





EFT interpretations (ATLAS/CMS) of single-Higgs and Higgs boson pair production

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Standard Model at the LHC - Rome, May 7-10 2024

Introduction & outline of the talk

- No evidence of New Physics at the LHC direct searches of BSM effects continue but focus is shifting to indirect exploration of Higgs sector
 - increasing number of Higgs EFT measurements in CMS and ATLAS
- EFT results focus in interpretation of unfolded spectrum in presence of EFT effects or extract coefficients with dedicated analyses using optimal observables
- This talk will focus on recent EFT interpretations results in ATLAS/CMS for Higgs and double Higgs analyses
 - EFT parametrisation models
 - EFT interpretation using STXS scheme and full combination of Run 2 single Higgs analyses
 - highlights on optimal observable analyses targeting constraints on SMEFT EFT or anomalous coupling basis ($H \rightarrow ZZ$, $H \rightarrow WW$, $H \rightarrow \tau\tau$)
 - extending constraints on EFT coefficients using SMEFT or HEFT basis in double Higgs analyses
 - example of global EFT fit using Higgs and EWK constraints

3

- Experimental profile of the Higgs boson with Run 1/2 data becoming very precise
- Precision measurement is key to look for deviations of SM couplings: achieved using low-energy approximation (EFT) to UV complete theory

$$\mathscr{L}_{SMEFT} = \mathscr{L}_{SM}^{4} + \sum_{i} \frac{c_{i}^{(5)}}{\Lambda} \mathcal{O}_{i}^{(5)} + \sum_{i} \frac{c_{i}^{(6)}}{\Lambda^{2}} \mathcal{O}_{i}^{(6)} + \sum_{i} \frac{c_{i}^{(7)}}{\Lambda^{3}} \mathcal{O}_{i}^{(7)} + \mathcal{O}(8) + \dots = \mathscr{L}_{BSM}$$

Under the assumption that the theory cut-off is much larger than the energy, SMEFT series can be truncated at dimension 6

- dimension 5 operators violates L number conservation hence not relevant for LHC studies
- Warsaw basis often used to derive set of operators which can be accessed and constrained by analyses
- Dimension 8 operators often neglected and contributing as Ι/Λ⁴

Wilson coefficients & EFT Lagrangian expansion

4

Expansion of SM lagrangian in I/Λ : observables EFT effects are parametrised

with linear term in WC's and a linear+quadratic term in WC's (both are dim-6 operators)



SMEFT [link] is a popular model for EFT interpretation using dim-6 operators

Some EFT contributions are CP-odd operators: access on those operators is relevant as non vanishing components indicate CP violation

- Fundamental to keep all operators in interpretation due to correlation effects
- No single measurement constraints all operators need for global approach
 - EFT interpretation of STXS fit using STXS categorisation for Higgs production modes no sensitivity to CP given lack of dedicated CP-sensitive observables





Assumption of EFT interpretation in STXS bins: I) no EFT effects on background components, 2) acceptance corrections in STXS bins to account for EFT effects





7

Combined all Higgs final states and performed principal component analysis on SMEFT Fisher information matrix in the limit of Gaussian STXS measurements

$$V_{\text{SMEFT}}^{-1} = P_{(i,k',X) \rightarrow (j)}^{T} V_{\text{STXS}}^{-1} P_{(i,k',X) \rightarrow (j)}$$

Eigenvalue decomposition of SMEFT Fisher information matrix leads to eigenvectors which are constrained in the measurement

sensitivity criteria observations defining a new rotated basis achieving fit stability, fit-parameter interpretability and no flat directions in the likelihood scans due to poor sensitivity





EFT Higgs fit - Linear model

Linearised SMEFT model (SM p-value 94.5%): contribution of each measured production/decay extracted using Fisher information matrix

Most of WC's are statistically dominated - some WC's entirely constrained by single channels



ATLAS EFT Higgs fit including quadratic terms

- Constraints on the model with quadratic terms in most cases stronger then linear-only model (mostly for weak impact of BSM/SM interference terms)
 - observed uncertainties smaller than expected for quadratic due to double minima structure
 - difference between linear/quadratic magnitude as indicator of missing SM dim-8 interference term ATLAS





- Along with STXS, SMEFT operators can also be constrained from fiducial crosssection measurements (unfolded to particle-level) using pt(H) in H→yy/ZZ
 - fine granularity of kinematic regions in differential observables + large sensitivity to EFT in high pt(H), however inclusive treatment of Higgs production modes (unlike in STXS approach)
 - expected drop in sensitivity using differential approach compared to STXS interpretation
 - probing EFT contact term in production and decay: rotation defines a new set of eigenvectors which are decorrelated and extracted simultaneously



CMS

Anomalous couplings - H->ZZ->41



Constraints HVV using Anomalous Coupling: extended to WC constraints (SMEFT)

Various hypotheses on combined AC fit

- fixing all couplings but
 one to SM expectations
 or all couplings profiled
- access sensitivity to CP structure in HZZ decay





Phys. Rev. D 104 (2021) 052004		Coupling	Observed	Expected
Performing fit to all Wils targeting HZZ couplings modes	on coefficients: using VBF and VH	$c_{H\Box} c_{HD}$	$\begin{array}{r} 0.04\substack{+0.43\\-0.45}\\-0.73\substack{+0.97\\-4.21}\end{array}$	$\begin{array}{c} 0.00\substack{+0.75\\-0.93}\\ 0.00\substack{+1.06\\-4.60}\end{array}$
 both linear and quadratic to 	erms considered	$ \rightarrow c_{HW} $ c_{HWB}	$\begin{array}{c} 0.01\substack{+0.18\\-0.17}\\ 0.01\substack{+0.20\\-0.18}\end{array}$	$\begin{array}{c} 0.00\substack{+0.39\\-0.28}\\ 0.00\substack{+0.42\\-0.31}\end{array}$
Iargest precision for c(HW precision on CP-odd EFT V), also access with good VC	$\xrightarrow{C_{HB}} c_{H\tilde{W}}$	$\begin{array}{c} 0.00\substack{+0.05\\-0.05}\\-0.23\substack{+0.51\\-0.52}\end{array}$	$\begin{array}{c} 0.00\substack{+0.03\\-0.08}\\ 0.00\substack{+1.11\\-1.11}\\\end{array}$
		с _{НŴВ} с _{HB̃}	$-0.25^{+0.56}_{-0.57}\\-0.06^{+0.15}_{-0.16}$	$\begin{array}{r} 0.00\substack{+1.21\\-1.21}\\ 0.00\substack{+0.33\\-0.33}\end{array}$
Also provided constraints for c(ZZ) and CP-odd c(ZZ) coupling components using results on Warsaw basis	10 -	137 fb ⁻¹ (13 TeV) 10 10 8 10 8 10 0 0 0 0 0 0 0 0 0 0 0 0 0	CMS — Observed — Expected — 95% C.L. — 68% C.L — 68% C.L	137 fb ⁻¹ (13 TeV)
		0.0		



EFT combination across channels & operators

- Simultaneous measurement of EFT operators, c(gg),~c(gg), kt,~kt impacting gluonfusion loop - common EFT approach for several channels with additional sensitivity to CP odd operators
 - gluon fusion in addition to ttH/tH ($\rightarrow \gamma \gamma / ZZ$) used to constrain EFT top couplings









EFT constraints and CP tests in H->ZZ

Search for CP violation using SMEFT framework in HZZ(4I) final state

- BSM CP-odd dim-6 couplings constrained using matrix element-based optimal observables
- VBF production mode sensitive to CP structure

• Results provided with freely-floating parameters to account for correlations

measurements in agreement with SM CP-even prediction

arXiv:2304.0961 (submitted to JHEP)





Extending anomalous HVV constraints to H->WW



Constraints on anomalous effects at the HVV and Hgg vertices following AC and SMEFT interpretations

- analysis split in categories targeting gluon fusion, VBF-like and VH-like topologies
- MELA kinematic discriminant: Higgs mode discriminator, SM couplings vs BSM, interference vs SM/BSM

Results provided under
two fitting hypotheses

POI's are fixed/floating

- AC and Higgs SMEFT Warsaw basis
- Significant improvement in sensitivity/analysis coverage compared to full Run I analysis

Coupling	Observed	Expected	
$c_{H\square}$	$-0.76^{+1.43}_{-3.43}$	$0.0^{+1.37}_{-1.84}$	
$c_{\rm HD}$	$-0.12\substack{+0.93 \\ -0.32}$	$0.0\substack{+0.43 \\ -0.30}$	
$c_{\rm HW}$	$0.08\substack{+0.43 \\ -0.87}$	$0.0\substack{+0.37 \\ -0.48}$	
$c_{\rm HWB}$	$0.17\substack{+0.88 \\ -1.79}$	$0.0\substack{+0.77\\-0.96}$	
$c_{\rm HB}$	$0.03\substack{+0.13 \\ -0.26}$	$0.0\substack{+0.11 \\ -0.14}$	
$\mathcal{C}_{\mathrm{H}\tilde{\mathrm{W}}}$	$-0.26\substack{+0.67\\-0.50}$	$0.0\substack{+0.48 \\ -0.52}$	
$c_{\mathrm{H}\tilde{\mathrm{W}}\mathrm{B}}$	$-0.54^{+1.37}_{-1.03}$	$0.0\substack{+0.99 \\ -1.07}$	
$\mathcal{C}_{H ilde{B}}$	$-0.08\substack{+0.20\\-0.15}$	$0.0\substack{+0.15 \\ -0.16}$	



CMS PAS HIG-22-008 (submitted to EPJC)





Parameter value



Targeting measurement of EFT vertices using SMEFT and anomalous couplings

- ▶ <u>VBF production</u>: HVV EFT vertex, <u>ggH production</u>: Hgg EFT vertex
- ► HVV vertex constrained using $H \rightarrow \tau \tau$ decay in VBF production while Hgg vertex uses combination of $H \rightarrow \tau \tau$ and $H \rightarrow ZZ \rightarrow 4I$

Pure CP-odd hypothesis for Higgs couplings to gluons excluded at 2.4σ



EFT in double Higgs analyses





EFT in double Higgs analyses (2)



- ► <u>HH→bbyy</u>: excluded 4 of the considered 7 HEFT $_{B}^{B}$ benchmark points
- $HH \rightarrow bb\tau\tau$: one dimensional constraints on SMEFT/HEFT coefficients
- similar approach for $HH \rightarrow bbbb$ analysis





LHC EFT constraints towards global EFT fit

21



Precision measurements is key to look for deviations on SM couplings: several Effective Field Theory interpretations of Higgs measurements in CMS and ATLAS

- EFT interpretation of STXS/differential results allows to probe EFT parameters using various Higgs production modes
 - EFT effects parametrised in STXS bins and dedicated acceptance corrections in analysis phasespace
- Several dedicated measurements of EFT effects in ATLAS/CMS single and double Higgs analyses using optimal observables to enhance EFT effects
- Developing Principal Component Analyses to tackle large combinations and simultaneous constraints on Wilson coefficients
 - very relevant for global EW+Higgs EFT combination and to select non flat directions in EFT space
 - Important to assess interpretability and compatibility of EFT results using PCA analyses
- Ongoing effort in CMS+ATLAS to provide common STXS+SMEFT parameterisation in the context of the LHC EFT WG [LHC EFT WG prediction note]
- Several more EFT interpretation using Run 2 results will be released soon stay tuned!

Additional slides



- Tested difference in central values and uncertainty employing full likelihoods and multivariate Gaussian approximations for EFT combinations
 - STXS interpretation model combined likelihood functions based on the STXS results
 - all non-Gaussian effects accounted for (experimental/theory uncertainty models)
 - complexity of the full likelihood combination
 - combined likelihoods based on simplified approach straightforward to implement

Generally good agreement in expected/observed parameter estimation for linear-only and linear+quadratic EFT extraction



Wilson coefficients & EFT Lagrangian expansion

25

b

Η

Z/W

Z/W

q

 \bar{q}

Example of $VH \rightarrow bb$ channel

Operator	Wilson coefficient	Lagrangian modification	Channels
${\cal O}^{(1)}_{Hq}=iH^\dagger \overleftrightarrow{D}_\mu H \overline{q} \gamma^\mu q$	cHj1	qqV vertex, HVqq contact term	ZII
${\cal O}^{(3)}_{Hq}=iH^{\dagger}\sigma^{i}\overleftrightarrow{D}_{\mu}H\overline{q}\sigma^{i}\gamma^{\mu}q$	cHj3	qqV vertex, HVqq contact term	ZII, WIv
${\cal O}_{Hu}=iH^\dagger \overleftarrow{D}_\mu H \overline{u}_R \gamma^\mu u_R$	cHu	qqV vertex, HVqq contact term	ZII
${\cal O}_{Hd}=iH^\dagger \overleftarrow{D}_\mu H \overline{d}_R \gamma^\mu d_R$	cHd	qqV vertex, HVqq contact term	ZII
${\cal O}_{HW}=H^{\dagger}HW^{i}_{\mu u}W^{i\mu u}$	cHW	HVV vertex	ZII, WIv
${\cal O}_{H ilde W} = H^\dagger H ilde W^i_{\mu u} W^{i \mu u}$	cHWtil	HVV vertex	ZII, WIv
${\cal O}_{HB}=H^{\dagger}HB_{\mu u}B^{\mu u}$	cHB	HVV vertex	ZII
${\cal O}_{H {\tilde B}} = H^\dagger H {\tilde B}_{\mu u} B^{\mu u}$	cHBtil	HVV vertex	ZII
${\cal O}_{HWB}=H^{\dagger}\sigma^{i}HW^{i}_{\mu u}B^{\mu u}$	cHWB	HVV vertex, Wlv vertex	ZII
${\cal O}_{H ilde W B} = H^\dagger \sigma^i H ilde W^i_{\mu u} B^{\mu u}$	cHWBtil	HVV vertex, Wlv vertex	ZII
${\cal O}_{H\square} = (H^\dagger H) \square (H^\dagger H)$	cHbox	HVV vertex, hbb coupling	ZII, WIv
${\cal O}_{HD} = (D^\mu H^\dagger H) (H^\dagger D_\mu H)$	cHDD	HVV vertex, hbb coupling, qqV vertex	ZII, WIv
${\cal O}_{bH}=(H^{\dagger}H)(\overline{q}bH)$	cbHRe + cbHIm	hbb coupling	ZII, WIv



arXiv:2402.05742 (submitted to JHEP)



EFT interpretation using STXS

Constraints on main WC's in STXS bins affecting following vertices

• EW+Higgs boson interactions, boson couplings to fermions and 4-fermion interactions

EW+Higgs interactions	Boson couplings to fermions		4-fermion interactions	
Wilson coefficient	Operator	Wilson coefficient	Operator	
$c_{H\square}$	$(H^\dagger H) \square (H^\dagger H)$	C _{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{H} G^A_{\mu\nu}$	
C _{HDD}	$\left(H^{\dagger}D^{\mu}H ight)^{*}\left(H^{\dagger}D_{\mu}H ight)$	C _{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \widetilde{H} W^I_{\mu\nu}$	
c _{HG}	$H^{\dagger}HG^{A}_{\mu u}G^{A\mu u}$	C _{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \widetilde{H} B_{\mu\nu}$	
c _{HB}	$H^{\dagger}H B_{\mu u}B^{\mu u}$		$(\bar{l}_p \gamma_\mu l_t)(\bar{l}_r \gamma^\mu l_s)$	
C _{HW}	$H^{\dagger}HW^{I}_{\mu u}W^{I\mu u}$	$c_{qq}^{\scriptscriptstyle (1)}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{q}_r \gamma^\mu q_s)$	
C _{HWB}	$H^{\dagger} au^{I} H W^{I}_{\mu u} B^{\mu u}$	$c_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r) (\bar{q}_s \gamma^\mu \tau^I q_t)$	
C _{eH}	$(H^{\dagger}H)(l_{p}e_{r}H)$	c _{qq}	$(\bar{q}_p \gamma_\mu q_t) (\bar{q}_r \gamma^\mu q_s)$	
c _{uH}	$(H^{\dagger}H)(\bar{q}_{p}u_{r}H)$	$c_{oldsymbol{q}oldsymbol{q}}^{\scriptscriptstyle{(31)}}$	$(\bar{q}_p \gamma_\mu \tau^I q_t) (\bar{q}_r \gamma^\mu \tau^I q_s)$	
C _{dH}	$(H^{\dagger}H)(\bar{q}_{p}d_{r}\widetilde{H})$	c _{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	
$c_{Hl}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}_{p}\gamma^{\mu}l_{r})$	$c_{uu}^{(1)}$	$(\bar{u}_p \gamma_\mu u_t)(\bar{u}_r \gamma^\mu u_s)$	
$c_{Hl}^{\scriptscriptstyle (3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}_{p}\tau^{I}\gamma^{\mu}l_{r})$	$c_{qu}^{\scriptscriptstyle (1)}$	$(\bar{q}_p \gamma_\mu q_t)(\bar{u}_r \gamma^\mu u_s)$	
c_{He}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}_{p}\gamma^{\mu}e_{r})$	$c_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r) (\bar{d}_s \gamma^\mu T^A d_t)$	
$c_{Hq}^{\scriptscriptstyle (1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{q}_{p}\gamma^{\mu}q_{r})$	$c_{oldsymbol{qu}}^{\scriptscriptstyle{(8)}}$	$(\bar{q}_p \gamma_\mu T^A q_r) (\bar{u}_s \gamma^\mu T^A u_t)$	
$c_{Hq}^{\scriptscriptstyle{(3)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{q}_{p}\tau^{I}\gamma^{\mu}q_{r})$	$c_{qd}^{\scriptscriptstyle (8)}$	$(\bar{q}_p \gamma_\mu T^A q_r) (\bar{d}_s \gamma^\mu T^A d_t)$	
c_{Hu}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{u}_{p}\gamma^{\mu}u_{r})$	c _W	$\epsilon^{IJK}W^{I\nu}_{\mu}W^{J ho}_{\nu}W^{K\mu}_{ ho}$	
c_{Hd}	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{d}_{p}\gamma^{\mu}d_{r})$	c_G	$f^{ABC}G^{A\nu}_{\mu}G^{B ho}_{\nu}G^{C\mu}_{ ho}$	



H->ZZ ATLAS





H->ZZ ATLAS







CP violation in EtH/EH

JHEP 07 (2023) 092

Targeting HWW and Hπ final states - CP even/odd operators

BDT sensitive to CP observables (ΔR_{jj} , $\Delta \eta_{jj}$, ...)

Results including combination of multi-lepton and yy/ZZ

- f(CP)=0 SM expectation; f(CP)=0.28; <0.55 at 68% CL interval
 - results compatible to CP even scenario: CP odd contribution not favoured at 3.7σ





EFT basis definition

$$f_{a3}^{ggH} = \frac{|a_3^{gg}|^2}{|a_2^{gg}|^2 + |a_3^{gg}|^2} \operatorname{sign}\left(\frac{a_3^{gg}}{a_2^{gg}}\right)$$

Hff couplings - CP-even

$$|f_{CP}^{Hff}| = \left(1 + 2.38 \left[\frac{1}{|f_{a_3}^{ggH}|}\right]\right)^{-1} = \sin^2 \alpha^{Hff}$$

Hff couplings - CP-odd

$$f_{a_3} = \frac{|a_3^{gg}|^2}{|a_2^{gg}|^2 + |a_3^{gg}|^2} \operatorname{sign}\left(\frac{a_3^{gg}}{a_2^{gg}}\right)$$

HVV couplings - gg couplings (only non zero contributions are a2 and a3)

$$\begin{split} c_{zz} &= -\frac{s_w^2 c_w^2}{2\pi \alpha} a_2, \\ \tilde{c}_{zz} &= -\frac{s_w^2 c_w^2}{2\pi \alpha} a_3. \end{split}$$

$$\begin{split} c_{gg} &= -\frac{1}{2\pi \alpha_S} a_2^{gg}, \\ \tilde{c}_{gg} &= -\frac{1}{2\pi \alpha_S} a_3^{gg}, \end{split}$$
EFT interpretation



ATLAS HH4b EFT interpretation





Phys. Rev. D 108, 052003







HEL parameters

HEL Parameters	Definition	Others profiled	Fix others to SM
$c_A imes 10^4$	$c_A = \frac{m_W^2}{g'^2} \frac{f_A}{\Lambda^2}$	$-1.03^{+1.53}_{-1.59} onumber \ \left(egin{array}{c} +1.59 \ -1.56 \end{array} ight)$	$-0.78^{+1.11}_{-1.16} \ \left(egin{array}{c} +1.10 \ -1.11 \end{array} ight)$
$c_G imes 10^5$	$c_G = \frac{m_W^2}{g_s^2} \frac{f_G}{\Lambda^2}$	$1.43^{+3.20}_{-3.00}$ $\begin{pmatrix} +3.13\\ -2.74 \end{pmatrix}$	$0.27^{+1.05}_{-1.05} \\ \begin{pmatrix} +1.03 \\ -1.01 \end{pmatrix}$
$c_u \times 10$	$c_u = -v^2 \frac{f_u}{\Lambda^2}$	$0.68^{+0.82}_{-0.83}$ $\begin{pmatrix} +0.83\\ -0.79 \end{pmatrix}$	$0.43^{+0.69}_{-0.69} \\ \left(\begin{smallmatrix} +0.68 \\ -0.67 \end{smallmatrix} \right)$
$c_d \times 10$	$c_d = -v^2 \frac{f_d}{\Lambda^2}$	$\begin{array}{c} 0.59^{+1.03}_{-1.13} \\ \left(\substack{+1.08 \\ -1.05} \right) \end{array}$	$-0.01^{+0.31}_{-0.28}$ $\begin{pmatrix} +0.30\\ -0.28 \end{pmatrix}$
$c_\ell imes 10$	$c_\ell = - v^2 rac{f_\ell}{\Lambda^2}$	$-0.57^{+0.74}_{-0.73}$ $\begin{pmatrix} +0.72\\ -0.77 \end{pmatrix}$	$-0.75^{+0.60}_{-0.64} onumber {}^{+0.60}_{-0.60}$
$c_{HW} imes 10^2$	$c_{HW} = rac{m_W^2}{2g} rac{f_{HW}}{\Lambda^2}$	$-1.45^{+4.72}_{-3.03}$ $\begin{pmatrix} +3.93\\ -3.27 \end{pmatrix}$	$0.77^{+0.84}_{-1.20} onumber {}^{+1.04}_{-1.38}$
$(c_{WW} - c_B) \times 10^2$	$c_{WW} = \frac{m_W^2}{g} \frac{f_{WW}}{\Lambda^2}, c_B = \frac{2m_W^2}{g'} \frac{f_B}{\Lambda^2}$	$2.16^{+2.84}_{-5.35}$ $\begin{pmatrix} +3.46\\ -5.00 \end{pmatrix}$	$\begin{array}{c} 0.62^{+1.06}_{-1.22} \\ \left(\substack{+1.09 \\ -1.23} \right) \end{array}$



HH->WW88 SM 95% CL exclusion limits

