# The Non-Linear Early Universe



# The Numerical Early Universe



## Lattice Techniques in Cosmology





### The Art of Simulating the Early Universe

### The Art of Simulating the Early Universe (When do we need

to simulate it ?)

### The Art of Simulating the Early Universe

When things get complicated: non-linear, strong coupling, non-perturbative, etc















#### Non-linear field dynamics

Curvature Fluctuations

Cosmic Defects



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Cosmic Defects

























~ 10-25 yrs













~ 5-10 yrs



~ 5-10 yrs



~ 10-25 yrs



~ 10-25 yrs



~ 3 yrs



~ 3 yrs [+ upcoming]
## **The Early Universe**



~ 3 yrs



Code Manual: arXiv: 2102.01031 (+100 pages)



Simulates scalar-gauge field dynamics [w. self-consistent expanding background]
[U(1) x SU(2)]

Links & plaquettes (~ lattice-QCD)





- Simulates scalar-gauge field dynamics [w. self-consistent expanding background]
- Written in C++, with modular structure separating <u>physics</u> (CosmoInterface library) and <u>technical details</u> (TempLat library).



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- Written in C++, with modular structure separating <u>physics</u> (CosmoInterface library) and <u>technical details</u> (TempLat library).
- > Parallellized in multiple spatial dimensions (but you write in serial !)
- ► Family of evolution algorithms, accuracy ranging from  $\delta O(\delta t^2) \delta O(\delta t^{10})$ [LeapFrog, Verlet, Runge-Kutta, Yoshida, ...]



Code Manual: arXiv: 2102.01031

## http://www.cosmolattice.net/



Code Manual: arXiv: 2102.01031

## http://www.cosmolattice.net/



A modern code for lattice simulations of scalar and gauge field dynamics in an expanding universe

Cosmo Lattice – Default Field Content

#### ➤ Matter content:

$$S = -\int d^4x \sqrt{-g} \left\{ \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + (D^A_\mu \varphi)^* (D^\mu_A \varphi) + \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + (D_\mu \Phi)^\dagger (D^\mu \Phi) + \frac{1}{2} \text{Tr} \{ G_{\mu\nu} G^{\mu\nu} \} + V(\phi, |\varphi|, |\Phi|) \right\}$$

**Cosmo Lattice** – Default Field Content

Matter content:



SU(2) gauge sector

**Cosmo Lattice** – Default Field Content

Matter content:



Background Metric:

$$ds^{2} = -dt^{2} + a^{2}(t) \,\delta_{ij} \,dx^{i} dx^{j} \,\langle$$

Self-consistent expansion (Friedmann equations)
 Fixed power-law background a(t) ~ t<sup>2</sup>/<sub>3(1+w)</sub>

**Cosmo Lattice** – Default Field Content

Matter content:



Background Metric:

$$ds^{2} = -dt^{2} + a^{2}(t) \,\delta_{ij} \,dx^{i} dx^{j} \begin{cases} \blacktriangleright \text{ Self-consistent expansion (Friedmann equations)} \\ \blacktriangleright \text{ Fixed power-law background } a(t) \sim t^{\frac{2}{3(1+w)}} \end{cases}$$

**Cosmo Lattice** – Equations of Motion

Hamiltonian scheme: coupled first-order differential equations



**Cosmo Lattice** – Equations of Motion

Hamiltonian scheme: coupled first-order differential equations



Scalar Fields and momenta are defined in the lattice sites



**Cosmo Lattice** – Equations of Motion

Hamiltonian scheme: coupled first-order differential equations



Scalar Fields and momenta are defined in the lattice sites



Gauge fields introduced via links and plaquettes (like in lattice-QCD)



Cosmo Lattice – Expansion Evolution

### > Algorithms use **second Friedmann equation** to **evolve the scale factor**.

> The **first Friedmann equation** is used to check the accuracy of the simulation.

$$\frac{a''}{a} = \frac{a^{2\alpha}}{3m_p^2} \left\langle (\alpha - 2)(K_{\phi} + K_{\varphi} + K_{\Phi}) + \alpha(G_{\phi} + G_{\varphi} + G_{\Phi}) + (\alpha + 1)V + (\alpha - 1)(K_{U(1)} + G_{U(1)} + K_{SU(2)} + G_{SU(2)}) \right\rangle$$

$$\left(\frac{a'}{a}\right)^2 = \frac{a^{2\alpha}}{3m_p^2} \left\langle K_{\phi} + K_{\varphi} + K_{\Phi} + G_{\phi} + G_{\varphi} + G_{\Phi} + K_{U(1)} + G_{U(1)} + K_{SU(2)} + G_{SU(2)} + V \right\rangle$$

 $\langle ... \rangle$  represents volume averaging

Cosmo Lattice – Expansion Evolution

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\_ \_ \_ \_ \_ \_ \_ \_ \_

 $\langle ... \rangle$  represents volume averaging

$$K_{\phi} = \frac{1}{2a^{2\alpha}} {\phi'}^2 \qquad G_{\phi} = \frac{1}{2a^2} \sum_i (\partial_i \phi)^2 \qquad K_{U(1)} = \frac{1}{2a^{2+2\alpha}} \sum_i F_{0i}^2 K_{\varphi} = \frac{1}{a^{2\alpha}} (D_0^A \varphi)^* (D_0^A \varphi) ; \qquad G_{\varphi} = \frac{1}{a^2} \sum_i (D_i^A \varphi)^* (D_i^A \varphi) ; \qquad K_{SU(2)} = \frac{1}{2a^{2+2\alpha}} \sum_{a,i} (G_{0i}^a)^2 K_{\Phi} = \frac{1}{a^{2\alpha}} (D_0 \Phi)^{\dagger} (D_0 \Phi) \qquad G_{\Phi} = \frac{1}{a^2} \sum_i (D_i \Phi)^{\dagger} (D_i \Phi) \qquad G_{SU(2)} = \frac{1}{2a^4} \sum_{i,j < i} F_{ij}^2 K_{ij}^2 K_{ij}^$$

**Cosmo Lattice** – Expansion Evolution

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$$\frac{a''}{a} = \frac{a^{2\alpha}}{3m_p^2} \left\langle (\alpha - 2)(K_{\phi} + K_{\varphi} + K_{\Phi}) + \alpha(G_{\phi} + G_{\varphi} + G_{\Phi}) + (\alpha + 1)V + (\alpha - 1)(K_{U(1)} + G_{U(1)} + K_{SU(2)} + G_{SU(2)}) \right\rangle$$

$$\left(\frac{a'}{a}\right)^2 = \frac{a^{2\alpha}}{3m_p^2} \left\langle K_{\phi} + K_{\varphi} + K_{\Phi} + G_{\phi} + G_{\varphi} + G_{\Phi} + K_{U(1)} + G_{U(1)} + K_{SU(2)} + G_{SU(2)} + V \right\rangle$$

$$\left\langle \dots \right\rangle \text{ represents volume averaging}$$

 $K_{\phi} = \frac{1}{2a^{2\alpha}} {\phi'}^2 \qquad G_{\phi} = \frac{1}{2a^2} \sum_i (\partial_i \phi)^2 \qquad K_{U(1)} = \frac{1}{2a^{2+2\alpha}} \sum_i F_{0i}^2 \\ K_{\varphi} = \frac{1}{a^{2\alpha}} (D_0^A \varphi)^* (D_0^A \varphi) ; \qquad G_{\varphi} = \frac{1}{a^2} \sum_i (D_i^A \varphi)^* (D_i^A \varphi) ; \qquad K_{SU(2)} = \frac{1}{2a^{2+2\alpha}} \sum_{a,i} (G_{0i}^a)^2 \\ K_{\Phi} = \frac{1}{a^{2\alpha}} (D_0 \Phi)^{\dagger} (D_0 \Phi) \qquad G_{\Phi} = \frac{1}{a^2} \sum_i (D_i \Phi)^{\dagger} (D_i \Phi) \qquad G_{SU(2)} = \frac{1}{2a^4} \sum_{a,i,j < i} F_{ij}^2 \\ (\text{Kinetic-Scalar}) \qquad (\text{Gradient-Scalar}) \qquad (\text{Electric \& Magnetic})$ 

# Cosmo Lattice – Output / Observables



# Cosmo Lattice – Output / Observables





# Cosmo Lattice – Output / Observables





### In summary ...



\* Init Conditions\* Eqs. of Motion

Field Objects
Field Algebra





\* Init Conditions\* Eqs. of Motion



\* Choose Lattice: dt, N, dx

- \* Choose Algorithm  $\mathcal{O}(\delta t^n)$
- \* Choose Param: g, m, ...
- \* Choose Observables







\* Init Conditions\* Eqs. of Motion



- \* Choose Lattice: dt, N, dx
- \* Choose Algorithm  $\mathcal{O}(\delta t^n)$
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### **Algorithms**

- Staggered LeapFrog (LF)
- Position-Verlet (PV2)
- Velocity-Verlet (VV2)
- Runge-Kutta (RK2, RK3, RK4)
- Yoshida (VV4, VV6, VV8, VV10)



PosmoLattice

\* Init Conditions\* Eqs. of Motion



\* Choose Lattice: dt, N, dx

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$$\lambda_1, \lambda_2, \dots, g_1^2, g_2^2, \dots$$
  
 $m_{\phi}^2, m_{\psi}^2, \dots, v^2, \Phi_*, \dots$ 

#Output
<pre>outputfile = ./</pre>
#Evolution
expansion = true
evolver = VV2
#Lattice
N = 32
dt = 0.01
kIR = 0.75
nBinsSpectra = 55
#Times
tOutputFreq = 0.1
toutputIntreq = 1
tmax = 300
#TC
#IC VCu+Off - 1 75
initial amplitudes - 7 /2675e18 0 # homogeneous amplitudes in GeV
initial momenta = $-6.2060e30.0$ # homogeneous amplitudes in GeV2
initiat_momenta = 012505050 0 # nomogeneous amptitudes in devz
#Model Parameters
lambda = 9e-14

(no need to re-compile !)



Cosmo Lattice

\* Init Conditions\* Eqs. of Motion



\* Choose Lattice: dt, N, dx

\* Choose Algorithm  $\mathcal{O}(\delta t^n)$ 

\* Choose Param: g, m, ...

\* Choose Observables

Output







### CL is a platform for field theories You choose the problem to solve !



Posmo Lattice

\* Choose Lattice: dt, N, dx

**Field Th. Problem** 

\* Init C^

roblen

- \* Choose Algorithm  $\mathcal{O}(\delta t^n)$
- \* Choose Param: g, m, ...
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### CL is a platform for field theories You choose the problem to solve !



Cosmo Lattice



- \* Choose Lattice: dt, N, dx
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### Output



### > CL so far (v1.0, Public):

- Global scalar field dynamics
- ► U(1) scalar-gauge dynamics
- ► SU(2) scalar-gauge dynamics



Cosmo Lattice

- \* Choose Lattice: dt, N, dx
- \* Choose Algorithm  $\mathcal{O}(\delta t^n)$
- \* Choose Param: g, m, ...
- \* Choose Observables

### Output



### > CL so far (v1.0, Public):

**Field Th. Problem** 

Vew Problem

- Global scalar field dynamics
- ► U(1) scalar-gauge dynamics
- ► SU(2) scalar-gauge dynamics

- > CL update (v2.0, to be released by ~2024)
  - Gravitational waves

$$\Box h_{ij} = 2\Pi_{ij}^{\mathrm{TT}}$$

- > Axion-like couplings  $\phi F_{\mu\nu} \tilde{F}^{\mu\nu}$
- ► Non-minimal coupling  $\xi \phi^2 R$
- Cosmic String Networks



Cosmo Lattice

- \* Choose Lattice: dt, N, dx
- \* Choose Algorithm  $\mathcal{O}(\delta t^n)$
- \* Choose Param: g, m, ...
- \* Choose Observables

### Output



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**Field Th. Problem** 

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- Global scalar field dynamics
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CL update (v2.0, to be released by ~2024)
 Gravitational waves □ h<sub>ij</sub> = 2Π<sup>TT</sup><sub>ij</sub>
 Axion-like couplings φF<sub>µν</sub>I Released in 2022/23 !
 Non-minimal coupling ξφ<sup>2</sup>
 Cosmic String Networks

## New Modules

- \* Magneto Hydro-dynamics (MHD)
- \* Axion-gauge interactions
- \* Cosmic string networks
- \* Non-minimal Grav. coupling

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## Magneto Hydro-dynamics (MHD)

$$T^{\mu\nu} = (p+\rho)U^{\mu}U^{\nu} - pg^{\mu\nu}$$

$$D_{\nu}T^{\mu\nu} = \partial_{\nu}T^{\mu\nu} + \Gamma^{\mu}_{\sigma\nu}T^{\sigma\nu} + \Gamma^{\nu}_{\nu\sigma}T^{\mu\sigma} = \frac{1}{\sqrt{-g}}\partial_{\nu}(\sqrt{-g}\ T^{\mu\nu}) = 0,$$

$$\partial_{\eta} \tilde{T}^{00} + \partial_{i} \tilde{T}^{0i} = \tilde{S}^{0} [\phi, A_{k}, \{\tilde{T}_{lk}\}], \partial_{\eta} \tilde{T}^{0i} + \partial_{j} \tilde{T}^{ij} = \tilde{S}^{i} [\phi, A_{k}, \{\tilde{T}_{lk}\}],$$

### Work in progress ... key to GWs from PhT's !

(w/ K. Marschall, A. Midiri, and A. Roper Pol)

## New Physics

- \* Magneto Hydro-dynamics (MHD)
- \* Axion-gauge interactions
- \* Cosmic string networks
- \* Non-minimal Grav. coupling
#### **Axion-gauge interactions**

$$\mathcal{S}_{ax} = -\int dx^4 \sqrt{-g} \left\{ \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + V(\phi) + \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{\alpha_\Lambda}{4} \frac{\phi}{m_p} F_{\mu\nu} \tilde{F}^{\mu\nu} \right\}$$

#### Used in Phys. Rev. Lett. 131 (2023) 15, 151003 (Topic I)

# New Physics

- \* Magneto Hydro-dynamics (MHD)
- \* Axion-gauge interactions
- \* Cosmic defect networks
- \* Non-minimal Grav. coupling

#### **Cosmic Defect Networks**

## **Particle & GW Emission**



Used in arXiv:2308.08456 ; Phys. Rev. Lett. (submitted) (Topic II)

# New Physics

- \* Magneto Hydro-dynamics (MHD)
- \* Axion-gauge interactions
- \* Cosmic string networks
- \* Non-minimal Grav. coupling

#### Non-minimal Grav. coupling

$$\mathcal{S}_{\rm NMC} = -\int d^4x \sqrt{-g} \left( \frac{1}{2} \xi R \phi^2 + \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + V(\phi, \{\varphi_{\rm m}\}) \right)$$

$$\begin{split} & [\text{Non-minimally coupled}] \left\{ \begin{array}{l} \phi' = a^{\alpha - 3} \pi_{\phi} \,, \\ & \pi'_{\phi} = a^{1 + \alpha} \, \nabla^2 \phi - a^{3 + \alpha} \left( \xi R \phi + V_{,\phi} \right) \,, \\ & \\ & \\ & \left[ \text{Expanding background} \right] \left\{ \begin{array}{l} a' = a^{\alpha - 1} \pi_a \,, \\ & \pi'_a = \frac{a^{2 + \alpha}}{6} R \,, \end{array} \right. \end{split}$$

with

 $R = \frac{1}{m_p^2} \left[ \frac{2\left(1 - 6\xi\right) \left(E_G^{\phi} - E_K^{\phi}\right) + 4\langle V \rangle - 6\xi \langle \phi V_{,\phi} \rangle + \left(\rho_{\rm m} - 3p_{\rm m}\right)}{1 + \left(6\xi - 1\right) \xi \langle \phi^2 \rangle / m_p^2} \right] \,,$ 

(w/ B. Stefanek, T. Opferkuch, and A. Florio)

# New Physics

- \* Magneto Hydro-dynamics (MHD)
- \* Axion-gauge interactions
- \* Cosmic string networks
- \* Non-minimal Grav. coupling



# New Physics

- \* Magneto Hydro-dynamics (MHD)
- \* Axion-gauge interactions
- \* Cosmic string networks
- \* Non-minimal Grav. coupling
- \* Grav. Pert. Th / Full GR

- 1) Non-linear inflation dynamics
- 2) GW from non-linear dynamics
- 3) Preheating & Equation of State after inflation
- 4) Cosmic string networks (axions, AH, ...)
- 5) Single string loop dynamics
- 6) Non-minimal gravitational Interactions
- 7) Phase transitions
- X) Your project !

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- 2) GW from non-linear dynamics
- 3) Preheating & Equation of State after inflation
- 4) Cosmic string networks (axions, AH, ...)
- 5) Single string loop dynamics
- 6) Non-minimal gravitational Interactions
- 7) Phase transitions
- X) Your project !

**I)** Non-linear inflation dynamics

**I)** Single string loop dynamics

- **I)** Non-linear inflation dynamics
- **II)** Single string loop dynamics



I) Non-linear inflation dynamics (e.g Axion-inflation)
 II) Single string loop dynamics

I) Non-linear inflation dynamics (e.g Axion-inflation)
 II) Single string loop dynamics (If time permits)

**Example I** 

(Non-Linear) Field Dynamics of Axion Inflation

#### **Strong Backreaction Regime in Axion Inflation**

Daniel G. Figueroa<sup>®</sup>,<sup>1,\*</sup> Joanes Lizarraga<sup>®</sup>,<sup>2,3,†</sup> Ander Urio<sup>®</sup>,<sup>2,3,‡</sup> and Jon Urrestilla<sup>®</sup>,<sup>2,3,§</sup> <sup>1</sup>Instituto de Física Corpuscular (IFIC), Consejo Superior de Investigaciones Científicas (CSIC) and Universitat de València, 46980, Valencia, Spain <sup>2</sup>Department of Physics, University of Basque Country, UPV/EHU, 48080, Bilbao, Spain

<sup>3</sup>EHU Quantum Center, University of the Basque Country UPV/EHU, 48940, Leioa Biscay, Spain

(Received 29 April 2023; accepted 8 September 2023; published 13 October 2023)

We study the nonlinear dynamics of axion inflation, capturing for the first time the inhomogeneity and full dynamical range during strong backreaction, till the end of inflation. Accounting for inhomogeneous effects leads to a number of new relevant results, compared to spatially homogeneous studies: (i) the number of extra efoldings beyond slow-roll inflation increases very rapidly with the coupling, (ii) oscillations of the inflaton velocity are attenuated, (iii) the tachyonic gauge field helicity spectrum is smoothed out (i.e., the spectral oscillatory features disappear), broadened, and shifted to smaller scales, and (iv) the nontachyonic helicity is excited, reducing the chiral asymmetry, now scale dependent. Our results are expected to impact strongly on the phenomenology and observability of axion inflation, including gravitational wave generation and primordial black hole production.

DOI: 10.1103/PhysRevLett.131.151003

(e-Print: 2303.17436 [astro-ph.CO])

#### **Strong Backreaction Regime in Axion Inflation**

Daniel G. Figueroa,<sup>1,\*</sup> Joanes Lizarraga,<sup>2,3,†</sup> Ander Urio,<sup>2,3,‡</sup> and Jon Urrestilla,<sup>2,3,§</sup> <sup>1</sup>Instituto de Física Corpuscular (IFIC), Consejo Superior de Investigaciones Científicas (CSIC) and Universitat de València, 46980, Valencia, Spain <sup>2</sup>Department of Physics, University of Basque Country, UPV/EHU, 48080, Bilbao, Spain <sup>3</sup>EHU Quantum Center, University of the Basque Country UPV/EHU, 48940, Leioa Biscay, Spain (Received 29 April 2023: accented & September 2023: published 13 October 2023)

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#### + Nico Loayza

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self-couplings the aquantage that small inflaton. const. <sup>Infla</sup> Smallness of Lishift t  $\mu \nu$  (review Pajer, MP teomytheayy  $\rightarrow \phi + G_n \otimes n_5 C \circ \mu p | i n g_s to other strates as the strates of the strate$ E COMPANY CONTRACTOR CONTRACTOR lings to other (ields/) V shut() atter (predictivit  $10^{16}$  GeV  $\sim 10^{16}$  GeV ekvala attopotentia · Flatness and gatus atmess sands grades lagaity  $\psi^{\nu_p} = \forall 1 \text{ GeV}^*$   $\psi^{\nu_p} = \forall 1 \text{ ocal NG}^{\alpha}$   $\psi^{\nu_p} = \psi^{\nu_p} = \forall 1 \text{ ocal NG}^{\alpha}$   $\psi^{\nu_p} = \psi^{\nu_p} = \psi^{\nu_p} = \psi^{\nu_p}$   $\psi^{\nu_p} = \psi^{\nu_p} = \psi^{\nu_p} = \psi^{\nu_p} = \psi^{\nu_p}$   $\psi^{\nu_p} = \psi^{\nu_p} =$ self-couplings Agreement with standard single field slow (Natural) Inflation Out a solution of the second state of the seco  $= (\phi, \phi, \phi, \phi) = (\phi, \phi, \phi) = (\phi, \phi)$ Freese Sheftarsymmetersy .on2couplings to oth Inflaton perturbations  $\delta \phi$  $s_{\rm Mift} \propto V_{\rm through inverse decay}$ ift sypaner, where the strain shire shire share the strain shire s elds <u>Liction ty</u> (highly non-Gaussian) se<sup>v</sup> <sup>∉</sup>Frienna<sub>f</sub>r2, Vanta of that  $V_{shift} = \frac{M_p^2}{2} \left( \frac{V_{,\phi}}{V} \right)$  $\ll 1$ ,  $\sqrt[\eta]{\eta} = MOt_V the OCD axion; reference values <math>f$ Freese, Frieman, Olinto 90; steen wighted high natural. A wordany, Roiffrons reheasing an (review Pajer, MP '13)  $\sqrt{5}$ Barnaby, Peloso '10<sup>th</sup> • Smallness of  $V_{\text{shift}}$  technically natural  $\Delta V V_{\text{shift}}$ xion; reference values  $f \sim 10^{16} \text{ GeV}$ ,  $m_{\phi} \geq 10^{13} \text{ GeV}$ 

self-couplings the advantage that he advantage that din P vales societs const. infla • Shir syxn Oyperkin yr GABit Smallness of Leshift technicady  $\mu\nu$  (review Pajer, MP  $\phi \rightarrow \phi + \operatorname{Growns} \operatorname{coupling}_{4\Lambda} \operatorname{to}_{\delta} \operatorname{coupling}_{4\Lambda} \operatorname{to}_{\delta} \operatorname{coupling}_{6\Lambda} \operatorname{to}_{\delta} \operatorname{coupling}_{6\Lambda} \operatorname{to}_{\delta} \operatorname{coupling}_{6\Lambda} \operatorname{to}_{\delta} \operatorname$ ilings to other fields ) hat the other fields ) hat the other fields () hat the other of the other other of the other o atter (predict  $10^{16}$  GeV  $\simeq 10^{15}$  GeV • Flatness and gatusathtess sandsefaussions ity restandard single field slow rob Natural) Inflation self-couplings Agreement wit \_ Applitude to othe sunded n2 couplings to othe s M  $V \propto V_{\rm shift}$ ordial Black blokes (PBH)  $\begin{array}{c} Value (1, 0, 0) \\ Value$  $f = \frac{1}{1}$ ,  $f = \frac{1}{1}$ , f =spectrum "(review Pajer, MP •  $S_{n}^{\text{mallness of } V_{\text{shift technically natural}}}$ xion: reference values  $f \sim 10^{16} \text{ GeV}$ 

he advantage that **A CALE STOR** const. infla • Shinsyxnoyperkinyy SABit Smallness of Eshift technicady  $\phi \rightarrow \phi + \Theta_n \otimes \Pi_5 C O \mu p | i n g_F to chart if end on the strain of t$ lings to other (ields/) atter (predicti  $\frac{10^{16} \mu^3}{10^{16} \text{ GeV}^3} \simeq 10^{18} \text{ GeV} \simeq 10^{18} \text{ GeV}$ ekvisianpotentian • Flatness and gatus at the same signads grades signade signades in the second rstandard single field slow row Natural Inflation \_self-couplings\_Agreement - Freese She han yon male to; y. On 200 yp in g y s for  $V \propto V_{
m shift}$ steen with the shear of the star of the shear of the shea if inflation GeV,  $m_{\phi} \simeq 10^{13} \text{ GeV}$  (review raje, .... if of inflation GeV,  $m_{\phi} \simeq 10^{13} \text{ GeV}$  matter if the dealine of V an important 'rece also ons' to the • Smallness of  $V_{\text{shift}}$  technically natural 2 Addres Same in  $5\psi + \frac{\alpha}{6} \frac{\delta}{F_{\mu\nu}} \frac{f_{\mu\nu}}{F^{\mu\nu}}$ • Smallness of  $V_{\text{shift}}$  technically natural 2 Addres Same in  $5\psi + \frac{\alpha}{6} \frac{\delta}{F_{\mu\nu}} \frac{f_{\mu\nu}}{F^{\mu\nu}}$ • Smallness of  $V_{\text{shift}}$  technically natural 2 Addres Same in Amber Sisting (12) Amber Sisting (12)

self-couplings the aquantage that infla const. Smallness of Ushift  $\frac{1}{2} = \frac{1}{2} \frac$ tecinytheayy  $\rightarrow \phi + \Theta_n \otimes \Omega_5 C O U D I I M G S to O D C S C O D S$ lings to other fields )  $10^{16} \text{ GeV}_{3,\text{GeV}} \simeq 10^{16} \text{ GeV}_{10} \text{ GeV}_{10} \approx 10^{16} \text{ GeV}_{10} \approx 1$ ervitiethpote Flatness and gaussamtess south graussion ity  $\frac{1}{2} = \frac{1}{2} = \frac{1}$ self-couplings Agre ment with standard single fier Slong Natural Inflation Applitude the system beinged n2COuppings to Inflaton perturbations  $\delta \phi$  $s_{\rm Mift} \propto V_{\rm through inverse decay}$ elds Hiction to the second s the OCD axion, reference values f steen (the review Pajer, MP 13) • Smallness of  $V_{\text{shift}}$  technically natural  $\Delta V V_{\text{shift}}$   $V_{5}\psi + \frac{\alpha}{F}\phi F_{\mu\nu}F^{\mu\nu}$ • Smallness of  $V_{\text{shift}}$  technically natural  $\Delta V V_{\text{shift}}$  ( Sign frequence values  $f_{1} \sim 10^{16} \text{ GeV}_{10}$  m/s  $2 \pm 10^{13} \text{ GeV}_{10}$ 

self-couplings the aquantage that **diases y d**els infla const. Smallness of Lishift t teciniticagy Pendentiere strained Sorbolet a  $\phi + \ell \omega \alpha 0 1 5 COUP I 1905$ lings to other fields ) atter (pred ekvitatopote  $\int_{M_{p}^{p}} 10^{16} \text{GeV}_{3,\text{GeV}} \simeq 10^{15} \text{GeV}_{1} = 10^{15}$ what here and attractive se some dserause ious ity restandard single fiels long Natural Inflation self-couplings Agreement wit A HERE Y HERE FRANK SALE A LOOPATE Amplitude the system ded n2COuplings to Inflaton perturbations  $\delta \phi$  $\frac{1}{2} = \frac{1}{2} + \frac{1}$ lds shift  $\equiv \frac{M_p^2}{2} \left(\frac{V_{,\phi}}{V}\right)$ <sup>#</sup>the OCD axion, reference values f Problem ly natural A wordanty confernations and an e<sup>4</sup>, pNGr, PBH, FGWr, very sensitive jeg finflatoon GeV  $5\psi + -\phi F_{\mu\nu}F^{\mu\nu}$  • Smallness of  $V_{\text{ohift}}$  technically nature blender  $f \sim 10^{16} \text{ GeV} h$   $m \to 2000 \text{ Totoles}$ ice to the president of the second seco

**PROBLEM:** PNG, GW and PBH **Approximations (e.g. Analytical)** 

# Let's have a look to the full problem !

$$\left(V(\phi) = \frac{1}{2}m^2\phi^2\right)$$

**PROBLEM:** PNG, GW and PBH  $\longrightarrow$  Approximations (e.g. Analytical)  $\pi_{\phi} \equiv \dot{\phi} , \quad E_i \equiv \dot{A}_i , \quad B_i \equiv \epsilon_{ijk} \partial_j A_k$  $\tilde{\pi}_{\phi} = a^3 \pi_{\phi} , \quad \tilde{\vec{E}} = a \vec{E} , \ \pi_a \equiv \dot{a}$ 

# Let's have a look to the full problem !

$$\left(V(\phi) = \frac{1}{2}m^2\phi^2\right)$$




















**PROBLEM:** PNG, GW and PBH **Approximations (e.g. Analytical)** 

#### Can we do better than homogeneous backreaction ?



#### Yes, we need a full lattice approach



**PROBLEM:** PNG, GW and PBH **Approximations (e.g. Analytical)** 

Yes, we need a full lattice approach



**PROBLEM:** PNG, GW and PBH **Approximations (e.g. Analytical)** 

#### Yes, we need a full lattice approach



**PROBLEM:** PNG, GW and PBH **Approximations (e.g. Analytical)** 

#### Let's "latticize" the system of EOM !



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DGF, Shaposhnikov 2017 Canivete, DGF 2018

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DGF, Shaposhnikov 2017 Canivete, DGF 2018

- 1. Lattice Gauge Inv:  $A_{\mu} \longrightarrow A_{\mu} + \Delta_{\mu}^{+} \alpha$
- 2. Cont. Limit to  $\mathcal{O}(dx^2)$

3. Lattice Bianchi Identities:  $\Delta_i^-(B_i^{(4)} + B_{i,+\hat{0}}^{(4)}) = 0, \dots$ 

4. Topological Term:  $(F_{\mu\nu}\tilde{F}^{\mu\nu})_L = \Delta^+_{\mu}K^{\mu}$  (CS current)  $\begin{bmatrix} F_{\mu\nu}\tilde{F}^{\mu\nu} = \partial_{\mu}K^{\mu} \end{bmatrix}$  Exact Shift Sym. on the lattice !

#### Let's "latticize" the system of EOM !

DGF, Shaposhnikov 2017 Canivete, DGF 2018

- 1. Lattice Gauge Inv:  $A_{\mu} \longrightarrow A_{\mu} + \Delta_{\mu}^{+} \alpha$
- 2. Cont. Limit to  $\mathcal{O}(dx^2)$

3. Lattice Bianchi Identities:  $\Delta_i^-(B_i^{(4)} + B_{i,+\hat{0}}^{(4)}) = 0, \dots$ 

4. Topological Term:  $(F_{\mu\nu}\tilde{F}^{\mu\nu})_L = \Delta^+_{\mu}K^{\mu}$  (CS current)  $\begin{bmatrix} F_{\mu\nu}\tilde{F}^{\mu\nu} = \partial_{\mu}K^{\mu} \end{bmatrix}$  Exact Shift Sym. on the lattice !

**PROBLEM:** PNG, GW and PBH **Approximations** (e.g. Analytical)

#### Let's "latticize" the system of EOM !

DGF, Shaposhnikov 2017 Canivete, DGF 2018

We show now our recent work

Phys.Rev.Lett. 131 (2023) 15, 151003

e-Print: 2303.17436 [astro-ph.CO]



CosmoLattice

# **Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 18\right)$ $\downarrow$ $\alpha_{\Lambda} \equiv \frac{m_p}{\Lambda}$

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Comparison to Homogeneous Backreaction

**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 15\right)$ 

#### **Comparison to GEF (Homogeneous backreaction)**

GEF (O. Sobol, R. von Eckardstein, K. Schmitz)



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**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 18\right)$ 



### **Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \Lambda = \frac{m_p}{\alpha}\right)$ $(\alpha = 15, 18, 20)$



## Gauge Amplification

**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 18\right)$ 













# UV sensitivity (convergence)













#### **UV convergence ?**



#### **UV convergence ?**



#### **UV convergence ?**



**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 15\right)$ 



Zoom

**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 15\right)$ 

ZOOM



**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 15\right)$ 



**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 15\right)$ 



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**Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \alpha_{\Lambda} = 15\right)$ 



### Chirality













### **Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \Lambda = \frac{m_p}{X}\right)$ (X = 15, 20, 25)

#### Summary

- \*  $\xi$  Controls the Gauge field excitation
- \* Linear change in  $\xi$  : exponential response in  $A_{\mu}$
- \* Predictions/constraints (PNG, PBH and GWs) depend crucially on  $\xi$  : we will re-assess real observability !
- \* Adding Schwinger pair production easy via  $\vec{J} = \sigma \vec{E}$
- \* Other phenomena: BAU, Magnetogenesis, ...



#### Summary

Phys. Rev. Lett. 131 (2023) 15, 151003

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Work in Progress ...

\* Predictions/constraints (PNG, PBH and GWs) depend crucially on  $\xi$  : we will re-assess real observability !

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## **Axion-Inflation** $\left(V(\phi) = \frac{1}{2}m^2\phi^2; \frac{\phi}{4\Lambda}F\tilde{F}; \Lambda = \frac{m_p}{X}\right)$ (X = 15, 20, 25)

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- \* Linear change in  $\xi$  : exponential response in  $A_{\mu}$
- \* Predictions/constraints (PNG, PBH and GWs) depend crucially on  $\xi$ : we will re-assess real observability ! Future work ... \* Adding Schwinger pair production easy via  $\vec{J} = \sigma \vec{E}$ \* Other phenomena: BAU, Magnetogenesis, ...

## **Example II**

## Cosmic String Loops (+ GW emission)

ArXiv:2308.08456 [astro-ph]

(Submitted to Phys. Rev. Lett.)

with J. Baeza-Ballesteros, E. Copeland, J. Lizarraga (PhD student)

Cosmic strings are one-dimensional topological defects





Symm. Breaking Fld ('Higgs')

**Different Vacua** 

Cosmic strings are one-dimensional topological defects





Symm. Breaking Fld ('Higgs')

**Different Vacua (at different locations)** 

Cosmic strings are one-dimensional topological defects



Symm. Breaking Fld ('Higgs')

**Different Vacua (at different locations)** 

Cosmic strings are one-dimensional topological defects



Symm. Breaking Fld ('Higgs')

**Different Vacua (at different locations)** 

#### Global (scalar) or Local (Scalar + Gauge fld.)

[Cosmic Strings: Global (scalar) or Local (Scalar + Gauge fld.)]

#### e.g. Global cosmic strings



# Cosmic Strings: Global (scalar) or the cal (Scalar + Gauge fld.) ] Intensity of magnetic energy density





\* Scaling dynamics

\* Infinitely thin

\* Inter-commutation

# **Cosmic String Networks** Scaling $H^{-1}$ $H^{-1}$

\* Scaling dynamics

\* Infinitely thin

\* Inter-commutation

Infinitely thin:  $H^{-1} \gg m^{-1}$ 



\* Scaling dynamics

\* Infinitely thin

\* Inter-commutation

#### Intercommutation





Loops !

Loops !



Loops are formed !



#### **Loops are formed** ! Vibrate under their tension !



#### Periodic Oscillations

#### Loops are formed !

#### Vibrate under their tension !



#### Gravitational Waves (GW) are emitted !

#### Loops are formed ! Vibrate under their tension !



#### Gravitational Waves (GW) are emitted !

**Superposition from many loop signals** 



**Gravitational Wave Background** 

Traditional picture --- Nambu-Goto approximation (zero width)

String networks = Infinite strings + Loops

Traditional picture --- Nambu-Goto approximation (zero width)

String networks = Infinite strings + Loops
Decay to loops

Traditional picture --- Nambu-Goto approximation (zero width)

► Decay to GWs (Vilenkin '81)

String networks = Infinite strings + Loops
Decay to loops

Traditional picture --- Nambu-Goto approximation (zero width)

> Loops decay via GWs radiated in all harmonic frequencies  $\nu_j$ 

$$P_j = \Gamma G \mu^2 \frac{j^{-q}}{\zeta(q)} \longrightarrow P_{GW} = \dot{E}_{GW} = \sum_{j=1}^{\infty} P_j = \Gamma G \mu^2$$

Traditional picture --- Nambu-Goto approximation (zero width)

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But ...

Field-theory strings can also decay via particle emission

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#### **Goal**: Particle and GW emission

Traditional picture --- Nambu-Goto approximation (zero width)

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But ...

Field-theory strings can also decay via particle emission

Goal: Particle and GW emission using lattice simulations

Traditional picture --- Nambu-Goto approximation (zero width)

> Loops decay via GWs radiated in all harmonic frequencies  $\nu_j$ 

$$P_j = \Gamma G \mu^2 \frac{j^{-q}}{\zeta(q)} \longrightarrow P_{GW} = \dot{E}_{GW} = \sum_{j=1}^{\infty} P_j = \Gamma G \mu^2$$

But ...

Field-theory strings can also decay via particle emission



GOAL Dynamics of an isolated loop and its particle & GW emission

GOAL Dynamics of an isolated loop and its particle & GW emission

Today's focus on ... Global Strings

[but Local String analysis coming !]

#### GOAL Dynamics of an isolated loop and its particle & GW emission

**Case I**: Nielsen-Olesen

**Case II** : Network



(following Vachaspati et al 2020)



(following Lizarraga et al 2020/21)

#### GOAL Dynamics of an isolated loop and its particle & GW emission

#### Case I: 'Artificial'

Isolate the inner loop —







**Boost 2 string pairs** 

Intersect -> 2 Loops Inner/Outer **Isolate Inner Loop** 

#### GOAL Dynamics of an isolated loop and its particle & GW emission

Case II: Network — Only one loop remains eventually —



## **Loop Resolution**














## **Decay of a Loop**



#### **Higgs isosurface**

String Core

















































































(Due to Particle Emission)



#### Network





(Due to Particle Emission)







(Due to Particle Emission)







(Due to Particle Emission)







(Due to Particle Emission)







(Due to Particle Emission)



#### Network

#### Artificial

$$P_{arphi} = -rac{\mathrm{d}E_{\mathrm{str}}}{\mathrm{d}t} pprox (11.2\pm1.6)v^2$$

$$P_{arphi} = -rac{\mathrm{d} E_{\mathrm{str}}}{\mathrm{d} t} pprox (3-8) v^2$$



(Due to Particle Emission)



#### Network

Artificial

$$P_{arphi} = -rac{\mathrm{d}E_{\mathrm{str}}}{\mathrm{d}t} pprox (11.2\pm1.6)v^2$$

$$P_{arphi} = -rac{\mathrm{d}E_{\mathrm{str}}}{\mathrm{d}t} pprox (3-8)v^2$$

# Independent\* of Resolution & Length !

[\*within the error]

#### **String Loop: GW emission**

$$\tilde{t} = \sqrt{\lambda} v t \qquad \qquad \tilde{l} = \sqrt{\lambda} v l \qquad \qquad \tilde{k} = k/\sqrt{\lambda} v$$

\_\_\_\_




















## **String Loop: GW emission**

$$\tilde{t} = \sqrt{\lambda} v t \qquad \qquad \tilde{l} = \sqrt{\lambda} v l \qquad \qquad \tilde{k} = k/\sqrt{\lambda} v$$

(Artificial Loop)





$$\tilde{L}/\tilde{L}_{1/4} = 2,3,4,...,8$$











### String Loop: GW emission







#### Network

**Artificial** 





#### Network

**Artificial** 





#### Network

**Artificial** 



## **GW Power Emitted**



(GW emission)

\_\_\_\_





Baeza-Ballesteros et al, 2023 (Global Strings)  $[L/w \simeq 80 - 1700]$ 

$$\frac{P_{\rm GW}}{P_{\phi}} \simeq \mathcal{O}(10) \left(\frac{v}{m_p}\right)^2 \ll 1$$

$$\left(\frac{v^2}{m_p^2} \lesssim 10^{-6} - 10^{-3} \text{ [Lopez-Eiguren, et al. (2017), Benabou, et al. (2023)]}\right)$$





Baeza-Ballesteros et al, 2023 (Global Strings)  $[L/w \simeq 80 - 1700]$ 

$$\frac{P_{\rm GW}}{P_{\phi}} \simeq \mathcal{O}(10) \left(\frac{v}{m_p}\right)^2 \iff 1$$

#### So what happens with Local Strings?

Results will impact on Real evaluation of GW emission Re-evaluation of PTA constraints (Pulsar Time Array)

Results will impact on Real evaluation of GW emission Re-evaluation of PTA constraints (Pulsar Time Array)

### Implications for

Dark Matter Axion string network Local (Abelian-Higgs) string network Comparison with Nambu-Goto GUT models

. . . .

Almost ... the End If you want to learn how to "latticesize" your problems ...







## 2nd CL School 2023: Sept 25-29

### https://www.youtube.com /@CosmoLattice/videos



CosmoLattice School 2023, Day 4: Practice 3 (Simulating Gravitational Waves)

17 views · 4 months ago



CosmoLattice School 2023, Day 4: Lecture 8 (Plotting Features of CosmoLattice)



CosmoLattice School 2023, Day 3: Lecture 7 [SU(2) Scalar-Gauge Theory Lattice...

#### Evolution of GWs modes non-local operations are computationally expensive! Solution: we define a set of unphysical tensor modes u's (1) Evolve equation of motion of u's) $\ddot{u}_{ij} + 3H\dot{u}_{ij} - \frac{\nabla^2}{a^2}u_{ij} = \frac{2}{m_{ij}^2a^2}\Pi_{ij}^2$ (2) When needed, (compute power spectrum energy density) we apply transformation $h_{ij}(k, t) = \Lambda_{ij,kl}(k)u_{kl}(k, t)$ 1:19:39

CosmoLattice School 2023, Day 3: Lecture 6 (Creation and Propagation of Grav. Waves)

12 views • 4 months ago



## Details for CL School 2024 TBA at:

### https://cosmolattice.net



# Thanks for your attention

## Merci pour votre attention

# Thanks for your attention

## Merci pour votre attention

## **Back Slides**



Constraints Applications Program Variables Axion Inflation self-couplings the aquantage that all cardeles const. <sup>Infla</sup> • Shi syx Opporter my Stati Smallness of Eshift technicady lings to other fields ) atter (predic  $\frac{10^{16} \text{ (M}^{3} \text{ (M}^{3}$ self-couplings, Agreement with standard single field slow now Natural Freese Sheftarsymmetery ... On 2000 proching Streed s V shift  $\propto V_{\rm shift}$ Example, GM torediet onstrains  $\begin{array}{l} \hline \mathcal{H}_{\mathcal{A}} \\ \hline \mathcal{H}_{\mathcal{H}} \\ \hline \mathcal{H}$ steen wird with nically natural. A wordant, Roiffrons reheasing an (review Pajer, MP '13) in the final tool GeV,  $m_{\phi} \simeq 10^{13} GeV$ an important role also  $\phi \psi \gamma^{\mu} \gamma_5 \psi + \frac{1}{2} \phi F_{\mu\nu} F^{\mu\nu}$ • Smallness of  $V_{\text{shift}}$  technically natural  $\Delta V = V_{\text{shift}} V_{\text{shift}} (\phi)$ • Smallness of  $V_{\text{shift}}$  technically natural  $\Delta V = V_{\text{shift}} V_{\text{shift}} (\phi)$ • Smallness of  $V_{\text{shift}} = 0.12$ 



## INFLATIONARY MODELS Axion-Inflation



Blue-Tilted + Chiral + Non-G GW background

Bartolo et al '16, 1610.06481

## INFLATIONARY MODELS Axion-Inflation



Bartolo et al '16, 1610.06481

## INFLATIONARY MODELS Axion-Inflation



Bartolo et al '16, 1610.06481

## Constraints



> The **first Friedmann equation** is used to check the accuracy of the simulation.




#### > The **first Friedmann equation** is used to check the accuracy of the simulation.





#### > The **first Friedmann equation** is used to check the accuracy of the simulation.



#### **Gauge theories: Gauss constraint**

Preservation of U(1) & SU(2) Gauss constraints (for all integrators!)



#### **Gauge theories: Gauss constraint**

Preservation of U(1) & SU(2) Gauss constraints (for all integrators!)



# Applications (papers)

Gravitational Wave Symphony from Oscillating Spectator Scalar Fields #	<b>‡1</b>
Yanou Cui (UC, Riverside), Pankaj Saha, Evangelos I. Sfakianakis (Barcelona, IFAE and Case Western Reserve U. (Oct 19, 2023)	.)
e-Print: 2310.13060 [hep-ph]	
Higher-form symmetry and chiral transport in real-time lattice $U(1)$ gauge theory	
Arpit Das, Adrien Florio, Nabil Iqbal, Napat Poovuttikul (Sep 25, 2023)	
e-Print: 2309.14438 [hep-th]	
·	

Gravitational Wave Symphony from Oscillating Spectator Scalar Fields #1	   
Yanou Cui (UC, Riverside), Pankaj Saha, Evangelos I. Sfakianakis (Barcelona, IFAE and Case Western Reserve U.) (Oct 19, 2023)	     
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Effects of Fragmentation on Post-Inflationary Reheating	] [
Marcos A.G. Garcia (Mexico U.), Mathieu Gross (IJCLab, Orsay), Yann Mambrini (IJCLab, Orsay), Keith A. Olive (Minnesota U., Theor. Phys. Inst.), Mathias Pierre (DESY) et al. (Aug 30, 2023)	1 1 1 1
e-Print: 2308.16231 [hep-ph] Gravitational Wave Emission from a Cosmic String Loop, I: Global Case	-  #
Jorge Baeza-Ballesteros (Valencia U., IFIC), Edmund J. Copeland (Nottingham U.), Daniel G. Figueroa (NU., IFIC), Joanes Lizarraga (Basque U., Bilbao and U. Basque Country, Leioa) (Aug 16, 2023)	Valencia

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e-Print: 2308.16231 [hep-ph]	
Gravitational Wave Emission from a Cosmic String Loop, I: Global Case	#5
Jorge Baeza-Ballesteros (Valencia U., IFIC), Edmund J. Copeland (Nottingham U.), Daniel G. Figueroa (Val U., IFIC), Joanes Lizarraga (Basque U., Bilbao and U. Basque Country, Leioa) (Aug 16, 2023)	lencia
e-Print: 2308.08456 [astro-ph.CO]	   
Reheating after Inflaton Fragmentation	
Marcos A.G. Garcia (Mexico U.), Mathias Pierre (DESY) (Jun 13, 2023)	
e-Print: 2306.08038 [hep-ph]	
On unitarity in singlet inflation with a non-minimal coupling to gravity	
Oleg Lebedev (Helsinki U.), Yann Mambrini (IJCLab, Orsay), Jong-Hyun Yoon (IJCLab, Orsay) (I	May 9, 2
Published in: JCAP 08 (2023) 009 • e-Print: 2305.05682 [hep-ph]	



Gravitational freeze-in dark matter from Higgs preheating
Ruopeng Zhang (Chongqing U.), Zixuan Xu (Chongqing U.), Sibo Zheng (Chongqing U.) (May 4, 2023)
Published in: JCAP 07 (2023) 048, JCAP 07 (2023) 048 • e-Print: 2305.02568 [hep-ph]

Gravitational freeze-in dark matter from Higgs preheating	   
R. Dissipative Genesis of the Inflationary Universe	
<ul> <li><sup>Pu</sup> Hiroki Matsui (Kyoto U., Yukawa Inst., Kyoto), Alexandros Papageorgiou (IBS, Daejeon), Fuminobu Takahashi (Tohoku U.), Takahiro Terada (IBS, Daejeon) (May 3, 2023)</li> <li>e-Print: 2305.02366 [gr-qc]</li> </ul>	

Gravitational freeze-in dark matter from Higgs preheating	   
R Dissipative Genesis of the Inflationary Universe	   
Pt Hi Dissipative Emergence of Inflation from Quasi-Cyclic Universe	۱ – – – ۲ – ۱ ۱
Ta Hiroki Matsui (Kyoto U., Yukawa Inst., Kyoto), Alexandros Papageorgiou (IBS, Daejeon, CTPU), Fur	ninobu
<sup>e</sup> - Takahashi (Tohoku U.), Takahiro Terada (IBS, Daejeon, CTPU) (May 3, 2023)	
e-Print: 2305.02367 [gr-qc]	ا ا ا ـ ـ ـ ـ ـ ـ

Gravitational freeze-in dark matter from Higgs preheating	
R Dissipative Genesis of the Inflationary Universe	
P <sup>I</sup> HI Dissipative Emergence of Inflation from Quasi-Cyclic Universe	     
Preheating in Einstein-Cartan Higgs Inflation: Oscillon formation	ה, CTPU), Fuminobu
Matteo Piani (Lisbon, CENTRA), Javier Rubio (Madrid U.) (Apr 25, 2023)	
e-Print: 2304.13056 [hep-ph]	L    

Gravitational freeze-in dark matter from Higgs preheating	
R Dissipative Genesis of the Inflationary Universe	
Pu Hi Dissipative Emergence of Inflation from Quasi-Cyclic Universe	1   
Ta Preheating in Einstein-Cartan Higgs Inflation: Oscillon formation <sup>م, CTPU), Fuminobu</sup>	1 1
Solving the domain wall problem with first-order phase transition	#13
e Yang Li (Beijing, GUCAS and Beijing, Inst. Theor. Phys.), Ligong Bian (Chongqing U. and Maryland U. an U., CHEP), Yongtao Jia (Chongqing U.) (Apr 11, 2023)	าd Peking
e-Print: 2304.05220 [hep-ph]	، ۱ ۱

Gravitational freeze-in dark matter from Higgs preheating		
R Dissipative Genesis of the Inflationary Universe		
Pu Hi Dissipative Emergence of Inflation from Quasi-Cyclic Universe		
Preheating in Einstein-Cartan Higgs Inflation: Oscillon formation <sup>a, CTPU), Fuminobu</sup>		
Solving the domain wall problem with first-order phase transition	#13	
e <sup>e</sup> Y Strong Backreaction Regime in Axion Inflation	#15	
<sup>U</sup> Daniel G. Figueroa (Valencia U., IFIC), Joanes Lizarraga (Basque U., Bilbao and U. Basque Country, Leioa), Ander <sup>e</sup> Urio (Basque U., Bilbao and U. Basque Country, Leioa), Jon Urrestilla (Basque U., Bilbao and U. Basque Country, Leioa) (Mar 30, 2023)		
Published in: <i>Phys.Rev.Lett</i> . 131 (2023) 15, 151003 • e-Print: 2303.17436 [astro-ph.CO]	   	

Gravitational freeze-in dark matter from Higgs preheating	
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<sup>U</sup> D: Oscillon formation from preheating in asymmetric inflationary potentials	ry, Leioa), Ander
Rafid Mahbub (Gustavus Adolphus Coll.), Swagat S. Mishra (Nottingham U.) (Mar 13, 2023)	3asque Country,
PuPublished in: <i>Phys.Rev.D</i> 108 (2023) 6, 063524 • e-Print: 2303.07503 [astro-ph.CO]	

Gravitational freeze-in dark matter from Higgs preheating		
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e Y Strong Backreaction Regime in Axion Inflation		#15
Oscillon formation from preheating in asymmetric inflationar	ry potentials	ry, Leioa), Ander
$\int_{L_{\epsilon}}^{C} \frac{1}{R}$ Misaligned, tilted and distorted: the hard life of audible axi	ons <sup>13, 2023)</sup>	3asque Country,
Pu Wolfram Ratzinger (Mainz U.) (Jan 26, 2023)	esis !!	





#### Just in 2023 ....

# Program Variables

Cosmo Lattice – Program variables

Choose:

 $\{\alpha, \omega_*, f_*\}$ 



Space and time



 $\widetilde{A}_{\mu} = \frac{A_{\mu}}{\Omega} \qquad \widetilde{B}_{\mu}^{a} = \frac{B_{\mu}^{a}}{\Omega}$ 



Cosmo Lattice – Program variables

Choose:

 $\left\{ \alpha, \omega_*, f_* \right\} \longrightarrow$ 

$$d\tilde{\eta} \equiv a^{-\alpha} \omega_* dt$$
$$d\tilde{x}^i \equiv \omega_* dx^i$$

Space and time



Gauge fields

Cosmo Lattice – Program variables

Choose:

 $\left\{ \alpha, \boldsymbol{\omega}_*, \boldsymbol{f}_* \right\}$ 

$$\begin{vmatrix} d\tilde{\eta} \equiv a^{-\alpha} \omega_* dt \\ d\tilde{x}^i \equiv \omega_* dx^i \end{vmatrix}$$

Space and time





Gauge fields

Cosmo Lattice – Program variables

Choose:

 $\left\{ \alpha, \omega_*, f_* \right\}$ 

$$d\tilde{\eta} \equiv a^{-\alpha} \omega_* dt$$
$$d\tilde{x}^i \equiv \omega_* dx^i$$

Space and time



 $\widetilde{A}_{\mu} = \frac{A_{\mu}}{\widehat{B}_{\mu}^{a}} = \frac{B_{\mu}^{a}}{\widehat{B}_{\mu}^{a}} = \frac{B_{\mu}^{a}}{\widehat{B}_{\mu}^{a}}$ 



Cosmo Lattice – Program variables

Choose:  $\{\alpha, \omega_*, f_*\}$ 

$$d\tilde{\eta} \equiv a^{-\alpha} \omega_* dt$$
$$d\tilde{x}^i \equiv \omega_* dx^i$$

Space and time





Gauge fields

#### How do I choose them?

Cosmo Lattice – Program variables

Choose:  $\left\{ \alpha, \omega_*, f_* \right\}$ 

$$d\tilde{\eta} \equiv a^{-\alpha} \omega_* dt$$
$$d\tilde{x}^i \equiv \omega_* dx^i$$

Space and time





Gauge fields

**Example:** 
$$\phi(t) \simeq \Phi_* \times f_{\rm osc}(t)$$



**Cosmo Lattice** – Program variables



 $\overline{\widetilde{A}_{\mu}} = \frac{A_{\mu}}{\overline{B}_{\mu}} \qquad \widetilde{B}_{\mu}^{a} = \frac{B_{\mu}^{a}}{\overline{B}_{\mu}}$ 

fields

Gauge

fields

Scalar

**Example:** 

 $\phi(t) \simeq \Phi_* \times f_{\rm osc}(t)$ 



Cosmo Lattice – Program variables



> Write scalar potential and first and second derivatives in one file (model.h)

$$\tilde{V}(\tilde{\phi}, |\tilde{\varphi}|, |\widetilde{\Phi}|) \equiv \frac{1}{f_*^2 \omega_*^2} V(f_* \tilde{\phi}, f_* |\tilde{\varphi}|, f_* |\widetilde{\Phi}|) \longrightarrow \frac{\partial \tilde{V}}{\partial \tilde{\phi}}, \frac{\partial \tilde{V}}{\partial |\tilde{\varphi}|}, \frac{\partial \tilde{V}}{\partial |\tilde{\Phi}|}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\phi}|^2}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\varphi}|^2}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\Phi}|^2}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\Phi}|^2}$$

Cosmo Lattice – Program variables



> Write scalar potential and first and second derivatives in one file (model.h)

$$\tilde{V}(\tilde{\phi}, |\tilde{\varphi}|, |\widetilde{\Phi}|) \equiv \frac{1}{f_*^2 \omega_*^2} V(f_* \tilde{\phi}, f_* |\tilde{\varphi}|, f_* |\widetilde{\Phi}|) \longrightarrow \frac{\partial \tilde{V}}{\partial \tilde{\phi}}, \frac{\partial \tilde{V}}{\partial |\tilde{\varphi}|}, \frac{\partial \tilde{V}}{\partial |\tilde{\Phi}|}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\phi}|^2}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\varphi}|^2}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\Phi}|^2}, \frac{\partial^2 \tilde{V}}{\partial |\tilde{\Phi}|^2}$$

Parameters passed via one file (input.txt) (no need to re-compile !)



Axion-inflation extra stuff

#### **Axion-Inflation**

**PROBLEM:** PNG, GW and PBH **Analytical approximations** !

#### Let's "latticize" the system of EOM !





$$\begin{array}{c} \text{Inflator}_{\text{EOM}} \Delta_{0}^{+}\left(a^{3}\pi_{\phi}\right) = a_{+\frac{0}{2}}\sum_{i}\Delta_{i}^{-}\Delta_{i}^{+}\phi_{+\frac{0}{2}} - a_{+\frac{0}{2}}^{3}m^{2}\phi_{+\frac{0}{2}} + \frac{1}{\Lambda}\sum_{i}\frac{1}{2}E_{i,+\frac{0}{2}}^{(2)}\left(B_{i}^{(4)} + B_{i,+0}^{(4)}\right) \\ \hline \\ \text{Gauge}_{\text{Fid}} \Delta_{0}^{-}\left(a_{+\frac{0}{2}}E_{i,+\frac{0}{2}}\right) = -\frac{1}{a}\sum_{j,k}\epsilon_{ijk}\Delta_{j}^{-}B_{k} - \frac{1}{2\Lambda}\left(\pi_{\phi}B_{i}^{(4)} + \pi_{\phi,+i}B_{i,+i}^{(4)}\right) \\ & \quad +\frac{1}{8\Lambda}(2 + dx\Delta_{i}^{+})\sum_{\pm}\sum_{j,k}\left\{\epsilon_{ijk}[(\Delta_{j}^{\pm}\phi)E_{k,\pm j}^{(2)}]_{+\frac{0}{2}} + [(\Delta_{j}^{\pm}\phi)E_{k,\pm j}^{(2)}]_{-\frac{0}{2}}\right\} \\ a_{+\frac{0}{2}}\sum_{i}\Delta_{i}^{-}E_{i,+\frac{0}{2}} = -\frac{1}{4\Lambda}\sum_{\pm}\sum_{i}\left(\Delta_{i}^{\pm}\phi_{+\frac{0}{2}}\right)\left(B_{i}^{(4)} + B_{i,+0}^{(4)}\right)_{\pm i}, \quad \text{(Gauss Law)} \\ 1 \text{ Lattice Gauge Inv: } A \longrightarrow A + A^{\pm}\alpha \end{array}$$

- 1. Lattice Gauge Inv:  $A_{\mu} \longrightarrow A_{\mu} + \Delta_{\mu}^{+} \alpha$
- 2. Cont. Limit to  $\mathcal{O}(dx^2)$
- 3. Lattice Bianchi Identities:  $\Delta_i^-(B_i^{(4)} + B_{i,+\hat{0}}^{(4)}) = 0, \dots$
- 4. Topological Term:  $(F_{\mu\nu}\tilde{F}^{\mu\nu})_L = \Delta^+_{\mu}K^{\mu}$  (CS current)  $\begin{bmatrix} F_{\mu\nu}\tilde{F}^{\mu\nu} = \partial_{\mu}K^{\mu} \end{bmatrix}$  Exact Shift Sym. on the lattice !

$$\begin{split} & \text{LATTICE FORMULATION of } \phi F \tilde{F} \\ & \text{Eom} \\ \Delta_{0}^{+} \left( a^{3} \pi_{\phi} \right) = a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} \Delta_{i}^{+} \phi_{+\frac{0}{2}} - a_{+\frac{0}{2}}^{3} m^{2} \phi_{+\frac{0}{2}} + \frac{1}{\Lambda} \sum_{i} \frac{1}{2} E_{i,+\frac{0}{2}}^{(2)} \left( B_{i}^{(4)} + B_{i,+0}^{(4)} \right) \\ & \Delta_{0}^{-} \left( a_{+\frac{0}{2}} E_{i,+\frac{0}{2}} \right) = -\frac{1}{a} \sum_{j,k} \epsilon_{ijk} \Delta_{j}^{-} B_{k} - \frac{1}{2\Lambda} \left( \pi_{\phi} B_{i}^{(4)} + \pi_{\phi,+i} B_{i,+i}^{(4)} \right) \\ & + \frac{1}{8\Lambda} (2 + dx \Delta_{i}^{+}) \sum_{\pm} \sum_{j,k} \left\{ \epsilon_{ijk} [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{+\frac{0}{2}} + [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{-\frac{0}{2}} \right\} \\ & a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} E_{i,+\frac{0}{2}} = -\frac{1}{4\Lambda} \sum_{\pm} \sum_{i} \left( \Delta_{i}^{\pm} \phi_{+\frac{0}{2}} \right) \left( B_{i}^{(4)} + B_{i,+0}^{(4)} \right)_{\pm i} , \quad \text{(Gauss Law)} \\ & \text{Def, Shaposhnikov 2017} \\ & \dot{\pi}_{\phi} = -3H \pi_{\phi} + \frac{1}{a^{2}} \vec{\nabla}^{2} \phi - m^{2} \phi + \frac{1}{a^{3}\Lambda} \vec{E} \cdot \vec{B}, \\ & \dot{\vec{E}} = -H \vec{E} - \frac{1}{a^{2}} \vec{\nabla} \times \vec{B} - \frac{1}{a\Lambda} \pi_{\phi} \vec{B} + \frac{1}{a\Lambda} \vec{\nabla} \phi \times \vec{E} \\ & \vec{\nabla} \cdot \vec{E} = -\frac{1}{a\Lambda} \vec{\nabla} \phi \cdot \vec{B} \text{(Gauss Law)} \\ & \text{EoM} \\ & \text{Continuum} \end{aligned}$$

$$\begin{split} & \text{LATTICE FORMULATION of } \phi F \tilde{F} \\ & \text{Eom} \\ \Delta_{0}^{+} \left( a^{3} \pi_{\phi} \right) = a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} \Delta_{i}^{+} \phi_{+\frac{0}{2}} - a_{+\frac{0}{2}}^{3} m^{2} \phi_{+\frac{0}{2}} + \frac{1}{\Lambda} \sum_{i} \frac{1}{2} E_{i,+\frac{0}{2}}^{(2)} \left( B_{i}^{(4)} + B_{i,+\hat{0}}^{(4)} \right) \\ & \Delta_{0}^{-} \left( a_{+\frac{0}{2}} E_{i,+\frac{0}{2}} \right) = -\frac{1}{a} \sum_{j,k} \epsilon_{ijk} \Delta_{j}^{-} B_{k} - \frac{1}{2\Lambda} \left( \pi_{\phi} B_{i}^{(4)} + \pi_{\phi,+i} B_{i,+i}^{(4)} \right) \\ & + \frac{1}{8\Lambda} (2 + dx \Delta_{i}^{+}) \sum_{\pm} \sum_{j,k} \left\{ \epsilon_{ijk} [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{+\frac{0}{2}} + [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{-\frac{0}{2}} \right\} \\ & a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} E_{i,+\frac{0}{2}} = -\frac{1}{4\Lambda} \sum_{\pm} \sum_{i} \left( \Delta_{i}^{\pm} \phi_{+\frac{0}{2}} \right) \left( B_{i}^{(4)} + B_{i,+\hat{0}}^{(4)} \right)_{\pm i} , \quad \text{(Gauss Law)} \\ & \text{DGF, Shaposhnikov 2017} \\ & \dot{\pi}_{\phi} = -3H \pi_{\phi} + \frac{1}{a^{2}} \vec{\nabla}^{2} \phi - m^{2} \phi + \frac{1}{a^{3}\Lambda} \vec{E} \cdot \vec{B}, \\ & \dot{\vec{E}} = -H \vec{E} - \frac{1}{a^{2}} \vec{\nabla} \times \vec{B} - \frac{1}{a\Lambda} \pi_{\phi} \vec{B} + \frac{1}{a\Lambda} \vec{\nabla} \phi \times \vec{E} \\ & \vec{\nabla} \cdot \vec{E} = -\frac{1}{a\Lambda} \vec{\nabla} \phi \cdot \vec{B} \text{(Gauss Law)} \\ & \text{EoM} \\ & \text{Continuum} \end{aligned}$$

$$\begin{split} & \text{LATTICE FORMULATION of } \phi F \tilde{F} \\ & \text{Eom} \\ \Delta_{0}^{+} \left( a^{3} \pi_{\phi} \right) = a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} \Delta_{i}^{+} \phi_{+\frac{0}{2}} - a_{+\frac{0}{2}}^{3} m^{2} \phi_{+\frac{0}{2}} + \frac{1}{\Lambda} \sum_{i} \frac{1}{2} E_{i,+\frac{0}{2}}^{(2)} \left( B_{i}^{(4)} + B_{i,+\hat{0}}^{(4)} \right) \\ & \Delta_{0}^{-} \left( a_{+\frac{0}{2}} E_{i,+\frac{0}{2}} \right) = -\frac{1}{a} \sum_{j,k} \epsilon_{ijk} \Delta_{j}^{-} B_{k} - \frac{1}{2\Lambda} \left( \pi_{\phi} B_{i}^{(4)} + \pi_{\phi,+i} B_{i,+i}^{(4)} \right) \\ & + \frac{1}{8\Lambda} (2 + dx \Delta_{i}^{+}) \sum_{\pm} \sum_{j,k} \left\{ \epsilon_{ijk} [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{+\frac{0}{2}} + [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{-\frac{0}{2}} \right\} \\ & a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} E_{i,+\frac{0}{2}} = -\frac{1}{4\Lambda} \sum_{\pm} \sum_{i} \left( \Delta_{i}^{\pm} \phi_{+\frac{0}{2}} \right) \left( B_{i}^{(4)} + B_{i,+\hat{0}}^{(4)} \right)_{\pm i} , \quad \text{(Gauss Law)} \\ & \text{Def, Shaposhnikov 2017} \\ & \dot{\pi}_{\phi} = -3H \pi_{\phi} + \frac{1}{a^{2}} \vec{\nabla}^{2} \phi - m^{2} \phi + \frac{1}{a^{3}\Lambda} \vec{E} \cdot \vec{B}, \\ & \dot{\vec{E}} = -H \vec{E} - \frac{1}{a^{2}} \vec{\nabla} \times \vec{B} - \frac{1}{a\Lambda} \pi_{\phi} \vec{B} + \frac{1}{a\Lambda} \vec{\nabla} \phi \times \vec{E} \\ & \vec{\nabla} \cdot \vec{E} = -\frac{1}{a\Lambda} \vec{\nabla} \phi \cdot \vec{B} \text{(Gauss Law)} \\ & \text{EoM} \\ & \text{Continuum} \\ \end{array}$$

$$\begin{split} & \text{LATTICE FORMULATION of } \phi F \tilde{F} \\ & \text{Eom} \\ \Delta_{0}^{+} \left( a^{3} \pi_{\phi} \right) = a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} \Delta_{i}^{+} \phi_{+\frac{0}{2}} - a_{+\frac{0}{2}}^{3} m^{2} \phi_{+\frac{0}{2}} + \frac{1}{\Lambda} \sum_{i} \frac{1}{2} E_{i,+\frac{0}{2}}^{(2)} \left( B_{i}^{(4)} + B_{i,+\hat{0}}^{(4)} \right) \\ & \Delta_{0}^{-} \left( a_{+\frac{0}{2}} E_{i,+\frac{0}{2}} \right) = -\frac{1}{a} \sum_{j,k} \epsilon_{ijk} \Delta_{j}^{-} B_{k} - \frac{1}{2\Lambda} \left( \pi_{\phi} B_{i}^{(4)} + \pi_{\phi,+i} B_{i,+i}^{(4)} \right) \\ & + \frac{1}{8\Lambda} (2 + dx \Delta_{i}^{+}) \sum_{\pm} \sum_{j,k} \frac{\left\{ \epsilon_{ijk} [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{+\frac{0}{2}} + [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{-\frac{0}{2}} \right\} \\ & a_{+\frac{0}{2}} \sum_{i} \Delta_{i}^{-} E_{i,+\frac{0}{2}} = -\frac{1}{4\Lambda} \sum_{\pm} \sum_{i} \left( \Delta_{i}^{\pm} \phi_{+\frac{0}{2}} \right) \left( B_{i}^{(4)} + B_{i,+\hat{0}}^{(4)} \right)_{\pm i} , \quad \text{(Gauss Law)} \\ & \text{DGF, Shaposhnikov 2017} \\ & \dot{\pi}_{\phi} = -3H \pi_{\phi} + \frac{1}{a^{2}} \vec{\nabla}^{2} \phi - m^{2} \phi + \frac{1}{a^{3}\Lambda} \vec{E} \cdot \vec{B}, \\ & \dot{\vec{E}} = -H \vec{E} - \frac{1}{a^{2}} \vec{\nabla} \times \vec{B} - \frac{1}{a\Lambda} \pi_{\phi} \vec{B} + \frac{1}{a\Lambda} \vec{\nabla} \phi \times \vec{E} \\ & \vec{\nabla} \cdot \vec{E} = -\frac{1}{a\Lambda} \vec{\nabla} \phi \cdot \vec{B} \text{(Gauss Law)} \\ & \text{EoM} \\ & \text{Continuum} \end{aligned}$$
## **LATTICE FORMULATION of** $\phi F \tilde{F}$ **Lattice Formulation**

EOIVI		
		$\mathbf{x} + (\mathbf{x}) \mathbf{x} = \mathbf{x} + (\mathbf{x}) \mathbf{x} + $
		$\Delta_{0}^{+}\left(a^{3}\pi_{\phi} ight)=a_{+rac{\hat{0}}{2}}\sum_{i}\Delta_{i}^{-}\Delta_{i}^{+}\phi_{+rac{\hat{0}}{2}}-a^{3}_{+rac{\hat{0}}{2}}m^{2}\phi_{+rac{\hat{0}}{2}}$
		$+rac{1}{\Lambda}\sum_i rac{1}{2}E^{(2)}_{i,+rac{\hat{0}}{2}}\left(B^{(4)}_i+B^{(4)}_{i,+\hat{0}} ight),$
	$\Delta_0^-$	$\left(a_{+rac{\hat{0}}{2}}E_{i,+rac{\hat{0}}{2}} ight) = -rac{1}{a}\sum_{j,k}\epsilon_{ijk}\Delta_{j}^{-}B_{k} - rac{1}{2\Lambda}\left(\pi_{\phi}B_{i}^{(4)} + \pi_{\phi,+i}B_{i,+i}^{(4)} ight)$
		$+\frac{1}{8\Lambda}(2+dx\Delta_{i}^{+})\sum_{\pm}\sum_{j,k}\left\{\epsilon_{ijk}[(\Delta_{j}^{\pm}\phi)E_{k,\pm j}^{(2)}]_{+\frac{\hat{0}}{2}}+[(\Delta_{j}^{\pm}\phi)E_{k,\pm j}^{(2)}]_{-\frac{\hat{0}}{2}}\right\}$
	$a_{+{\hat 0\over 2}}$	$\sum_{i} \Delta_{i}^{-} E_{i,+\frac{\hat{0}}{2}} = -\frac{1}{4\Lambda} \sum_{\pm} \sum_{i} \left( \Delta_{i}^{\pm} \phi_{+\frac{\hat{0}}{2}} \right) \left( B_{i}^{(4)} + B_{i,+\hat{0}}^{(4)} \right)_{\pm i} ,  \text{(Gauss Law)}$

Expansion		
	$\left(\Delta_{0}^{+}a_{-\hat{0}/2} ight)^{2}$ $\Delta_{0}^{-}\Delta_{0}^{+}a_{+\hat{0}/2}$	$egin{aligned} & e^2 = rac{a^2}{3m_{ m pl}^2}  ho_L, \ & e^2 = -rac{a_{+\hat{0}/2}}{6m_{ m pl}^2} ( ho_L + 3p_L)_{+\hat{0}/2}. \end{aligned}$

## **LATTICE FORMULATION of** $\phi FF$ **Lattice Formulation**

<b>EoM</b> $\Delta_{0}^{+}(a^{3}\pi_{\phi}) = a_{\pm 0} \sum \Delta_{i}^{-} \Delta_{i}^{+} \phi_{\pm 0} - a_{\pm 0}^{3} m^{2} \phi_{\pm 0}$	Expansion
$ \begin{split} & =_{0} \left( a^{\pm} u \phi \right)^{-} = a_{\pm \frac{0}{2}} \sum_{i} -i -i + \frac{1}{2} \sum_{i} a_{\pm \frac{0}{2}} e^{-i + \frac{1}{2}} + \frac{1}{2} \\ & + \frac{1}{\Lambda} \sum_{i} \frac{1}{2} E_{i,\pm \frac{0}{2}}^{(2)} \left( B_{i}^{(4)} + B_{i,\pm 0}^{(4)} \right) , \\ & \Delta_{0}^{-} \left( a_{\pm \frac{0}{2}} E_{i,\pm \frac{0}{2}} \right) = -\frac{1}{a} \sum_{j,k} \epsilon_{ijk} \Delta_{j}^{-} B_{k} - \frac{1}{2\Lambda} \left( \pi_{\phi} B_{i}^{(4)} + \pi_{\phi,\pm i} B_{i,\pm i}^{(4)} \right) \\ & + \frac{1}{8\Lambda} (2 + dx \Delta_{i}^{+}) \sum_{\pm} \sum_{j,k} \left\{ \epsilon_{ijk} [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{\pm \frac{0}{2}} + [(\Delta_{j}^{\pm} \phi) E_{k,\pm j}^{(2)}]_{-\frac{0}{2}} \right\} \\ & a_{\pm 0} \sum_{i} \Delta_{i}^{-} E_{i,\pm 0} = -\frac{1}{44} \sum_{i} \sum_{i} \sum_{i} \left( \Delta_{i}^{\pm} \phi_{\pm 0} \right) \left( B_{i}^{(4)} + B_{i,\pm 0}^{(4)} \right) ,  \text{(Gauss Law)} \end{split} $	$\begin{split} \left( \Delta_0^+ a_{-\hat{0}/2} \right)^2 &= \frac{a^2}{3m_{\rm pl}^2} \rho_L , \\ \Delta_0^- \Delta_0^+ a_{+\hat{0}/2} &= -\frac{a_{+\hat{0}/2}}{6m_{\rm pl}^2} (\rho_L + 3p_L)_{+\hat{0}/2} \end{split}$

$$\begin{split} \rho_L &= \bar{H}^{\rm kin} + \frac{1}{a^2} \frac{1}{2} (\bar{H}^{\rm grad}_{-\hat{0}/2} + \bar{H}^{\rm grad}_{+\hat{0}/2}) + \frac{1}{2} (\bar{H}^{\rm pot}_{-\hat{0}/2} + \bar{H}^{\rm pot}_{+\hat{0}/2}) + \frac{1}{a^2} \frac{1}{2} (\bar{H}^E_{-\hat{0}/2} + \bar{H}^E_{+\hat{0}/2}) + \frac{1}{a^4} \bar{H}^B \,, \\ (\rho_L + 3p_L)_{+\hat{0}/2} &= 2 (\bar{H}^{\rm kin} + \bar{H}^{\rm kin}_{+\hat{0}}) - 2 \bar{H}^{\rm pot}_{+\hat{0}/2} + \frac{2}{a^2_{+\hat{0}/2}} \bar{H}^E + \frac{1}{a^4_{+\hat{0}/2}} (\bar{H}^B + \bar{H}^B_{+\hat{0}}) \,, \end{split}$$

$$\begin{pmatrix} \bar{H}^{\rm kin} = \left\langle \frac{1}{N^3} \sum_{\vec{n}} \frac{\pi_{\phi}^2}{2} \right\rangle \qquad \bar{H}^{\rm grad} = \left\langle \frac{1}{N^3} \sum_{\vec{n}} \frac{1}{2} \sum_{i} (\Delta_i^+ \phi_{+\frac{\hat{0}}{2}})^2, \quad \bar{H}^{\rm pot} = \left\langle \frac{1}{N^3} \sum_{\vec{n}} \frac{1}{2} m^2 \phi_{+\frac{\hat{0}}{2}}^2 \right\rangle \\ \bar{H}^E = \left\langle \frac{1}{N^3} \sum_{\vec{n}} \sum_{i} \frac{1}{2} E_{i,+\frac{\hat{0}}{2}}^2 \right\rangle \quad \bar{H}^B = \left\langle \frac{1}{N^3} \sum_{\vec{n}} \sum_{i,j} \frac{1}{4} (\Delta_i^+ A_j - \Delta_j^+ A_i)^2 \right\rangle$$

#### **Gauss Constraint**





#### **Hubble Constraint**



0

 $N_e$ 

2

4

-2

4

 $10^{-17}$ 

 $V(\phi) = \frac{1}{2}m^2\phi^2 \; ; \; \frac{\phi}{4\Lambda}F\tilde{F} \; ; \; \Lambda = \frac{m_p}{15}$ 









Homogeneous ( — ) In-Homogeneous ( — )



#### **Example III**

## Non-minimally coupled Scalar fields in the Jordan Frame

#### with A. Florio, T. Opferkuch and B. Stefanek

SciPost, accepted ; 2112.08388 [astro-ph.CO]

#### Set-up

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_p^2}{2} R + \mathcal{L}_{\chi} + \mathcal{L}_{inf} + \frac{1}{2} \xi_{\phi} \phi^2 R \right]$$
  
or  $\frac{1}{2} \xi_{\chi} \chi^2 R$ 

#### $\bullet$ Inflaton $\phi$

 $\mathcal{L}_{\rm inf} = -\frac{1}{2} g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi - V_{\rm inf}(\phi)$ 

 $\bullet$  Spectator field  $\chi$ 

 $\mathcal{L}_{\chi} = -\frac{1}{2} g^{\mu\nu} \partial_{\mu} \chi \partial_{\nu} \chi - V(\chi, \phi)$ 

- Non minimal coupling to gravity  $\xi_{\phi}$  or  $\xi_{\chi}$
- Stay in Jordan frame

Set-up

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_P^2}{2} R + \mathcal{L}_{\chi} + \mathcal{L}_{inf} + \frac{1}{2} \xi_{\chi} \chi^2 R \right]$$
  
or  $\frac{1}{2} \xi_{\chi} \chi^2 R$  Spectator fld  
non-min. Coupled

 $\bullet$  Inflaton  $\phi$ 

 $\mathcal{L}_{\rm inf} = -\frac{1}{2} g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi - V_{\rm inf}(\phi)$ 

 $\bullet$  Spectator field  $\chi$ 

 $\mathcal{L}_{\chi} = -\frac{1}{2} g^{\mu\nu} \partial_{\mu} \chi \partial_{\nu} \chi - V(\chi, \phi)$ 

- Non minimal coupling to gravity  $\xi_{\phi}$  or  $\xi_{\chi}$
- Stay in Jordan frame



 $\nabla_{\mu}\nabla^{\mu}\chi + \frac{\partial V}{\partial\chi} + \xi_{\chi}\chi R = 0$ 



$$\nabla_{\mu}\nabla^{\mu}\chi + \frac{\partial V}{\partial\chi} + \xi_{\chi}\chi R = 0$$

$$R = F(\chi) \left( (1 - 6\xi) \langle \partial^{\mu} \chi \partial_{\mu} \chi \rangle + 4 \left( \langle V \rangle - \frac{3\xi}{2} \langle \chi V_{,\chi} \rangle \right) - \langle \rho_m \rangle - 3 \langle \rho_m \rangle \right)$$

$$F(\chi) = \frac{1}{M_{P}^{2} \left[1 + (6\xi - 1)\xi \langle \chi^{2} \rangle / M_{P}^{2}\right]}$$

Standard inflaton

 $V_{inf} \propto anh^4( ilde{\phi})$ 



#### **Curvature Oscillates !** (sourced by Inflaton Oscillations)

End

Standard inflaton

 $V_{inf} \propto anh^4( ilde{\phi})$ 



Geometric Preheating [Basset & Liberati '99]

 $\xi \chi^2 R$ ,  $\xi = 10,50,100$ 

The preheat field is excited exponentially

#### **Curvature Oscillates !** (sourced by Inflaton Oscillations)

### Geometric Preheating [Basset & Liberati '99]

$$\xi \chi^2 R$$
,  $\xi = 10,50,100$ 



#### How is the preheat field excited?

## Geometric Preheating [Basset & Liberati '99] $\xi \chi^2 R$ , $\xi = 10,50,100$ 10Ricci scalar $R/H^2$ and of inflatic 5

0

Time [e-folds post inflation]

2

4

How is the preheat field excited?



every time R < 0

Transparency worked out with Adrien Florio

0

-5

-4

Tachyonic growth

-2





#### Geometric Preheating [Basset & Liberati '99] How is the preheat field excited? $\xi \chi^2 R$ $\xi = 10,50,100$ **Non-Linear Regime Back-reaction** 15 $10^{2}$ $V_{\rm inf} = \Lambda^4 \tanh^4 \left( \frac{\alpha \phi}{M_{\rm e}} \right)$ $V_{\rm inf} = \Lambda^4 \tanh^4$ $= 100 = 10^{-5}$ Lattice simulation = 100Linear analysis = 50 10 $10^{-1}$ $\Delta_{\chi}(k,N)/m_p^2$ $R/H^2$ $10^{-4}$ 5 $10^{-7}$ 0 Lattice simulation achyoni $10^{-10}$ -51.50.51.02.02.5 $10^{-13}$ $10^{-2}$ e-folds post inflation $10^{-1}$

Transparency worked out with Adrien Florio

 $k/H_i$ 

#### **Geometric Preheating** [Basset & Liberati '99] $\xi \chi^2 R$ , $\xi = 10,50,100$

# How is the preheat field excited?

#### **Non-Linear Regime**



#### Transparency worked out with Adrien Florio

#### **Back-reaction**



#### **Geometric Preheating excitation (e.g. p = 4)**



#### **Geometric Preheating excitation (e.g. p = 4)**



#### Geometric Preheating excitation (e.g. p = 4)



#### **Geometric Preheating excitation (e.g. p = 4)**



#### We can do it for any **p**

 $V(|\phi|^p) \propto |\phi|^p$ 

#### **Full non-linear Geometric Preheating**



#### **Full non-linear Geometric Preheating**



#### We can compare Jordan vs Einstein fram