# Triboson production in the SMEFT

EC, Gauthier Durieux, Ken Mimasu, Eleni Vryonidou [240X.XXXXX]



SM@LHC 8/5/24, Roma, Italy

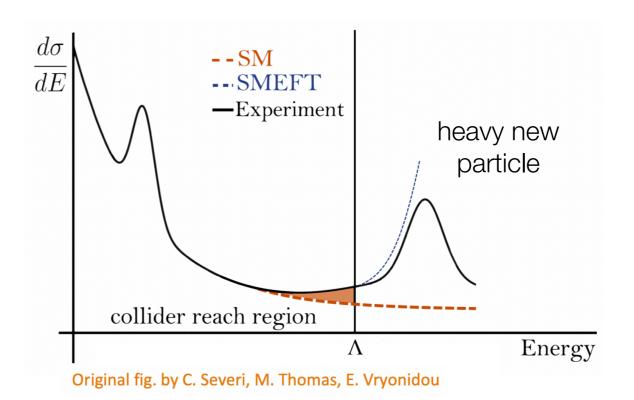


**Eugenia Celada University of Manchester** 





#### The SMEFT



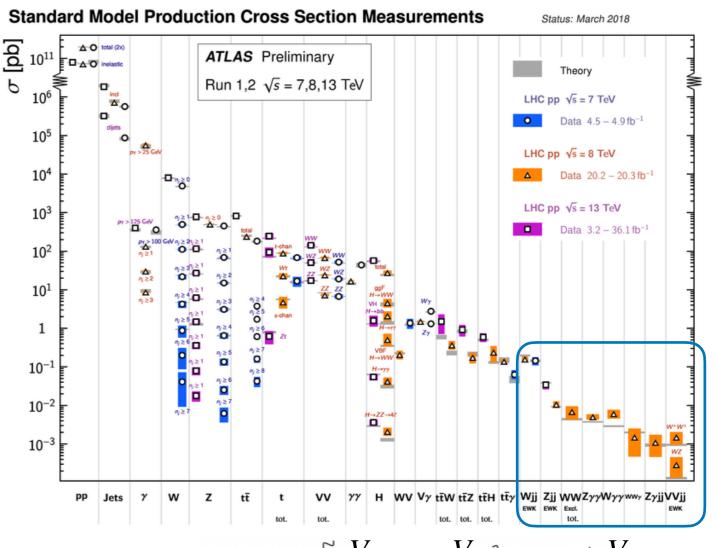
Dimension-6 operators Warsaw basis

$$\mathcal{L}_{ ext{SMEFT}} = \mathcal{L}_{ ext{SM}} + \sum_i rac{c_i^{(6)}}{\Lambda^2} O_i^{(6)} + \mathcal{O}(\Lambda^{-3})$$

$$\sigma \sim |\mathcal{M}_{\mathrm{SM}}|^2 + \frac{1}{\Lambda^2} \left( \sum c^{(6)} \, 2 \mathrm{Re} \left[ \mathcal{M}_{\mathrm{SM}}^* \mathcal{M}_{\mathrm{EFT}}^{(6)} \right] \right) + \frac{1}{\Lambda^4} \left( \sum c^{(6)} \mathcal{M}_{\mathrm{EFT}}^{(6)} \right)^2$$

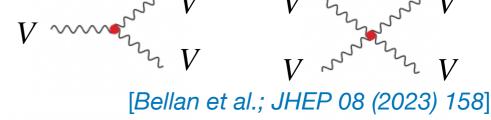
# Triboson production at the LHC

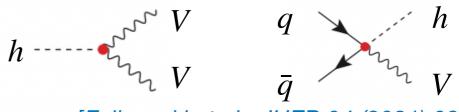
 Triboson have small cross sections, only accessible with LHC run 2 (total rates, mainly fully leptonic)



#### Why triboson?

- Tree-level access to TGC and QGC
- Interplay with the Higgs sector





[Falkowski et al.; JHEP 04 (2021) 023]

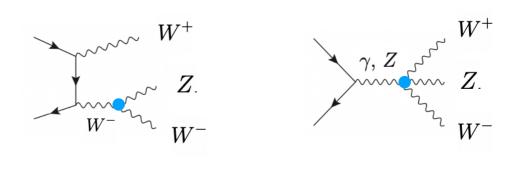
Operator	Definition		
bosonic			
$\mathcal{O}_{\phi D}$	$(\phi^\dagger D^\mu \phi)^\dagger (\phi^\dagger D_\mu \phi)$		
$\mathcal{O}_{\phi WB}$	$(\phi^\dagger  au_I \phi) B^{\mu  u} W^I_{\mu  u}$		
$\mathcal{O}_{WWW}$	$\epsilon_{IJK}W^I_{\mu\nu}W^{J,\nu\rho}W^{K,\mu}_{\rho}$		
two-fermion			
$\mathcal{O}_{\phi q}^{(1)}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu \phi)(\bar{q} \gamma^\mu q)$		
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four-fermion			
$\mathcal{O}_{\ell\ell}$	$(ar{\ell}\gamma_{\mu}\ell)(ar{\ell}\gamma^{\mu}\ell)$		

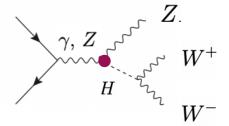
Subset of 11 EW&Higgs operators

• flavour universality,  $U(3)^5$ 

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$$pp \to W^+W^-Z$$





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$$pp \rightarrow W^{+}W^{-}Z$$

$$W^{+}$$

$$Z$$

$$W^{-}$$

$$W^{-}$$

$$W^{-}$$

$$W^{+}$$

$$W^{-}$$

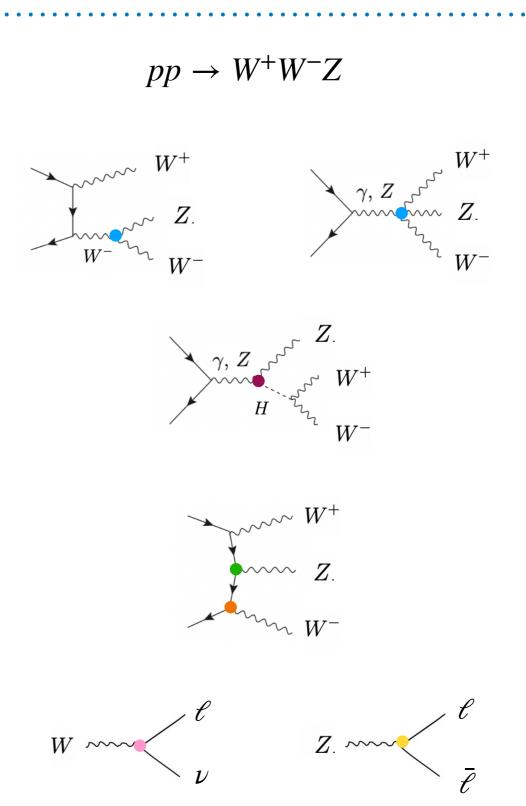
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$$W^{+}$$

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1	four-fermion	
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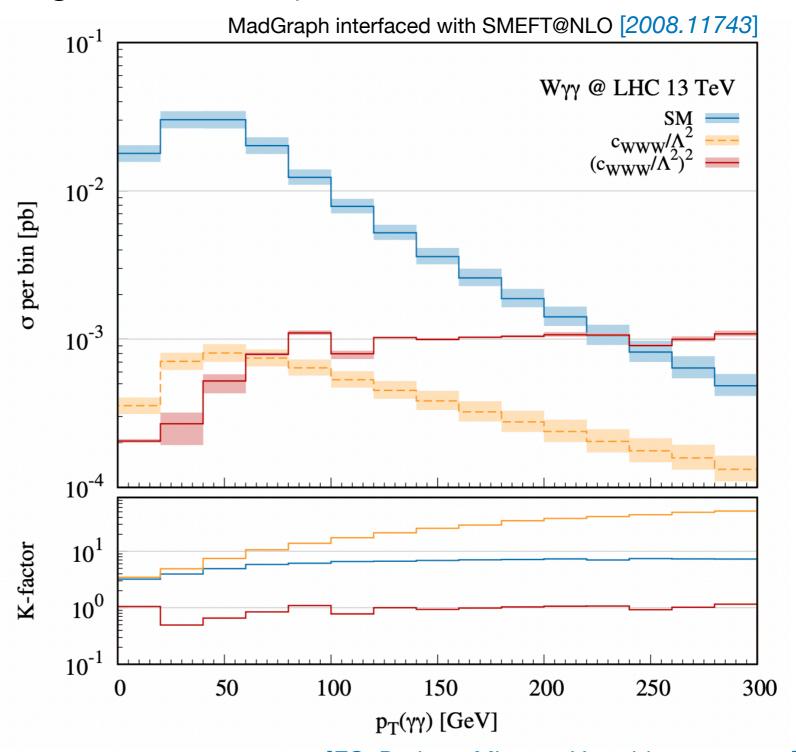


# Going NLO

NLO QCD corrections are large in triboson processes [Degrande et al.; 2008.11743]

$W\gamma\gamma$				
$\sigma({ m fb})$	K-factor			
$\sigma_{ m SM}$	4.84			
$\sigma_{\phi D}$	4.86			
$\sigma_{\phi D,\phi D}$	4.86			
$\sigma_{\phi WB}$	4.70			
$\sigma_{\phi WB,\phi WB}$	1.47			
$\sigma_{WWW}$	12.24			
$\sigma_{WWW,WWW}$	0.79			
$\sigma_{\phi\ell^{(3)}}$	4.85			
$\sigma_{\phi\ell^{(3)},\phi\ell^{(3)}}$	4.85			
$\sigma_{\phi q^{(3)}}$	4.80			
$\sigma_{\phi q^{(3)},\phi q^{(3)}}$	4.80			
$\sigma_{\ell\ell}$	4.82			
$\sigma_{\ell\ell,\ell\ell}$	4.82			

$$K = \frac{\sigma_{\rm NLO}}{\sigma_{\rm LO}}$$



#### Operators and observables

EWPOs and 
$$\alpha_{\rm EW} \sqrt{s} = m_{\rm Z}$$

$$\Gamma_Z, \sigma_{\mathrm{had}}^0, R_\ell^0, A_{FB}^\ell, A_\ell(\mathrm{SLD}), R_b^0, R_c^0, A_{FB}^b, A_{FB}^c, A_b, A_c$$
 [LEP; 0509008]  $\alpha_{EW}(m_Z)$ 

**LEP** 
$$WW \sqrt{s} = 183 - 209 \text{ GeV}$$

**LEP** 
$$WW\sqrt{s} = 183 - 209$$
 GeV  $\sigma(WW \to \ell\nu\ell\nu, qqqq)$   $\frac{d\sigma}{d\cos(\theta)}(WW \to \ell\nu qq)$ 

[LEP: 1302.3415]

LHC VV 
$$\sqrt{s} = 13 \text{ TeV}$$

$$\frac{d\sigma}{dm_{e\mu}}(WW \to e\nu\mu\nu)$$

[ATLAS; 1905.04242]

$$\frac{d\sigma}{dp_T^Z}(WZ\to\ell\nu\ell\nu)$$

[ATLAS; 1902.05759]

$$\frac{d\sigma}{d\Delta\phi_{ii}}(Zjj\to\ell\ell jj)$$

[ATLAS: 2006.15458]

LHC VVV 
$$\sqrt{s} = 13 \text{ TeV}$$

 $\sigma(WWW, WWZ, WZZ, WZ\gamma, WW\gamma, W\gamma\gamma)$ 

[ATLAS; 2201.13045, 2305.16994, 2308.03041] [CMS: 2006.11191, 2310.05164, 2105.12780]

# Operators and observables

Operator	Definition	EWPOs	LEP WW	LHC VV	$VVV, VV\gamma, V\gamma\gamma$
	bosonic				
$\mathcal{O}_{\phi D}$	$(\phi^\dagger D^\mu \phi)^\dagger (\phi^\dagger D_\mu \phi)$	✓	✓	✓	✓
$\mathcal{O}_{\phi WB}$	$(\phi^\dagger  au_I \phi) B^{\mu  u} W^I_{\mu  u}$	✓	✓	✓	✓
$\mathcal{O}_{WWW}$	$\epsilon_{IJK}W^I_{\mu\nu}W^{J,\nu\rho}W^{K,\mu}_\rho$		✓	✓	✓
	two-fermion				
$\mathcal{O}_{\phi q}^{(1)}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu \phi)(\bar{q} \gamma^\mu q)$	✓		✓	✓
$\mathcal{O}_{\phi q}^{(3)}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu  au_I \phi)(ar{q} \gamma^\mu  au^I q)$	✓	✓	✓	✓
$\mathcal{O}_{\phi u}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu \phi)(\bar{u}\gamma^\mu u)$	✓		✓	✓
$\mathcal{O}_{\phi d}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu \phi)(\bar{d}\gamma^\mu d)$	✓		✓	✓
$\mathcal{O}_{\phi\ell}^{(1)}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu \phi)(ar{\ell} \gamma^\mu \ell)$	✓	✓	✓	$\checkmark$
$\mathcal{O}_{\phi\ell}^{(3)}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu  au_I \phi)(ar{\ell} \gamma^\mu  au^I \ell)$	✓	$\checkmark$	✓	$\checkmark$
$\mathcal{O}_{\phi e}$	$i(\phi^\dagger \overleftrightarrow{D}_\mu \phi)(\bar{e} \gamma^\mu e)$	✓	✓	✓	✓
1	four-fermion				
$\mathcal{O}_{\ell\ell}$	$(ar{\ell}\gamma_{\mu}\ell)(ar{\ell}\gamma^{\mu}\ell)$	✓	✓	✓	✓

#### Fit results

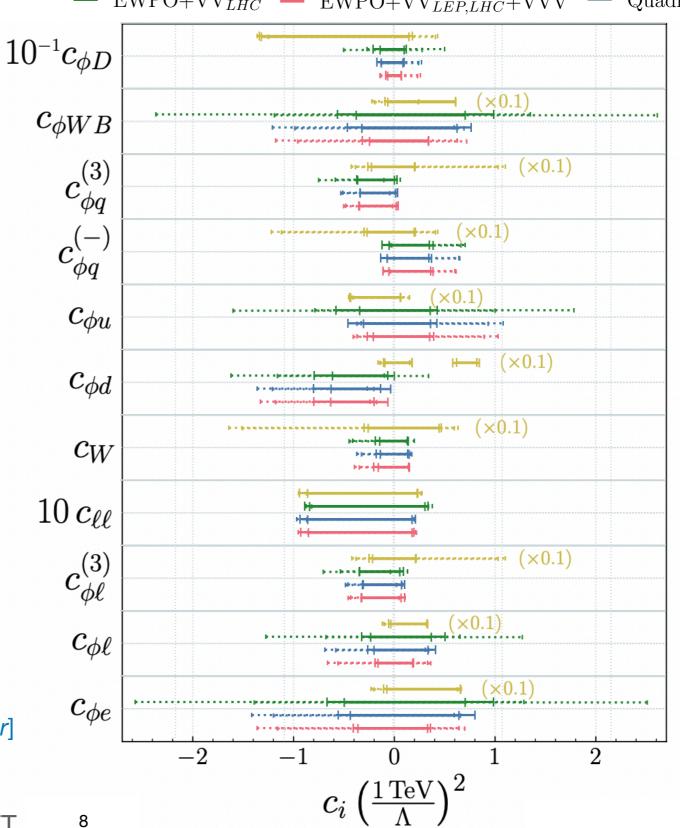
Fitmaker [Ellis et al.; 2012.02779]

Marginalised 95% C.I.

- EWPO+VV $_{LEP}$  EWPO+VV $_{LEP,LHC}$  ... Linear
- EWPO+VV<sub>LHC</sub> EWPO+VV<sub>LEP,LHC</sub>+VVV Quadratic

- LHC VV & VVV appear to improve significantly the bounds from EWPOs & LEP VV
- Quadratic fit: 50% improvement from VVV wrt VV on  $c_{\phi D}$ ,  $c_{\phi WB}$ ,  $c_{\phi \ell}$ ,  $c_{\phi e}$
- Bounds dominated by quadratic

[EC, Durieux, Mimasu, Vryonidou; to appear]



#### Interpretation

• Three EWPOs unconstrained directions:  $w_B, w_W + c_W$ 

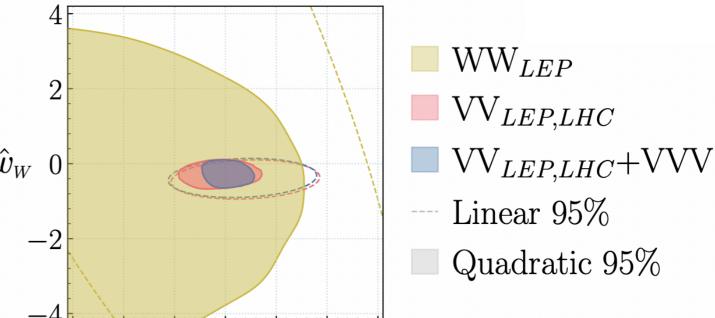
$$\begin{split} g_1^2 \, w_B &= g_1^2 \frac{\bar{v}_T^2}{\Lambda^2} \left( -\frac{1}{3} C_{Hd} - C_{He} - \frac{1}{2} C_{Hl}^{(1)} + \frac{1}{6} C_{Hq}^{(1)} + \frac{2}{3} C_{Hu} + 2 C_{HD} - \frac{1}{2 t_{\hat{\theta}}} C_{HWB} \right), \\ g_2^2 \, w_W &= g_2^2 \frac{\bar{v}_T^2}{\Lambda^2} \left( \frac{C_{Hq}^{(3)} + C_{Hl}^{(3)}}{2} - \frac{t_{\bar{\theta}}}{2} C_{HWB} \right). \end{split}$$
 [Brivio and Trott; 1701.06424]

- 3/11 directions unconstrained in a EWPOs only fit
- additional data is needed (multiboson)
- 2 possible origins of the improvement
- 1. constraints in EWPOs blind space + marginalisation
- 2. genuine effect of higher sensitivity in all directions

# Where do VV & VVV help?

Three EWPOs unconstrained parameters:  $\hat{w}_{R}$ ,  $\hat{w}_{W}$ ,  $c_{W}$ 

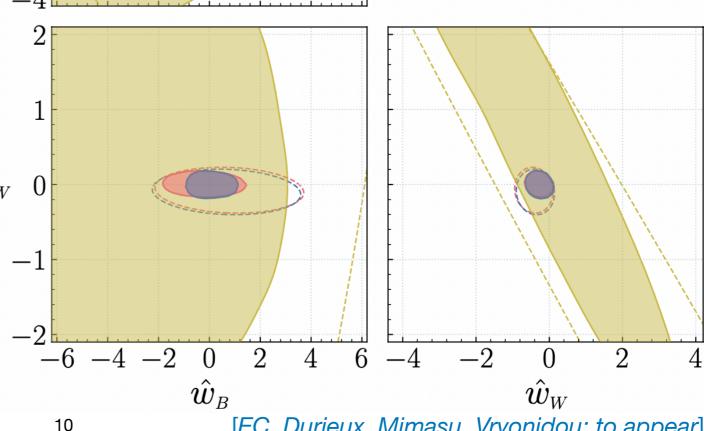
• Large  $\mathcal{O}(\Lambda^{-4})$  effect (also for  $\hat{w}_{\scriptscriptstyle W}$   $_0$ LEP VV!)



LHC VV dominates over LEP

•  $\vee\vee\vee$  at  $\mathcal{O}(\Lambda^{-2})$  doesn't help

• VVV constrains  $\hat{w}_R$  at  $\mathcal{O}(\Lambda^{-4}) c_W$  0

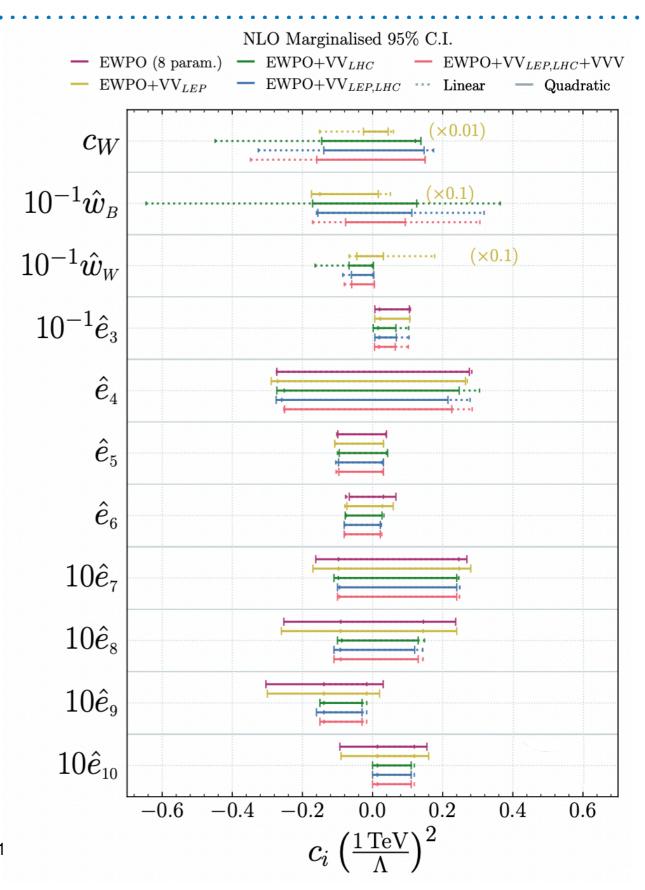


#### What about the other directions?

Does multiboson help EWPOs in the directions orthogonal to  $\{\hat{w}_B, \hat{w}_W, O_{WWW}\}$ ?

 in general, EWPOs constraints are dominant



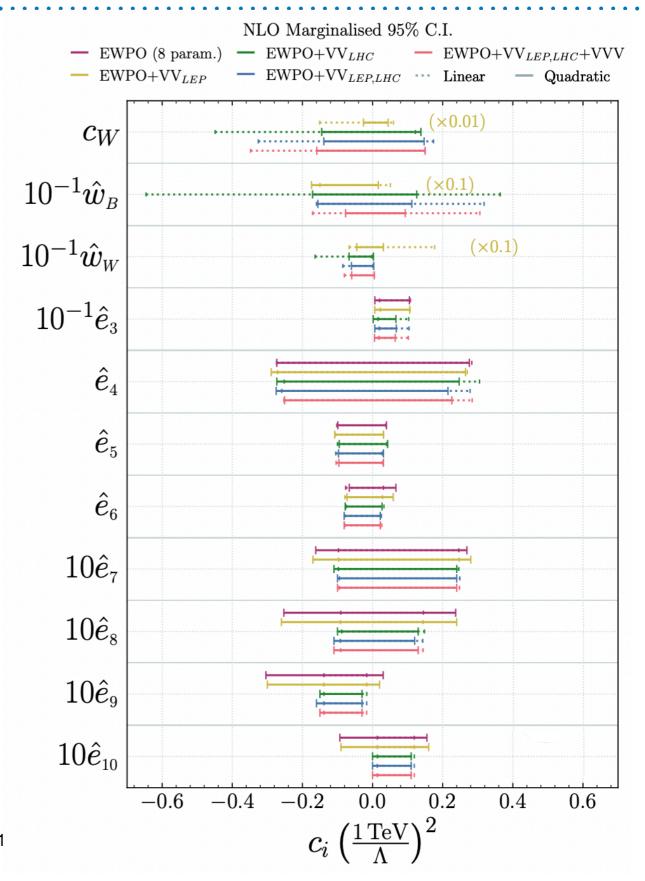


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Does multiboson help EWPOs in the directions orthogonal to  $\{\hat{w}_B, \hat{w}_W, O_{WWW}\}$ ?

- in general, EWPOs constraints are dominant
- mild improvement from quadratics (even EWPOs) on some directions

[EC, Durieux, Mimasu, Vryonidou; to appear]

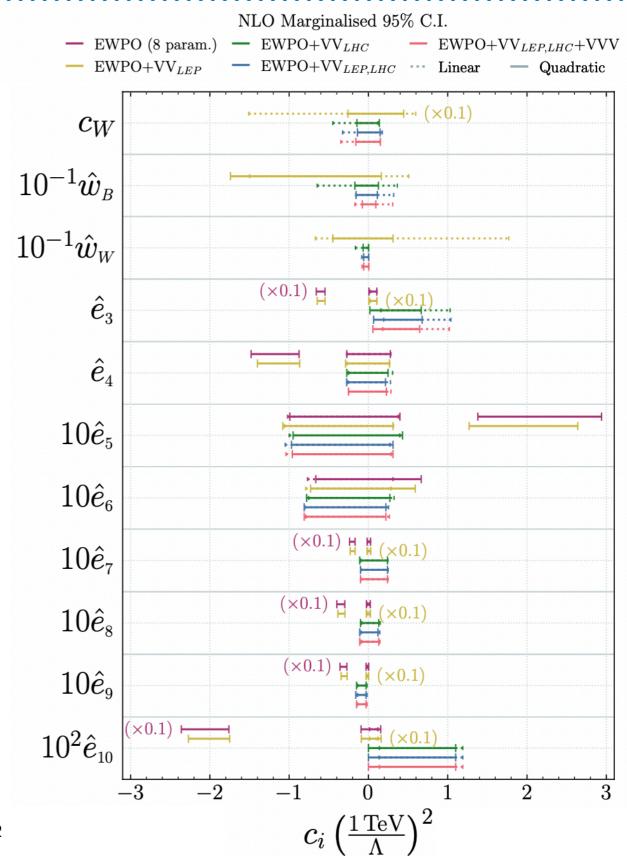


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- in general, EWPOs constraints are dominant
- mild improvement from quadratics (even EWPOs) on some directions
- secondary minima in EWPOs+LEP VV lifted by LHC VV

[EC, Durieux, Mimasu, Vryonidou; to appear]



#### Summary & conclusions

- QCD corrections have significant impact in LHC VV&VVV
- Multiboson production is complementary to EWPOs: triboson improves the bounds in EWPOs flat space
- Quadratic EFT contributions are sizeable for all the processes, from EWPO leading to secondary minima, to LEP diboson, and the LHC VV&VVV