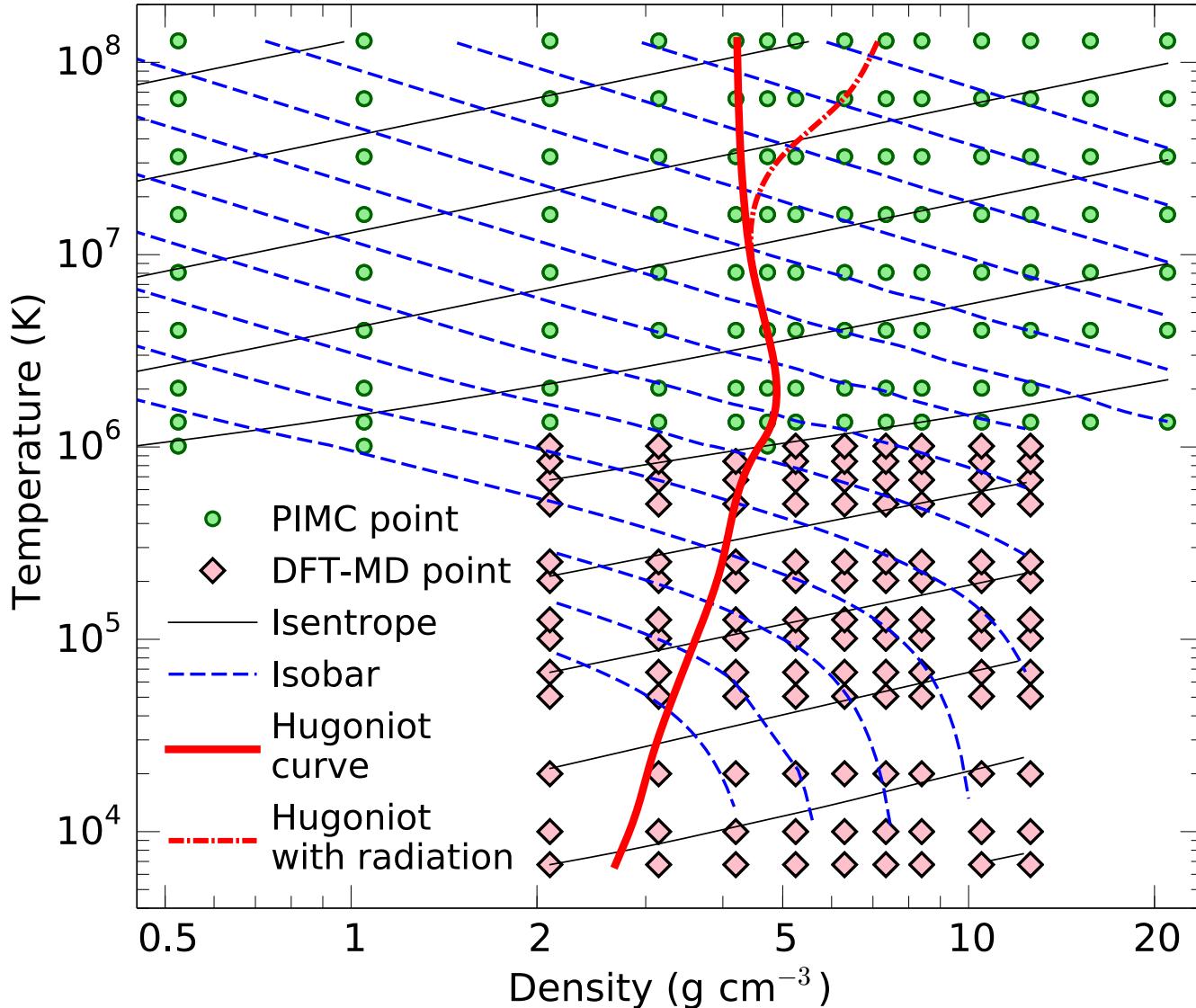
# Inertial confinement fusion experiments with plastic coated spheres of liquid H<sub>2</sub>

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(Graphics: Bachmann et al. LLNL)

# PIMC and DFT-MD simulations performed for $\text{C}_2\text{H}$ , $\text{CH}$ , $\text{C}_2\text{H}_3$ , $\text{CH}_3$ and $\text{CH}_4$ .

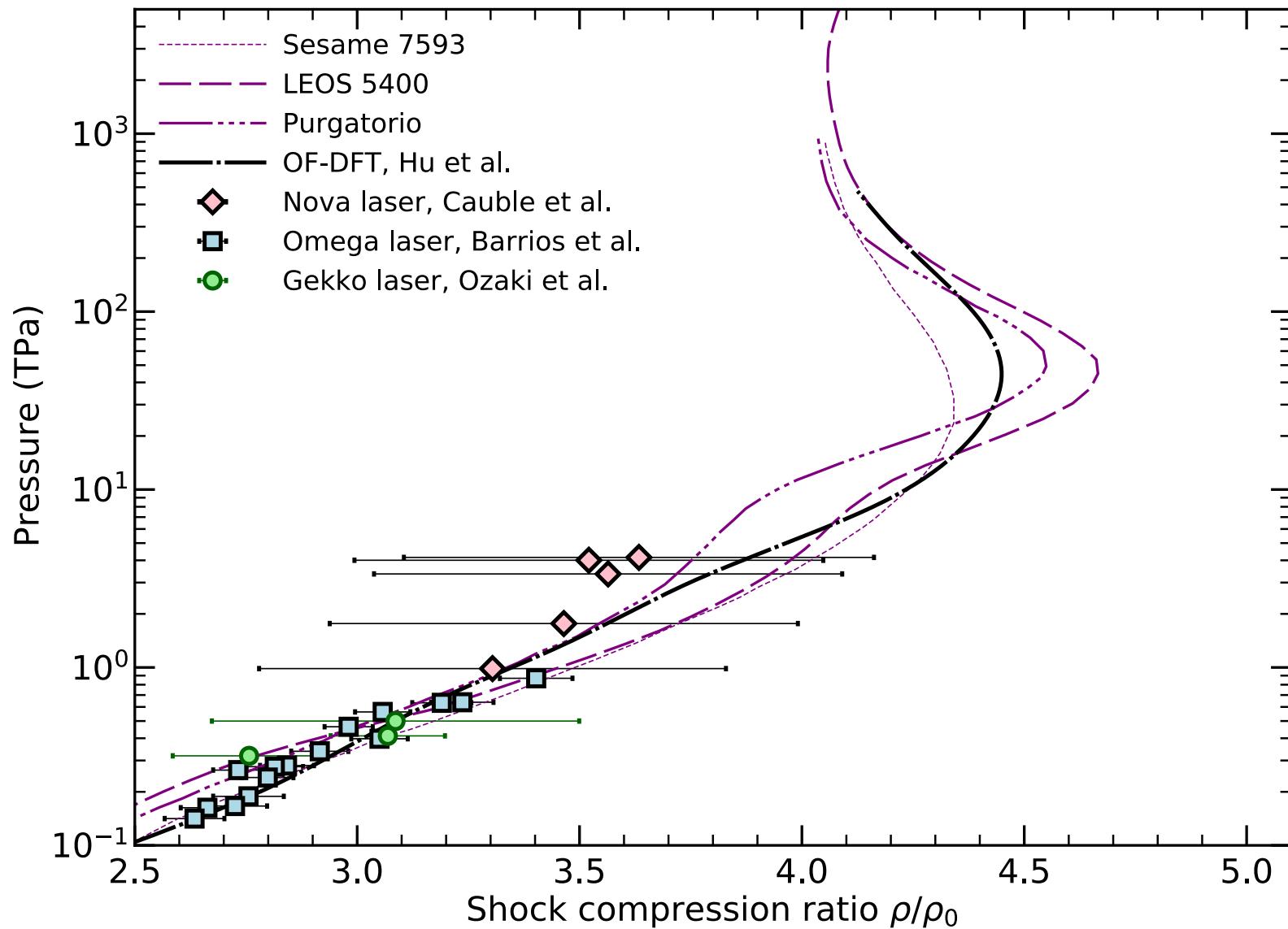


All calculations  
performed on

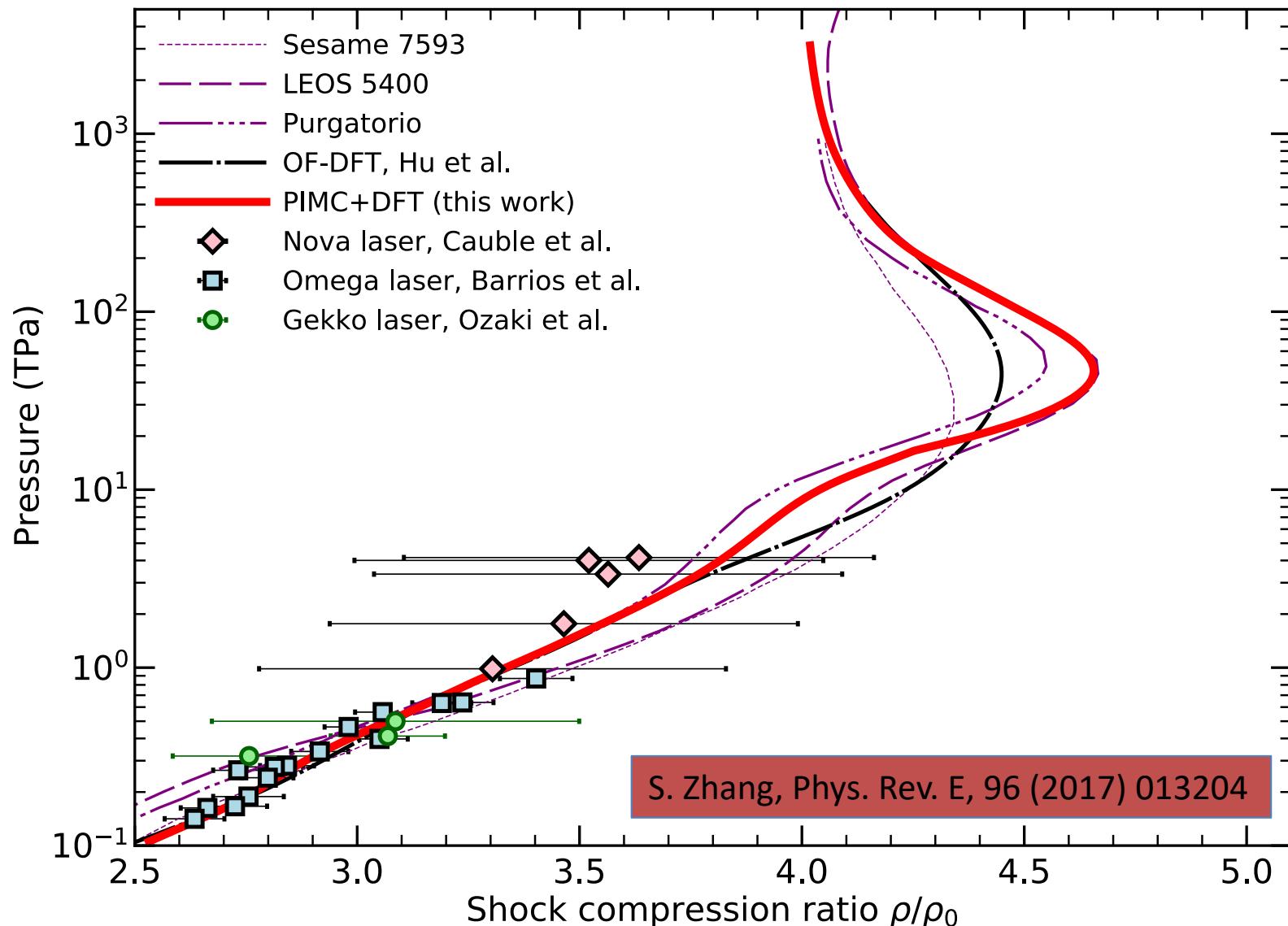
**BLUE WATERS**  
SUSTAINED PETASCALE COMPUTING



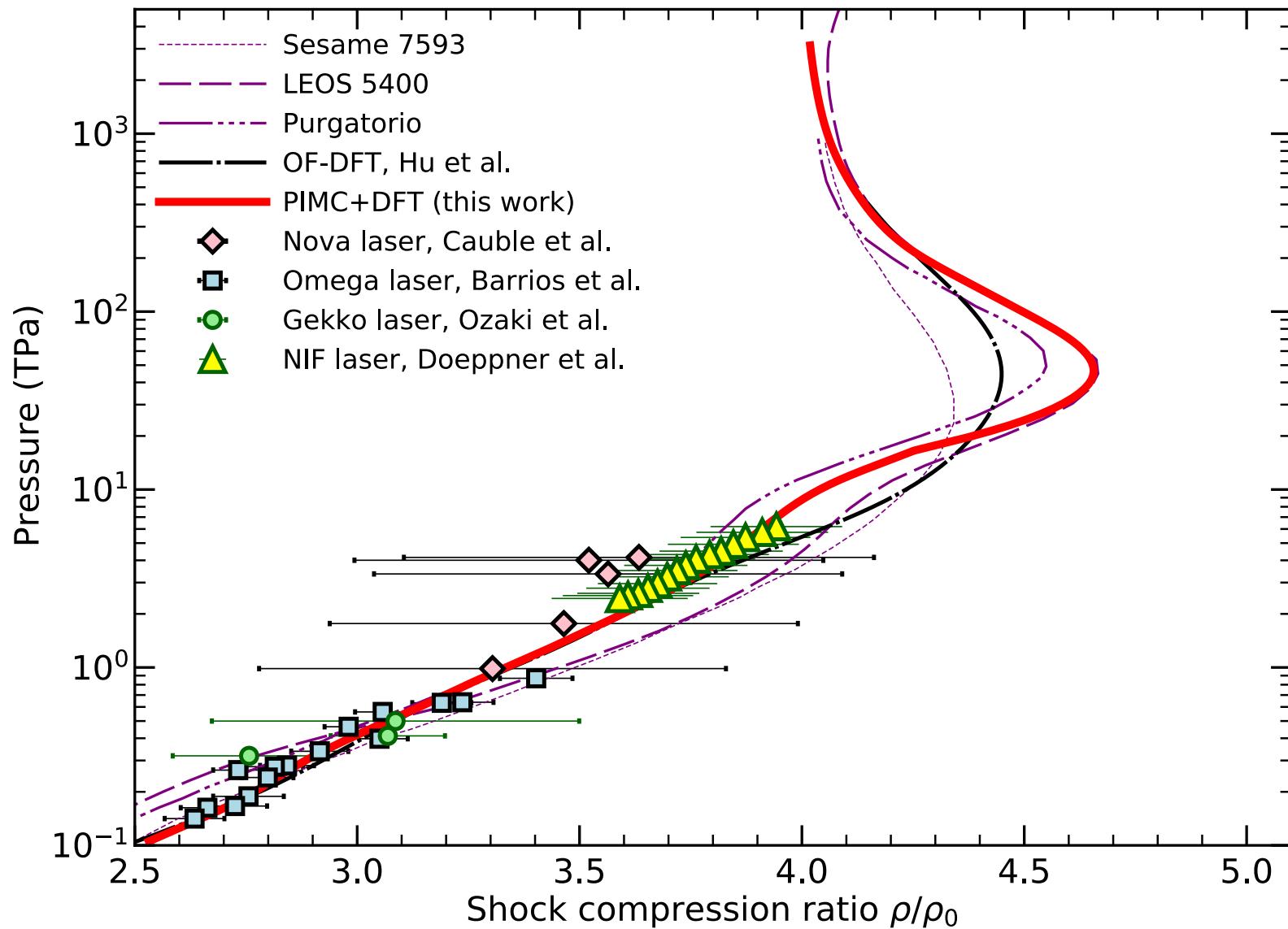
# CH Shock Hugoniot Curves: Comparison of Theory and Experiments



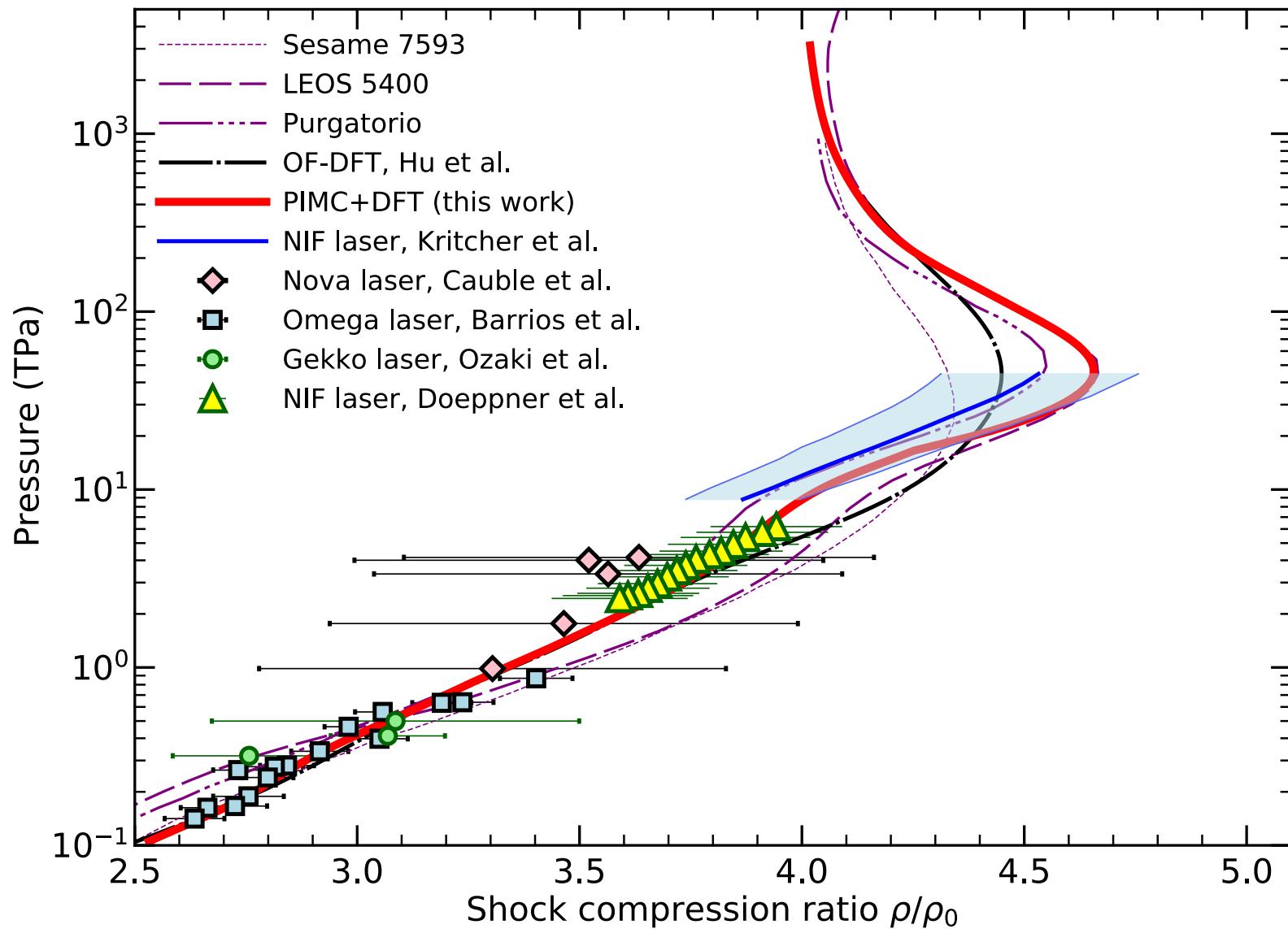
# CH Shock Hugoniot Curves: Comparison of Theory and Experiments



# CH Shock Hugoniot Curves: Comparison of Theory and Experiments

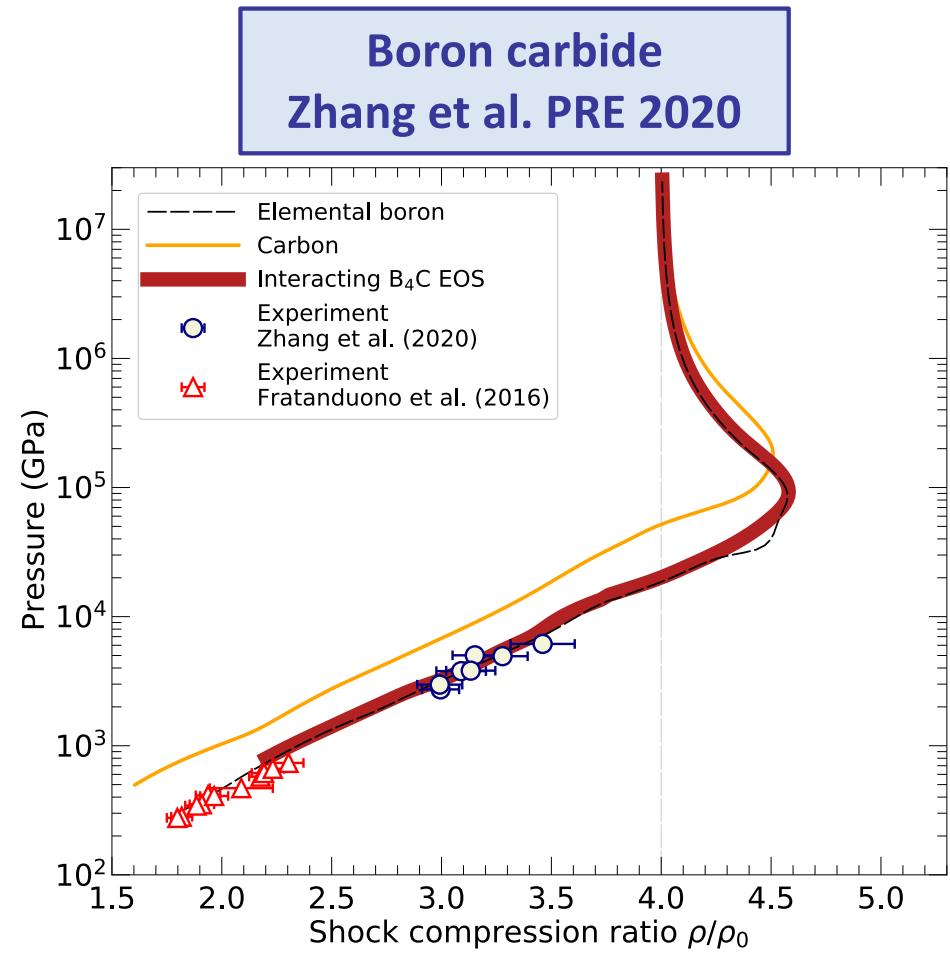
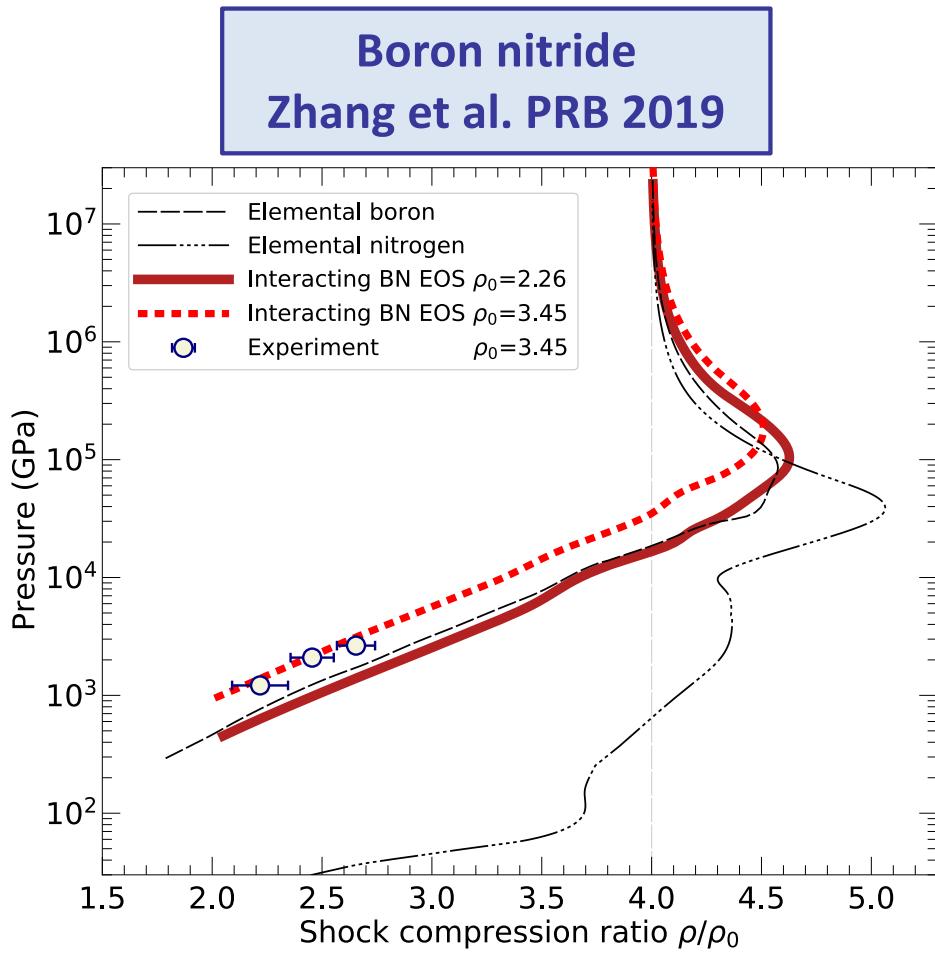


# CH Shock Hugoniot Curves: Comparison of Theory and Experiments



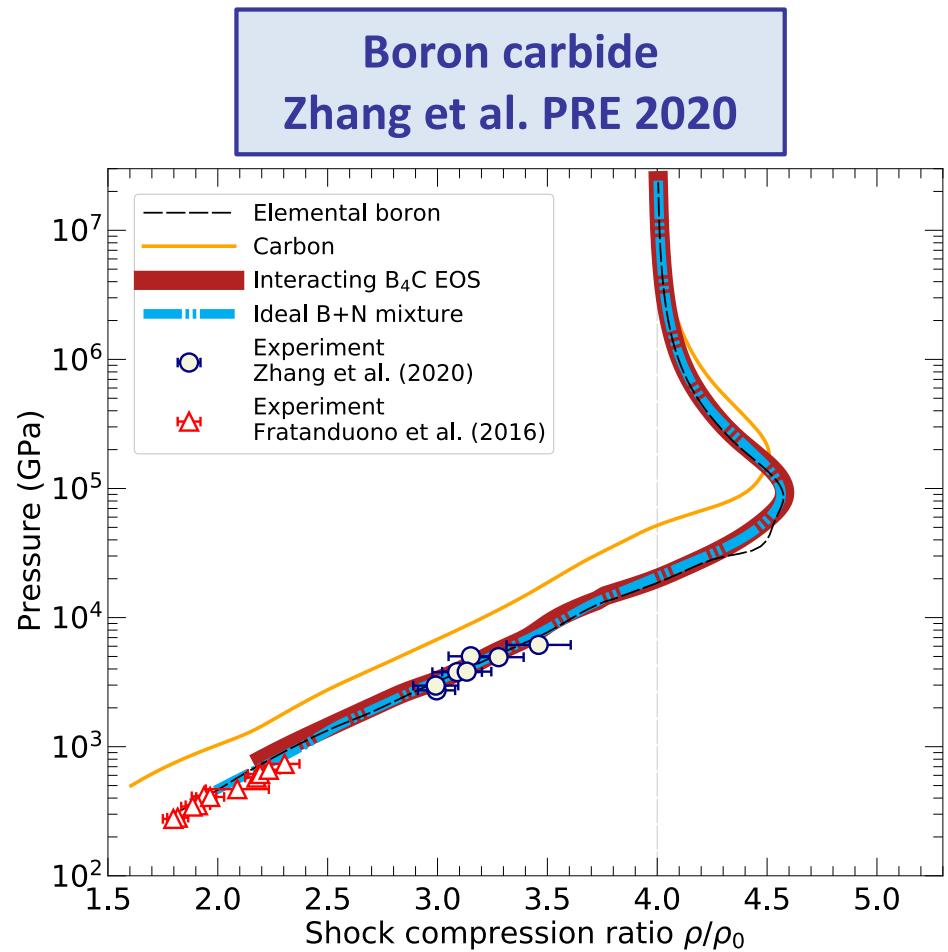
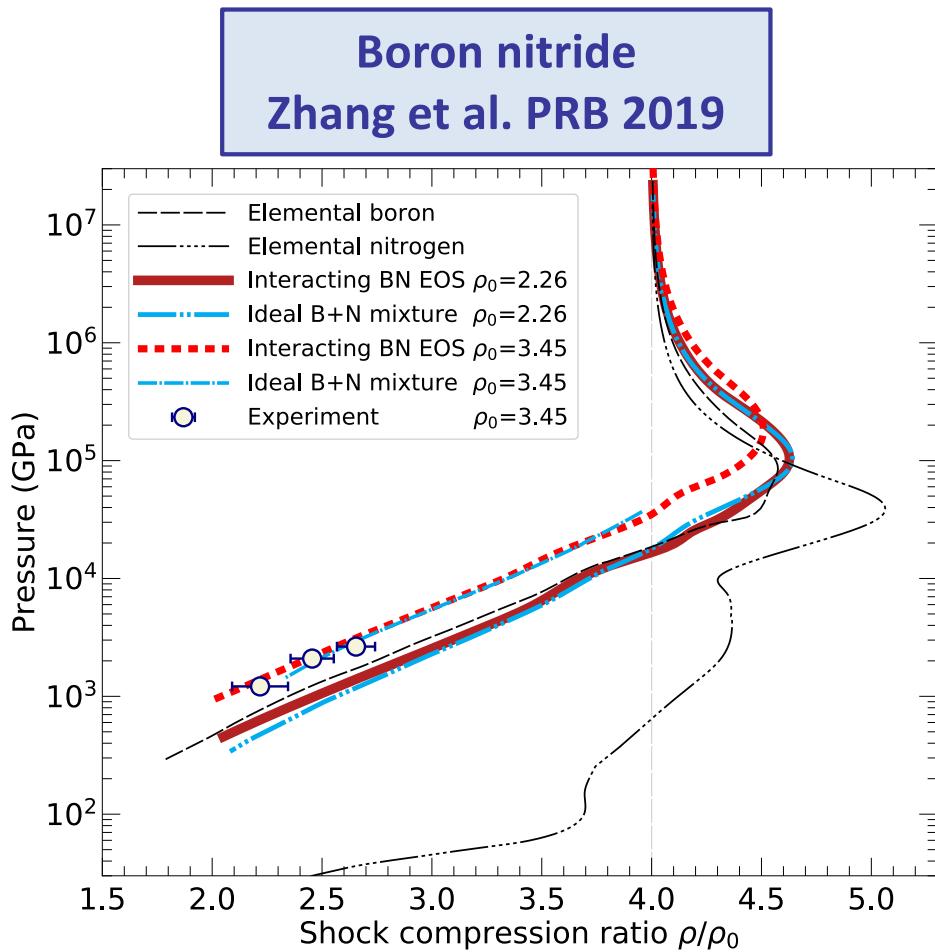
# Hugoniot Curves of BN and $B_4C$

Fully interacting EOS and Linear Mixing agree quite well.



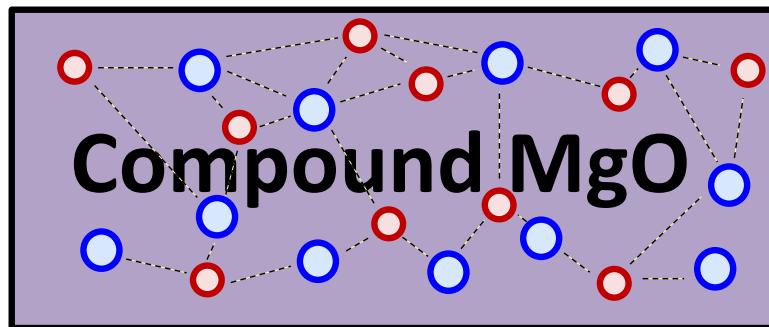
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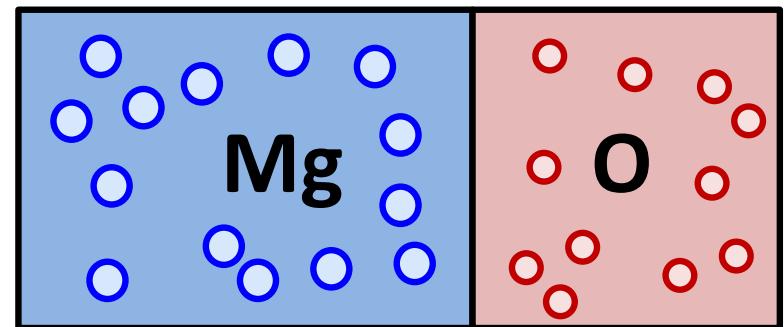


# Linear Mixing at Constant P and T

(Also called additive volume rule)



$\approx$



$$V_{\text{mix}} = N_1 V_1 + N_2 V_2 ,$$

$$m_{\text{mix}} = N_1 m_1 + N_2 m_2 ,$$

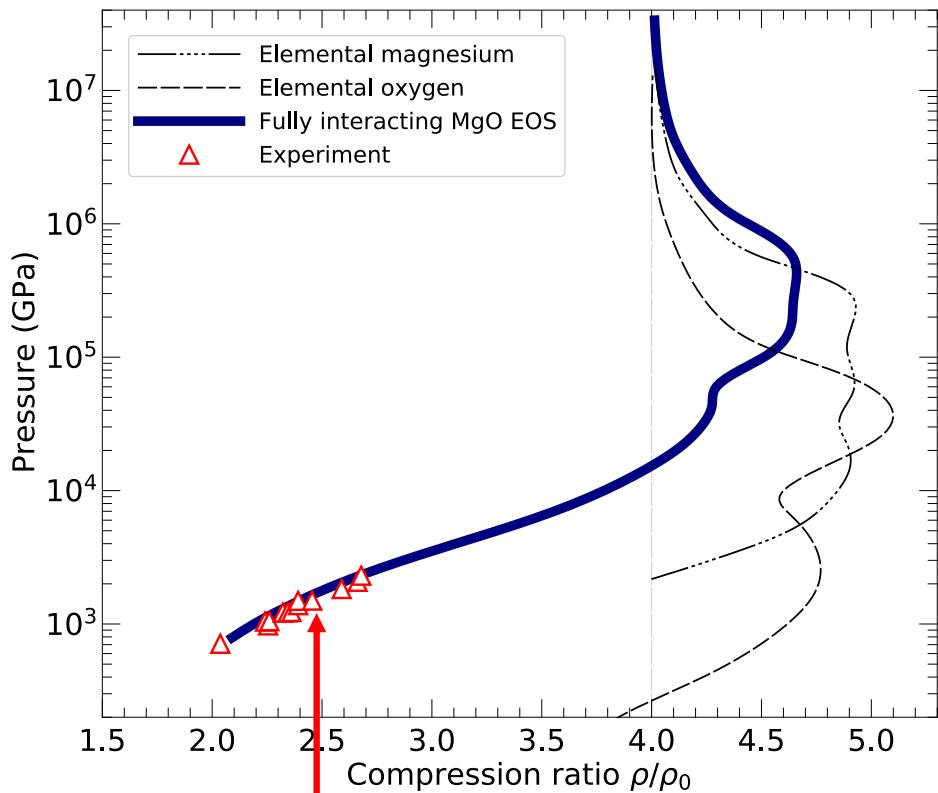
$$E_{\text{mix}} = N_1 E_1 + N_2 E_2 ,$$

$$\rho_{\text{mix}} = m_{\text{mix}} / V_{\text{mix}}$$

# Hugoniot Curves of MgO and MgSiO<sub>3</sub>

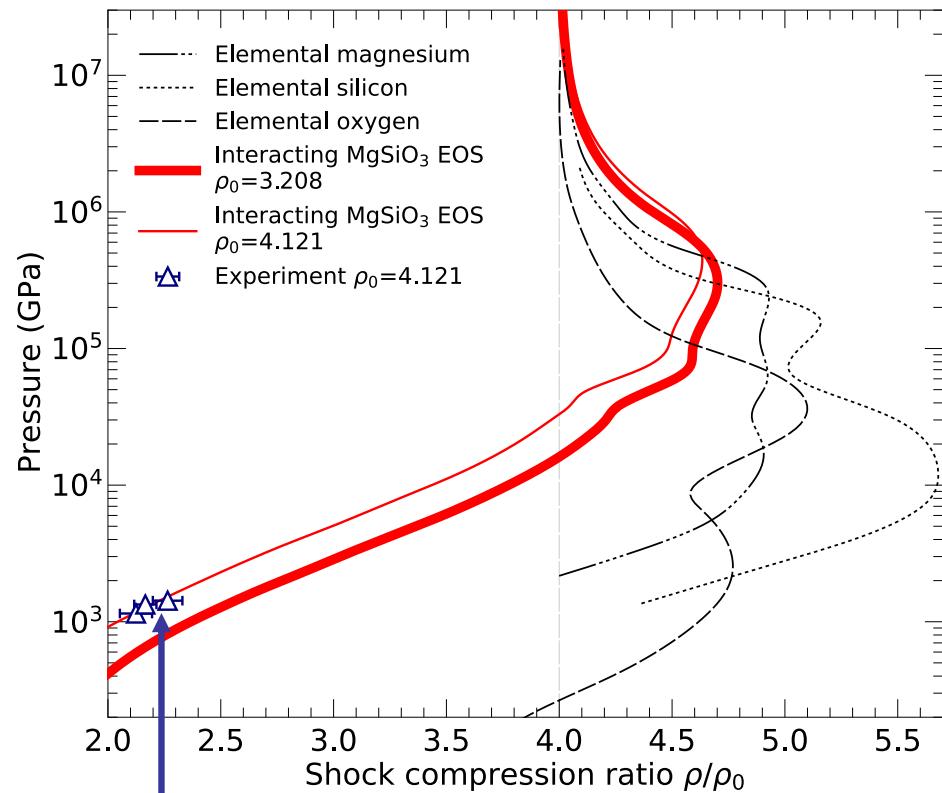
## Results from fully interacting EOS and experiment.

Soubiran et al. JCP 2019



McCoy et al.  
PRB 2019

Gonzalez et al. PRB 2020

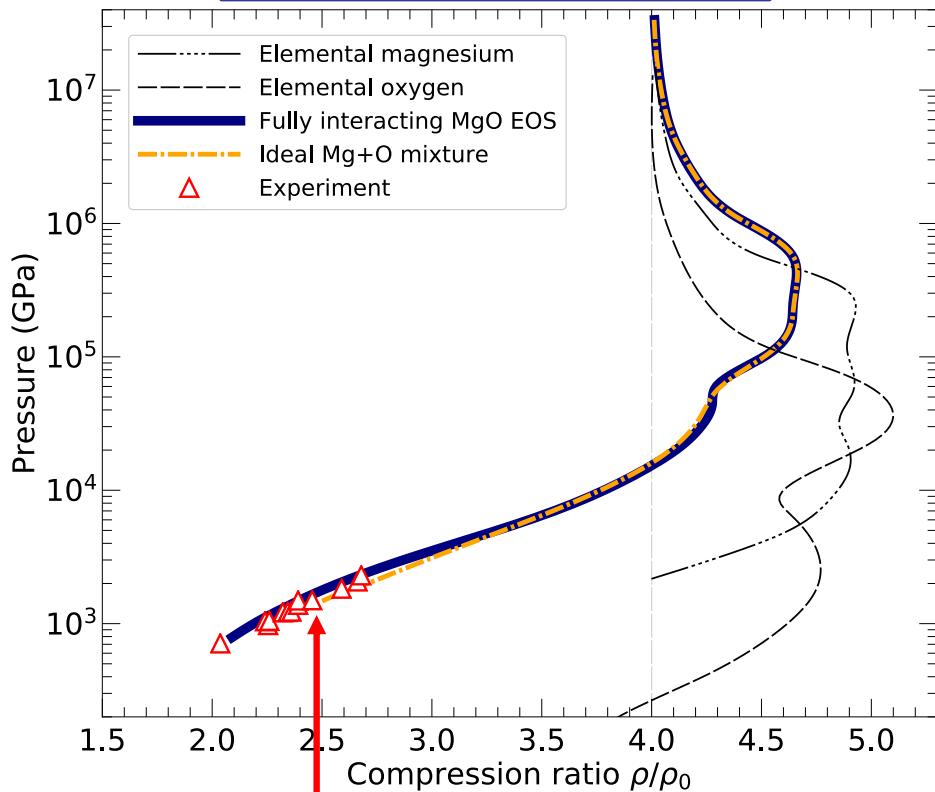


Millot et al.  
GRL 2020

# Hugoniot Curves of MgO and MgSiO<sub>3</sub>

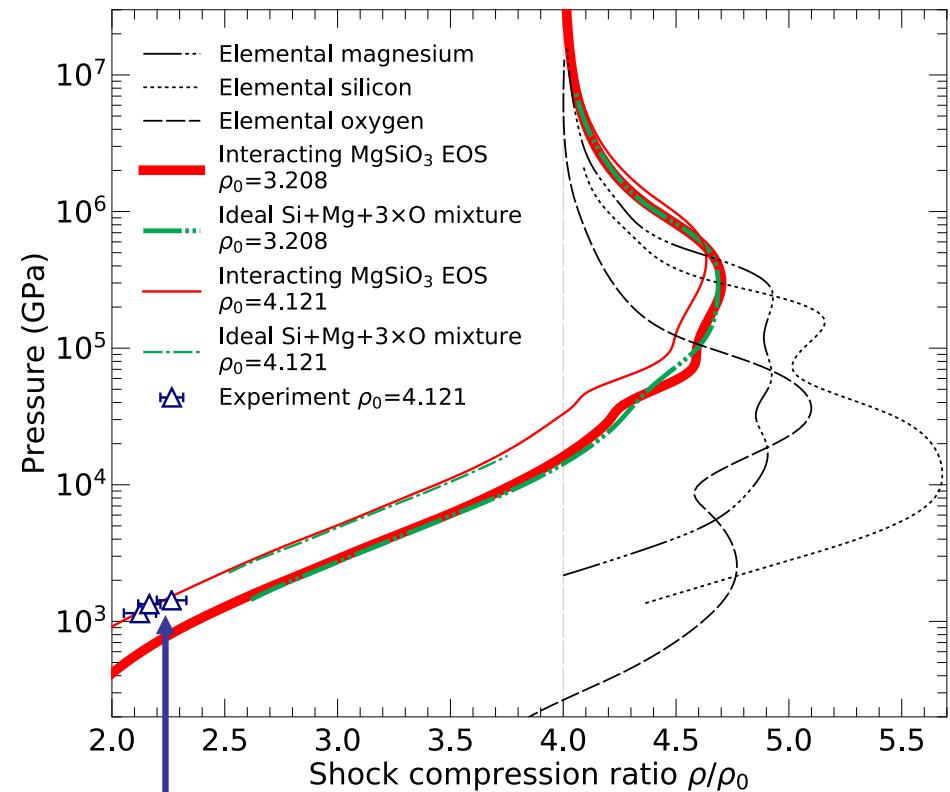
## Fully interacting EOS and Linear Mixing agree quite well.

Soubiran et al. JCP 2019



McCoy et al.  
PRB 2019

Gonzalez et al. PRB 2020



Millot et al.  
GRL 2020

Linear mixing works well for  $T \gtrsim 2 \times 10^5$  K and  $\rho/\rho_0 \gtrsim 3.2$

# Nonideal mixing effects in warm dense matter studied with first-principles computer simulations

Cite as: J. Chem. Phys. 153, 184101 (2020); doi: [10.1063/5.0023232](https://doi.org/10.1063/5.0023232)

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Published Online: 9 November 2020



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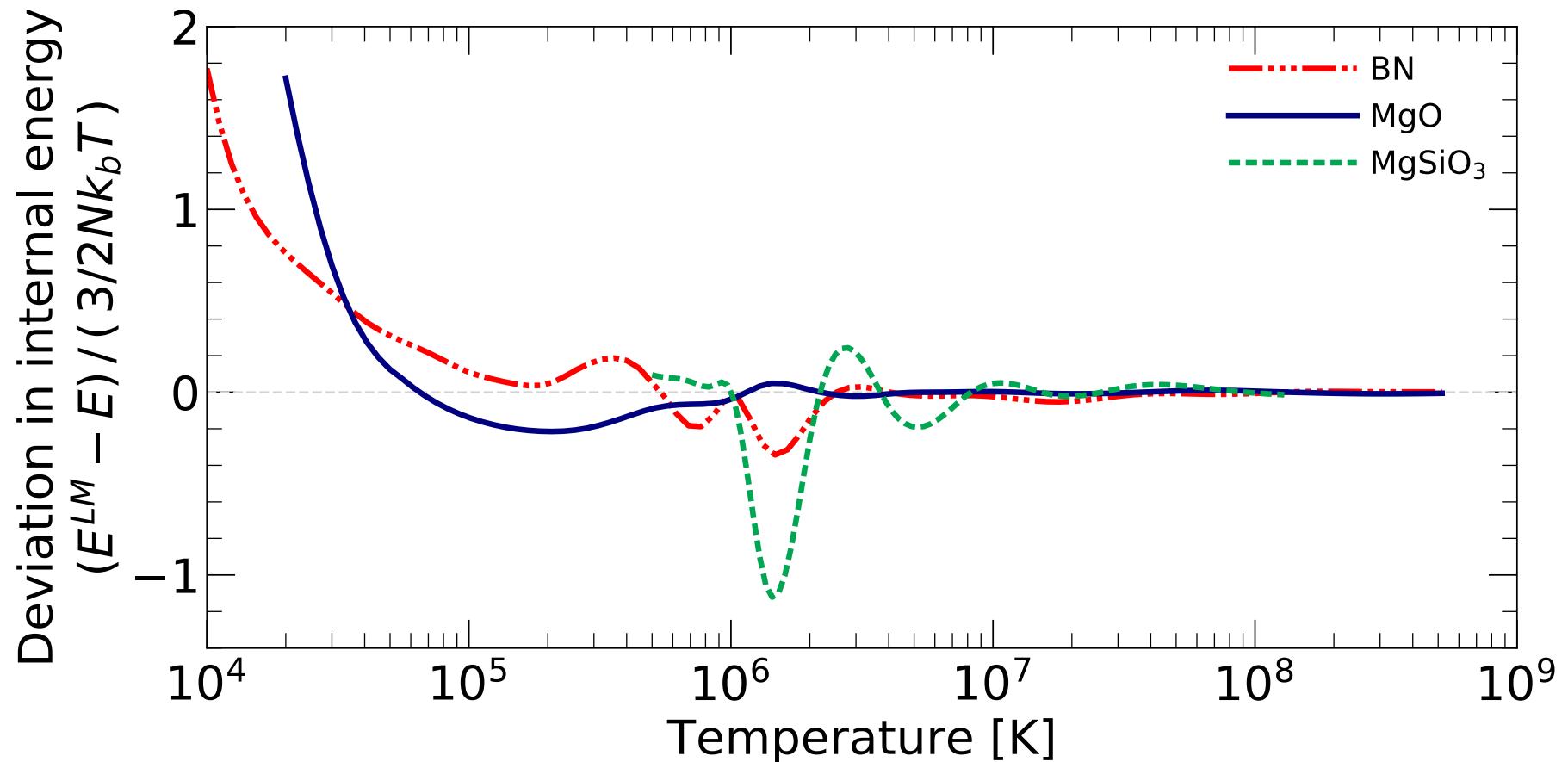


[CrossMark](#)

Burkhard Militzer,<sup>1,2,a)</sup>  Felipe González-Cataldo,<sup>1</sup>  Shuai Zhang,<sup>3</sup>  Heather D. Whitley,<sup>4</sup>  Damian C. Swift,<sup>4</sup> and Marius Millot<sup>4</sup> 

# Nonlinear Mixing Effects in $\text{MgSiO}_3$

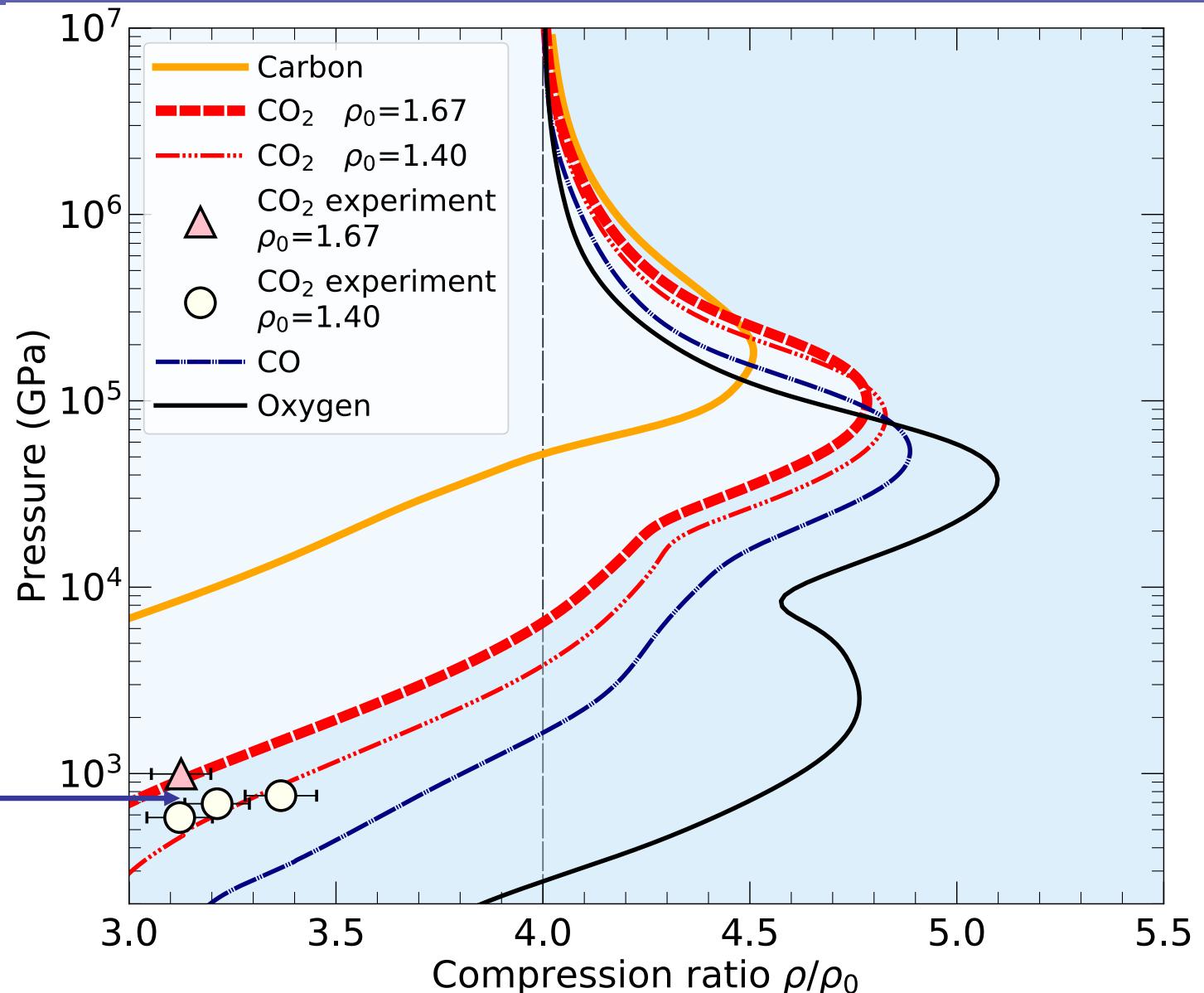
Fully interacting EOS and Linear Mixing agree quite well.



Linear mixing works well for  $T \gtrsim 2 \times 10^5 \text{ K}$  and  $\varrho/\varrho_0 \gtrsim 3.2$

# Hugoniot Curves of CO and CO<sub>2</sub>

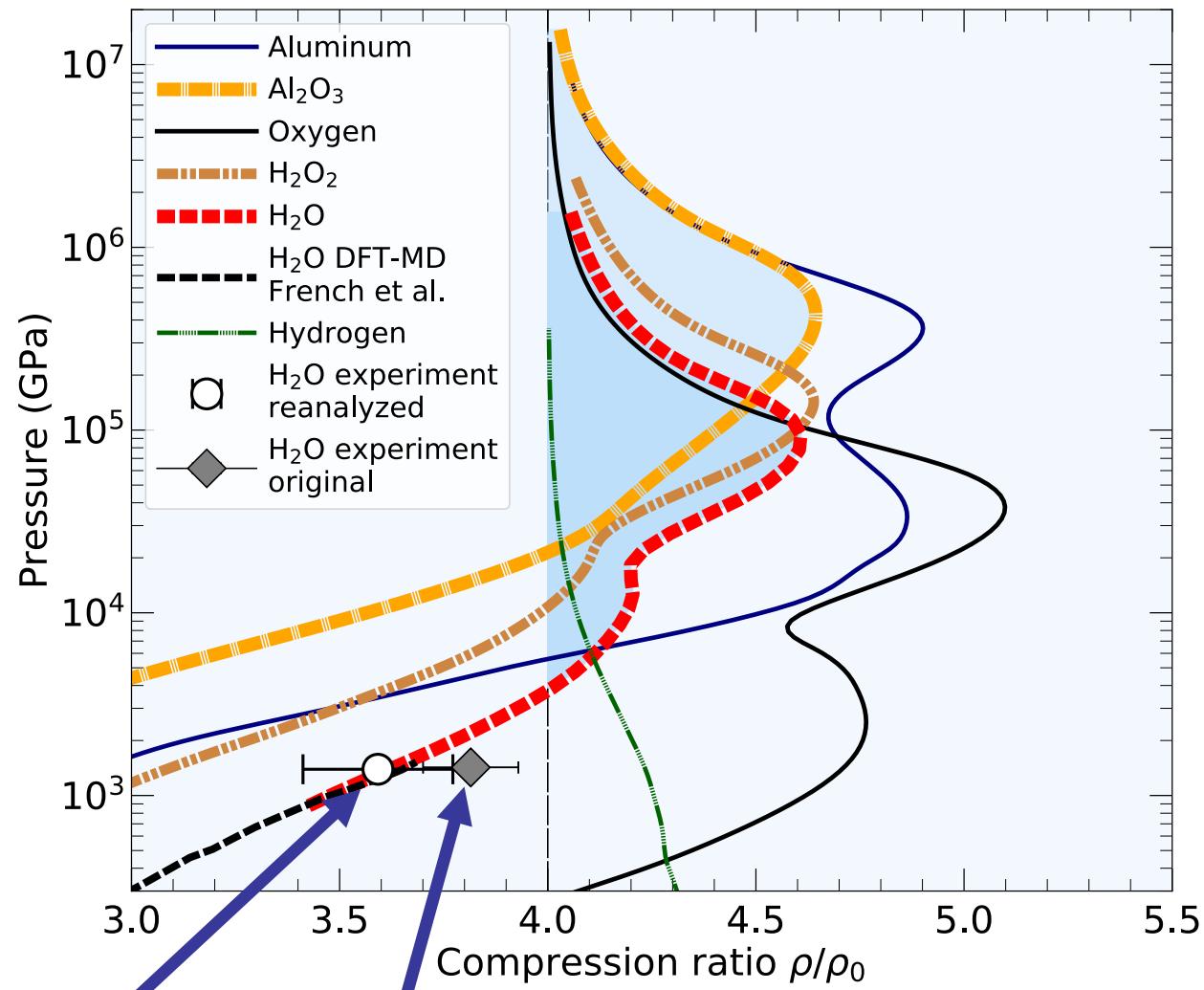
Experimental CO<sub>2</sub> Hugoniot agree with Linear Mixing result



Crandall et al.  
PRL 2020

# Hugoniot Curves of $\text{H}_2\text{O}$ , $\text{H}_2\text{O}_2$ , and $\text{Al}_2\text{O}_3$

## Experimental $\text{H}_2\text{O}$ Hugoniot agree with Linear Mixing result



Reanalyzed by  
Knudson et al.

Experiments by  
Podurets et al. 1972

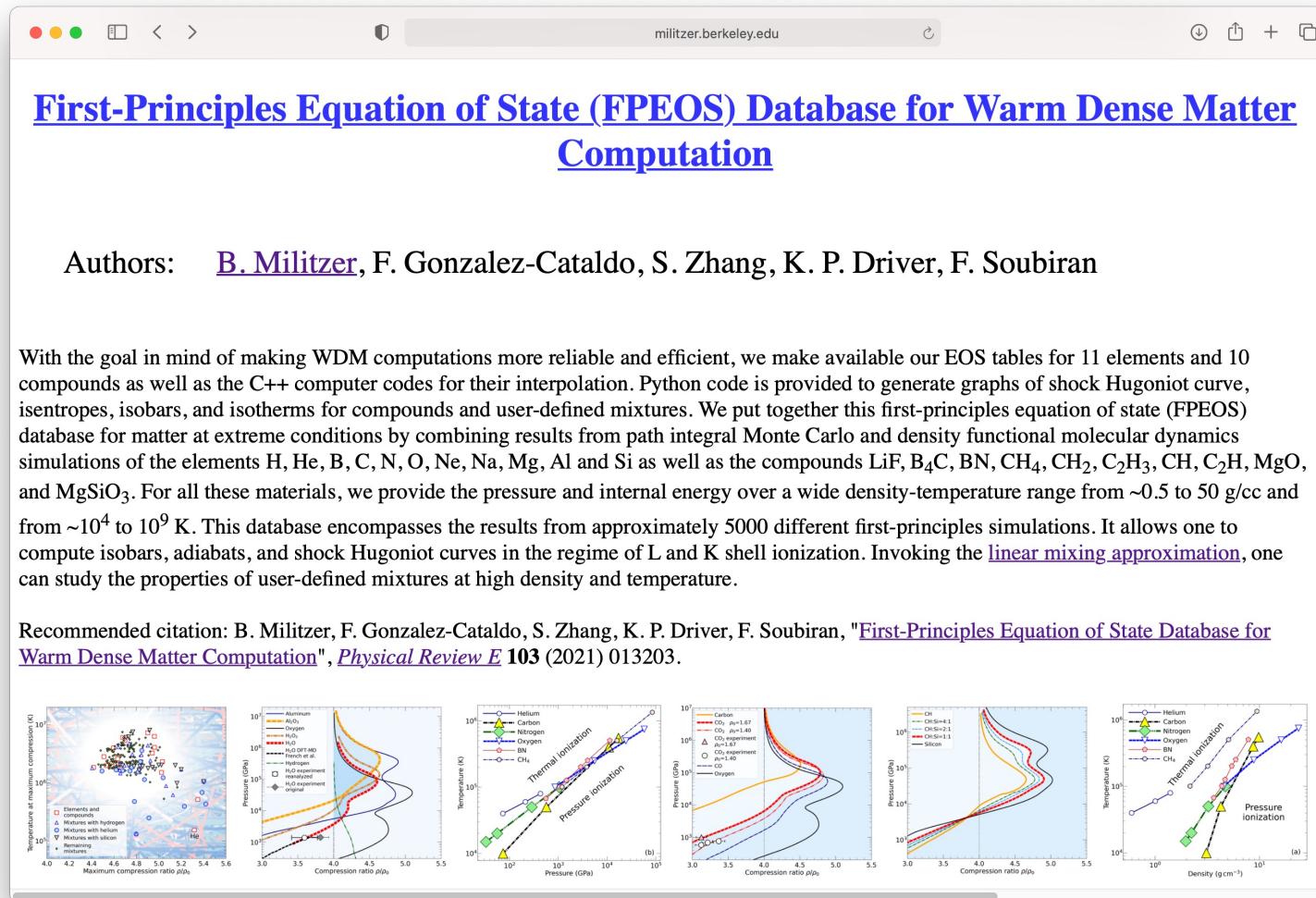


# FPEOS: 11+10 Available Tables



Material	Number of isochores	Minimum density [g cm <sup>-3</sup> ]	Maximum density [g cm <sup>-3</sup> ]	Minimum temperature [K]	Maximum temperature [K]	Number of EOS points	References
Hydrogen	33	0.001	798.913	15625	$6.400 \times 10^7$	401	[69–74]
Helium	9	0.387	10.457	500	$2.048 \times 10^9$	228	[75, 76]
Boron	16	0.247	49.303	2000	$5.174 \times 10^8$	314	[77]
Carbon	9	0.100	25.832	5000	$1.035 \times 10^9$	162	[78, 79]
Nitrogen	17	1.500	13.946	1000	$1.035 \times 10^9$	234	[80]
Oxygen	6	2.486	100.019	10000	$1.035 \times 10^9$	76	[81]
Neon	4	0.895	15.026	1000	$1.035 \times 10^9$	67	[82]
Sodium	9	1.933	11.600	1000	$1.293 \times 10^8$	193	[83, 84]
Magnesium	23	0.431	86.110	20000	$5.174 \times 10^8$	371	[85]
Aluminum	15	0.270	32.383	10000	$2.156 \times 10^8$	240	[86]
Silicon	7	2.329	18.632	50000	$1.293 \times 10^8$	85	[87, 88]
LiF	8	2.082	15.701	10000	$1.035 \times 10^9$	91	[89]
B <sub>4</sub> C	16	0.251	50.174	2000	$5.174 \times 10^8$	291	[90]
BN	16	0.226	45.161	2000	$5.174 \times 10^8$	311	[91]
CH <sub>4</sub>	16	0.072	14.376	6736	$1.293 \times 10^8$	247	[92, 93]
CH <sub>2</sub>	16	0.088	17.598	6736	$1.293 \times 10^8$	248	[92, 93]
C <sub>2</sub> H <sub>3</sub>	16	0.097	19.389	6736	$1.293 \times 10^8$	247	[92, 93]
CH	16	0.105	21.000	6736	$1.293 \times 10^8$	248	[92, 93]
C <sub>2</sub> H	16	0.112	22.430	6736	$1.293 \times 10^8$	245	[92, 93]
MgO	19	0.357	71.397	20000	$5.174 \times 10^8$	286	[94]
MgSiO <sub>3</sub>	16	0.321	64.158	6736	$5.174 \times 10^8$	284	[95, 96]

# First-Principles Equation of State Database online <http://militzer.berkeley.edu/FPEOS>

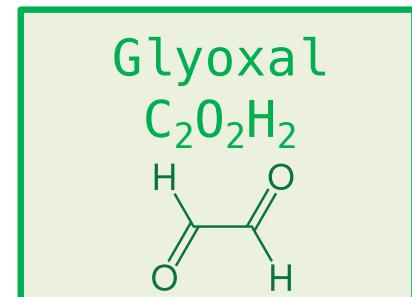
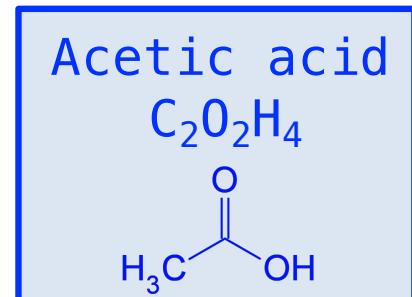
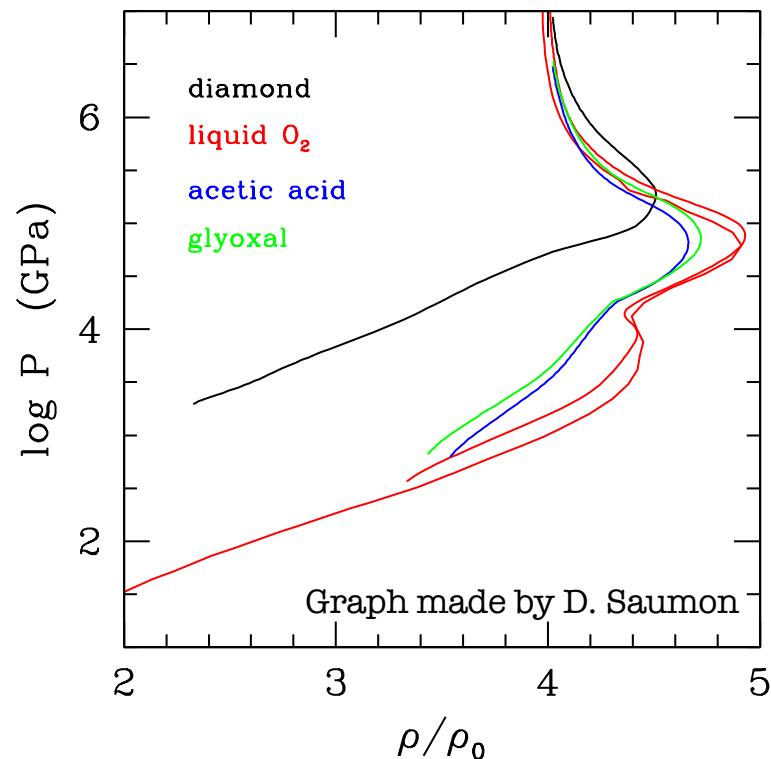
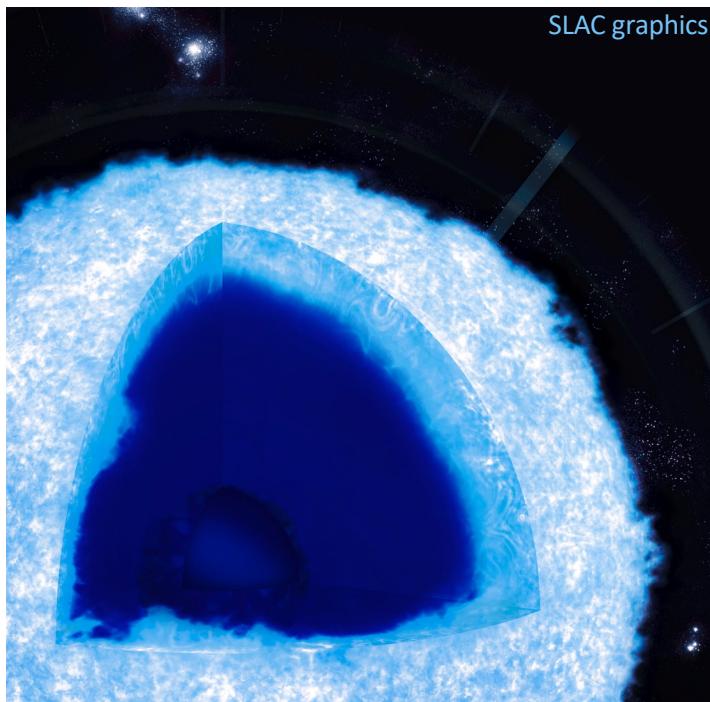


5000 first-principles calculations have been combined into our FPEOS database. So anyone can predict shock Hugoniot curves for a variety of compounds and mixtures. This will make warm dense matter calculations more reliable and efficient.

# NIF Gbar Experiment: Equations of State of C-O Mixtures in White Dwarf Stars

PI: D. Saumon (LANL), Blouin, Glenzer, Swift, Kritch, Doppner, Whitley, Lazicki, Falcone, Militzer

We propose to make EOS measurements along the Hugoniot with the Gbar platform of carbon-oxygen rich materials that resemble conditions in White Dwarf stars.



Glyoxal  $C_2O_2H_2$

comp:~/fpeos> **fpeos** binaryMixture EOS1=6 2.0 EOS2=18 2.0 rho0=1.27 E0=-227.8

Acetic acid  $C_2O_2H_4$

comp:~/fpeos> **fpeos** binaryMixture EOS1=6 2.0 EOS2=16 2.0 rho0=1.049 E0=-229.0

**FPEOS**

**demo**

**The End**