

# Charm, $\gamma$ and $\phi_s$ at LHCb

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on behalf of the  
LHCb Collaboration

Standard Model @ LHC 2024

Roma  
Consiglio nazionale delle ricerche

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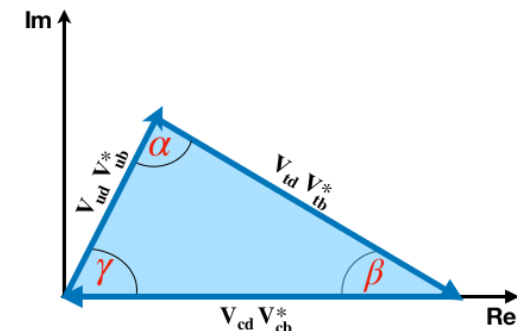
- The CKM matrix describes the quark charged current weak interactions

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \sim \begin{pmatrix} |V_{ud}| & |V_{us}| & |V_{ub}|e^{-iy} \\ -|V_{cd}| & |V_{cs}| & |V_{cb}| \\ |V_{td}|e^{-i\beta} & -|V_{ts}|e^{i\beta s} & |V_{tb}| \end{pmatrix}$$

- The unitarity of this matrix leads to

$$V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$$

- It can be visualized as a triangle in the complex plane

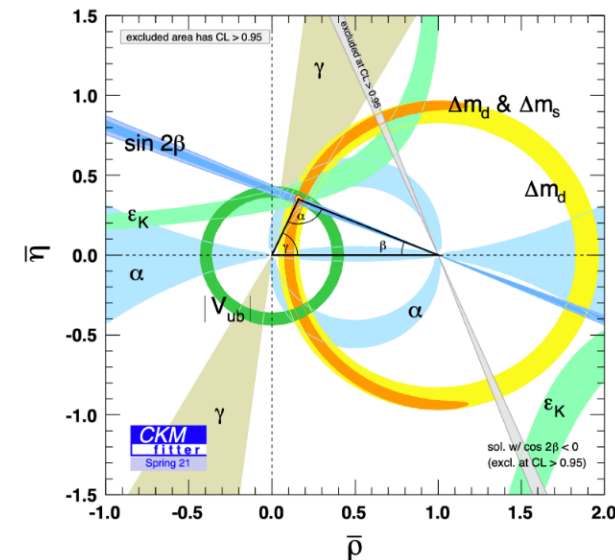


- The key test of the SM is the check of the unitarity of the CKM matrix

- A single complex phase in the CKM matrix is the only measured source of CP violation (CPV)

- Not enough to explain the matter-antimatter asymmetry in the universe.

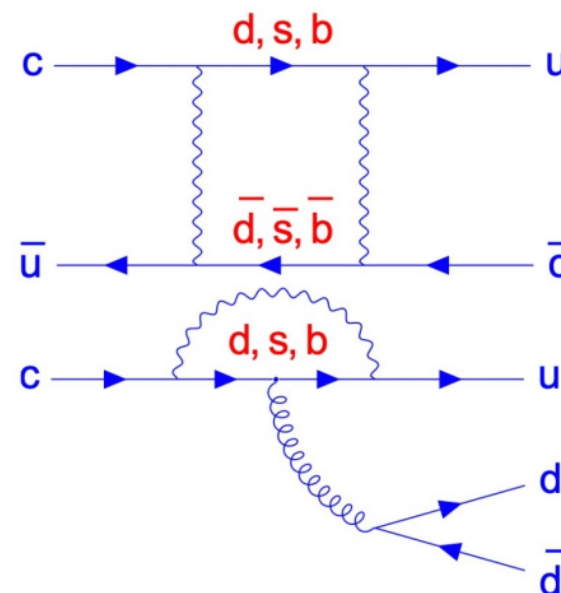
- B and D meson decays are a great laboratory to probe CP violation and to test the unitarity of CKM matrix



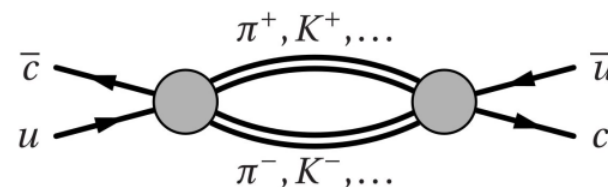
- Charm sector is an unique laboratory to study CPV in up-type quark decays
- CPV in charm is highly suppressed in the SM
  - beauty loop are suppressed by smallness of CKM elements

$$CPV \propto \text{Im} \left( \frac{V_{cb} V_{bu}^*}{V_{cs} V_{su}^*} \right) \approx -6 \times 10^{-4}$$

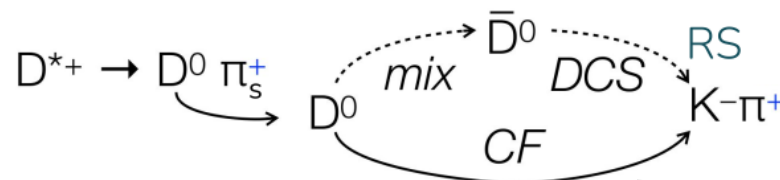
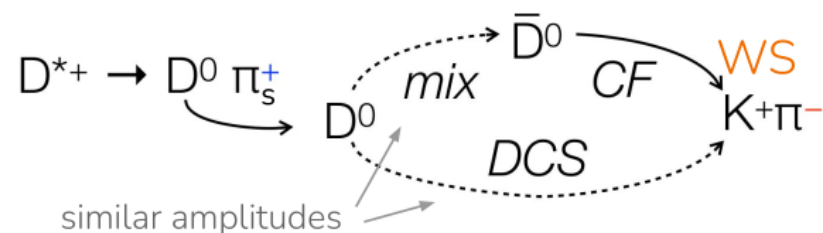
- strange-down loops suppressed by GIM cancellation broken by  $b$  quark
- CPV in charm small  $O(10^{-4}) \rightarrow$  sensitive to NP
- Theory predictions complicated by QCD effects, large and difficult to compute
- First observation of CPV in decay in  $D^0 \rightarrow h^+ h^-$  @LHCb( $5.3\sigma$ )  
[PRL122(2019)211803]
- First evidence for direct CP violation in specific  $D^0$  decays:  
[Phys. Rev. Lett. 131 (2023) 091802]
  - $3.8\sigma$  in  $D^0 \rightarrow \pi^- \pi^+$
  - $1.4\sigma$  in  $D^0 \rightarrow K^- K^+$



- $D^0 \rightarrow K^+\pi^-$  decays allows to simultaneously measure the mixing and all types of CPV.
- The SM prediction for the  $D^0$  mixing amplitude governed by two contributions:
  - Short distance: suppressed by CKM  $b$  coupling and GIM mechanism
  - Long distances: low energy QCD through on-shell resonances  $\rightarrow$  theoretical prediction of  $x$  and  $y$  very challenging



- Full run 2 LHCb analysis
- Time dependent analysis of  $D^0 \rightarrow K^+\pi^-$  decays
- Reconstruct  $D^0$  from  $D^{*+} \rightarrow D^0\pi^-$  decays
- Distinguish two processes
  - Wrong sign (WS)
  - Right sign (RS)



[LHCb-PAPER-2024-008, in preparation]

normalization ch.  
dominated by CF

- In order to reduce the dependence on the mixing and CPV quantities a time dependent fit of WS/RS yield is performed:

$$R_{K\pi}^+(t) \equiv \frac{\Gamma(D^0(t) \rightarrow K^+\pi^-)}{\Gamma(\bar{D}^0(t) \rightarrow K^+\pi^-)} \quad R_{K\pi}^-(t) \equiv \frac{\Gamma(\bar{D}^0(t) \rightarrow K^-\pi^+)}{\Gamma(D^0(t) \rightarrow K^-\pi^+)}$$

- Since  $x_{1,2}$  and  $y_{1,2} \ll 1$  this ratio can be expanded as:

$$R_{K\pi}^\pm(t) = R_{K\pi} (1 \pm A_{K\pi}) + \sqrt{R_{K\pi} (1 \pm A_{K\pi})} (c_{K\pi} \pm \Delta c_{K\pi}) t/\tau_{D^0} + (c'_{K\pi} \pm \Delta c'_{K\pi}) (t/\tau_{D^0})^2$$

$$R_{K\pi} = \frac{1}{2} \left( \left| \frac{A_{\bar{f}}}{\bar{A}_{\bar{f}}} \right|^2 + \left| \frac{\bar{A}_f}{A_f} \right|^2 \right),$$

$$A_{K\pi} = \frac{|A_{\bar{f}}/\bar{A}_{\bar{f}}|^2 - |\bar{A}_f/A_f|^2}{|A_{\bar{f}}/\bar{A}_{\bar{f}}|^2 + |\bar{A}_f/A_f|^2} \approx a_{\text{DCS}}^d, \quad \leftarrow \text{CPV in decay}$$

$$c_{K\pi} \approx y_{12} \cos \phi_f^\Gamma \cos \Delta_f + x_{12} \cos \phi_f^M \sin \Delta_f,$$

$$\Delta c_{K\pi} \approx x_{12} \sin \phi_f^M \cos \Delta_f - y_{12} \sin \phi_f^\Gamma \sin \Delta_f, \quad \leftarrow \text{CPV in the interference}$$

$$c'_{K\pi} \approx \frac{1}{4} (x_{12}^2 + y_{12}^2),$$

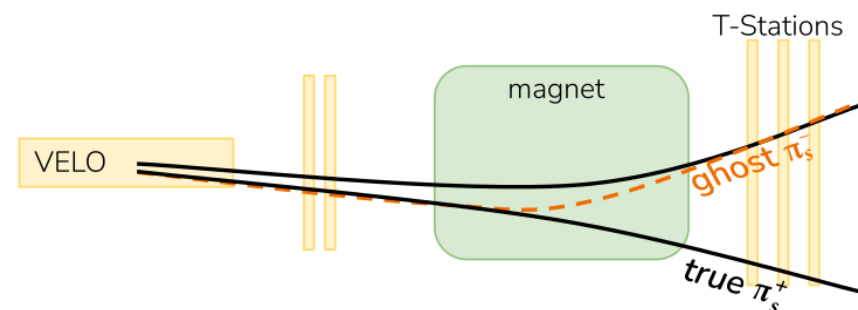
$$\Delta c'_{K\pi} \approx \frac{1}{2} x_{12} y_{12} \sin(\phi_f^M - \phi_f^\Gamma). \quad \leftarrow \text{CPV in mixing}$$

$$x_{12} \equiv 2|M_{12}|/\Gamma$$

$$\phi_2^M \sim \arg(M_{12}), \quad \phi_2^\Gamma \sim \arg(\Gamma_{12}) \quad y_{12} \equiv |\Gamma_{12}|/\Gamma$$

- CP violation measurements
- $A_{K\pi}$ : rigorous null test of SM since  $c \rightarrow uds$  doesn't receive any contribution from QCD nor chromomagnetic dipole operators
- $\Delta_f$ : improve the knowledge on  $SU(3)_F$  breaking and rescattering effects at energy scale of the charm mass

- Sample is divided between
  - $D^0$  final state ( $K\pi^+, K^+\pi^-$ ),
  - 18  $D^0$  decay-time intervals
  - 3 data-taking period (2015-16, 2017 and 2018)
- Extracted values:
  - average  $D^0$  decay time
  - WS-to-RS ratio,  $R$ , fitting  $D^*$  mass to disentangle signal from combinatorial and ghost backgrounds
  - ghost backgrounds: misassociation of correctly-identified hits in VELO with hits in T-Stations from different particles)
- Correction them from the known systematic effects
  - bias to the ratio
  - bias to asymmetry
  - bias to  $D^0$  decay-time
- Experimental challenges:
  - Backgrounds
  - Nuisance asymmetry
- Time dependence is fitted  $\rightarrow$  extract mixing and CPV parameters

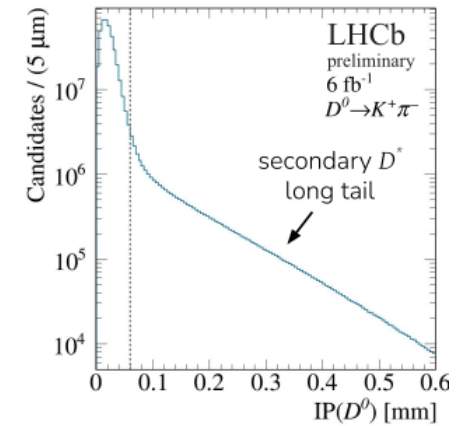
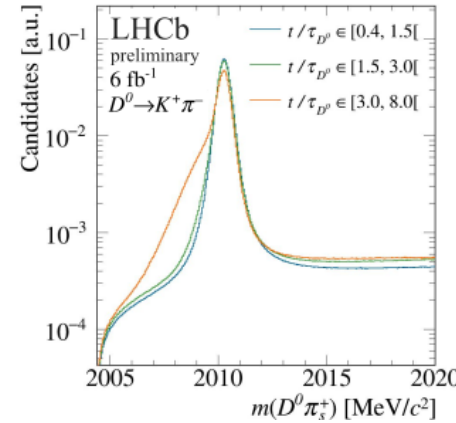
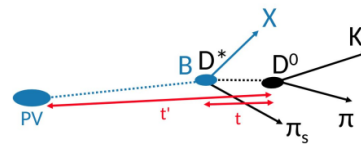


[LHCb-PAPER-2024-008]

- **Bias of the decay time:**
  - Poor  $D^*$  vertex resolution (1cm)  $\rightarrow$  request to  $D^*$  to originate from primary vertex

$\rightarrow$  contamination from secondary  $D^*$

- bias towards higher values
- deformed  $D^*$  shape



- **Bias of charge asymmetry:**

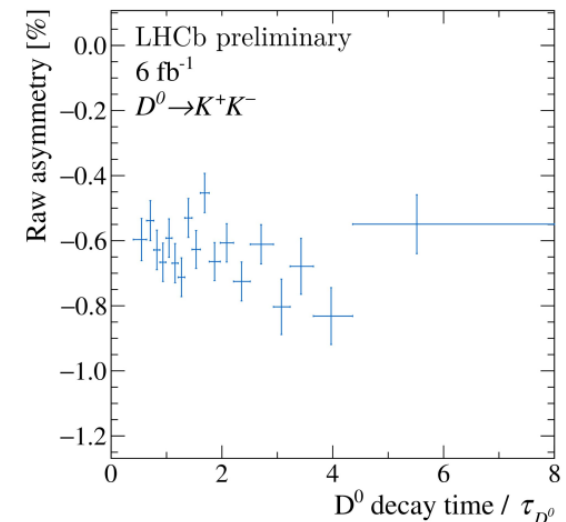
- Originate by the differences in reconstruction efficiency between WS and RS, may mimic CPV

$$\widetilde{R}'^{\pm} = R'^{\pm} \frac{\int [1 \pm A_P(D^*)] \epsilon(\pi_s^{\pm}) \epsilon(K^{\pm} \pi^{\mp}) \rho d\vec{p}_{D^0} d\vec{p}_{\pi_s}}{\int [1 \mp A_P(D^*)] \epsilon(\pi_s^{\mp}) \epsilon(K^{\pm} \pi^{\mp}) \rho d\vec{p}_{D^0} d\vec{p}_{\pi_s}} \simeq R'^{\pm} \frac{1 \pm [A_D(\pi_s) + A_P(D^*)]}{1 \mp [A_D(\pi_s) + A_P(D^*)]}$$

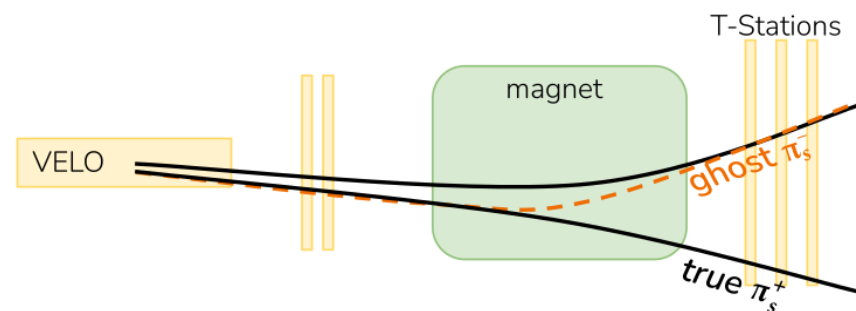
- Determinated with  $D^0 \rightarrow K^+ K^-$  control mode
- To extract the raw asymmetry  $D^{*-}$  and  $D^{*+}$  are fitted simultaneously

$$A_D(\pi_s) + A_P(D^*) = A^{raw}(KK) - a_{KK}^d - \Delta Y \langle t \rangle$$

direct CP asymmetry in  $D^0 \rightarrow K^+ K^-$ 
time dependent CP asymmetry in  $D^0 \rightarrow K^+ K^-$

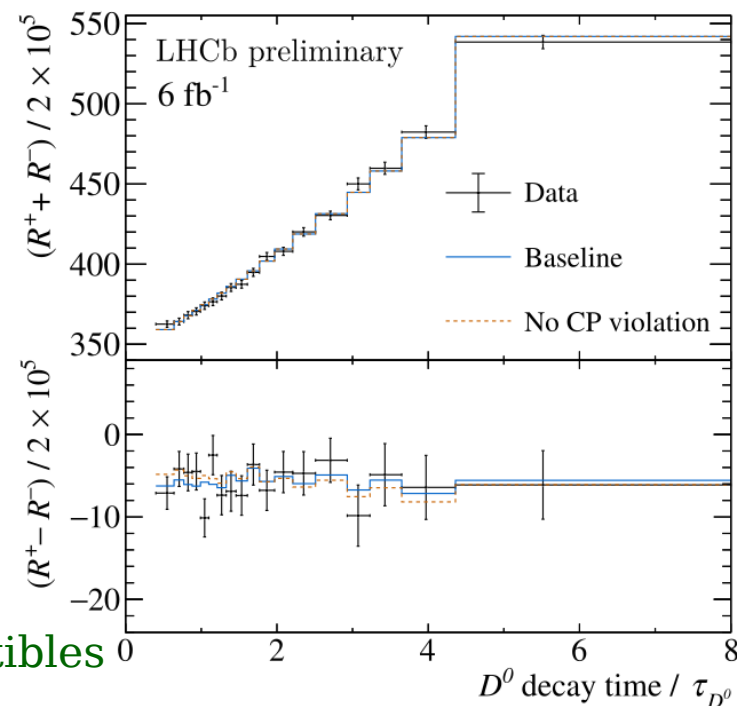
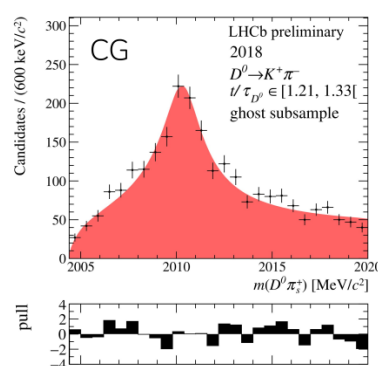
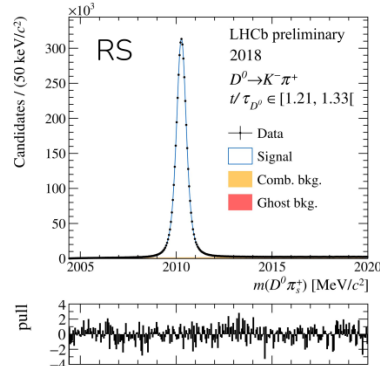
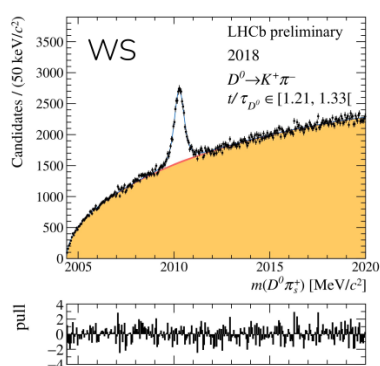


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- Bias of ratio:
  - due to the contamination of doubly misidentified  $D^0$  decays
  - removal of common candidates from the sample of WS decays
    - $D^0$  candidates used to reconstruct both WS  $D^{*+}$  and RS  $D^{*-}$ 
      - RS  $D^{*-}$  kept but this removes a small fraction of WS biasing the ratio
- A binned fit is performed simultaneously to  $D^*$  mass distribution of WS, RS and common ghost



- The main systematic sources are  $D^*$  mass fit model and ghost bkg pdf

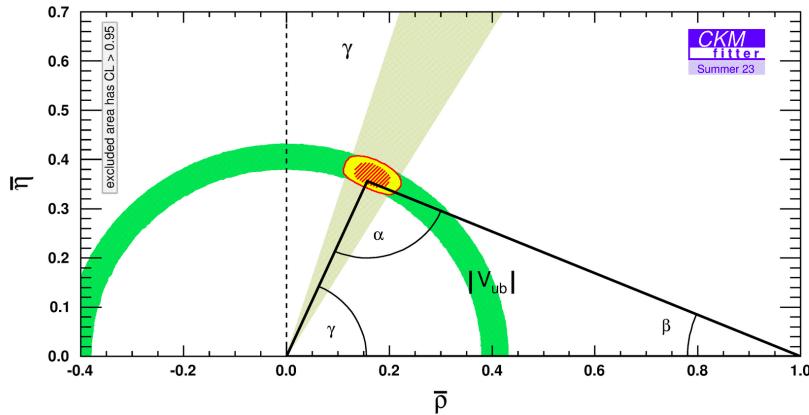
Parameters

$R_{K\pi}$	$(342.7 \pm 1.9) \times 10^{-5}$
$c_{K\pi}$	$(52.8 \pm 3.3) \times 10^{-4}$
$c'_{K\pi}$	$(12.0 \pm 3.5) \times 10^{-6}$
$A_{K\pi}$	$(-6.6 \pm 5.7) \times 10^{-3}$
$\Delta c_{K\pi}$	$(2.0 \pm 3.4) \times 10^{-4}$
$\Delta c'_{K\pi}$	$(-0.7 \pm 3.6) \times 10^{-6}$

- 20% improvement on  $\Phi_2^M$
- Run 1 and run 2 results compatibles
- Total uncertainty improved by 1.6 with respect Run 1
- No evidence CPV neither in decays, mixing nor interference

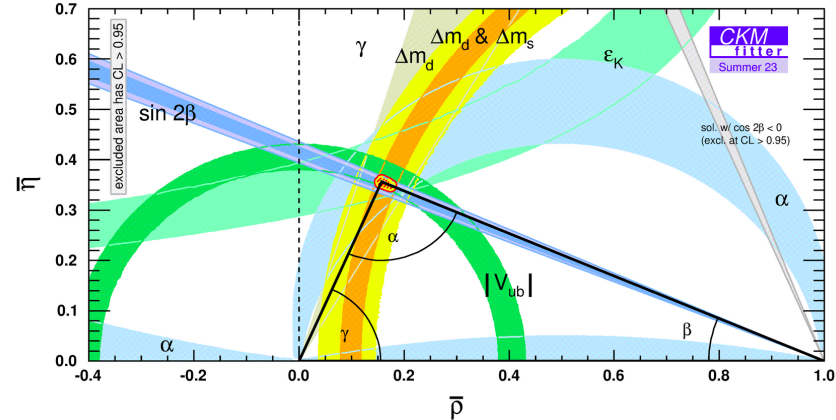
# CKM $\gamma$ angle

- $\gamma$  is the phase difference between  $b \rightarrow c$  and  $b \rightarrow u$  quark transition
  - measurable in purely tree level or indirectly
  - negligible theoretical uncertainty  $\sim 10^{-7}$  [Zupan & Brod 1308.5663]
  - current experimental uncertainty is  $< 4^\circ$



Tree level direct measurements

$$\gamma_{\text{direct}} = (66.2^{+3.4}_{-3.6})^\circ$$

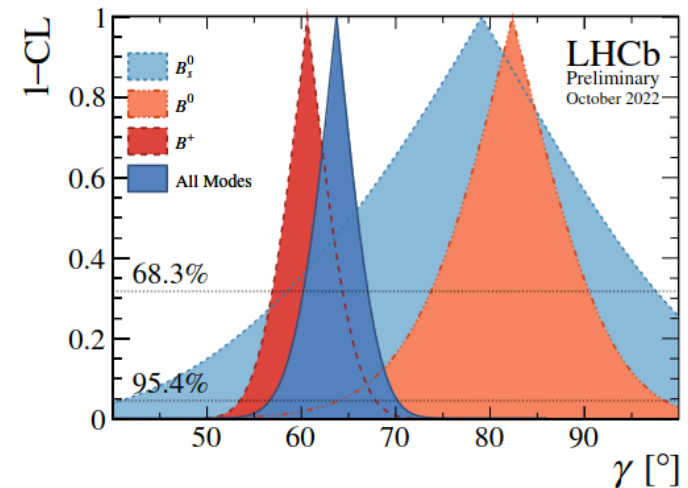


Loop-level indirect measurements

$$\gamma_{\text{indirect}} = (65.6^{+0.9}_{-2.7})^\circ$$

- LHCb combination:
  - simultaneous fit of  $\gamma$  and  $D^0$  mixing parameters

$$\gamma = (63.8^{+3.5}_{-3.7})^\circ \quad [\text{JHEP } 12(2021)141]$$



- It is typically measured in B decays such as  $B^\pm \rightarrow Dh^\pm$  (where  $D = D^0, \bar{D}^0$  and  $h = K, \pi$ )

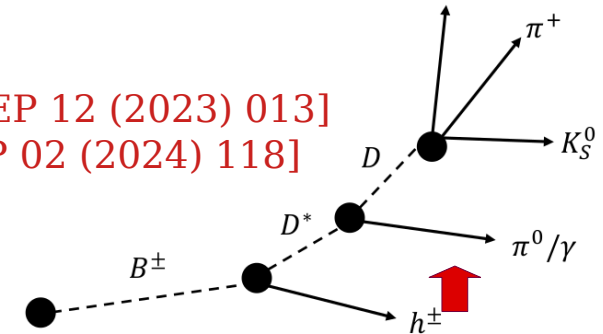
$$|A(B^-)|^2 \propto A_D^2 + r_B^2 A_{\bar{D}}^2 + 2A_D A_{\bar{D}} r_B \cos(\delta_B - \gamma)$$

- Measurement technique depends on D-decay mode

$$|A(B^+)|^2 \propto A_D^2 + r_B^2 A_{\bar{D}}^2 + 2A_D A_{\bar{D}} r_B \cos(\delta_B + \gamma)$$

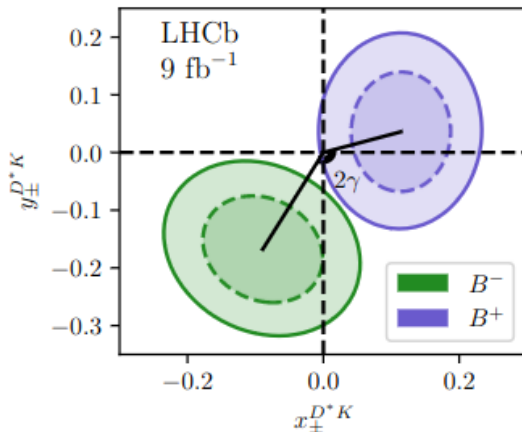
## LHCb $\gamma$ measurements with multibody D decays

- $B^\pm \rightarrow D^* K^\pm$  (full reconstructed  $\rightarrow$  better control of backgrounds) [JHEP 12 (2023) 013]
- $B^\pm \rightarrow D^* K^\pm$  (partially reconstructed  $\rightarrow$  higher signal efficiency) [JHEP 02 (2024) 118]  
with  $D^* \rightarrow D^0 \pi$  or and  $D^0 \rightarrow K^0_s h^+ h^-$



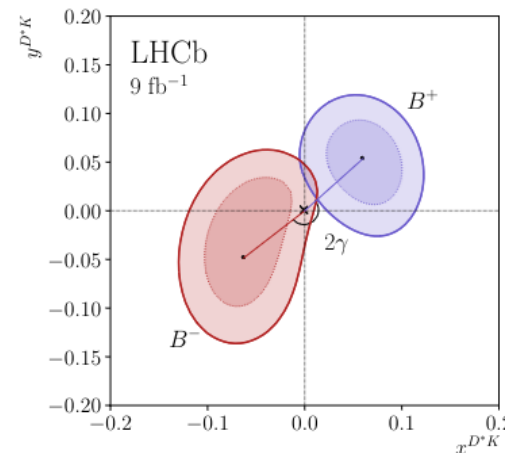
- The measurements are performed by analyzing the signal yields variation across the D decay phase space

- They are independent of any amplitude model
- direct measurement of D strong phase from BESIII and CLEO [JHEP05(2021)164]



[JHEP 12 (2023) 013]

$$\begin{aligned} \gamma &= (69_{-14}^{+13})^\circ \\ r_B^{D^*K} &= 0.15 \pm 0.03 \\ \delta_B^{D^*K} &= (311 \pm 15)^\circ \end{aligned}$$



[JHEP 02 (2024) 118]

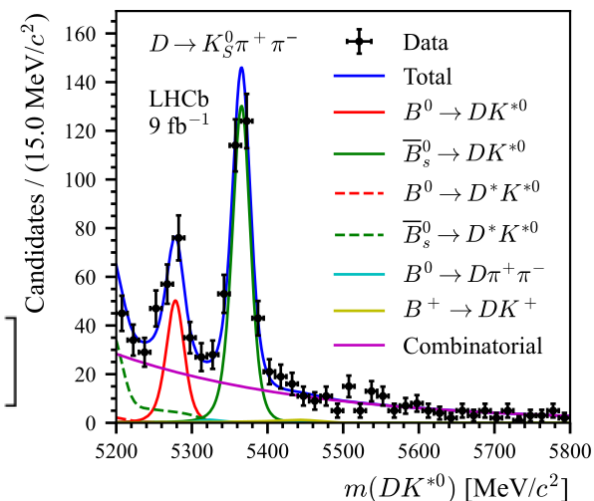
$$\begin{aligned} \gamma &= (92_{-17}^{+21})^\circ \\ r_B^{D^*K} &= 0.080_{-0.023}^{+0.022} \\ \delta_B^{D^*K} &= (310_{-20}^{+15})^\circ \end{aligned}$$

- Gain complementary info from  $B^0 \rightarrow DK^*(892)^0$ :
  - interference 3 times larger that for  $B^\pm \rightarrow Dh^\pm$
- $B^0 \rightarrow DK^*(892)^0$  with  $D \rightarrow K^0_s h^+h^-$  [Eur. Phys. J. C 84 (2024)]**
  - The  $\gamma$  angle is determined by examining the distributions of signal decays in phase space bins of  $D^0 \rightarrow K^0_s h^+h^-$

$$N_i(B^0) = h^{B^0} \left[ F_{-i} + (x_+^2 + y_+^2)F_i + 2\kappa\sqrt{F_i F_{-i}}(x_+c_i - y_+s_i) \right]$$

$$x_\pm \equiv r_{B^0} \cos(\delta_{B^0} \pm \gamma)$$

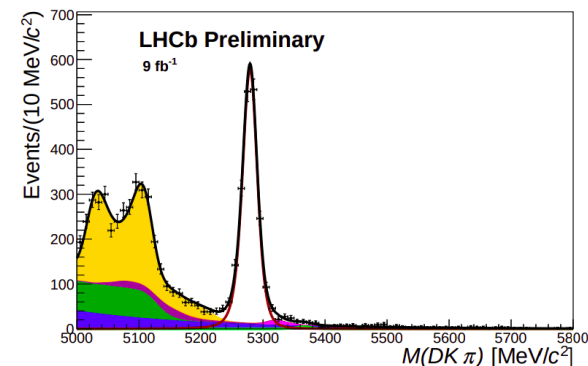
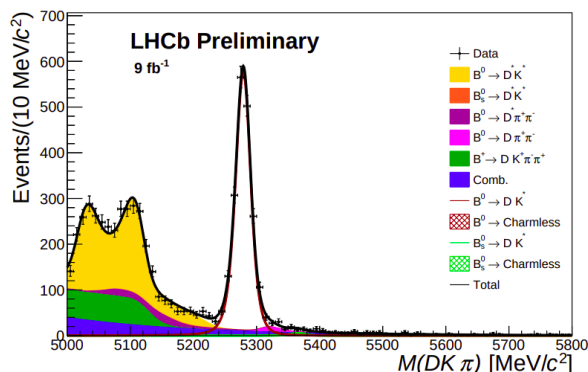
$$y_\pm \equiv r_{B^0} \sin(\delta_{B^0} \pm \gamma)$$



- $B^0 \rightarrow D^0K^*(892)^0$  with the ADS and GLW D-decays final states [arXiv:2401.17934]**

Fit to selected data through a simultaneous unbinned extended maximum-likelihood fit of the  $B^0$  candidate reconstructed mass on each flavour of each  $D^0$  final state

- Statistical precision on the CP-violating observables has improved by around 60% comparison to the previous results

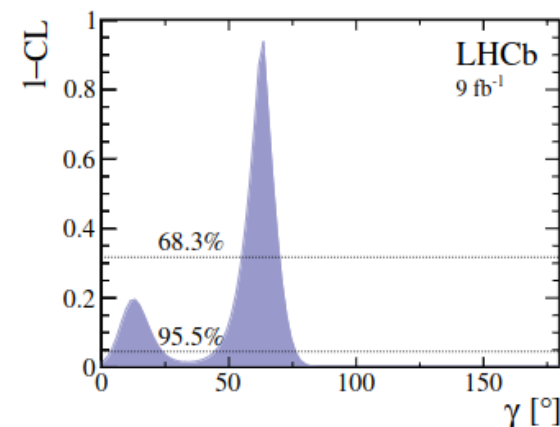


Combination of the last two results  
Results from  $D \rightarrow K^0_s h^+h^-$  broke the degeneracy

$$\gamma = (63.3 \pm 7.2)^\circ$$

$$r_{B^0}^{DK^*} = 0.233 \pm 0.016$$

$$\delta_{B^0}^{DK^*} = (191.8 \pm 6.0)^\circ$$



# Decay-time dependent $\gamma$ measurement

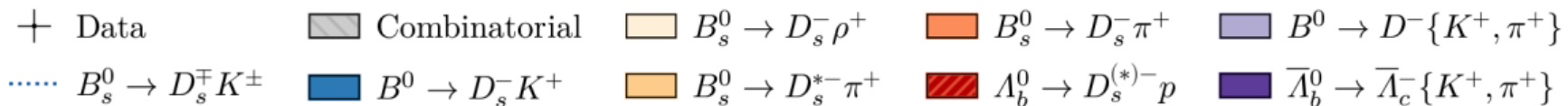
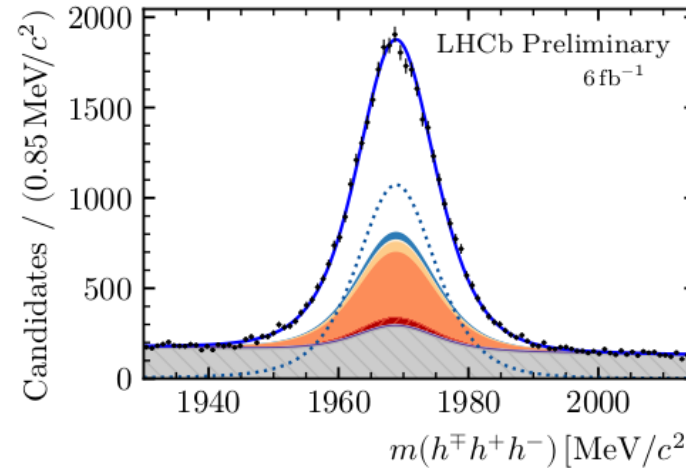
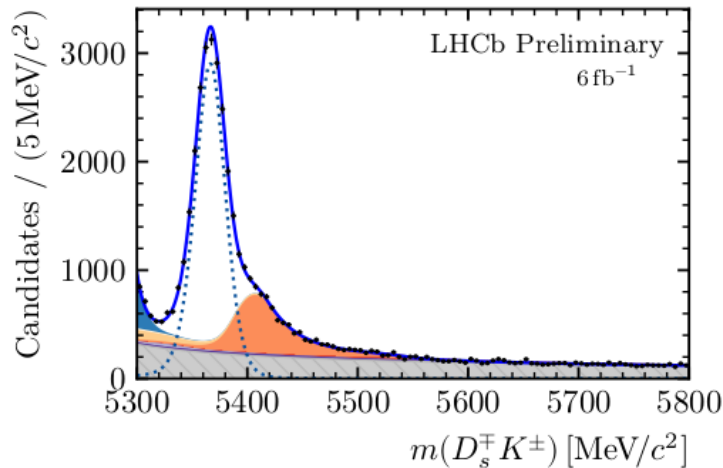
- **Measurement of CP asymmetry in  $B_s^0 \rightarrow D_s^- K^+$  [LHCb-CONF-2023-004]**
- With  $D_s^- \rightarrow K^+ \pi^+ \pi^-$ ,  $D_s^- \rightarrow K^- K^+ \pi^-$ ,  $D_s^- \rightarrow \pi^+ \pi^- \pi^-$
- Time dependent measurement of  $\gamma$

$$\Gamma(B_s^0(t) \rightarrow f\bar{f}) \sim e^{-\Gamma_s t} \left( \cosh\left(\frac{\Delta\Gamma_s}{2} t\right) + C_{f\bar{f}} \cos(\Delta m_s t) + A_{f\bar{f}}^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma_s}{2} t\right) - S_{f\bar{f}} \sin(\Delta m_s t) \right)$$

$$C_f = C_{\bar{f}} = \frac{1 - r_{D_s K}^2}{1 + r_{D_s K}^2} \quad A_f^{\Delta\Gamma} = \frac{-2 r_{D_s K} \cos(\delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2} \quad S_f = \frac{2 r_{D_s K} \sin(\delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}$$

$$A_{\bar{f}}^{\Delta\Gamma} = \frac{-2 r_{D_s K} \cos(\delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2} \quad S_{\bar{f}} = \frac{2 r_{D_s K} \sin(\delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}$$

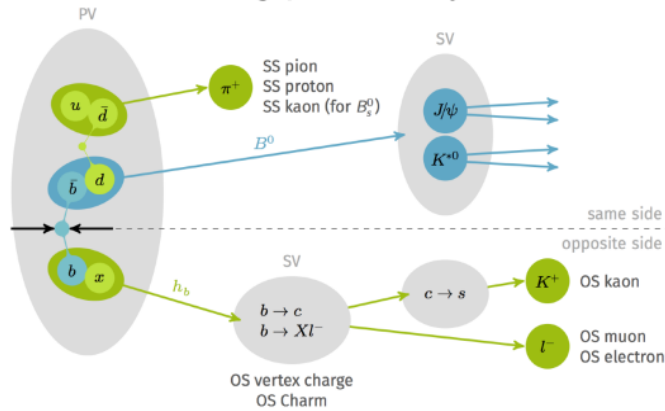
- Simultaneous fit of all modes and run 2 years
- Two dimensional invariant mass fit:  $20950 \pm 180$  candidates



- The decay time fit has to be corrected for:

To separate  $B_s^0/\bar{B}_s^0$  candidates  $\longrightarrow$  flavour-tagging

- estimates initial flavour
- exploits various fragmentation processes
- MVA-based mistag probability



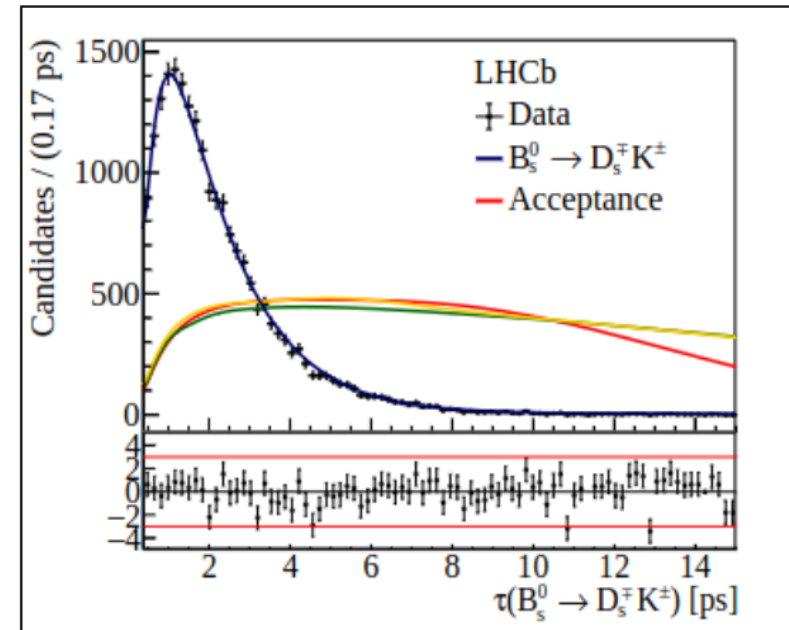
## Decay-time resolution

- Finite decay-time resolution in the detector leads to a dilution of the observed oscillation
- The prompt sample of  $D_s^- \pi^+$  was exploited

$$D_{res} \approx e^{-\frac{1}{2}\Delta m_s^2 \sigma^2}$$

## Decay-time acceptance

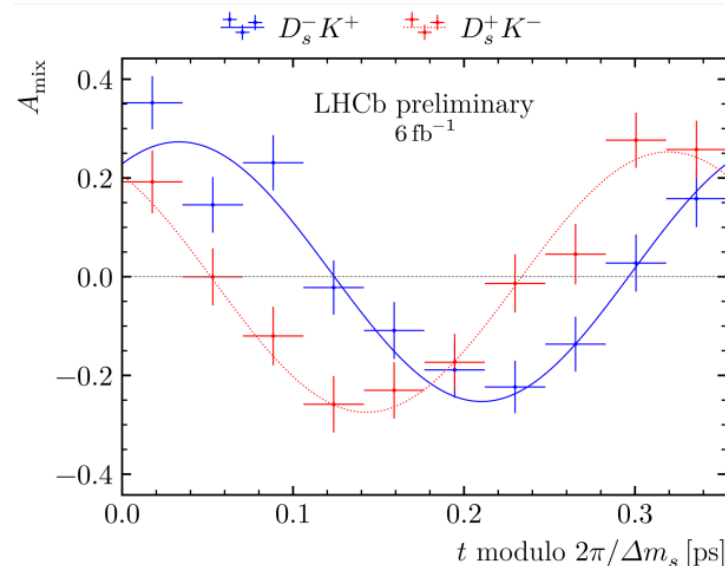
- Decay-time distorted by selection requirements
- Heavily correlated with the CP observable  $D_f, D_{\bar{f}}$
- Acceptance fixed to  $B^0 \rightarrow D_s^- \pi^+$  fit and corrected by the ratio of the decay-time acceptances of  $D_s^- K^+$  and  $D_s^- \pi^+$



# Decay-time dependent $\gamma$ measurement

- External input:  $\Delta\Gamma_s, \Gamma_s$ , detection asymmetry
- Input from  $B^0 \rightarrow D_s^- \pi^+$  [Nature Physics 18, (2022) 1-5]
  - Resolution calibration
  - Decay time acceptance
  - Tagging calibration
  - Production asymmetry
  - $\Delta m_s$
- External input:  $\phi_s = -2\beta_s$  [arXiv:2308.01468]
- Significant CP violation in the interference ( $8.8\sigma$ )
- Measured CPV observables

[LHCb-CONF-2023-004]



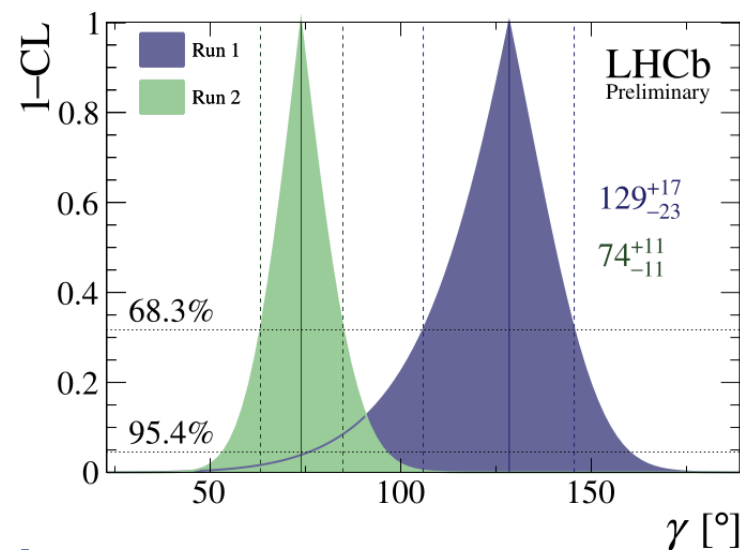
Parameter	Value
$C_f$	$0.791 \pm 0.061 \pm 0.022$
$A_f^{\Delta\Gamma}$	$-0.051 \pm 0.134 \pm 0.037$
$A_{\bar{f}}^{\Delta\Gamma}$	$-0.303 \pm 0.125 \pm 0.036$
$S_f$	$-0.571 \pm 0.084 \pm 0.023$
$S_{\bar{f}}$	$-0.503 \pm 0.084 \pm 0.025$

Source	$C_f$	$A_f^{\Delta\Gamma}$	$A_{\bar{f}}^{\Delta\Gamma}$	$S_f$	$S_{\bar{f}}$
$\Delta m_s$	0.007	0.004	0.004	0.108	0.103
Detection asymmetry	—	0.079	0.083	0.006	0.007
Multivariate fit	0.045	0.095	0.121	0.088	0.112
Flavour tagging	0.256	0.026	0.028	0.012	0.070
Decay-time resolution model	0.195	0.002	0.003	0.058	0.167
Decay-time bias	0.062	0.027	0.046	0.188	0.167
Decay-time acceptance, $\Gamma_s, \Delta\Gamma_s$	0.006	0.225	0.231	0.003	0.003
Decay-time acceptance ratios	0.001	0.018	0.018	—	—
Neglecting correlations	0.137	0.081	0.054	0.135	0.043
Total	0.358	0.273	0.285	0.278	0.294

$$\gamma = (74 \pm 11)^\circ,$$

$$\delta = (346.9 \pm 6.6)^\circ,$$

$$r_{D_s K} = 0.327 \pm 0.038,$$



- Compatible with run 1 @  $1.3\sigma$
- Combination between run 1 and run 2 ongoing

# Measurement of $\phi_s$ in $B_s^0 \rightarrow J/\psi K^+ K^-$

[PRL132(2024)051802]

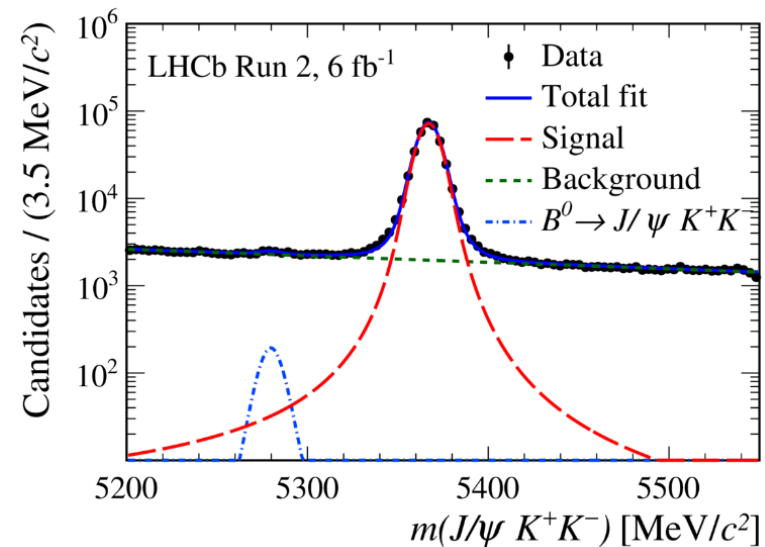
- B decays to CP eigenstates allow to probe the mixing phase  $\beta_s = \phi_s/2$  through the interference between decays with and without mixing with a  $c\bar{c}$  resonance in the final state.

$$\beta_s \equiv \arg \left[ - (V_{ts} V_{tb}^*) / (V_{cs} V_{cb}^*) \right]$$

- Unique opportunity to LHC due to Lorentz boost in p-p collisions
- $B_s^0 \rightarrow J/\psi K^+ K^-$  channel, in the vicinity of  $\phi(1020)$  resonance

- Full run II analysis

- The data sample is divided in 48 independent subsamples:
  - 6 bins in the of invariant mass in the  $[990, 1050]$  MeV/c<sup>2</sup>
  - 2 trigger categories,
  - 4 years of data taking



- Extended maximum likelihood fit to extract the signal yields

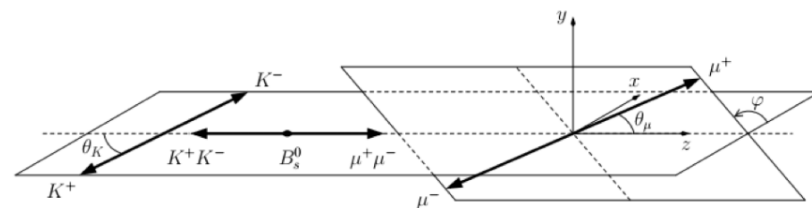


# Measurement of $\phi_s$ in $B_s^0 \rightarrow J/\psi K^+ K^-$

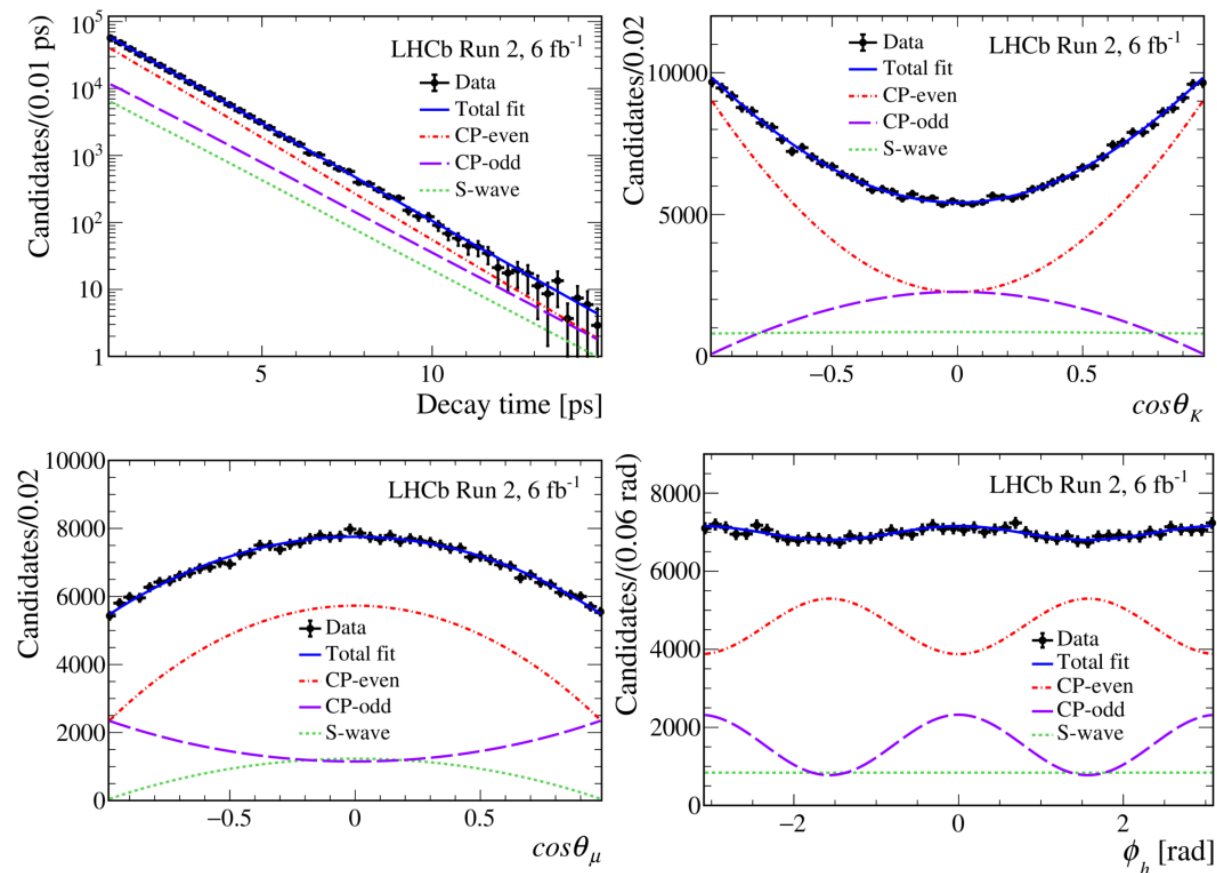
[PRL132(2024)051802]

- To extract  $\phi_s$ , CP even and CP-odd decay amplitude need to be disentangled since they depend on angular momentum between  $J/\psi$  and the kaons pair

→ A weighted simultaneous fit to decay time distribution and decays angles ( $\cos\theta_K$ ,  $\cos\theta_\mu$ ,  $\phi_h$ ) in the helicity basis is performed



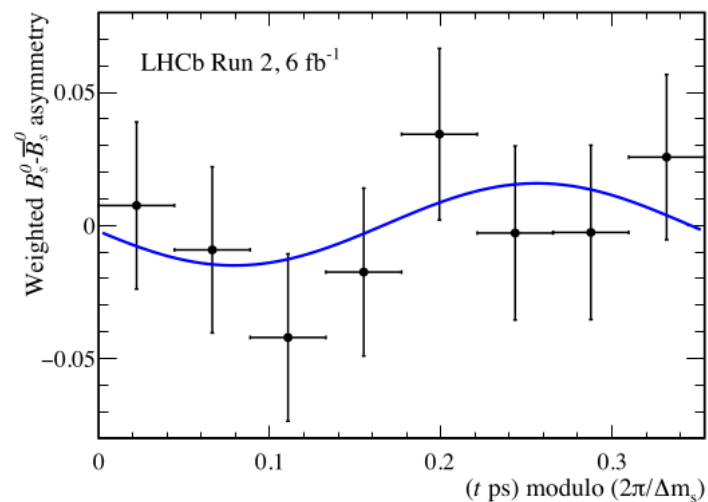
- The fit function account for :
  - Decay time resolution calibrated with prompt fake signals →  $\sigma \sim 42\text{ps}$
  - Flavour tagging calibrated with  $B^+ \rightarrow J/\psi K^+$  and  $B_s \rightarrow D_s \pi^+ \rightarrow \epsilon \sim 4\%$
  - Angular efficiencies



# Measurement of $\phi_s$ in $B_s^0 \rightarrow J/\psi K^+ K^-$

[PRL132(2024)051802]

Parameter	Values		
$\phi_s$ [rad]	$-0.039$	$\pm 0.022$	$\pm 0.006$
$ \lambda $	$1.001$	$\pm 0.011$	$\pm 0.005$
$\Gamma_s - \Gamma_d$ [ $\text{ps}^{-1}$ ]	$-0.0056$	$^{+0.0013}_{-0.0015}$	$\pm 0.0014$
$\Delta\Gamma_s$ [ $\text{ps}^{-1}$ ]	$0.0845$	$\pm 0.0044$	$\pm 0.0024$
$\Delta m_s$ [ $\text{ps}^{-1}$ ]	$17.743$	$\pm 0.033$	$\pm 0.009$
$ A_\perp ^2$	$0.2463$	$\pm 0.0023$	$\pm 0.0024$
$ A_0 ^2$	$0.5179$	$\pm 0.0017$	$\pm 0.0032$
$\delta_\perp - \delta_0$ [rad]	$2.903$	$^{+0.075}_{-0.074}$	$\pm 0.048$
$\delta_\parallel - \delta_0$ [rad]	$3.146$	$\pm 0.061$	$\pm 0.052$



[PRL 114 (2015) 041801, EPJC 79 (2019) 706, EPJC81 (2021) 1026]  
 [PLB 736 (2014) 186, PLB 797 (2019) 134789, PLB 762 (2016) 253, PRL 113 (2014) 211801]

- Most precise measurement to date and consistent with SM
- $|\lambda|$ : consistent with no direct CPV
- $\Gamma_s - \Gamma_d$ : consistent with HQE expectation [JHEP12 (2017) 068]
- No polarization dependence

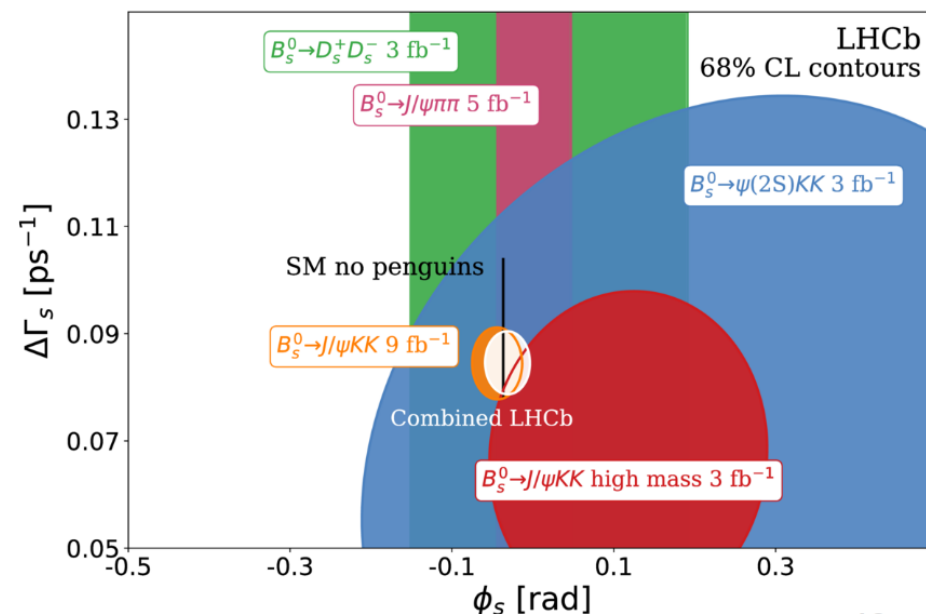
- Combination with run 1:

$$\phi_s = -0.044 \pm 0.020 \text{ rad}$$

$$|\lambda| = 0.990 \pm 0.010$$

- LHCb combination:

$$\phi_s = -0.031 \pm 0.018 \text{ rad}$$



A lot of results are still being produced with LHCb Run1 + Run2 sample

- World leading measurements of mixing phases of neutral B mesons
- New measurements of  $B \rightarrow Dh$  decays continuously improving the constraints on the  $\gamma$  angle
- LHCb is still exploiting its enormous charm data sample to look for CP violation in this sector
  - No evidence of discrepancies is observed with respect to SM expectations
- LHCb Upgrade I is going to start to collect data with the potential to more than double its sample in the next two years
- Complementarity and cross-check with Belle II will be fundamental as well



# Measurement of $\Delta\Gamma_s$ in $B_s^0 \rightarrow J/\psi \pi^+ \pi^-$ and $B_s^0 \rightarrow J/\psi \eta'$

- The decay-width difference between the light and heavy mass eigenstates
- $\Delta\Gamma_s$  can be determined from the decay-width difference between a CP-odd and a CP-even  $B_s^0$  mode.

If CP violation is negligible:

$$\Gamma(B_s^0(t) \rightarrow f) \propto e^{-\Gamma_s t} \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + \eta_{CP} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) \right]$$

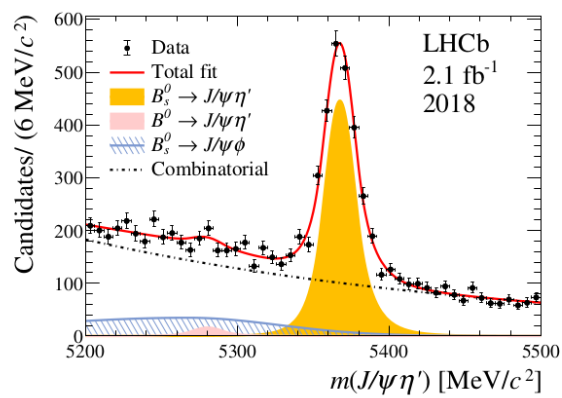
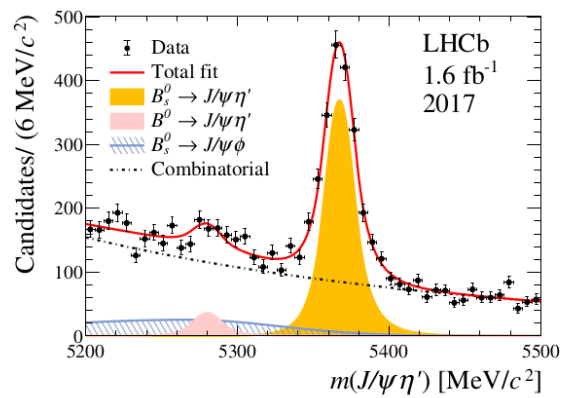
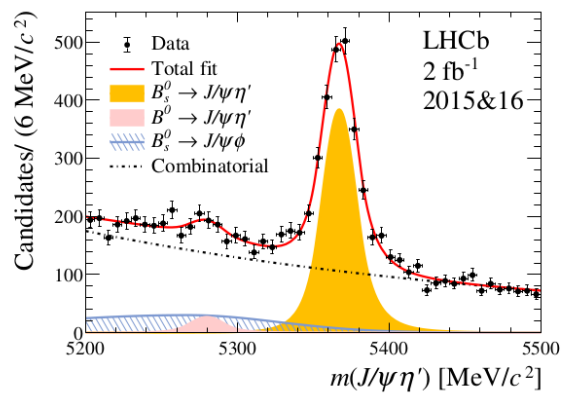
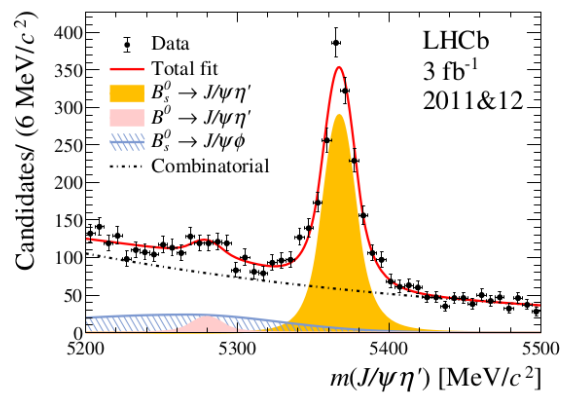
integrating over a time bin

$$R_i = \frac{N_L}{N_H} \propto \frac{[e^{-\Gamma_s t(1+y)}]_{t_1}^{t_2}}{[e^{-\Gamma_s t(1-y)}]_{t_1}^{t_2}} \cdot \frac{(1-y)}{(1+y)} \quad \longrightarrow \quad R_i = A_i \cdot \frac{N_L^{\text{RAW}}}{N_H^{\text{RAW}}}$$

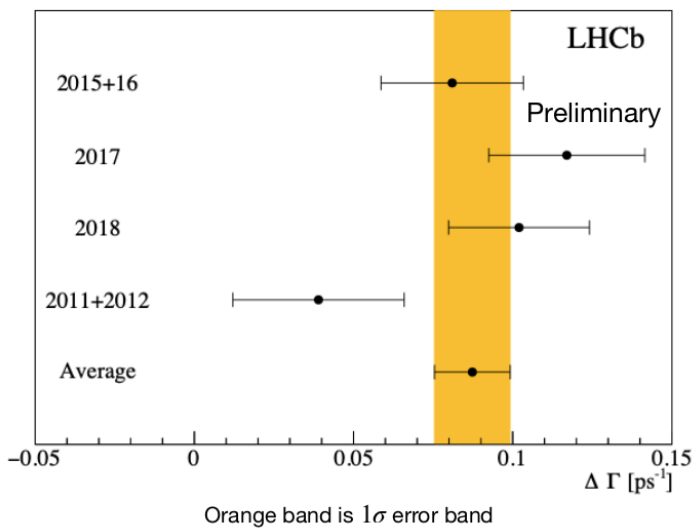
$$y = \Delta\Gamma_s / 2\Gamma_s$$

$N_{L(H)}$ : CP-even(odd) modes

efficiency in each decay time bin



$$\Delta\Gamma_s = 0.087 \pm 0.012 \pm 0.009 \text{ ps}^{-1}$$



- 1<sup>st</sup> measurement using  $\eta'$  channel
- In agreement with the SM

[LHCb-PAPER-2023-025]

# Measurement of $\sin(2\beta)$

- B decays to CP eigenstates allow to probe the mixing phase  $\beta$  through the interference between decays with and without mixing

$$\mathcal{A}^{CP}(t) = \frac{\Gamma(\bar{B}^0(t) \rightarrow f) - \Gamma(B^0(t) \rightarrow f)}{\Gamma(\bar{B}^0(t) \rightarrow f) + \Gamma(B^0(t) \rightarrow f)} = \frac{S \sin(\Delta m_d t) - C \cos(\Delta m_d t)}{\cosh(\frac{1}{2}\Delta\Gamma_d t) + \mathcal{A}_{\Delta\Gamma} \sinh(\frac{1}{2}\Delta\Gamma_d t)}$$

where

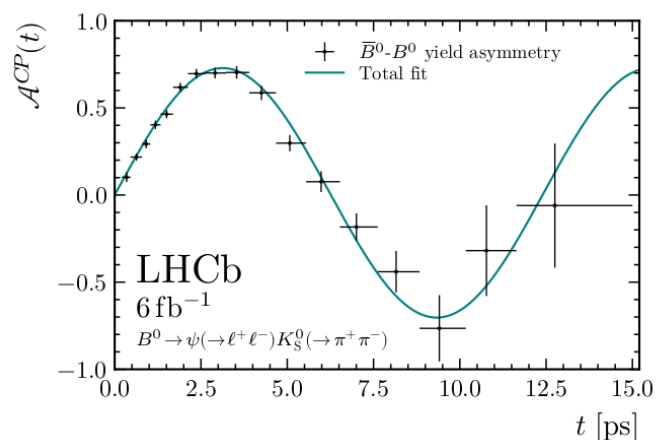
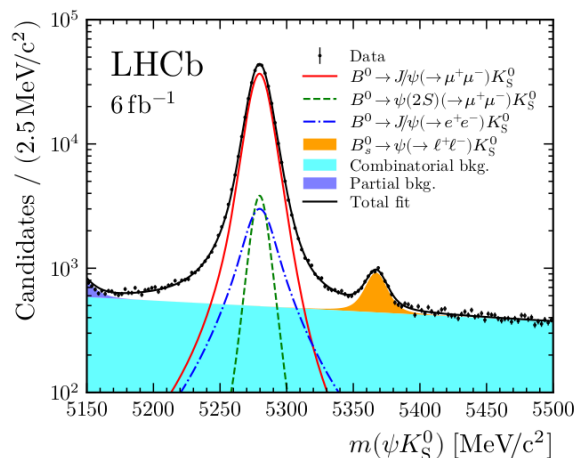
$$S = \sin(2\beta + \Delta\phi_d + \Delta\phi_d^{\text{NP}})$$

- Decay channels:  $B^0_s \rightarrow J/\psi(\rightarrow \mu^+\mu^-)K^0_s$ ,  $B^0_s \rightarrow \psi(2S)(\rightarrow \mu^+\mu^-)K^0_s$ ,  $B^0_s \rightarrow J/\psi(\rightarrow e^+e^-)K^0_s$  with  $K^0_s \rightarrow \pi\pi$

$$P(t, d, \eta) \propto [1 + d(1 - 2\omega^+(\eta))]P_{B^0}(t) + [1 + d(1 - 2\omega^-(\eta))]P_{\bar{B}^0}(t)$$

$$P_{B^0(\bar{B}^0)}(t) \propto \{(1 \mp A_P)(1 \mp \Delta\epsilon_{tag})e^{-\Gamma_d t'}(1 \mp S \sin(\Delta m_d t') \pm C \cos(\Delta m_d t'))\} \otimes R(t - t') \cdot \epsilon(t)$$

- Simultaneous fit of all channels
- Combination run2 and run1 data



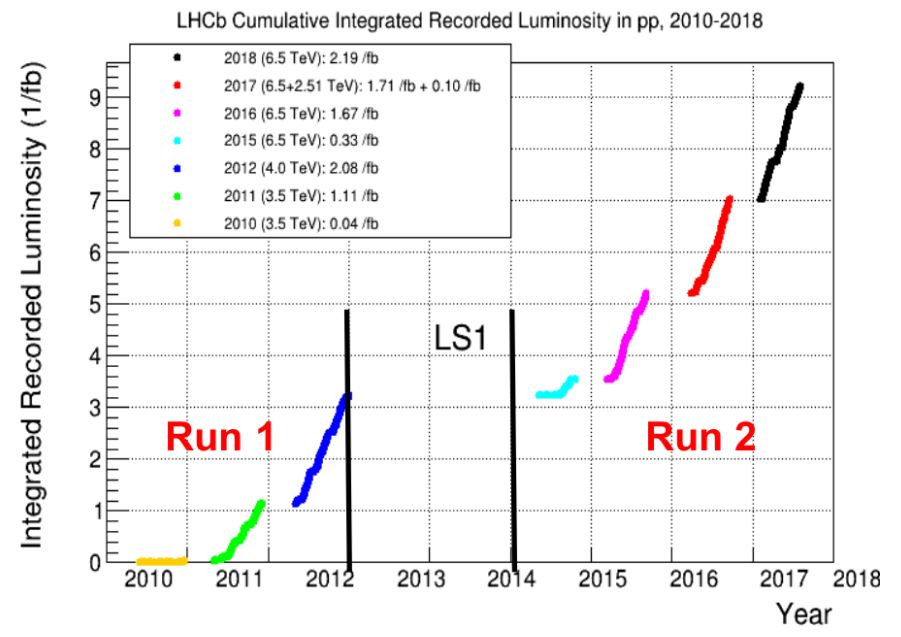
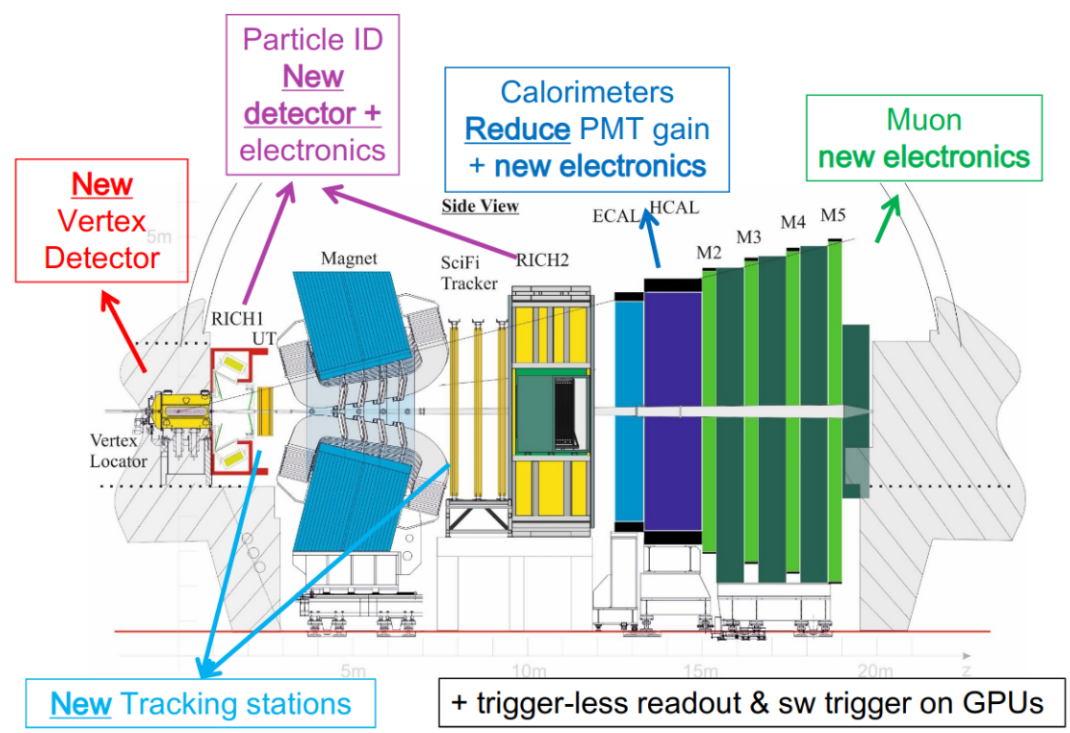
$$S_{J/\psi K_S^0}^{\text{Run 1\&2}} = 0.726 \pm 0.014 \text{ (stat+syst)}$$

$$C_{J/\psi K_S^0}^{\text{Run 1\&2}} = 0.010 \pm 0.012 \text{ (stat+syst)}$$

- Most precise single measurement
- Agreement with CKMfitter predictions

[LHCb-PAPER-2023-013]

- LHCb was originally designed for CP violation and rare beauty & charm decays
- But now it is a general purpose detector: *exotic spectroscopy, EW precision physics, heavy ions, fixed target program...*



Run1: 3 fb<sup>-1</sup> @  $\sqrt{s} = 7-8$  TeV  
Run2: 6 fb<sup>-1</sup> @  $\sqrt{s} = 13$  TeV

- LHCb is a spectrometer in the forward direction ( $2 < \eta < 5$ )
- Excellent vertexing, tracking and particle identification
- Low trigger threshold on hadrons, muons and photons
- Production of all types of *b* and *c* hadrons