

Accelerator based dark matter probes

Kate Pachal TRIUMF

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How we might think about framing this to the community

Today, largely talking about cases with dark matter and/or associated new particles over $\sim 1~\text{GeV}$ in mass

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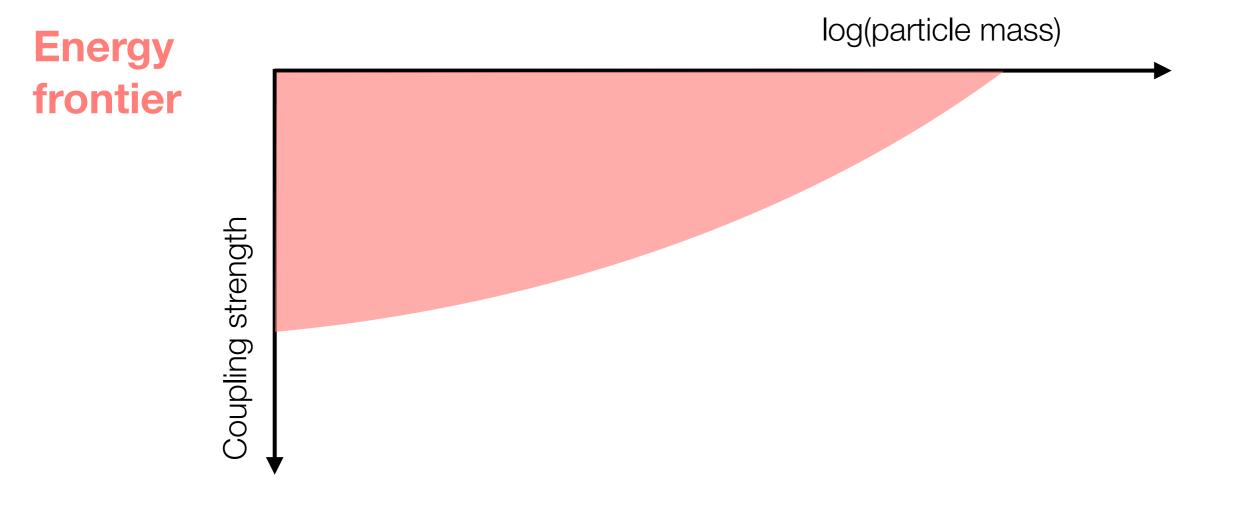
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Does not mean "vanilla" WIMPs only! But light dark matter, and with it light mediators, are easier at intensity frontier experiments

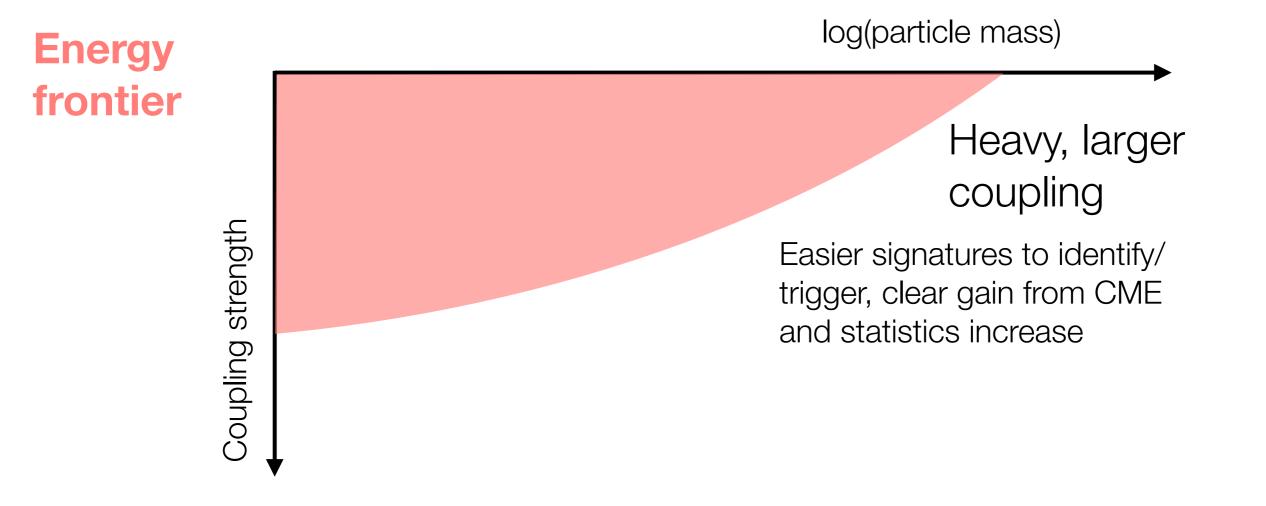
log(particle mass)

Coupling strength

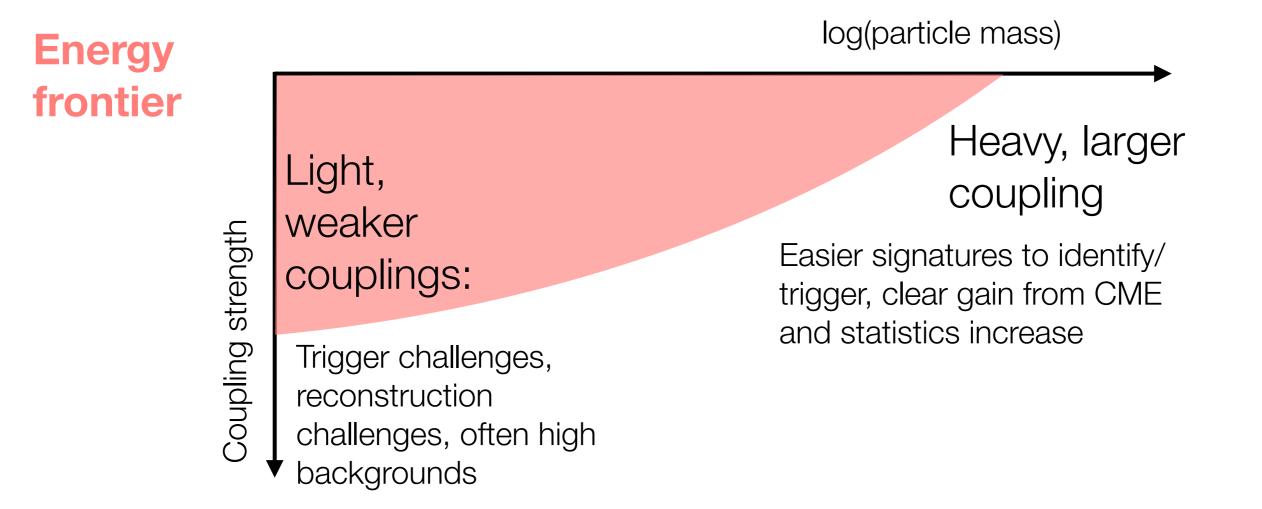
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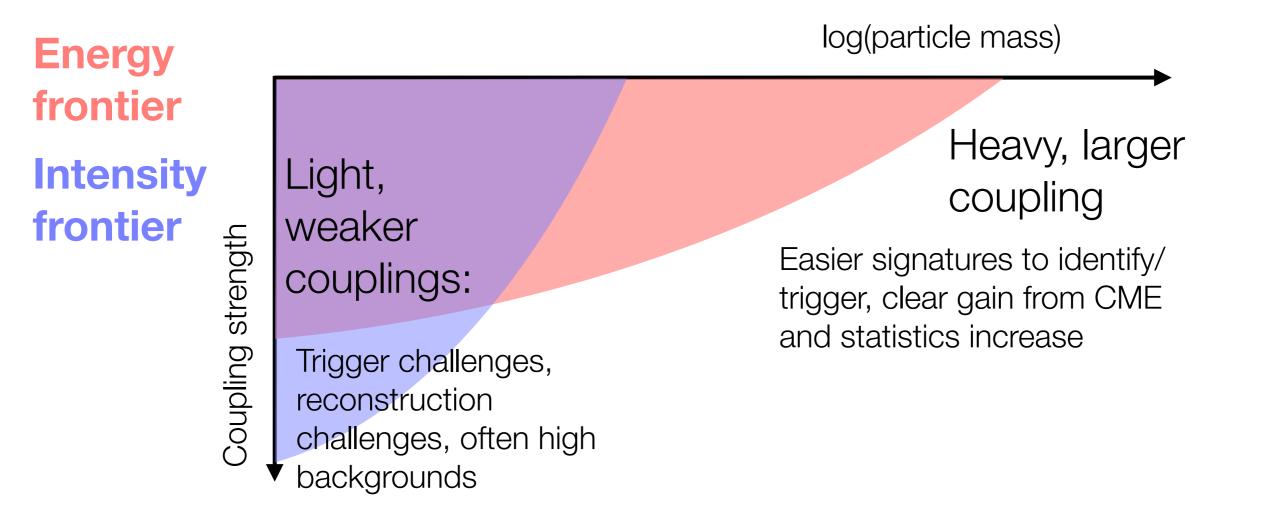
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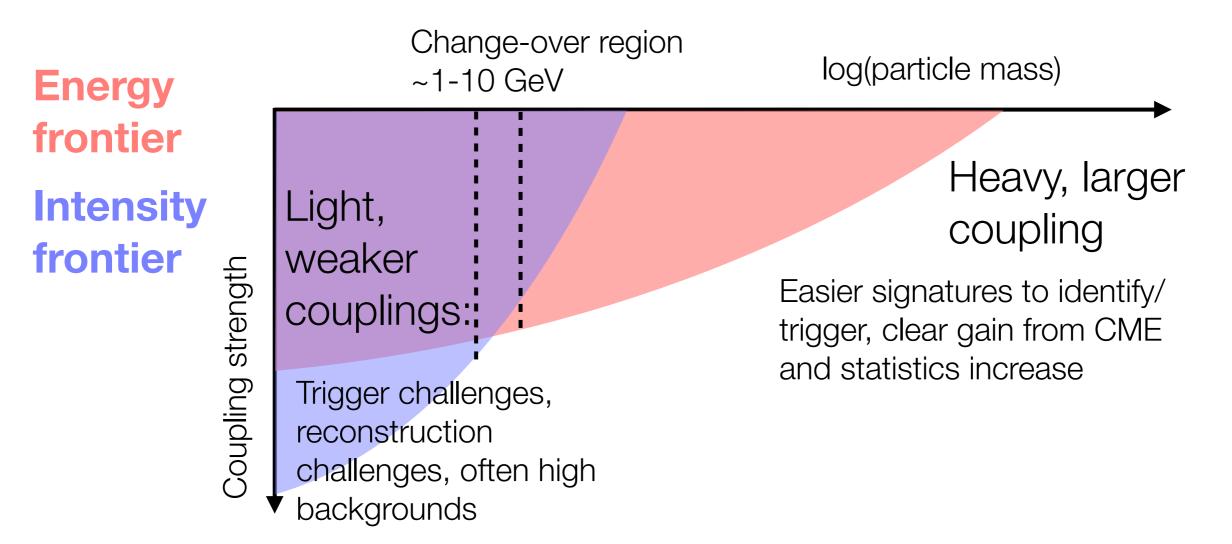
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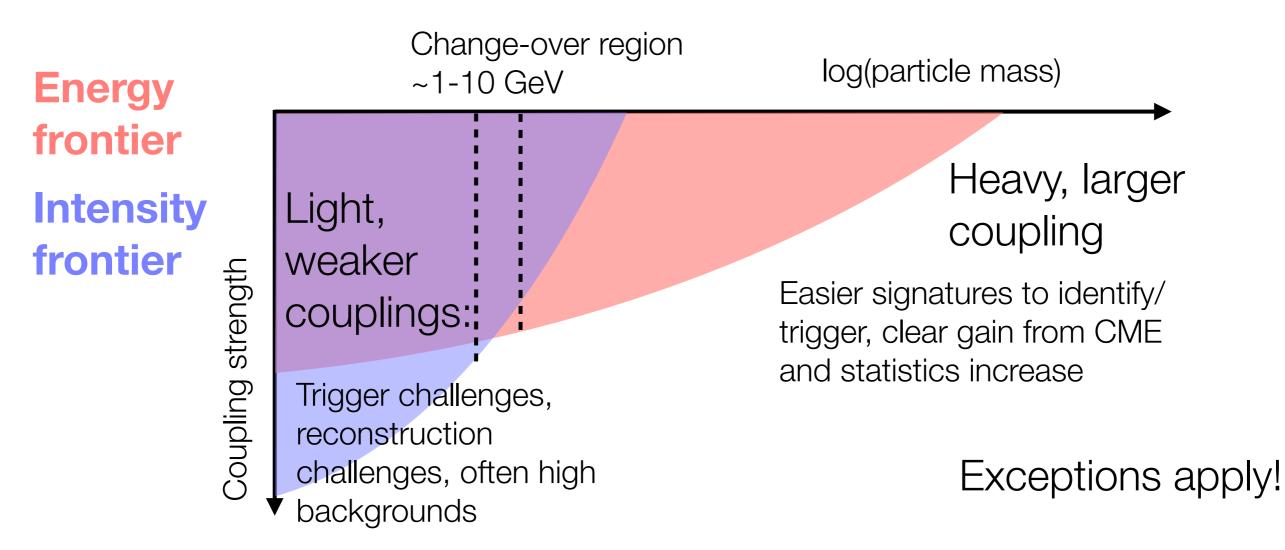
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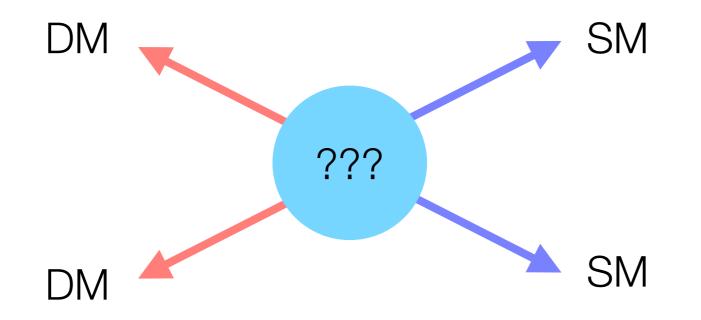


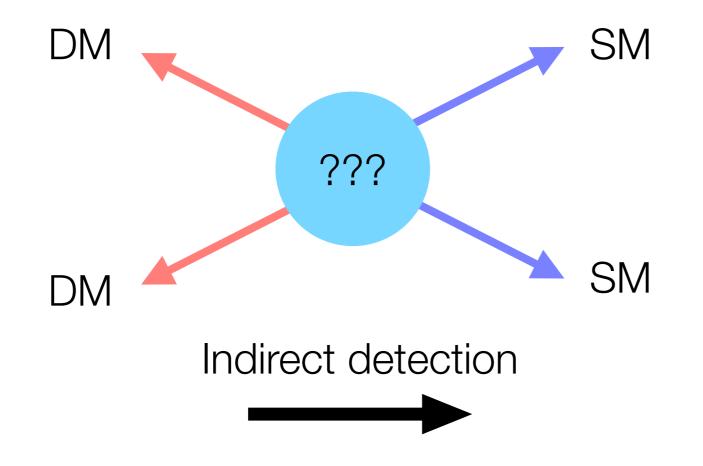
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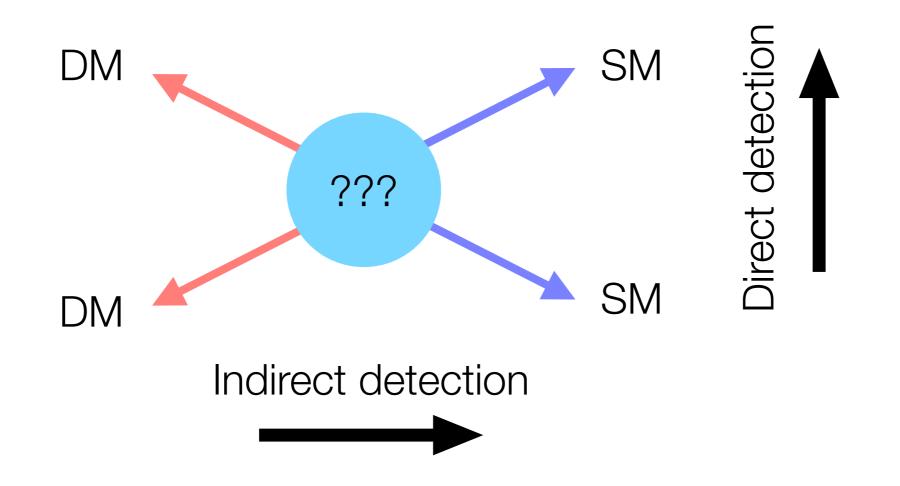


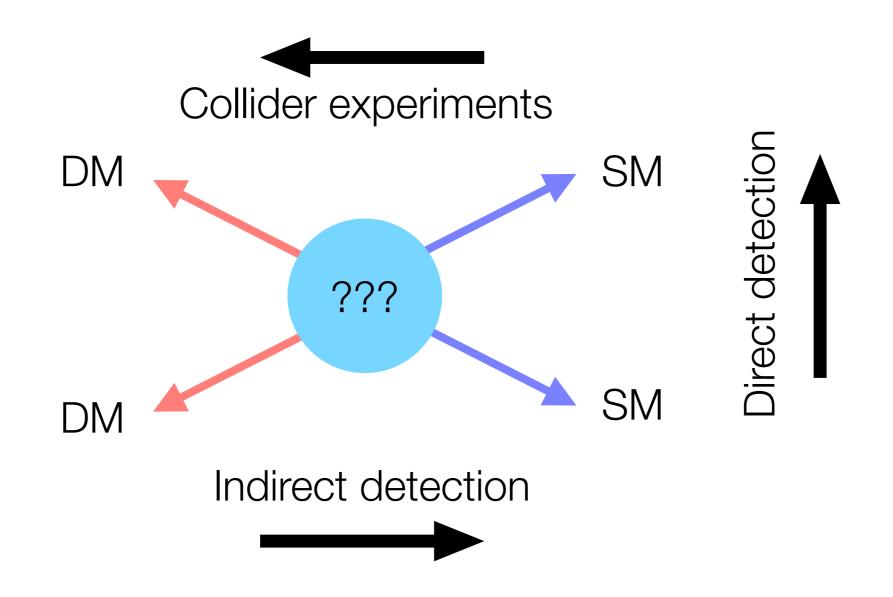
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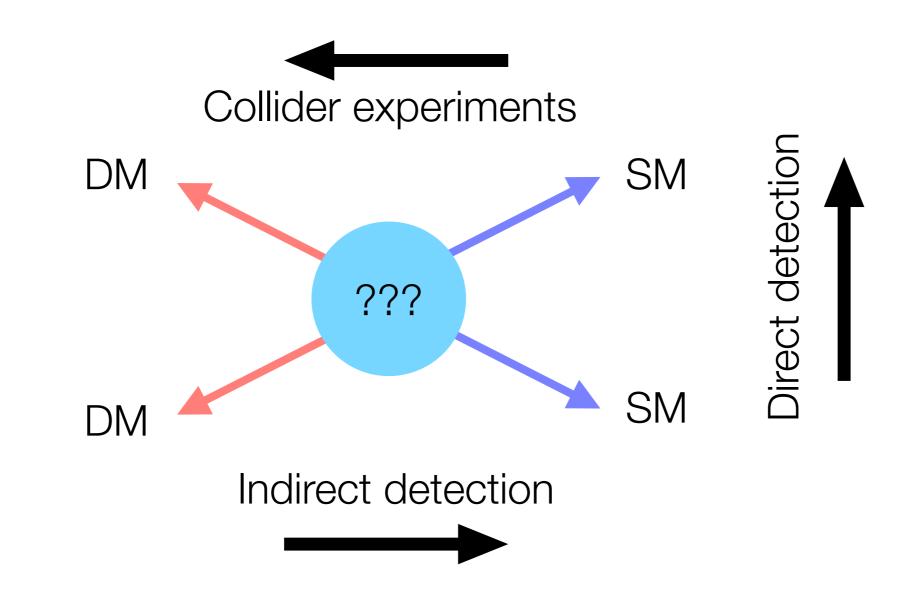




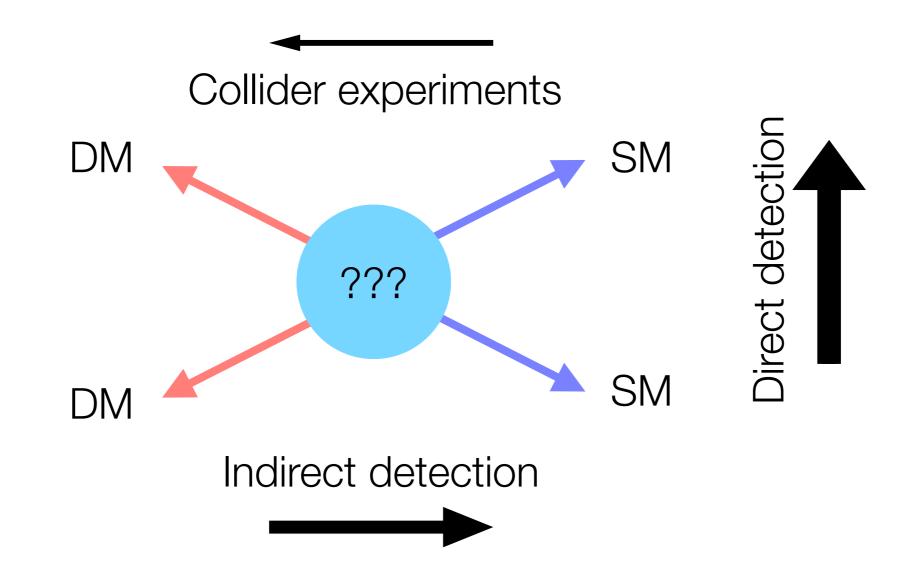






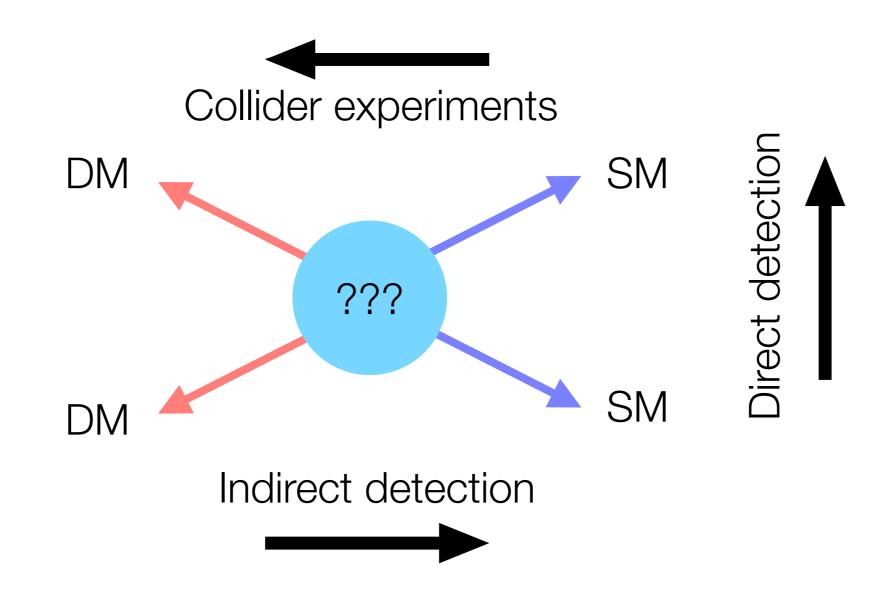


Tired: all three approaches are probing the same thing (interchangeable)



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Wired: different DM scenarios may be accessible to only one or two of the three approaches



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Wired: different DM scenarios may be accessible to only one or two of the three approaches

Inspired: the future of the field needs all three to ensure success

Benchmarks

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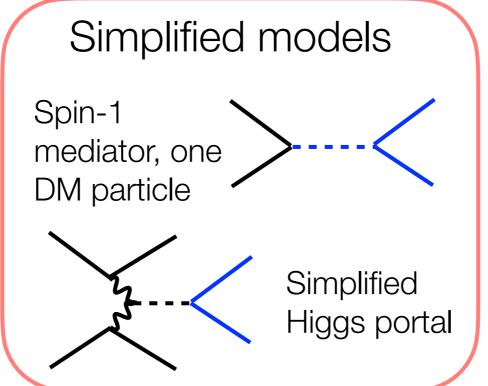
Dark matter models at the LHC



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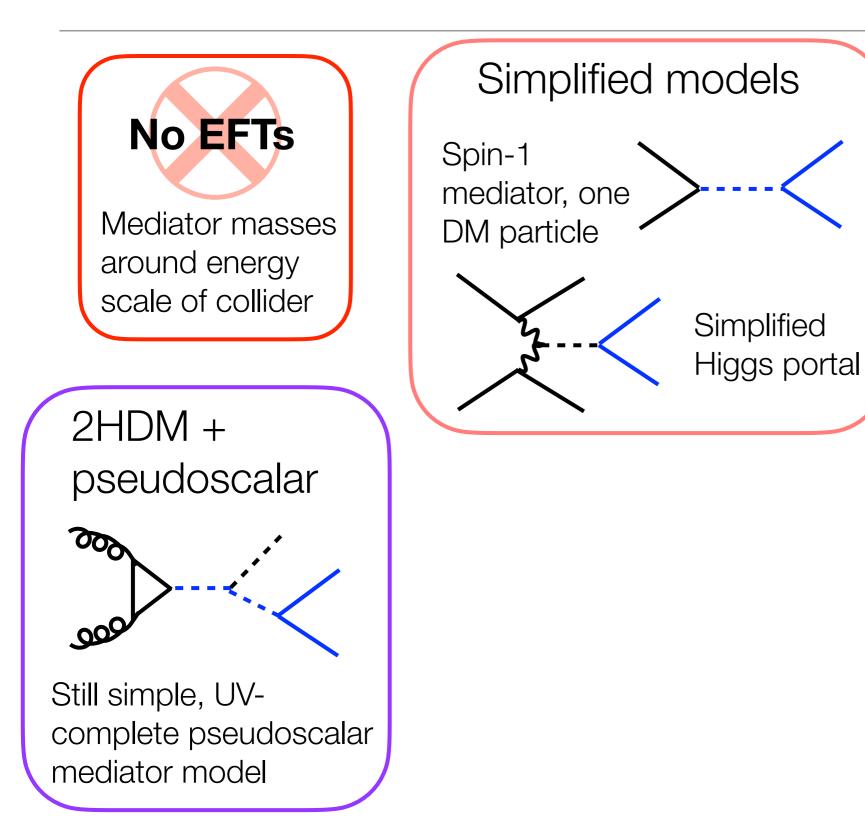
Standard Model: black BSM: blue

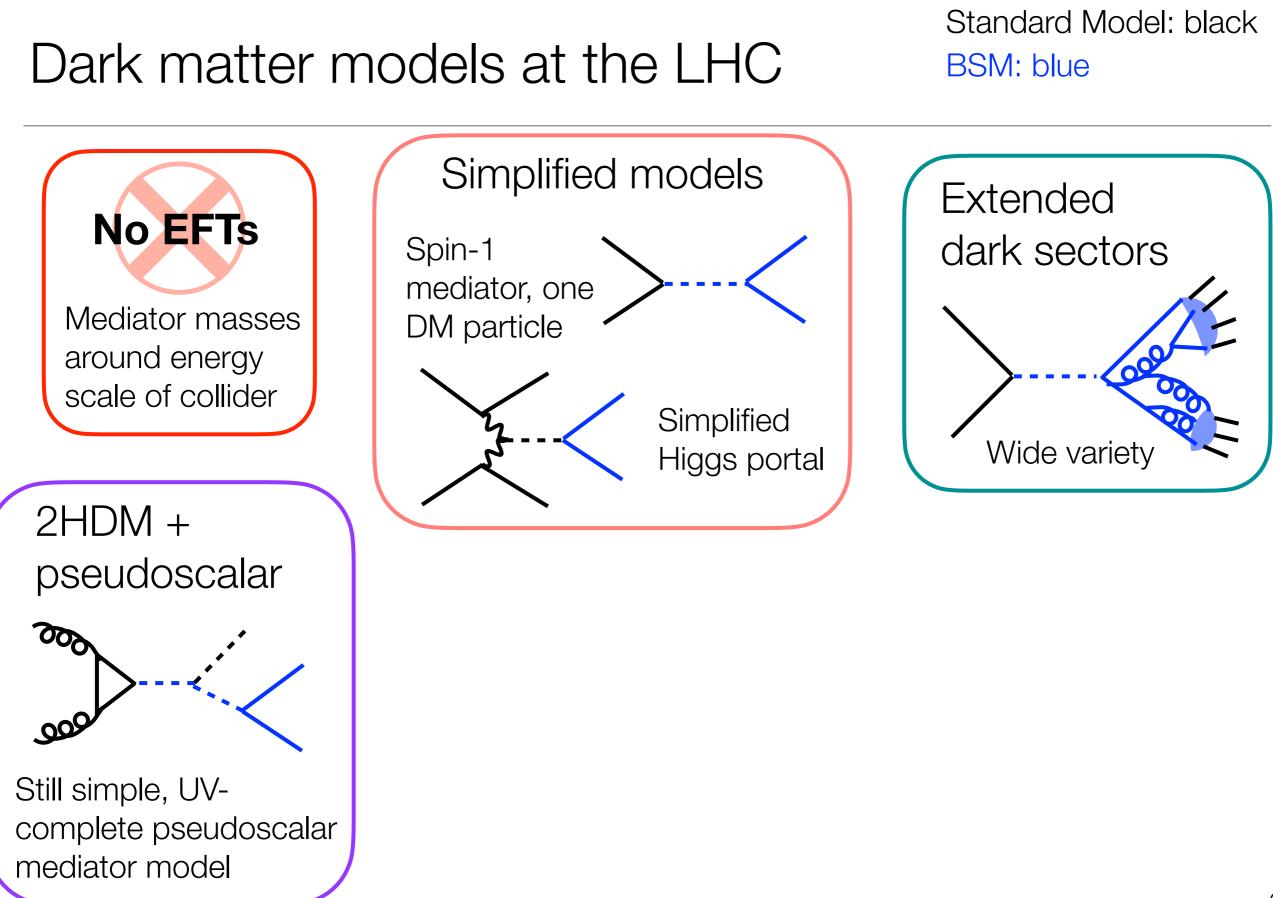
NO EFTS Nediator masses around energy scale of collider

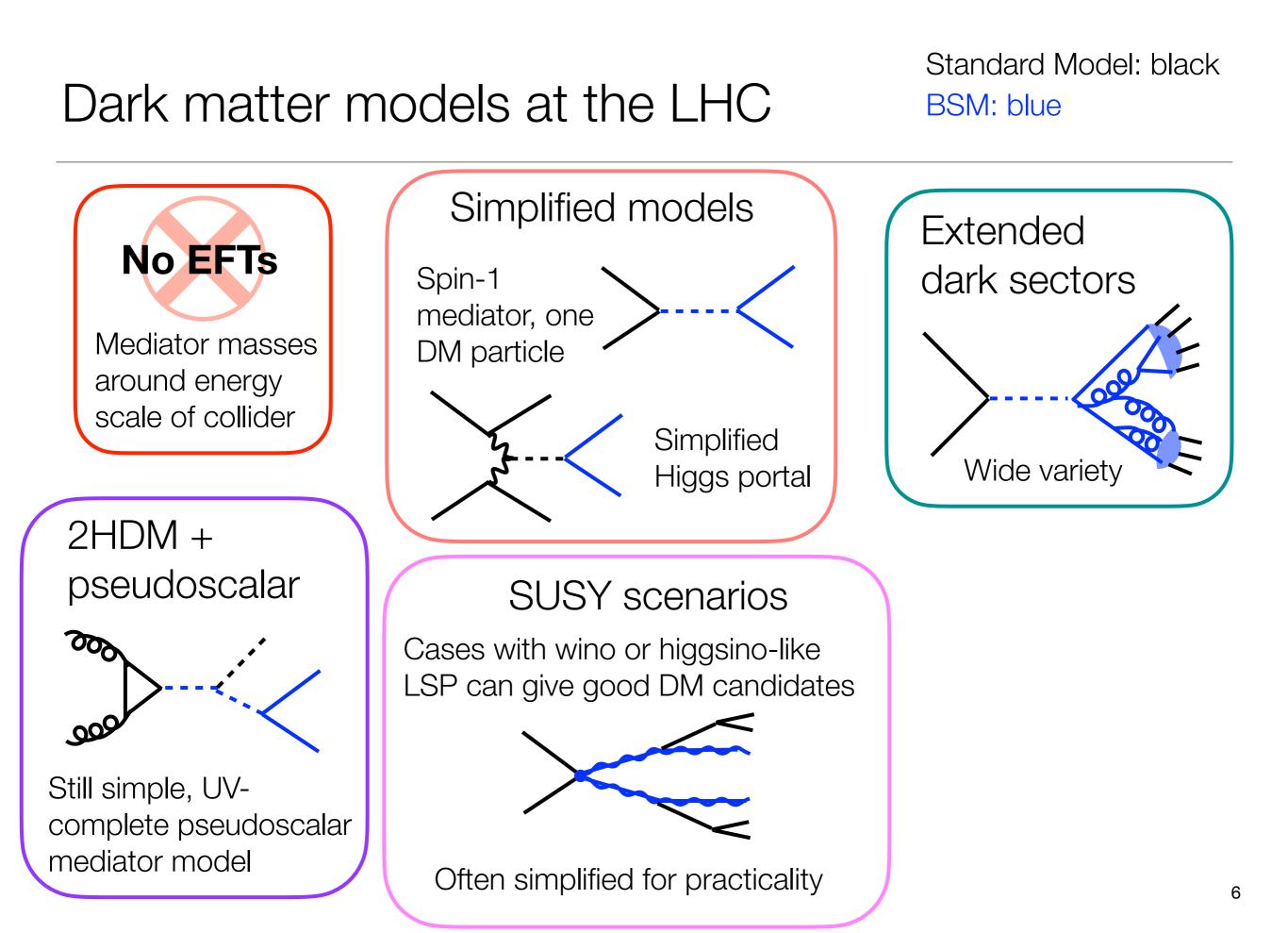


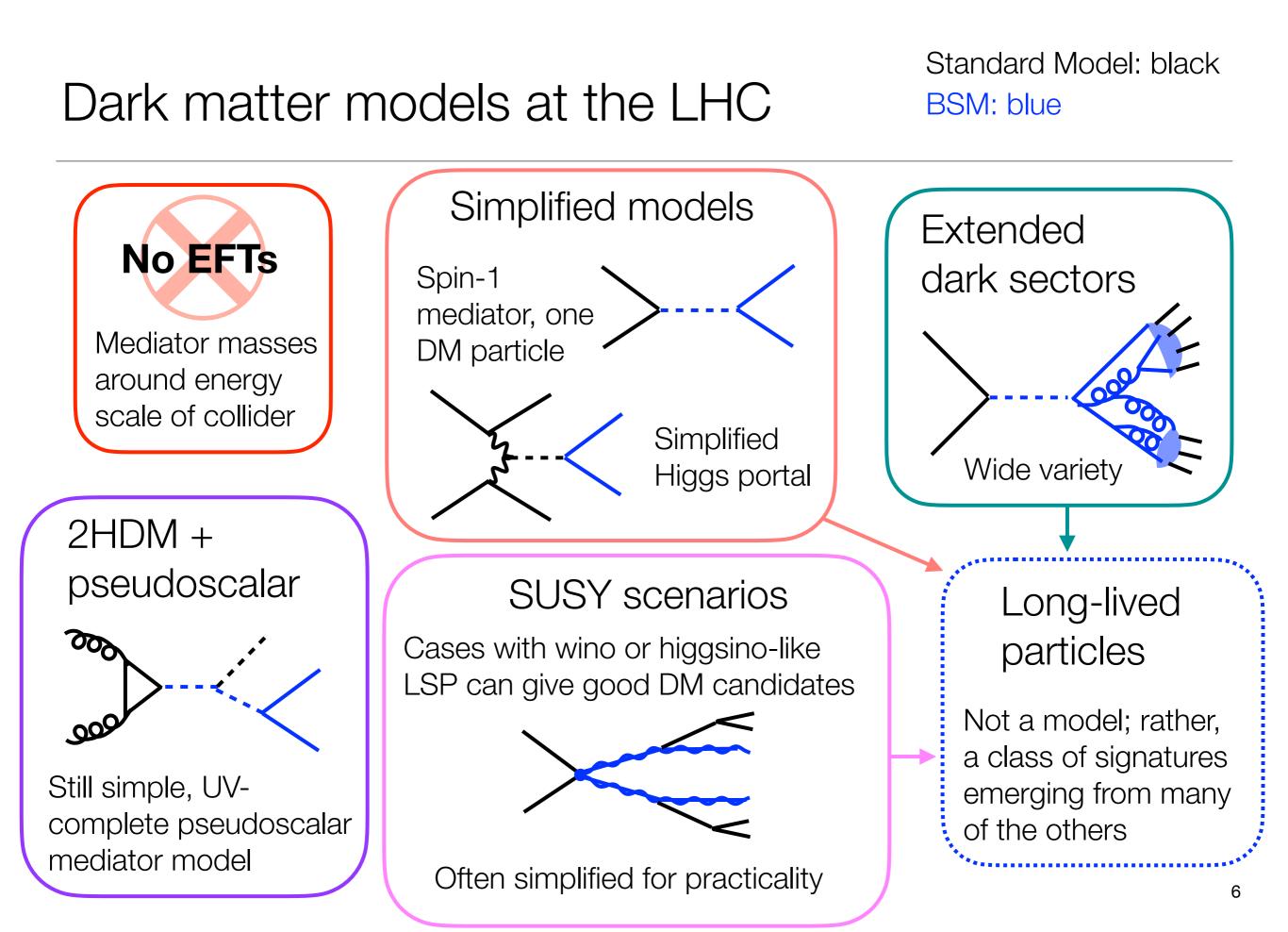
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* speaking as an experimentalist

Pros and cons of different benchmarks*

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Complete/ complex models

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No single answer. ATLAS & CMS lean on simplified models for comparisons; use complex models on analysis-by-analysis basis and for smaller comparison use cases

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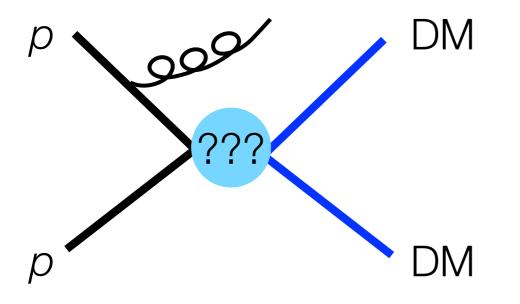
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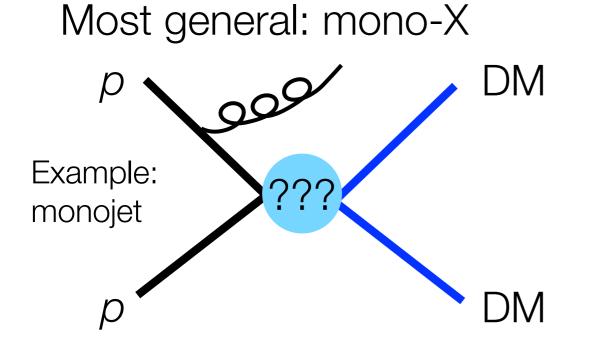
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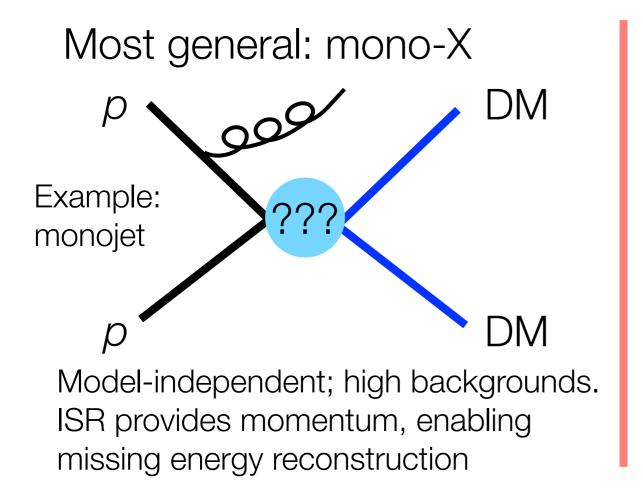
Could say a model is excluded once relic prediction reached

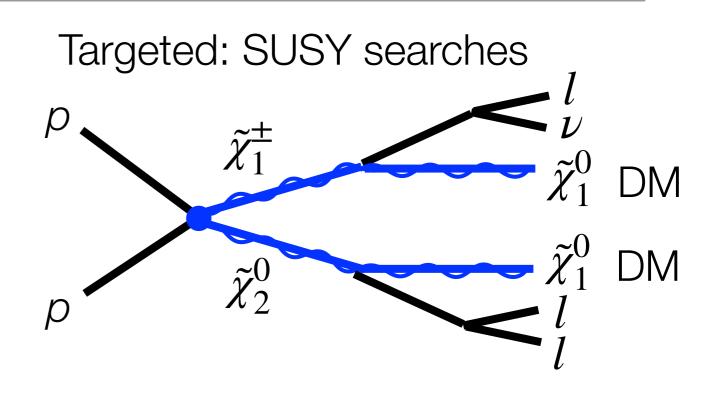
DM at the LHC

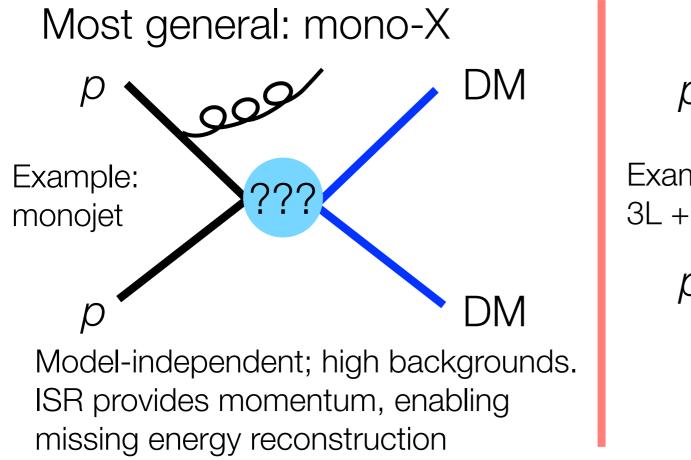
Most general: mono-X

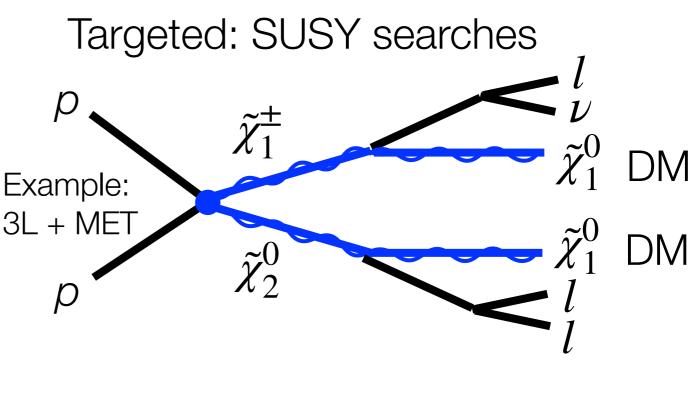


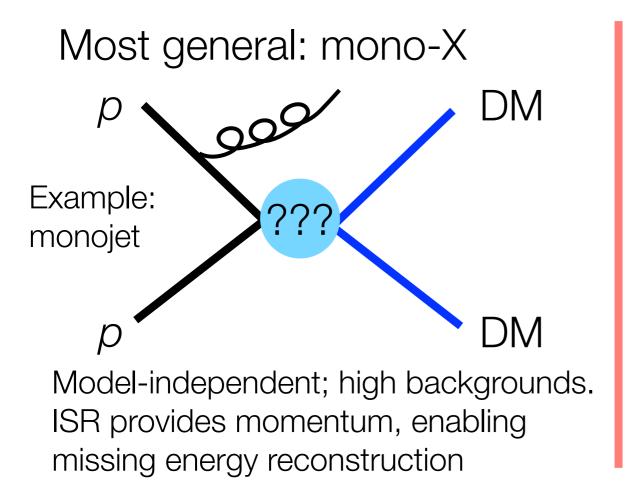


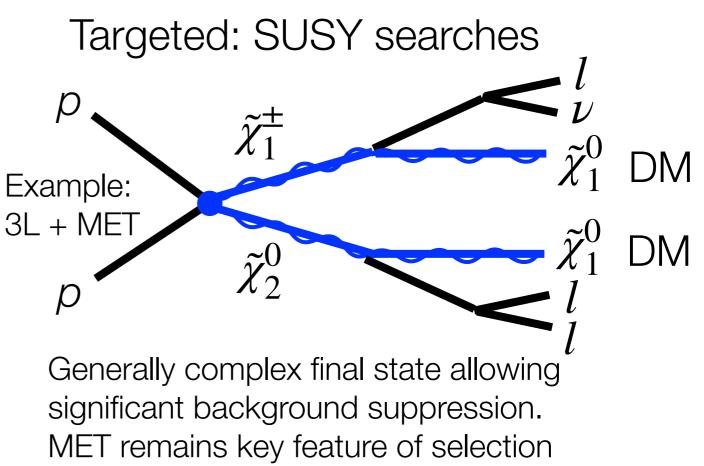


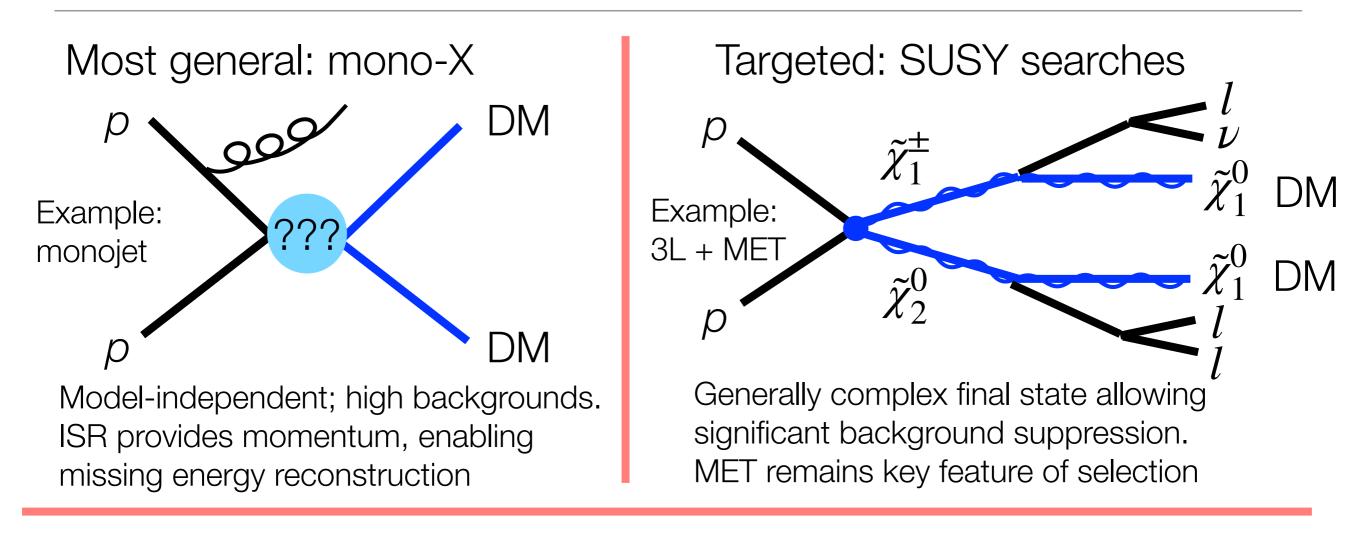




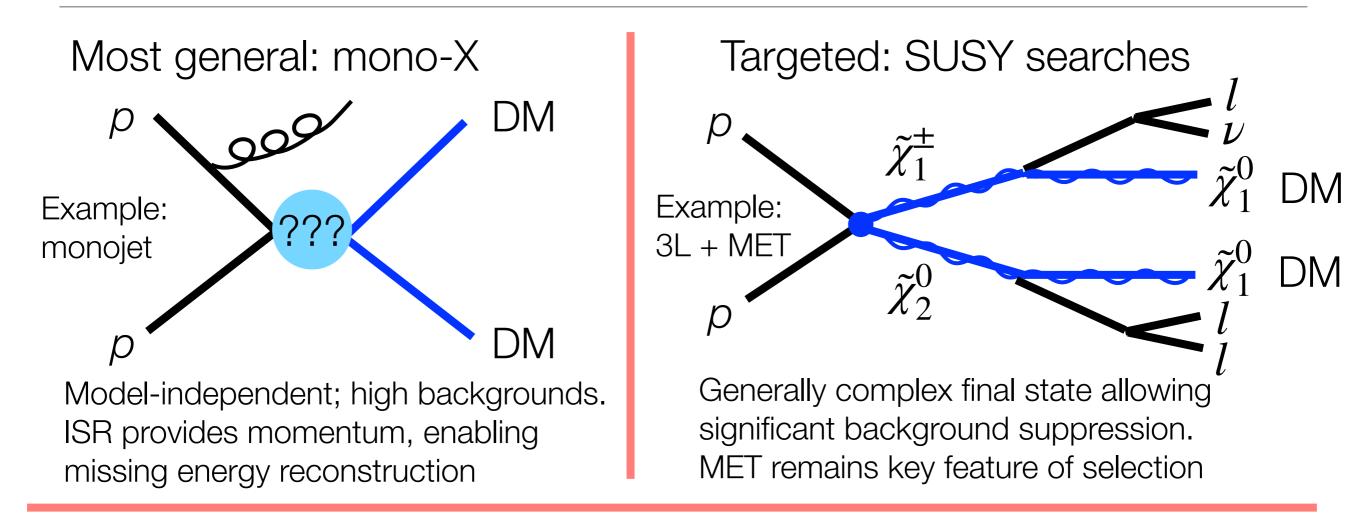






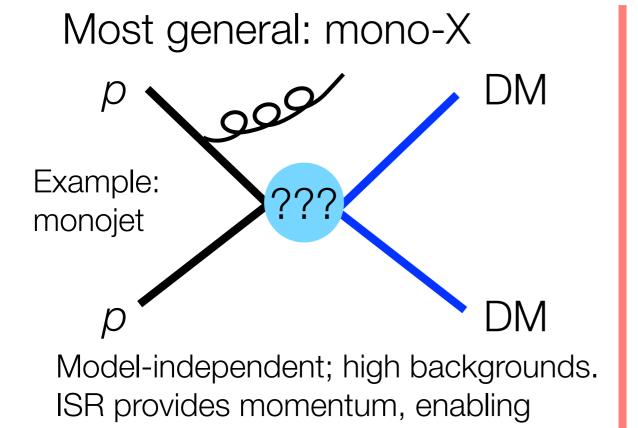


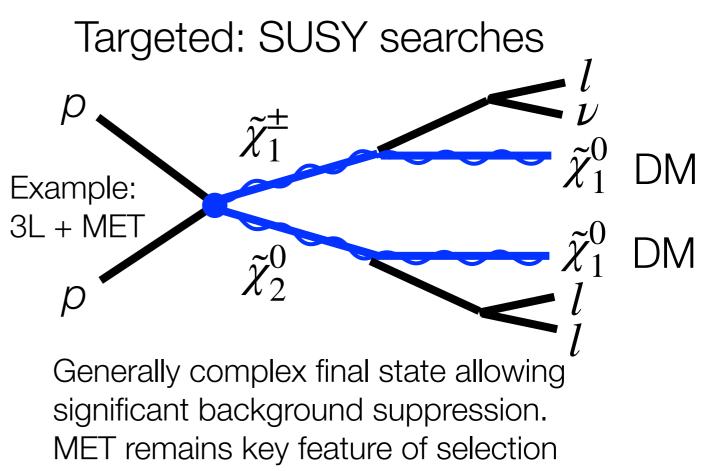
Non-MET-focused



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Various searches target models with dark matter implications, but that do not rely on MET in final state. Extended dark sectors, direct mediator searches, LLPs

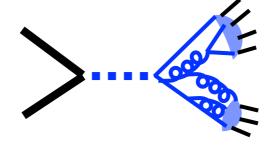




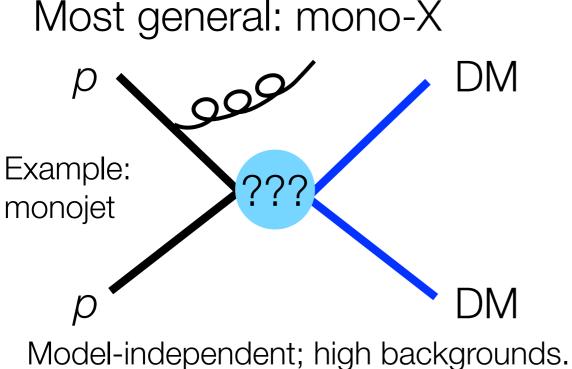
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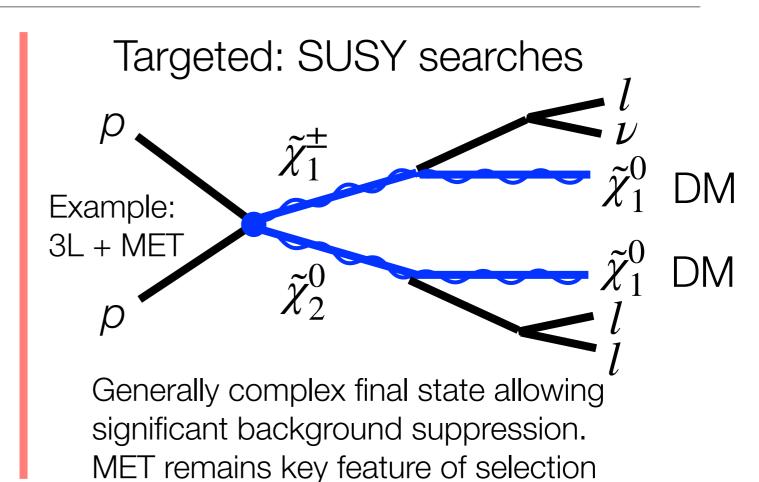
missing energy reconstruction



QCD final states with distinctive features

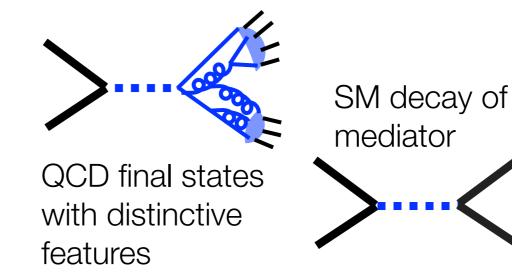


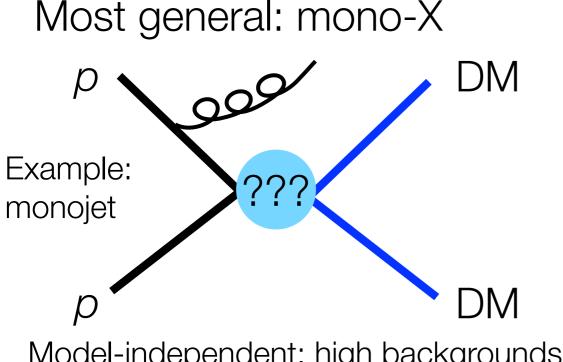
ISR provides momentum, enabling missing energy reconstruction



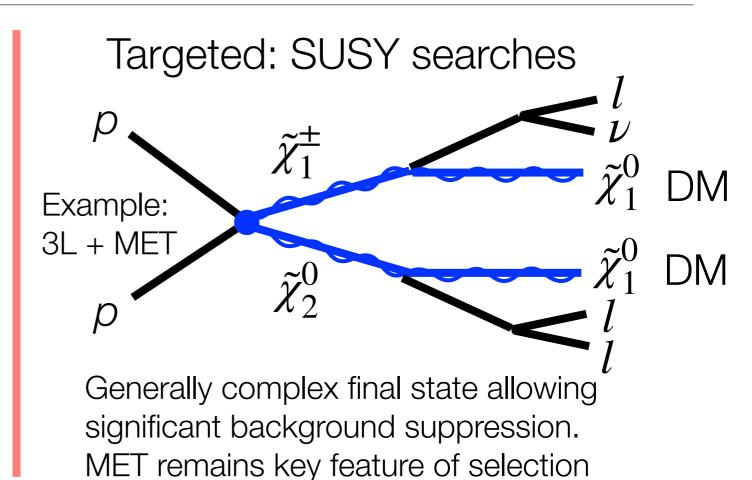
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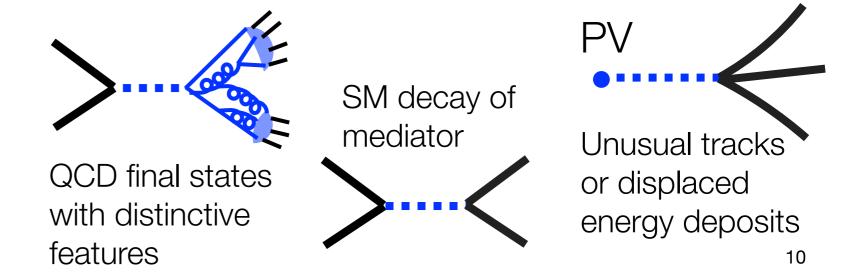


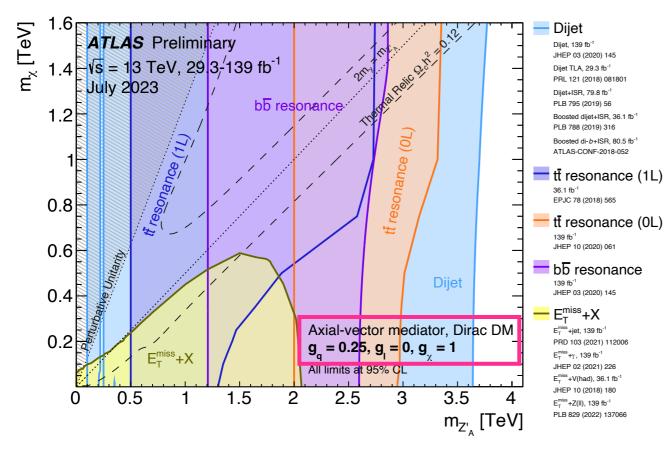
Model-independent; high backgrounds. ISR provides momentum, enabling missing energy reconstruction

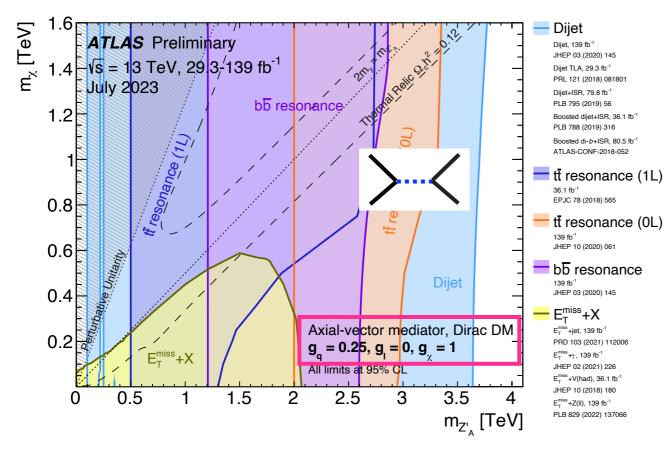


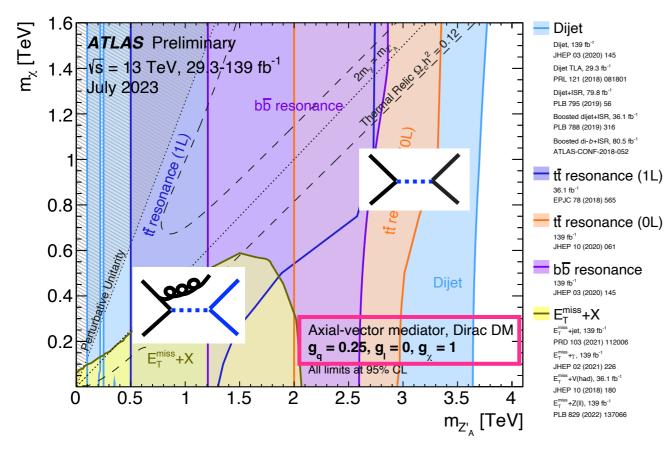
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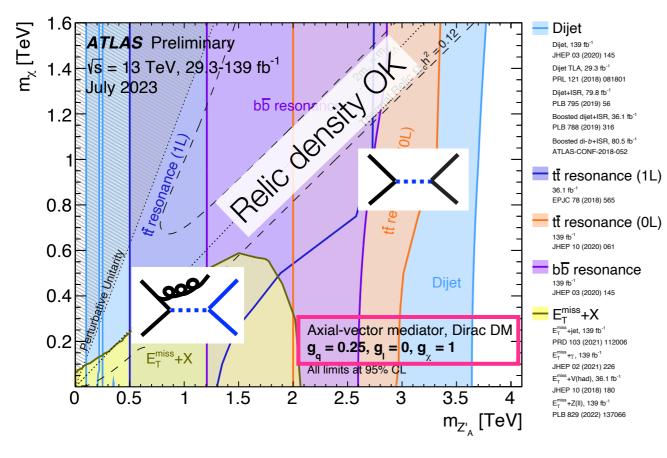
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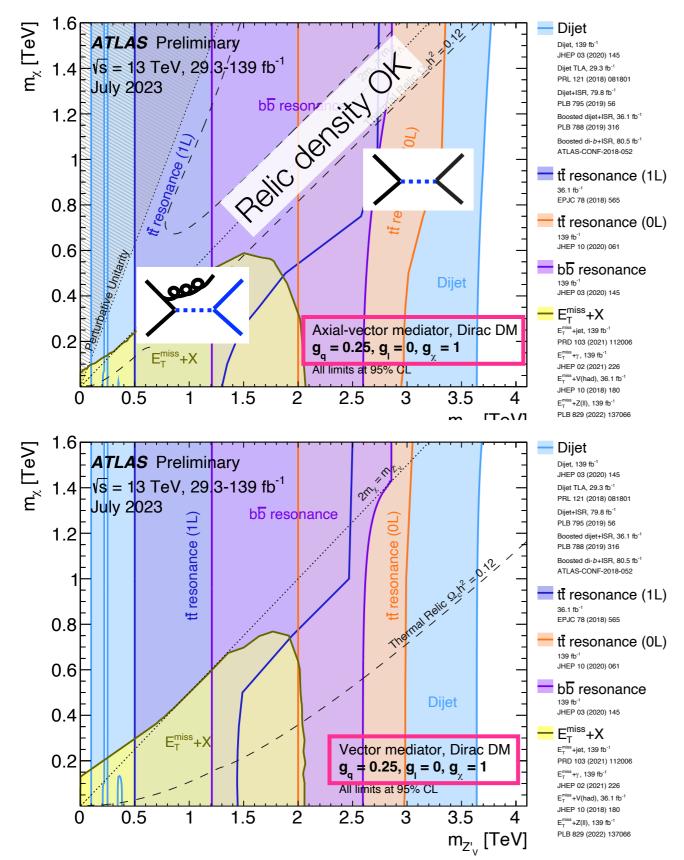


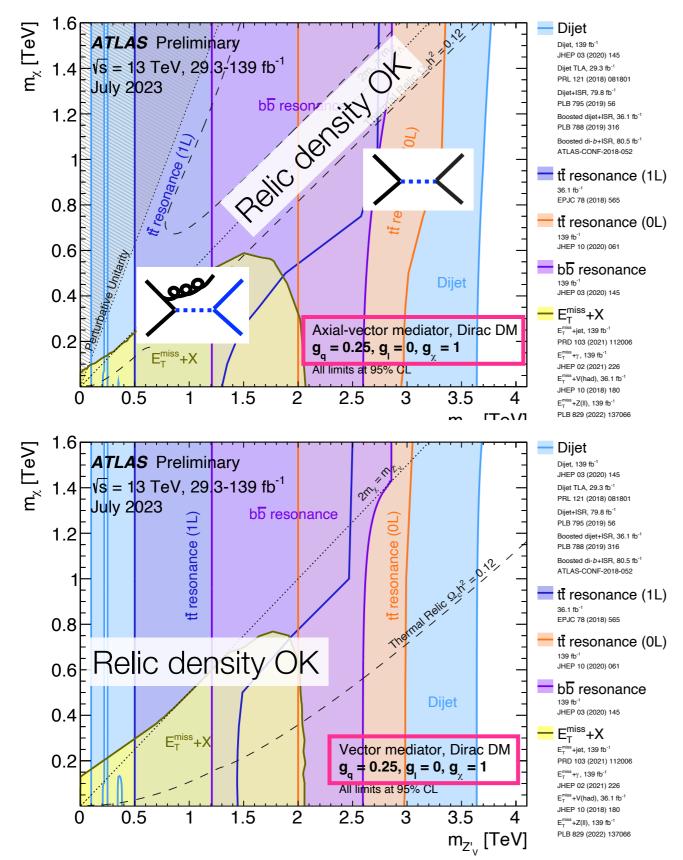


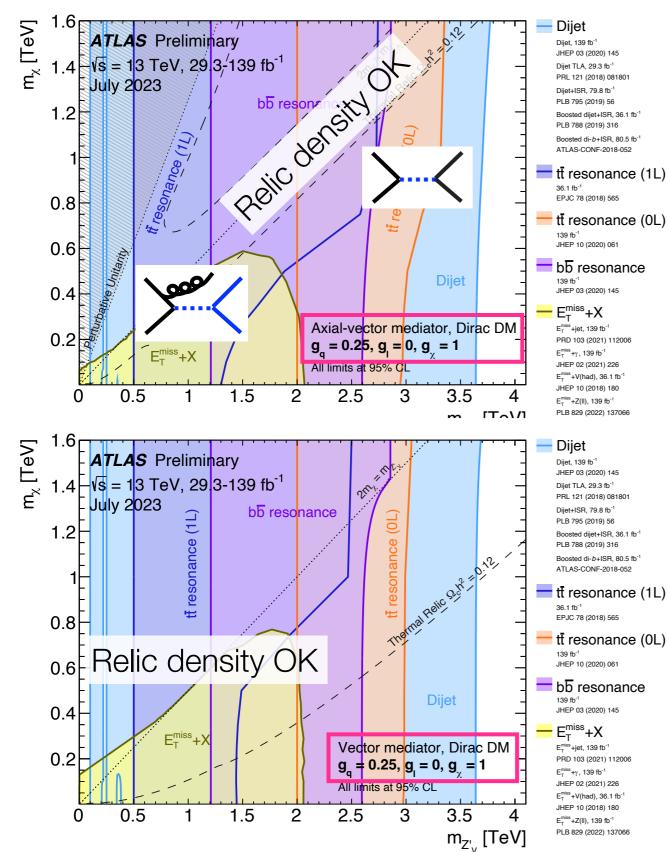


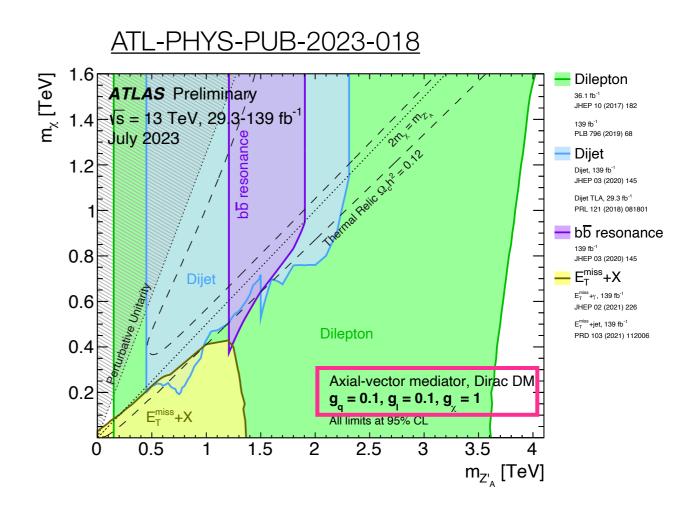


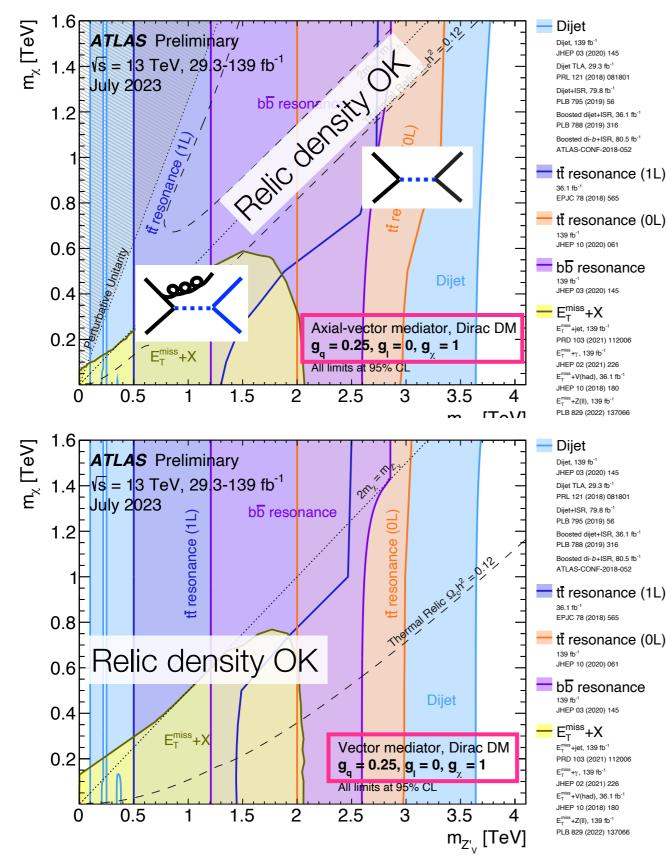


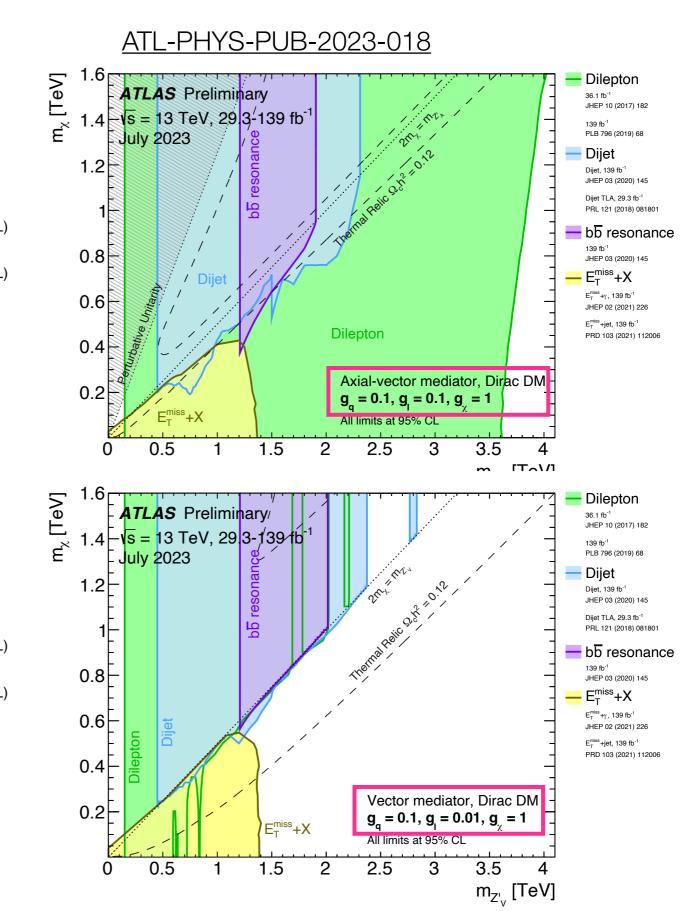


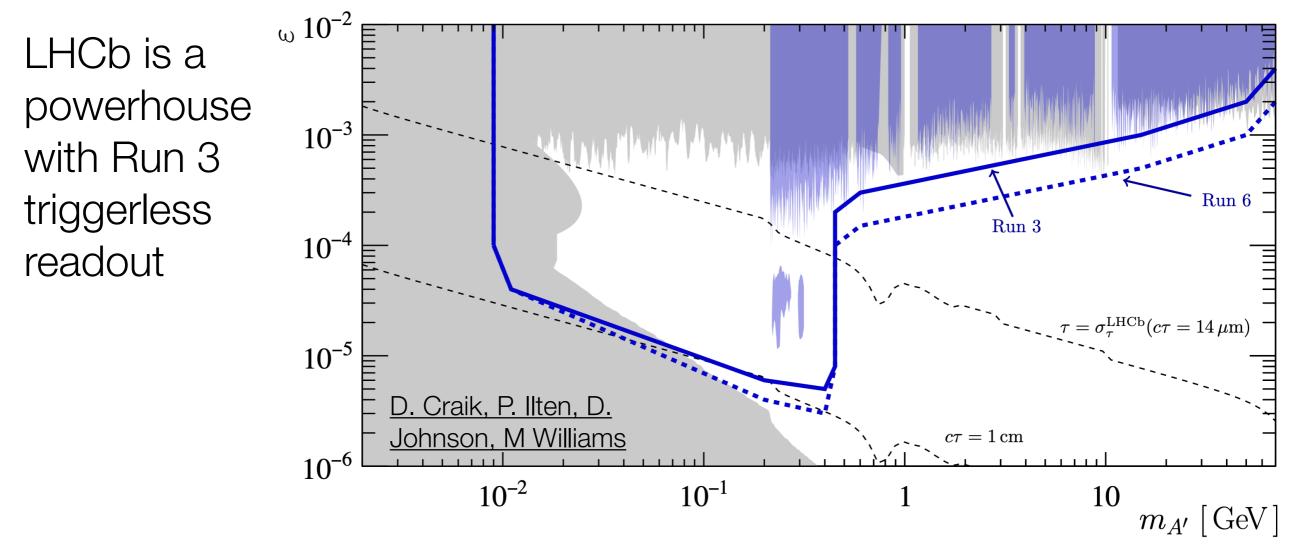


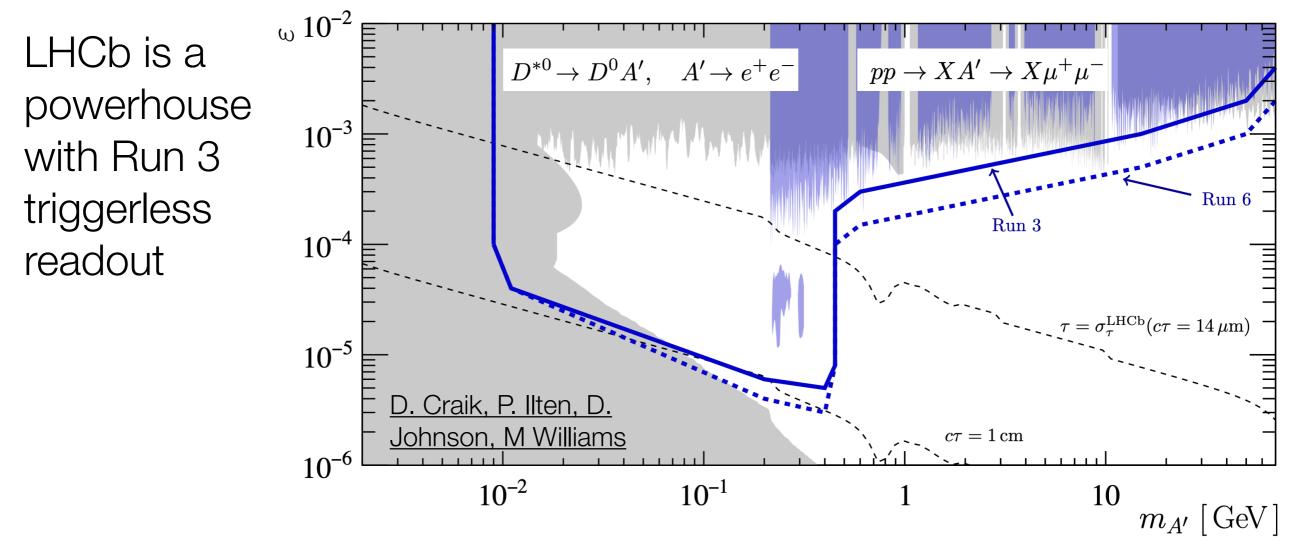


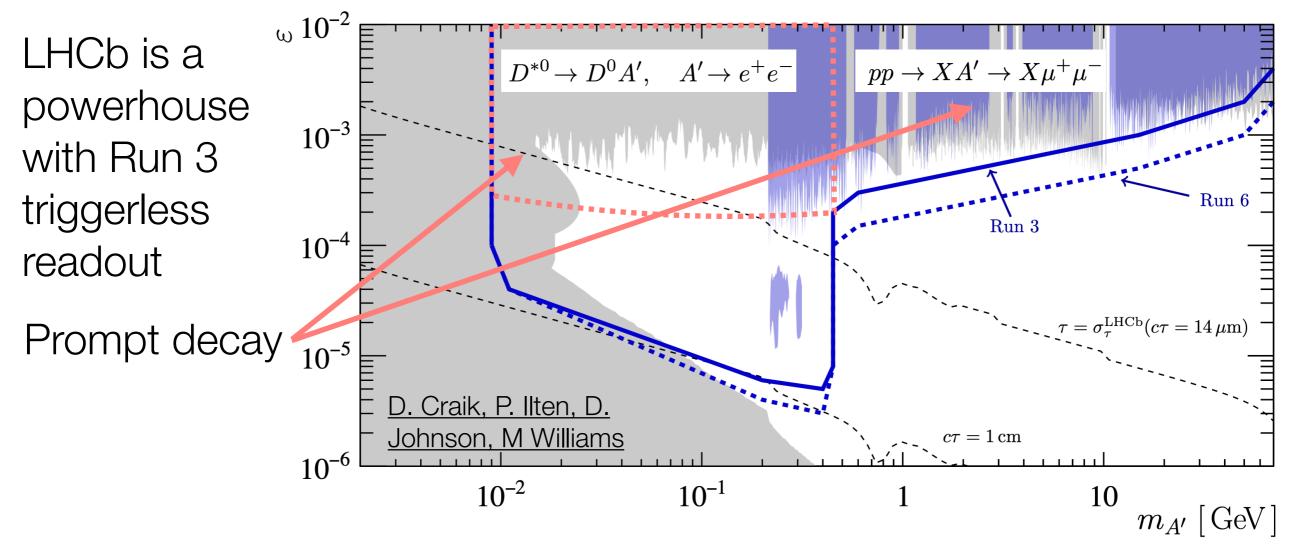


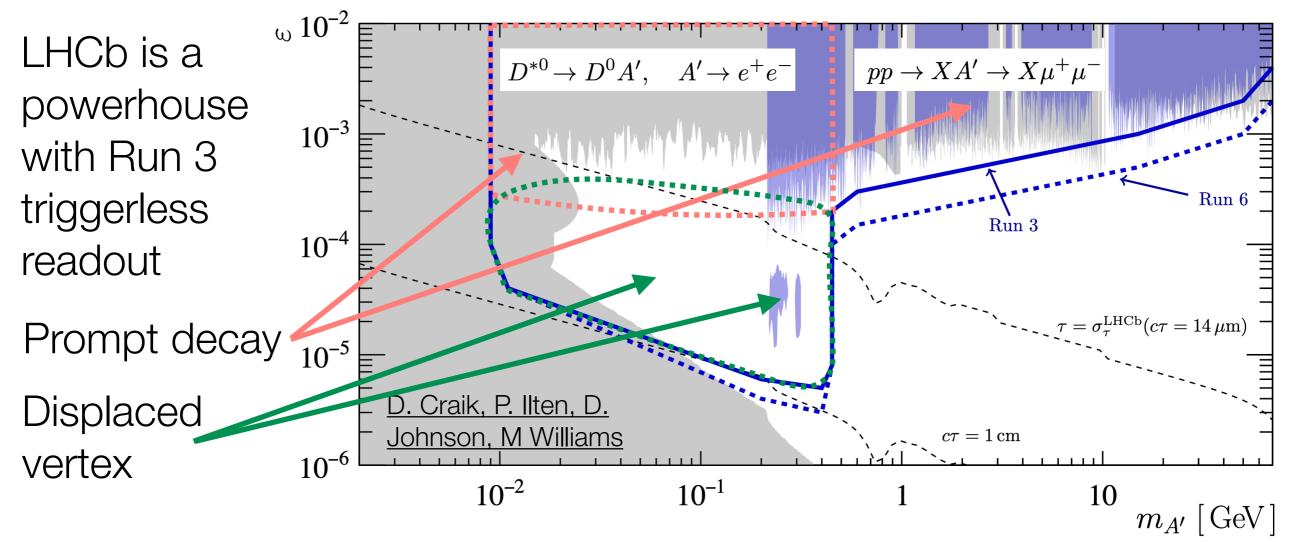




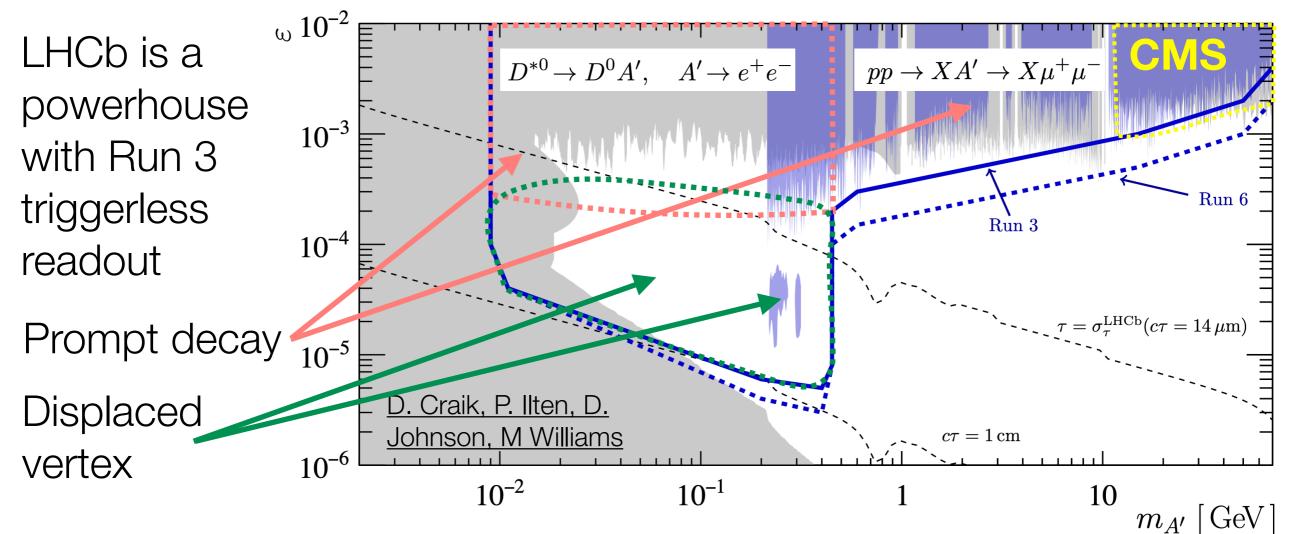








Very popular spin-1 vector benchmark, especially with intensity frontier and physics beyond colliders community



ATLAS & CMS can contribute at higher masses. Trigger poses a challenge. Simplified spin-1 limits translate fairly directly, but this is not currently a standard interpretation.

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4.9 fb⁻¹ (7 TeV), 19.7 fb⁻¹ (8 TeV), 140 fb⁻¹ (13 TeV) In Higgs portal models, the 10⁻³³ $\sigma_{\rm DM-nucleon}^{\rm SI}$ (cm²) Higgs decays to DM, **Higgs portal models CMS** 10^{-35} 90% CL limits creating a MET signature Majorana fermion DM 10^{-37} $B(H \rightarrow inv) < 0.14$ Scalar DM 10⁻³⁹ Vector DM UV-comp Vector DM radiative m₂ = 100 GeV **10**⁻⁴¹ Vector DM radiative m₂ = 65 GeV 10⁻⁴³ 10⁻⁴⁵ 10^{-47} **Direct-detection** 10⁻⁴⁹ KENON1T-MIGDAL 10⁻⁵¹ DarkSide-50 PandaX-4T 10⁻⁵³ LUX-ZEPLIN 10⁻⁵⁵ 10² 10^{3} **10**⁻¹

CMS-HIG-21-007

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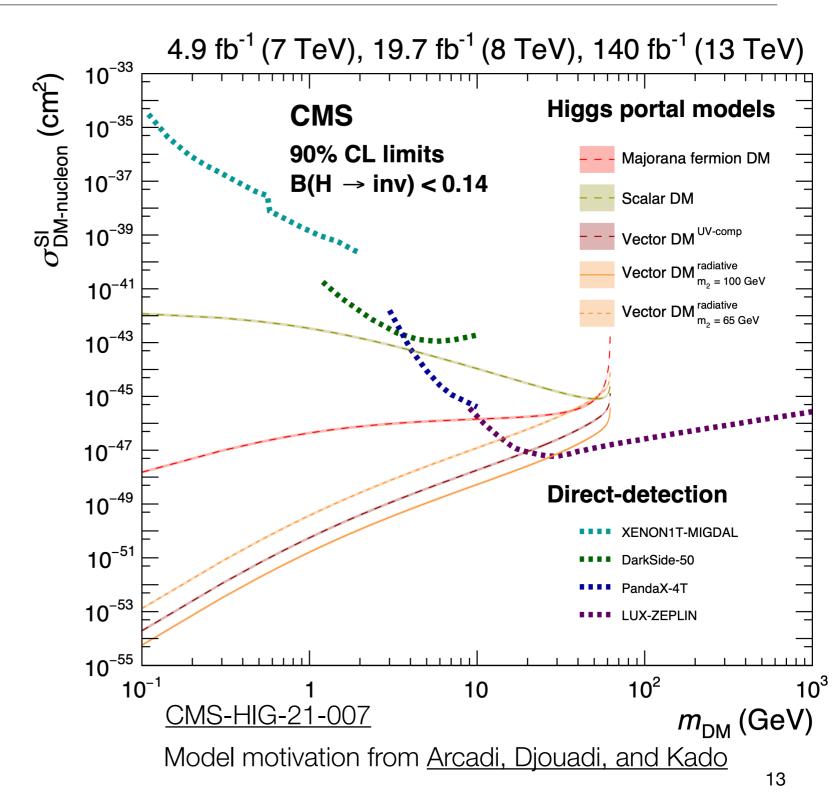
 $m_{\rm DM}\,({\rm GeV})$

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Model motivation from Arcadi, Djouadi, and Kado

In Higgs portal models, the Higgs decays to DM, creating a MET signature

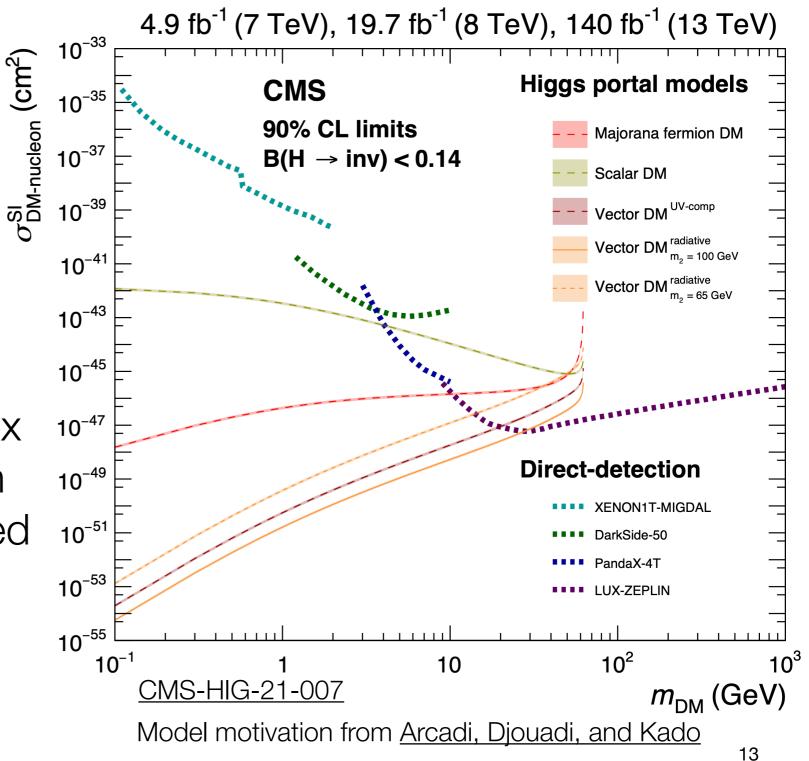
Possible UV-complete SM extension with just one DM particle if DM is a scalar



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For vector DM, more complex scenario with dark Higgs can still be appropriately estimated via this EFT approach (<u>ref</u>.)

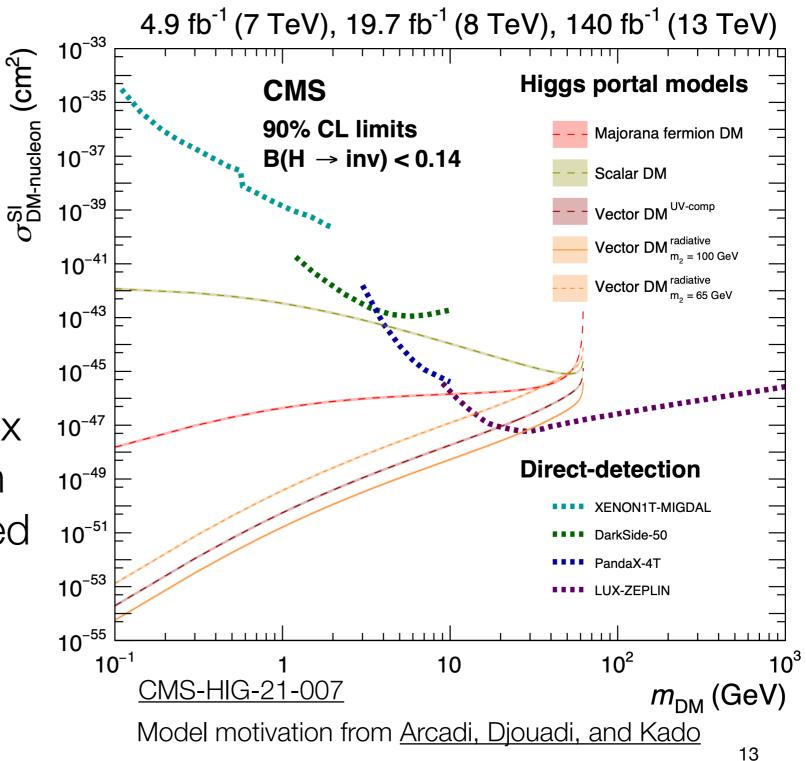


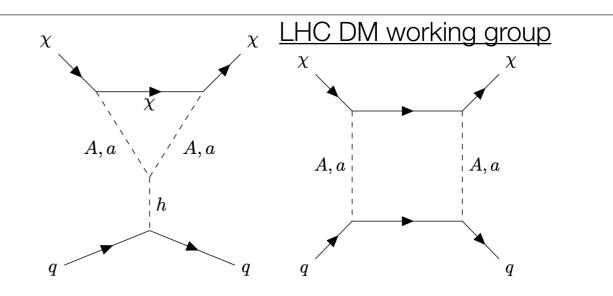
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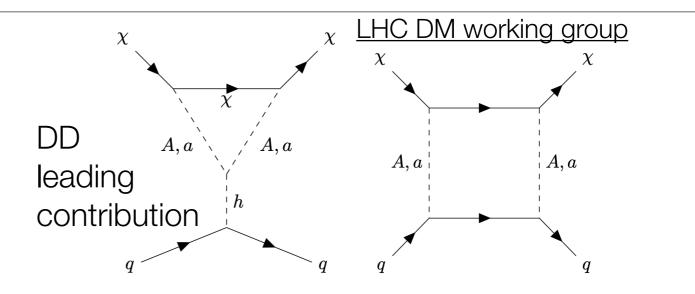
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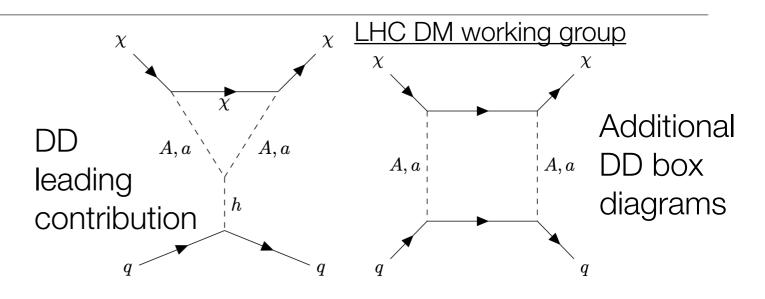
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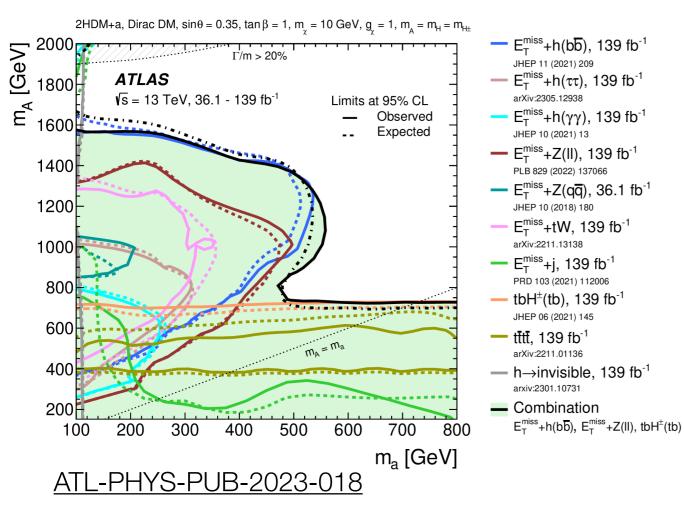
Current upper limits on $BR(h \rightarrow inv) \sim 0.11 (ATLAS)$

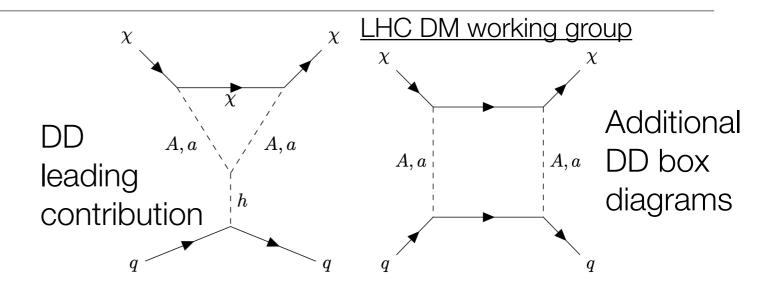


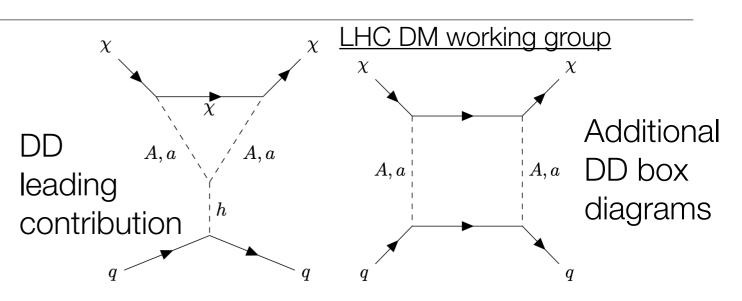


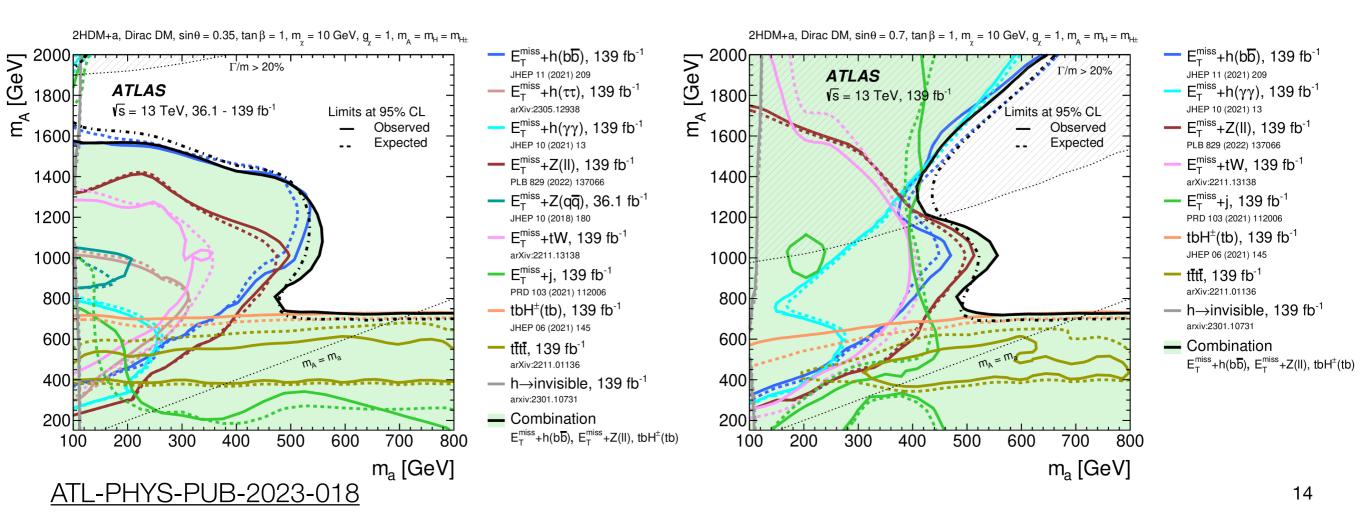


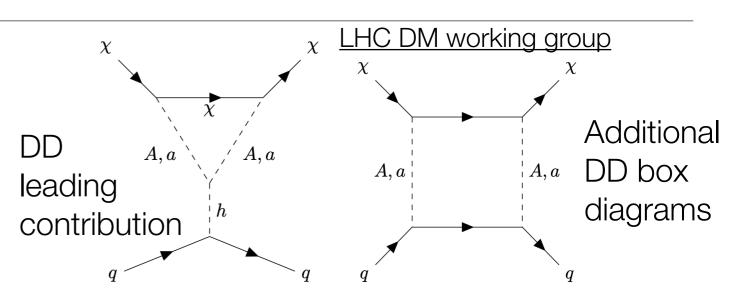


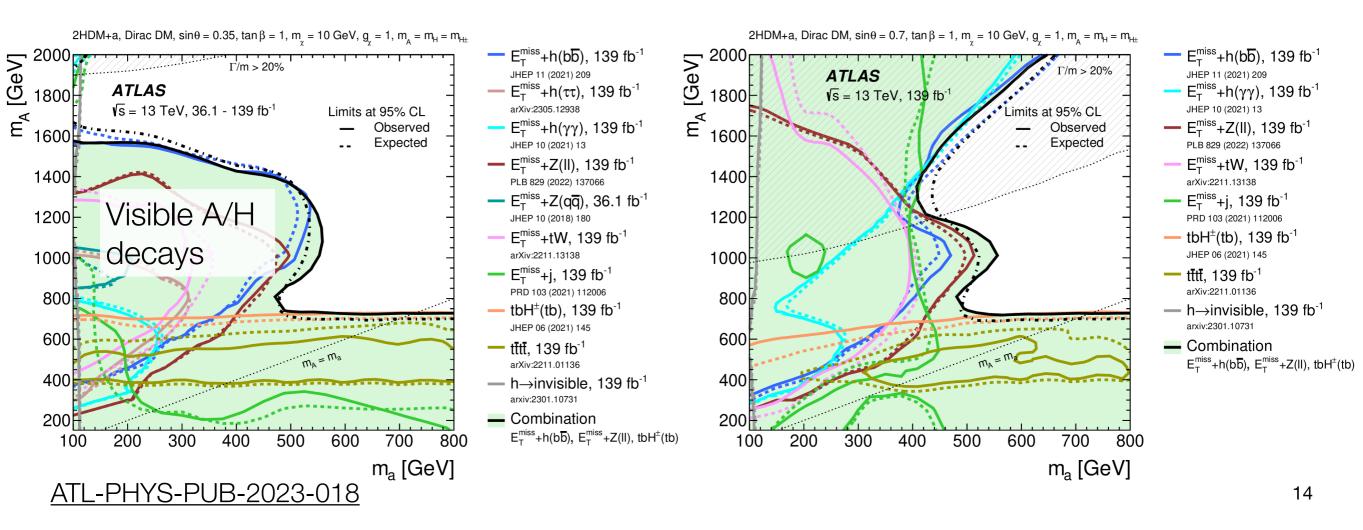


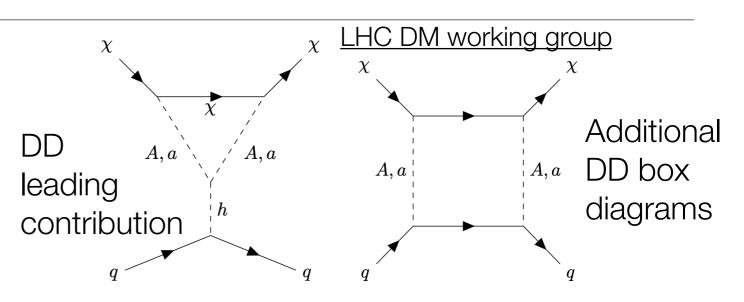


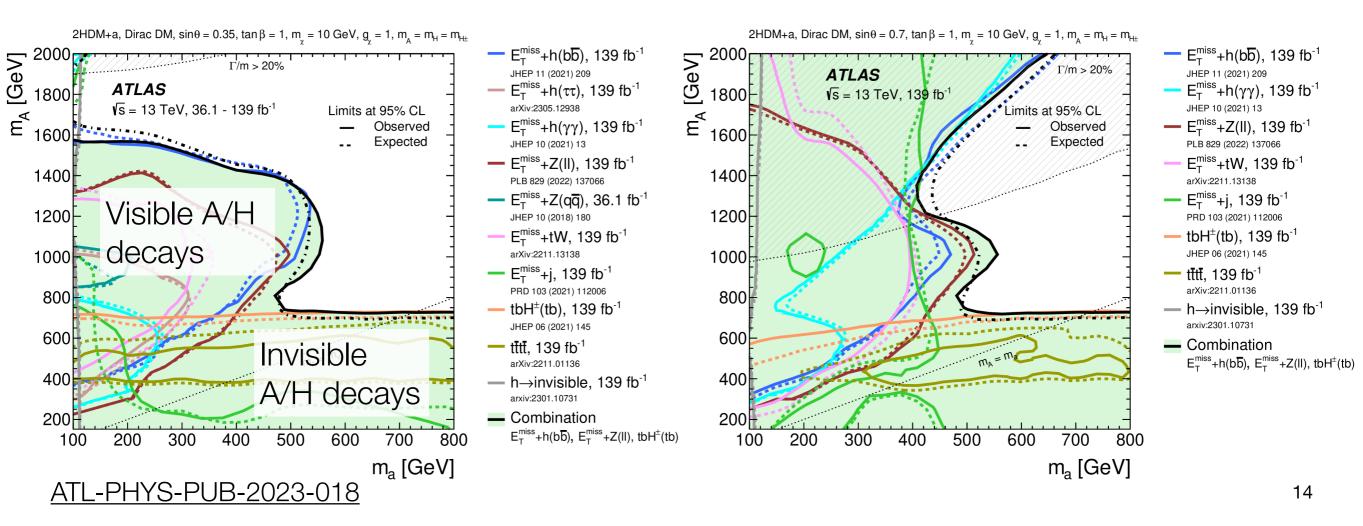










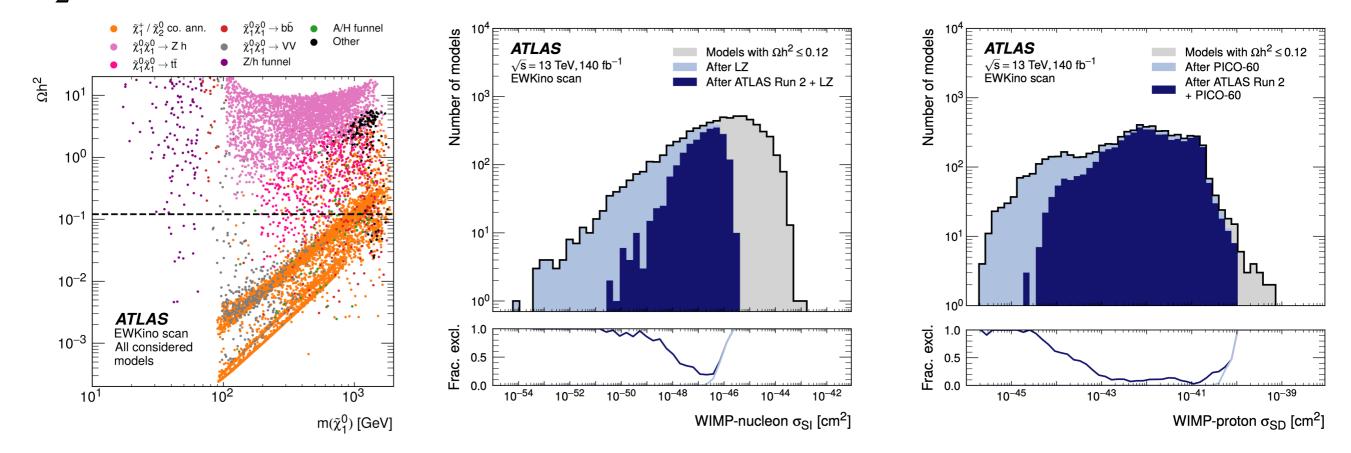


Let's look at pMSSM scan of DM candidates

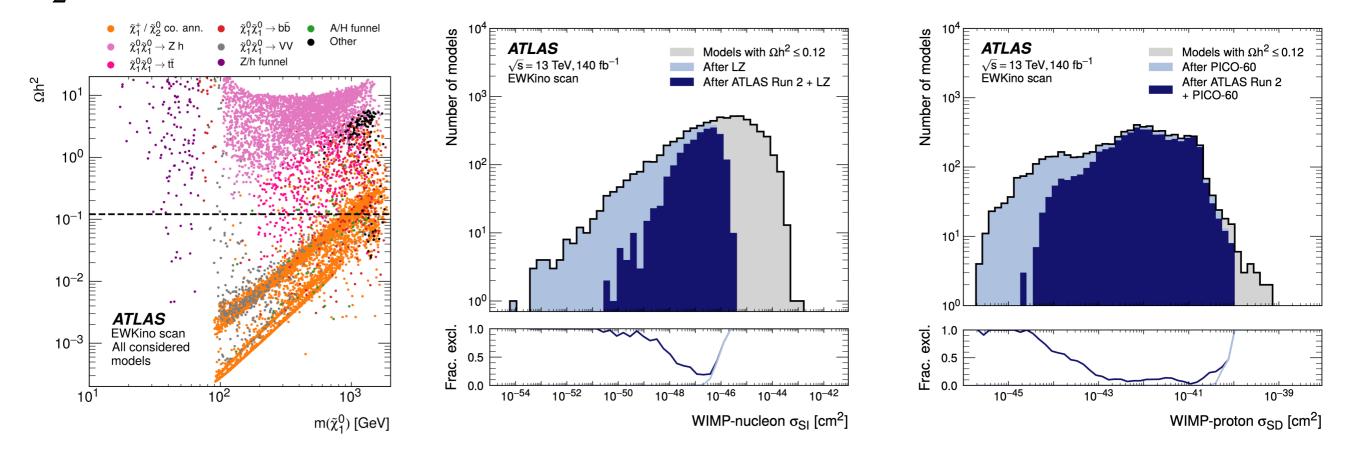
ATLAS CERN-EP-2024-021

Let's look at pMSSM scan of DM candidates $\Delta TLAS CERN-EP-2024-021$ Co-annihilation with small mass splitting from wino/higgsino-like $\tilde{\chi}_1^{\pm}$ and $\tilde{\chi}_2^0$ to LSP gives most of the viable candidates explored here

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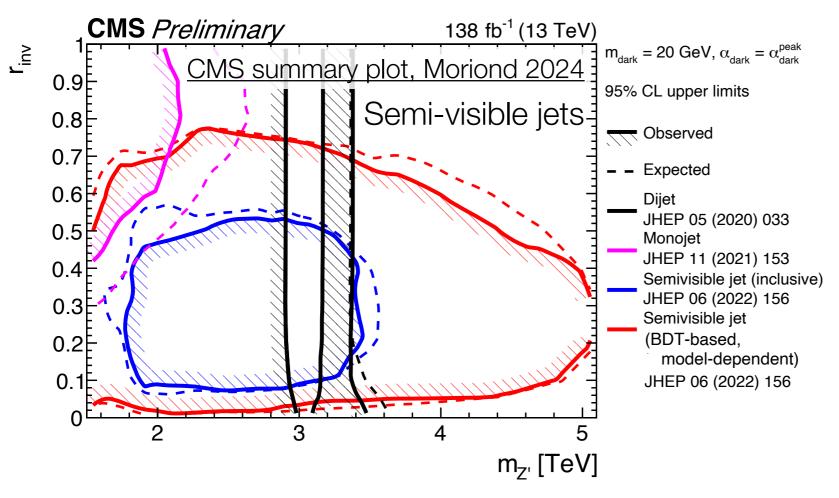
Can see 1) there is considerable space left for SUSY DM candidates in hard-to-reach electroweak signatures, and 2) there is good complementarity between LHC and direct detection reach

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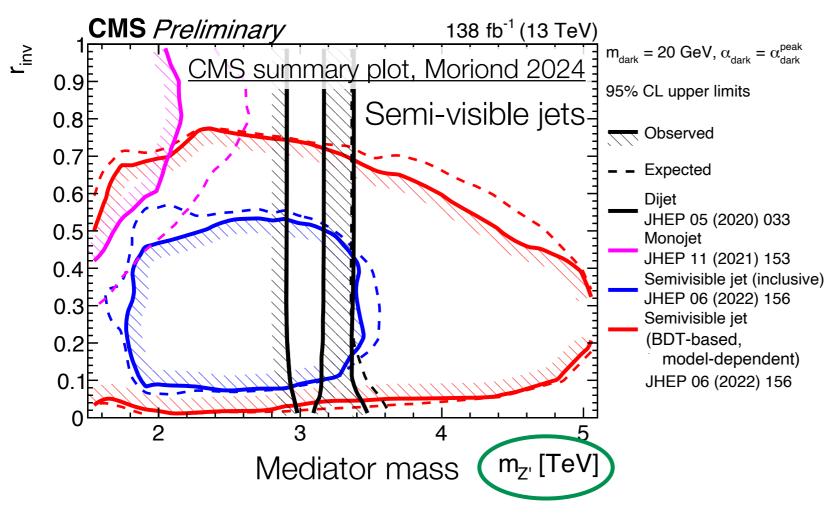
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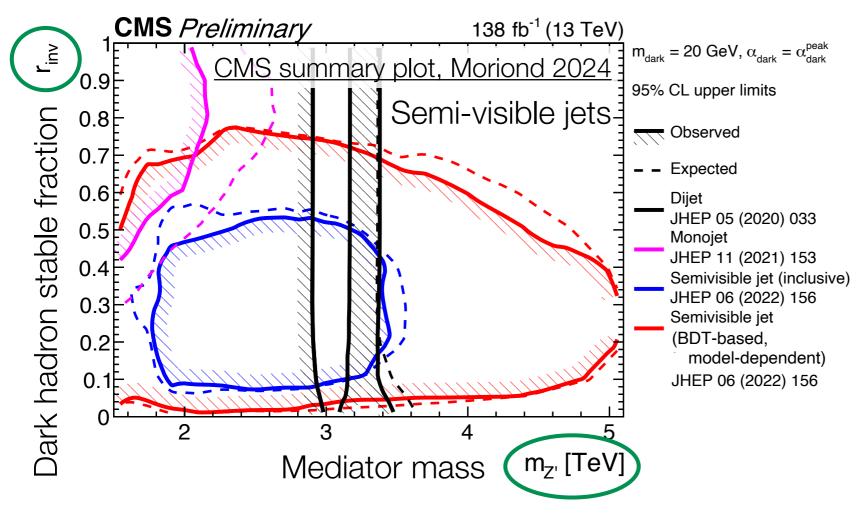
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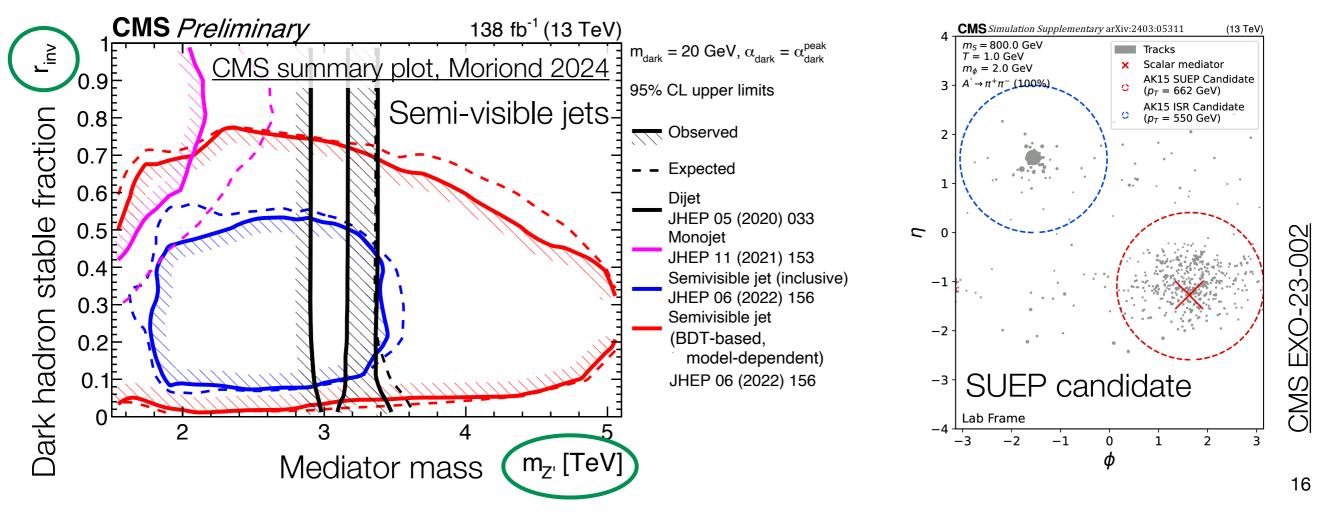
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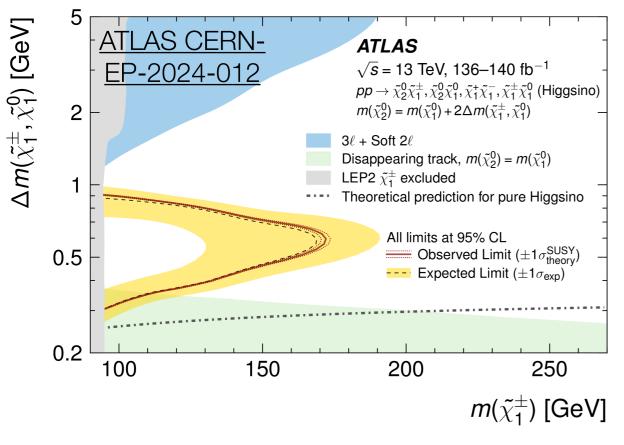


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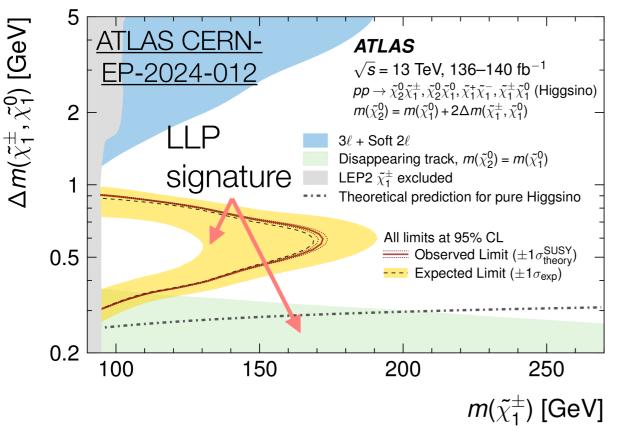


Saw one case already: displaced decays in dark photons with small ε. Other important examples:

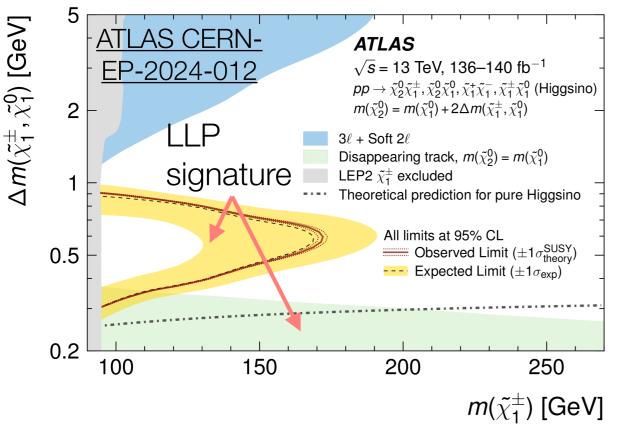
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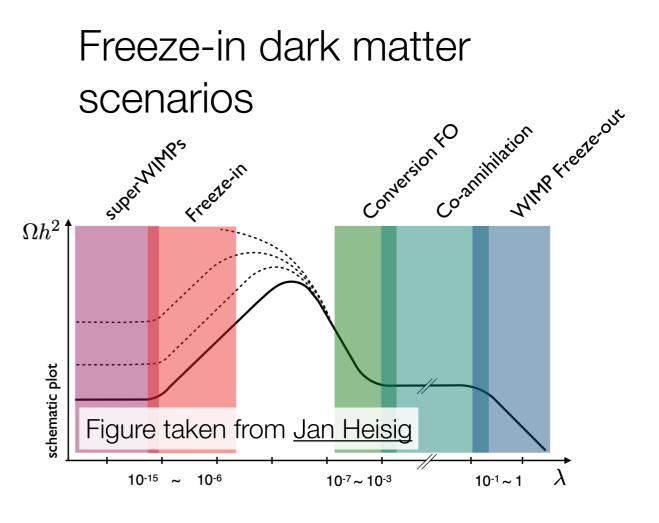


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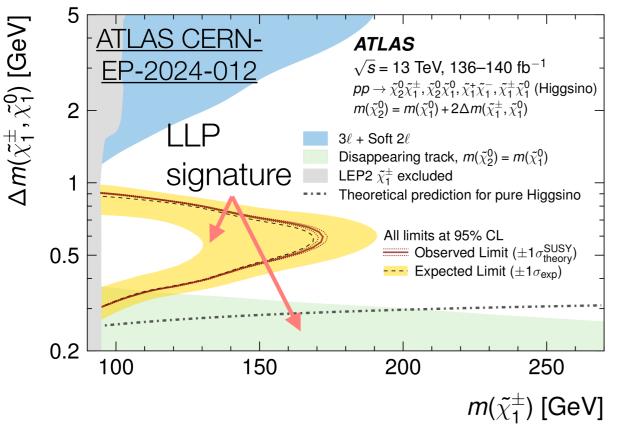


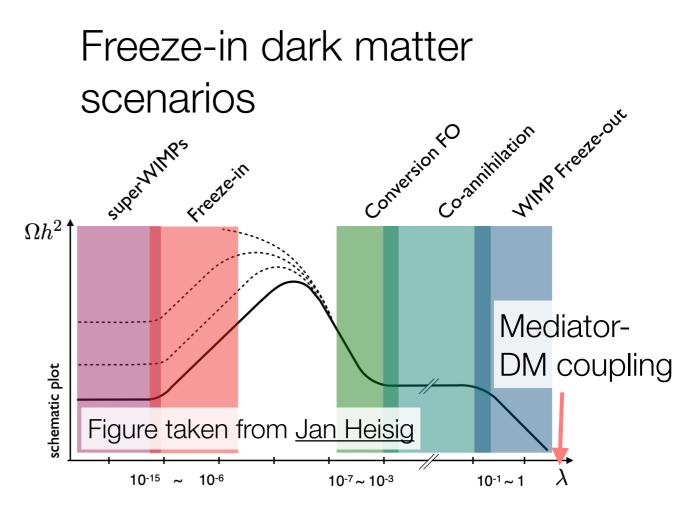
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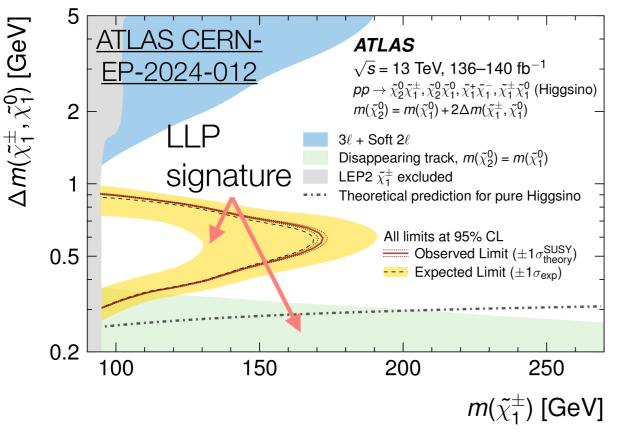


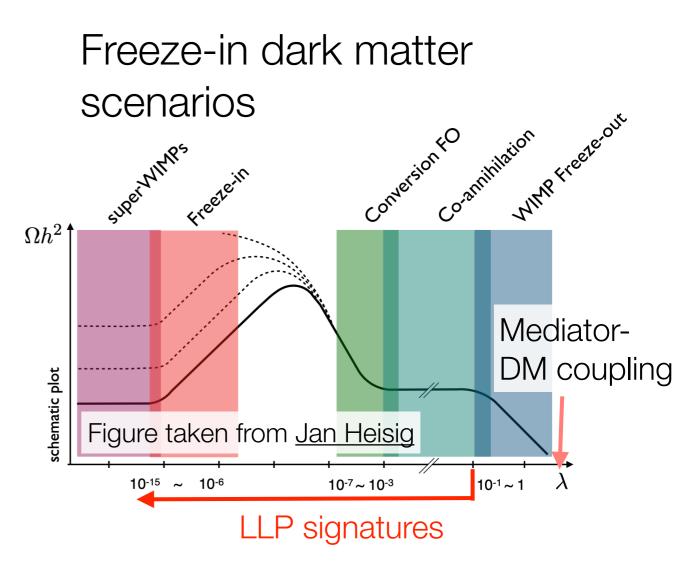
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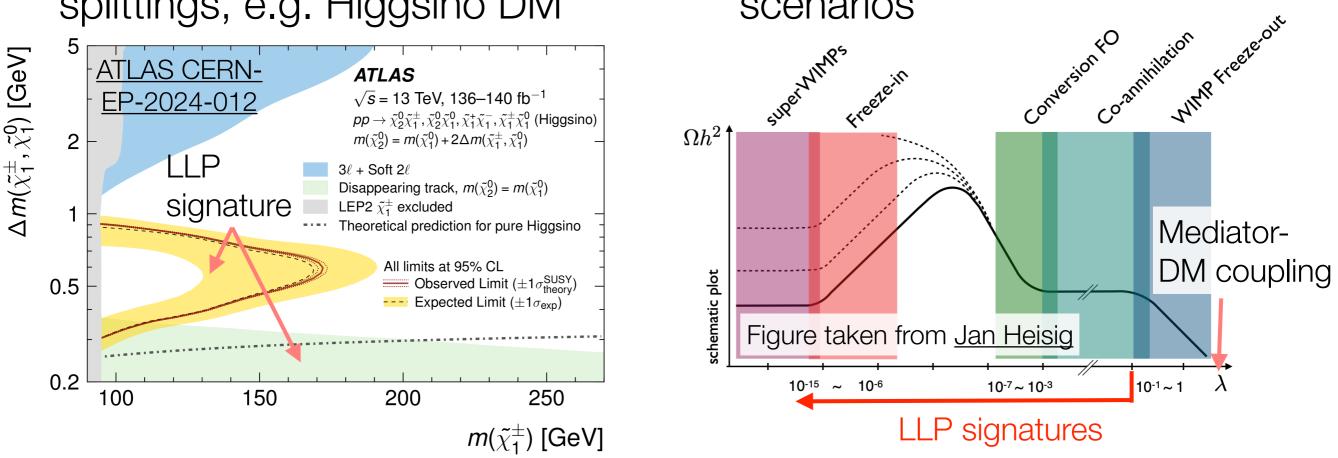


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Freeze-in dark matter

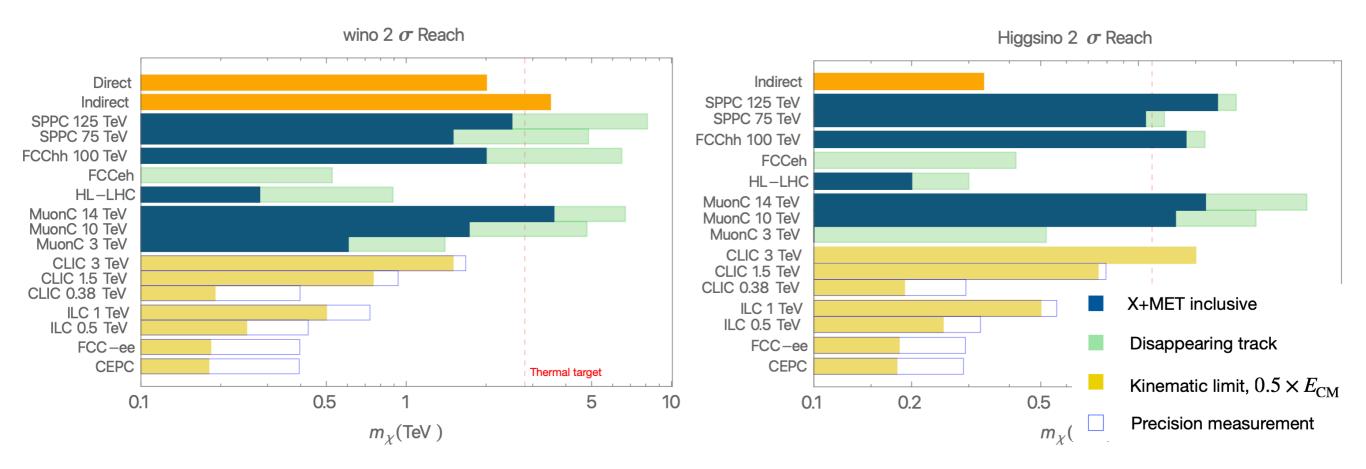
scenarios

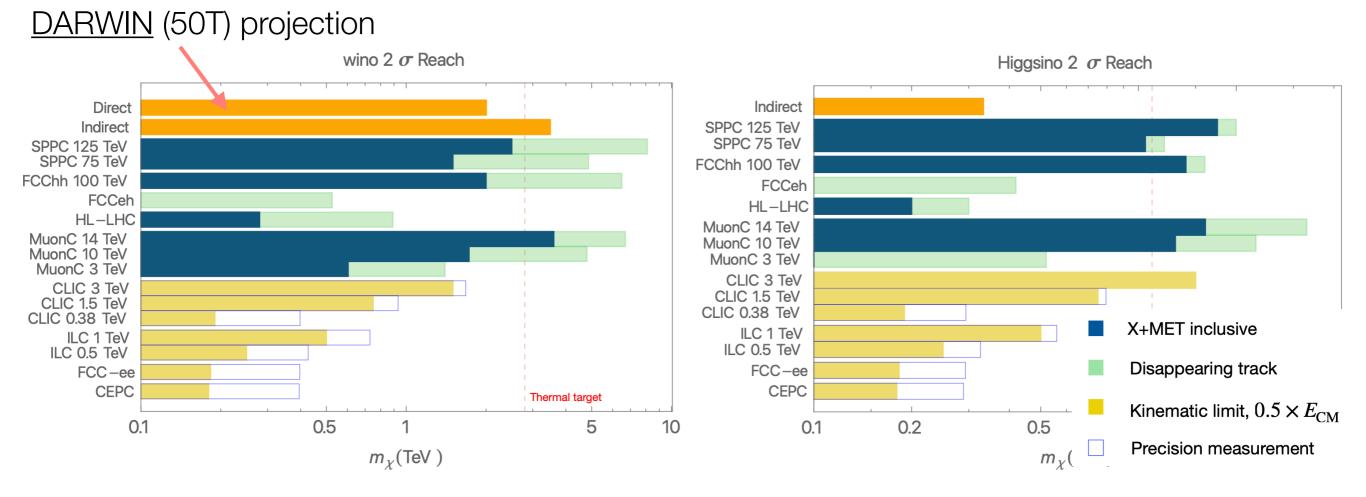
Models with very small mass splittings, e.g. Higgsino DM

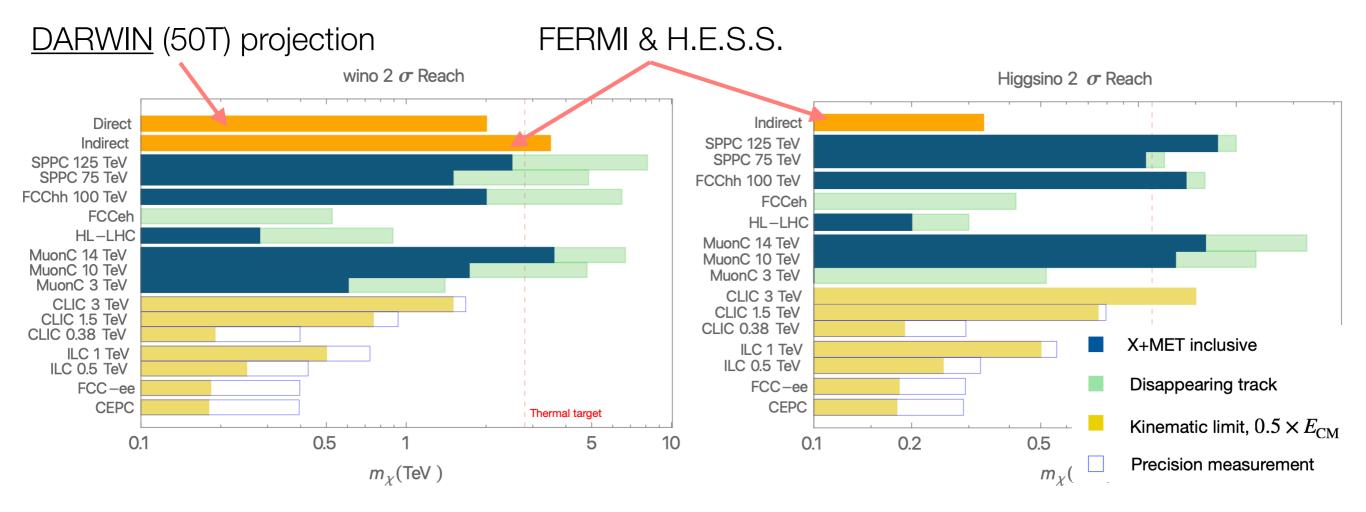


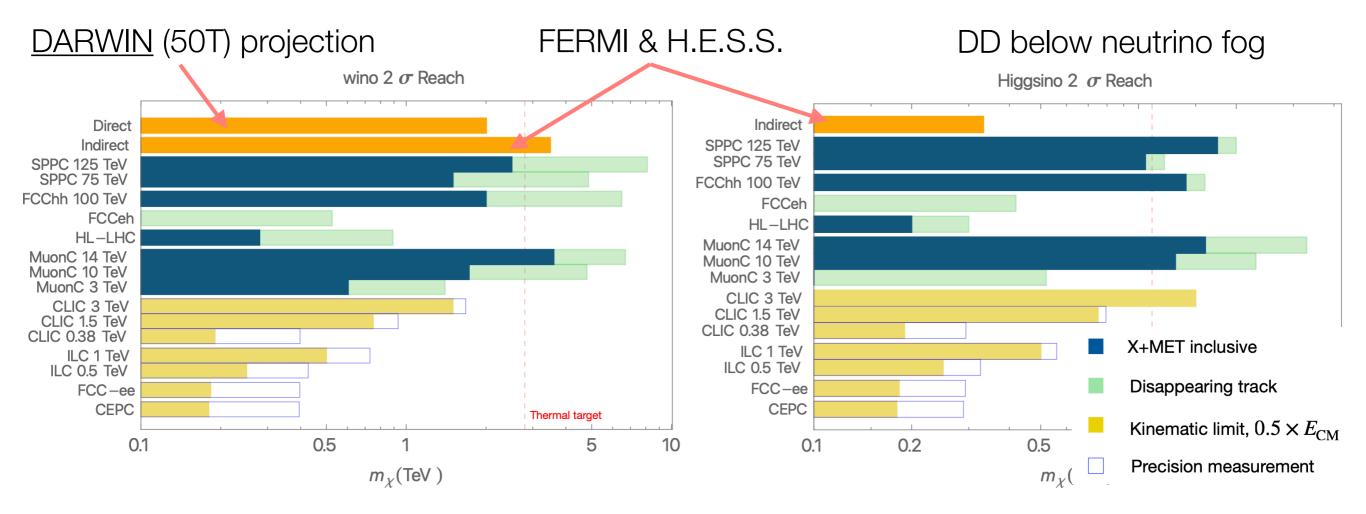
Can get LLPs from small mass splittings or small couplings, and turn up frequently in asymmetric, freeze-in, & SUSY DM

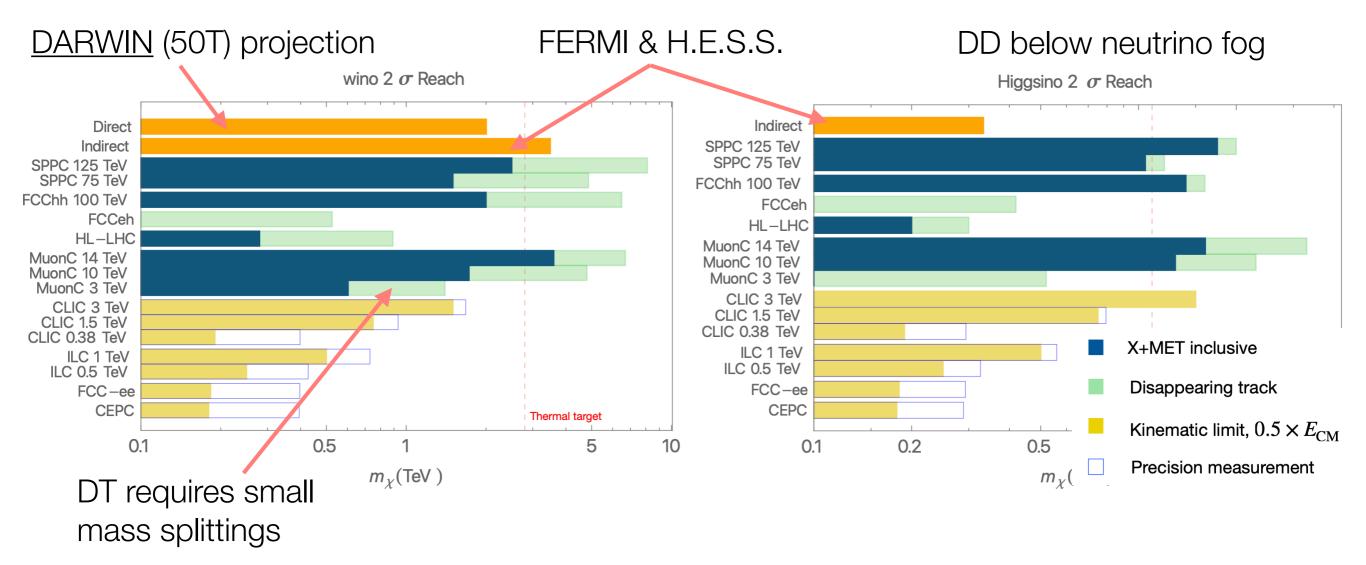
DM at HL-LHC and future colliders

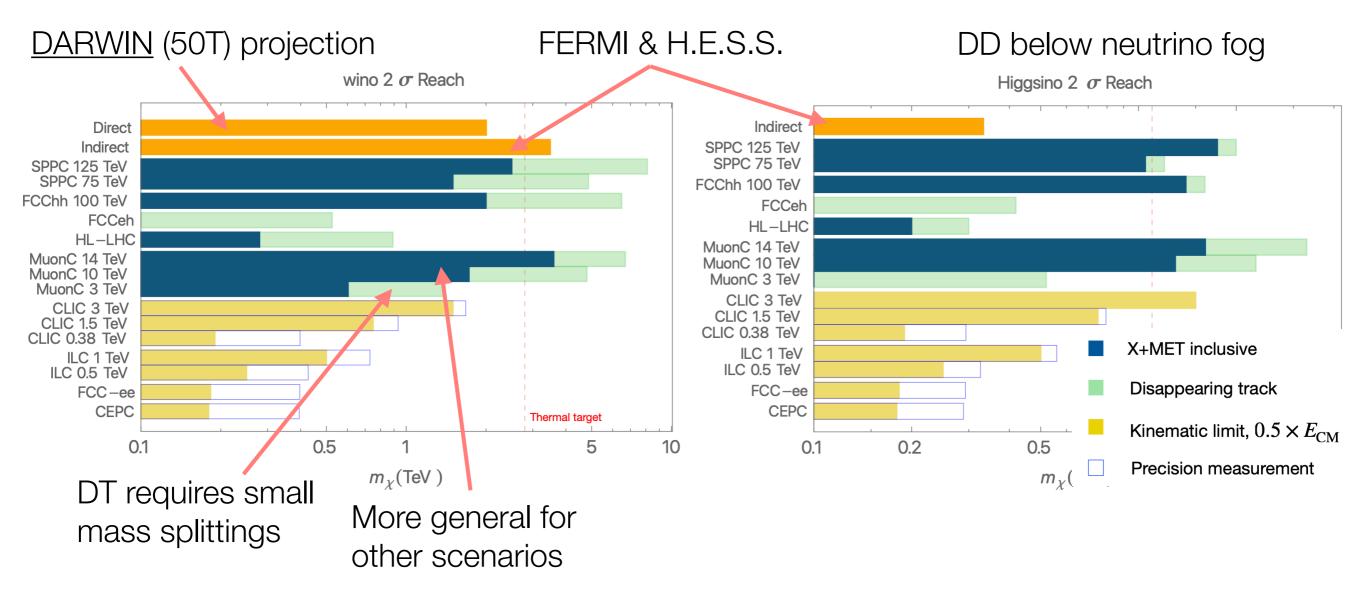


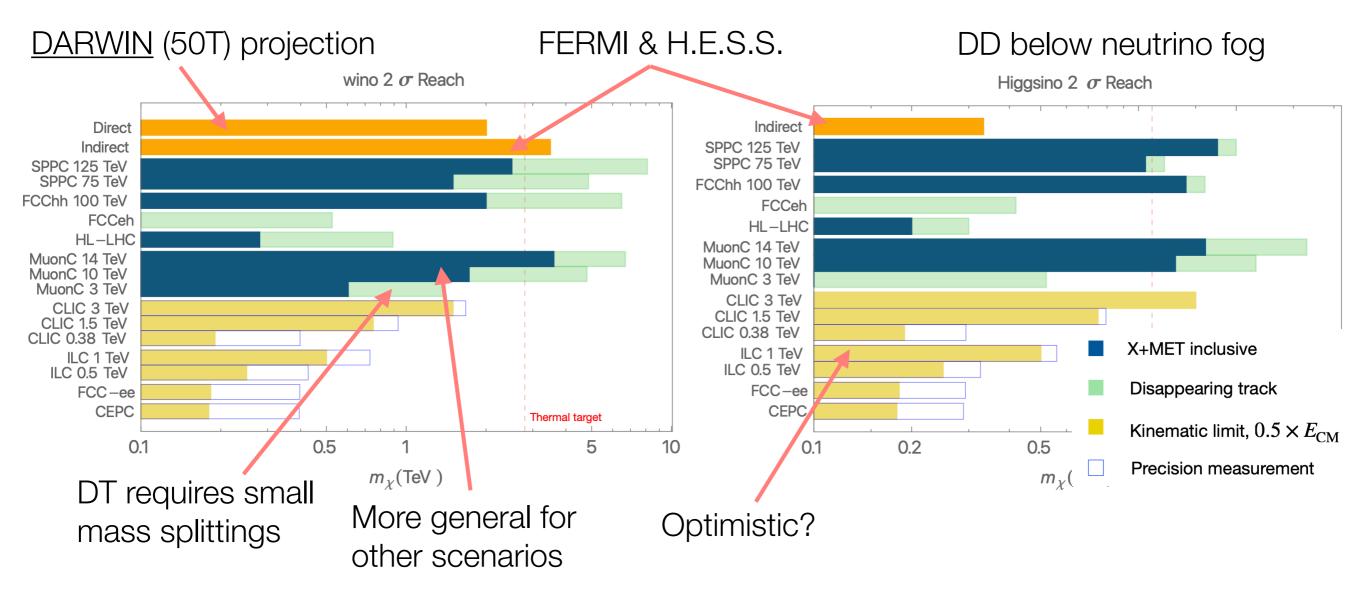






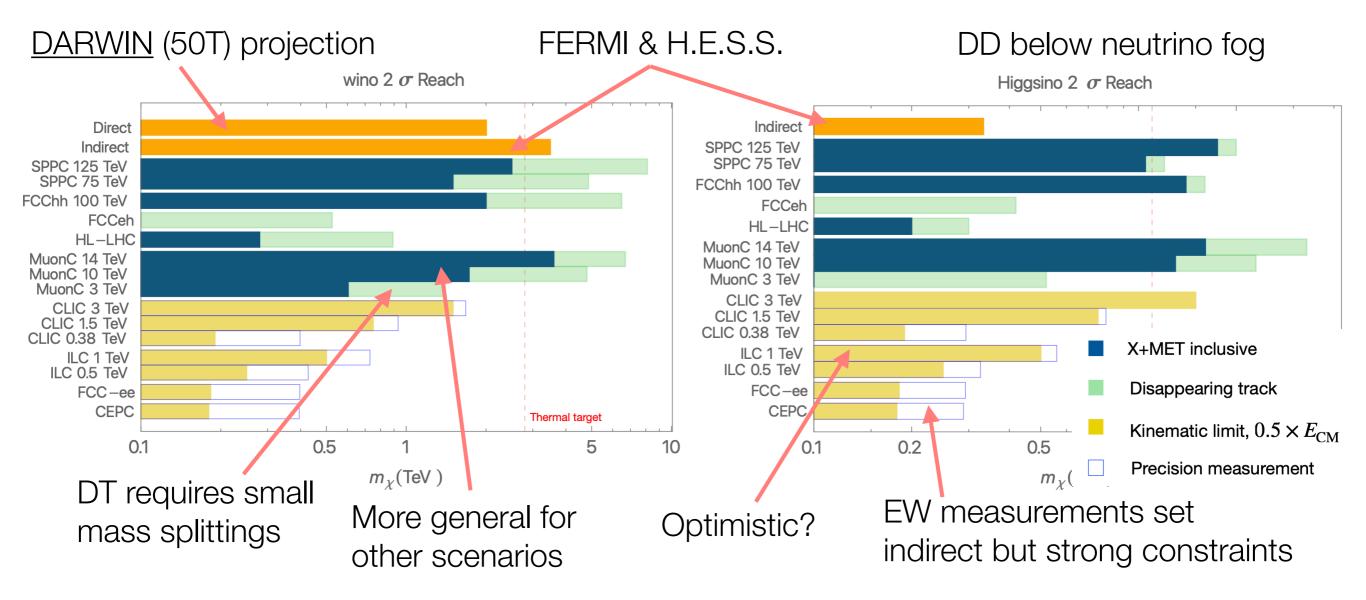






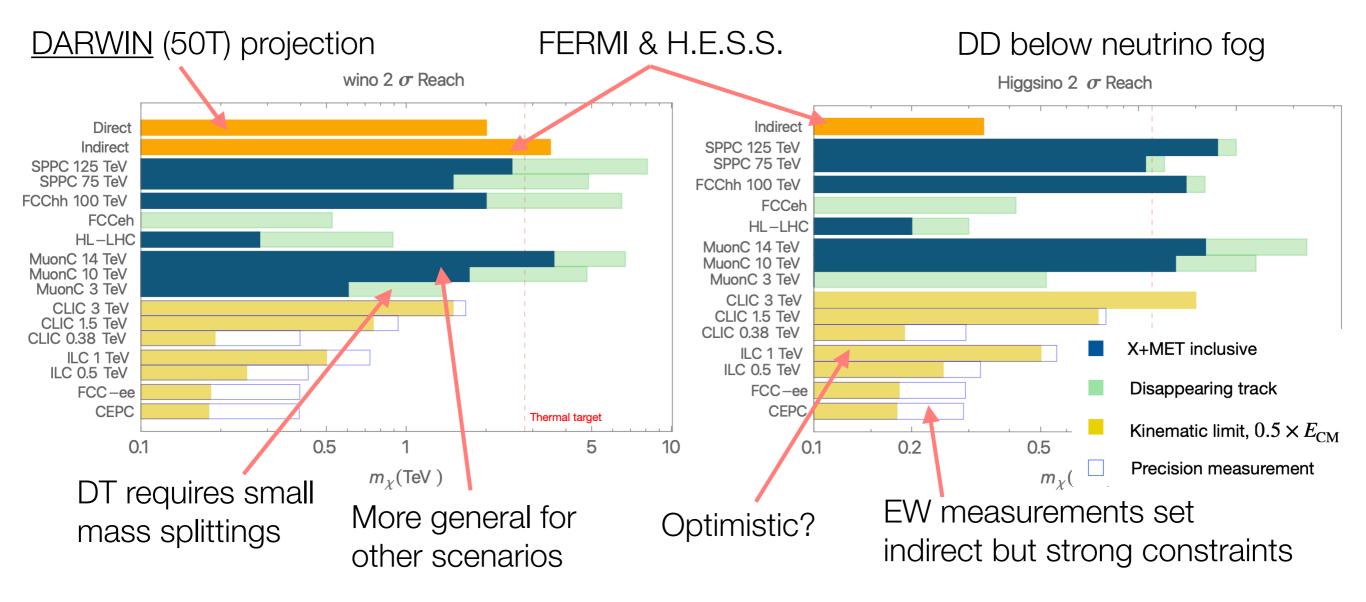
Opportunities at future colliders: SUSY DM

Minimal EW multiplet scenario: SM gauge couplings fix interactions so mass is only free parameter and thermal DM predictions simple.



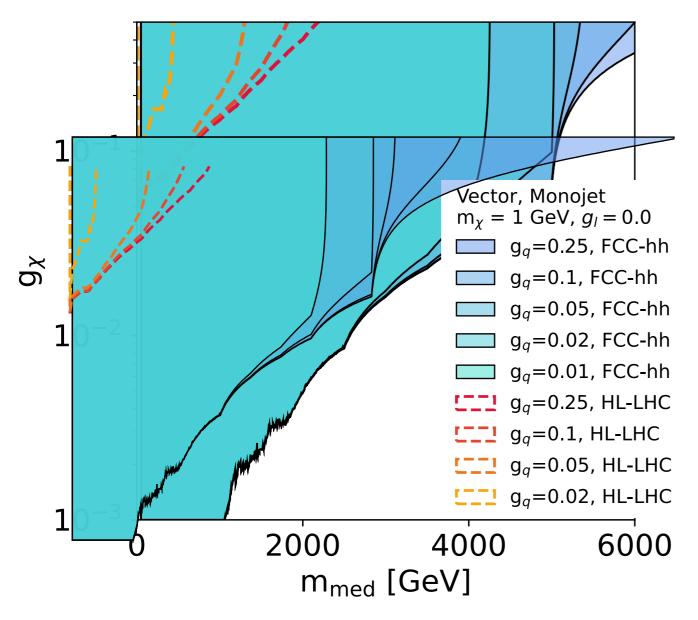
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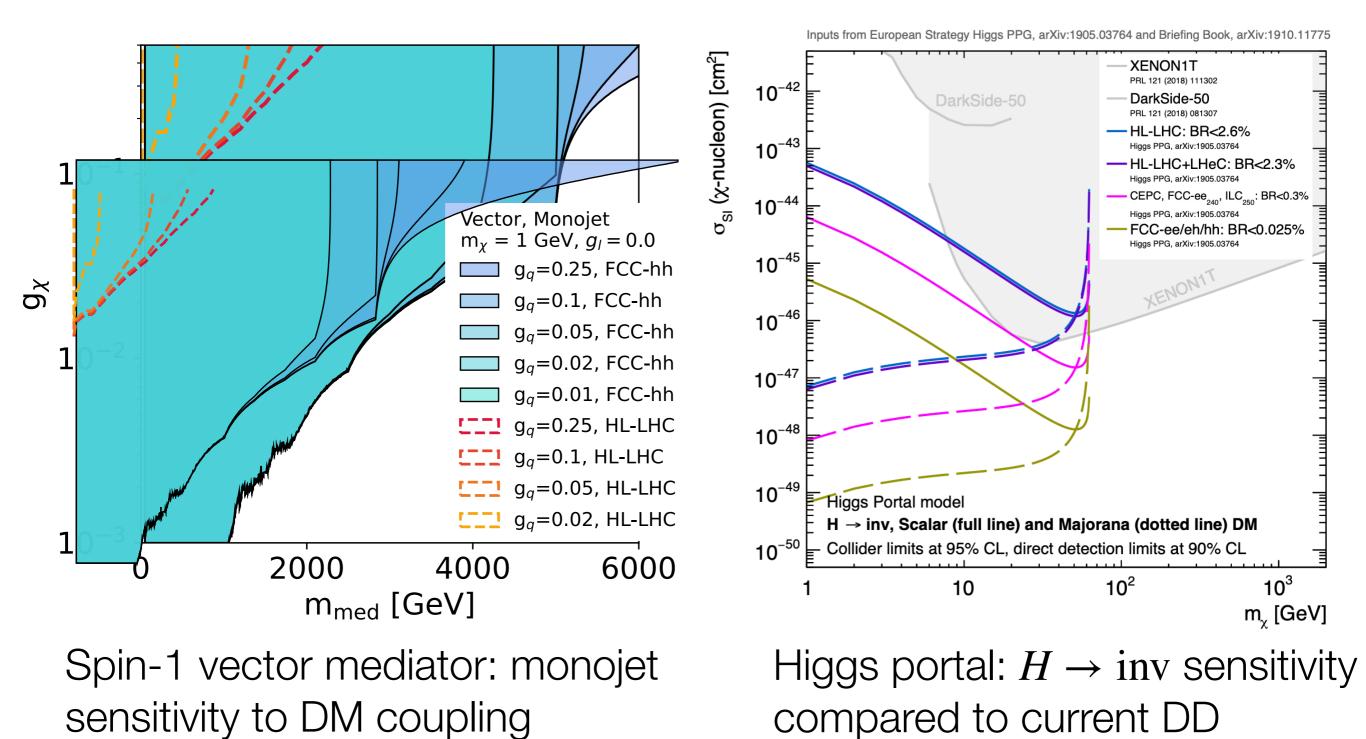
Reaching thermal target is not easy, but possible at some colliders

Opportunities at future colliders: non-SUSY DM



Spin-1 vector mediator: monojet sensitivity to DM coupling

Opportunities at future colliders: non-SUSY DM



20

Parasitic experiments at future colliders

Future colliders have possibilities beyond collision point detectors

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Dedicated LLP

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Valuable when LLP signature is trigger limited

Limited use at e+e- machines but useful at hadron & probably muon machines

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Beam dump experiments

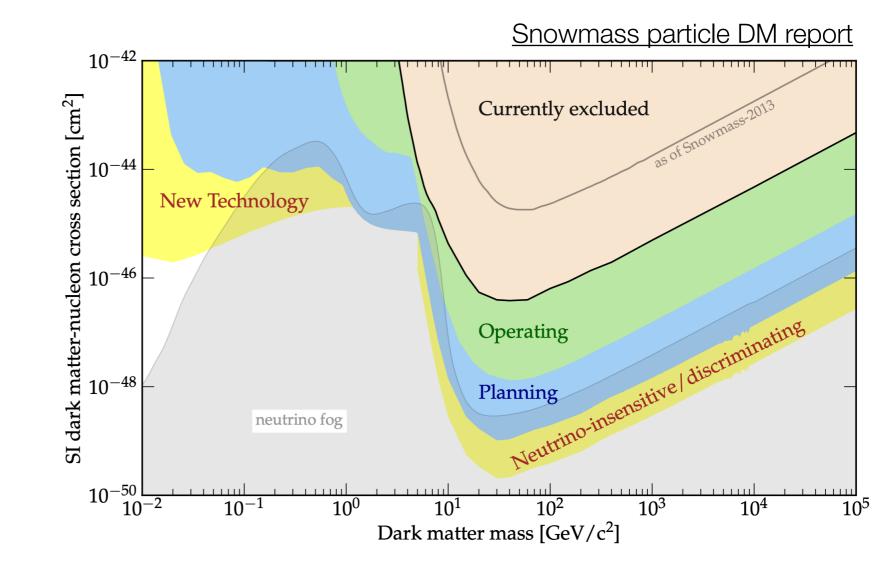
Missing energy/mass experiments not possible at EF machines

Could probably do a re-scattering experiment here but I've not seen it talked about

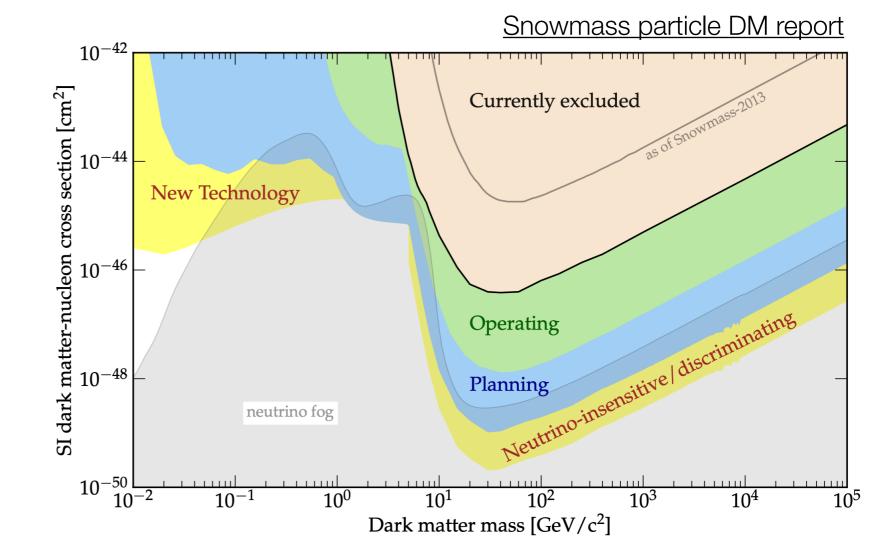
Visible decay searches are well suited and could be added to future colliders (examples 1, 2)

Discussing complementarity

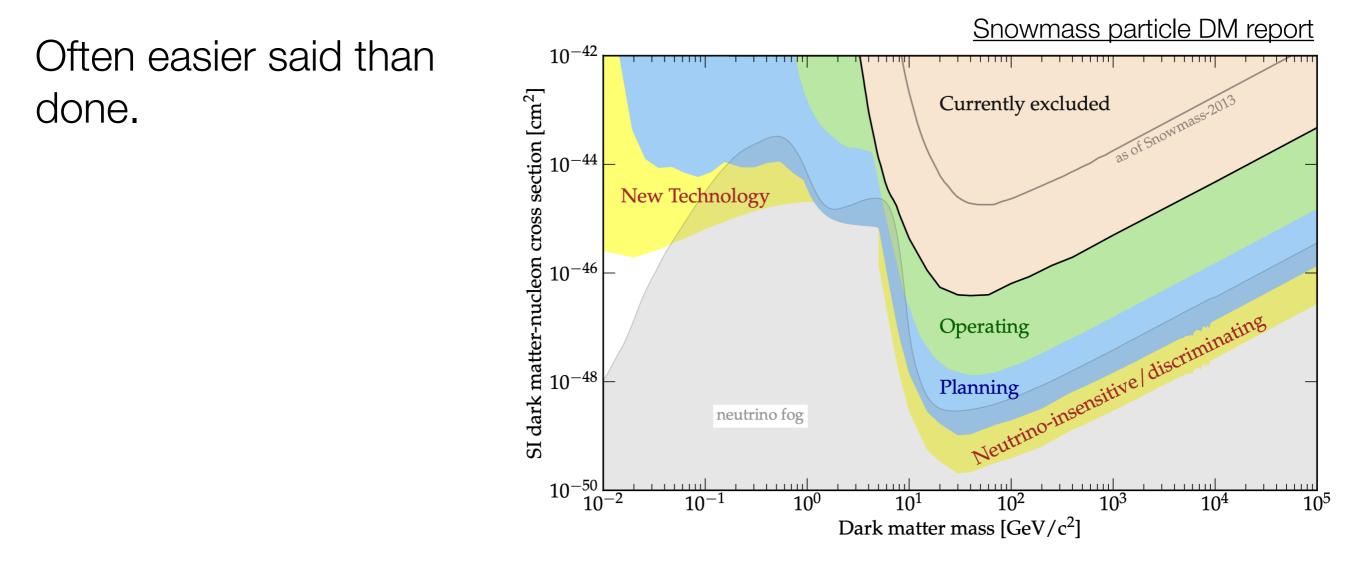
3



This will be key to building the field we want to see



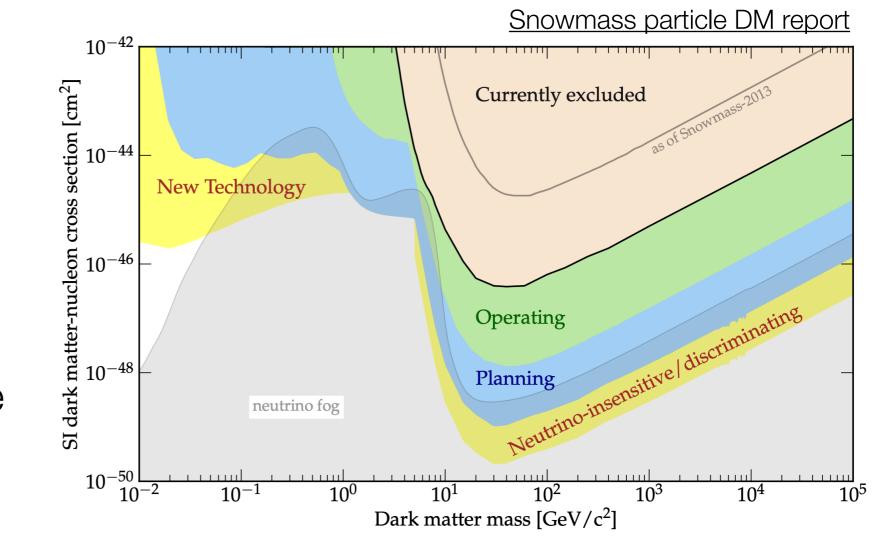
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Often easier said than done.

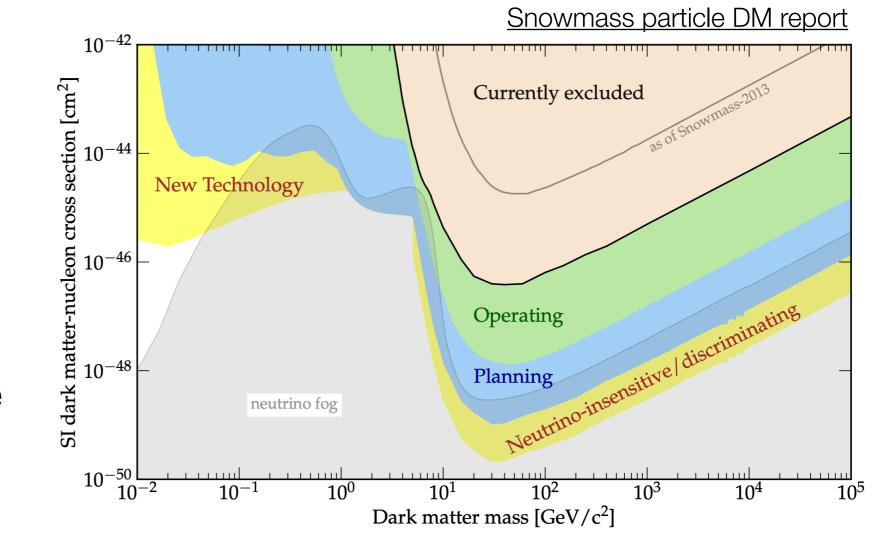
DD limits can use EFT; EF searches require model assumptions. Reducing problem dimensions to 2D plane usually needs extra assumptions



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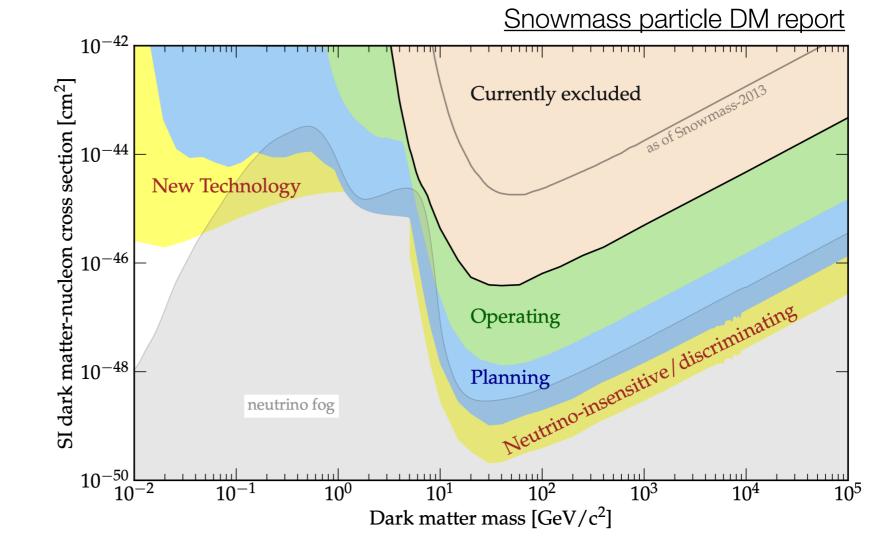


Show example I know best: LHC DMWG spin-1 simplified model

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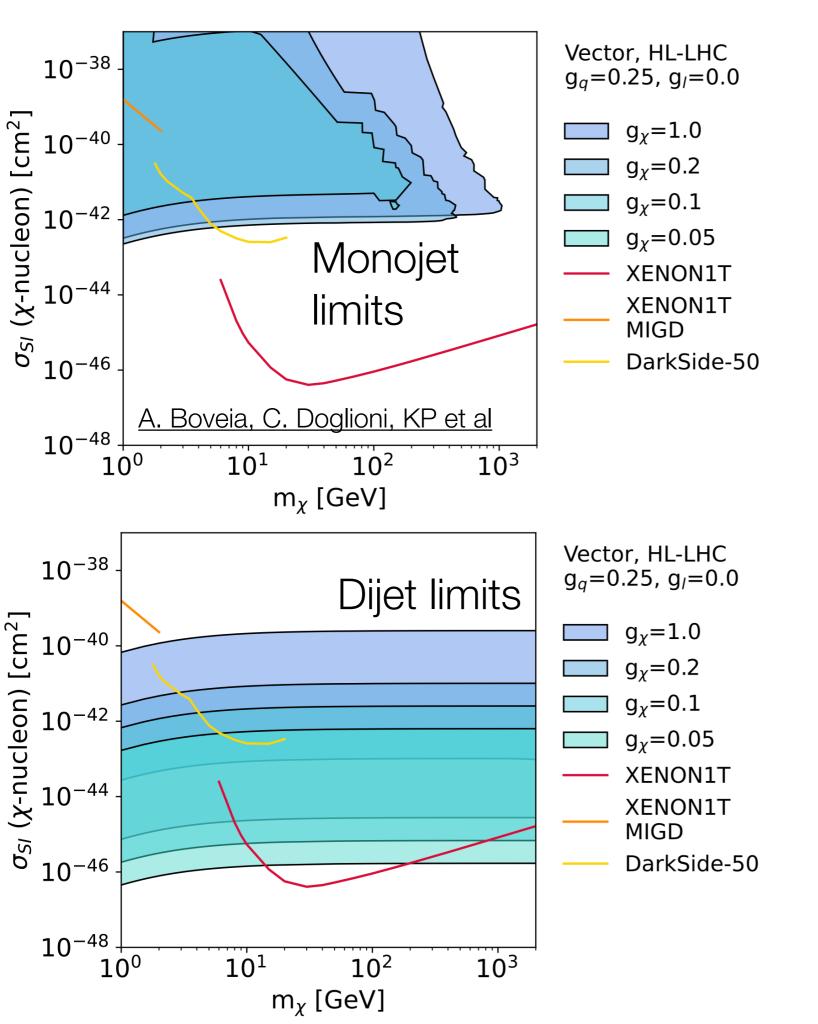
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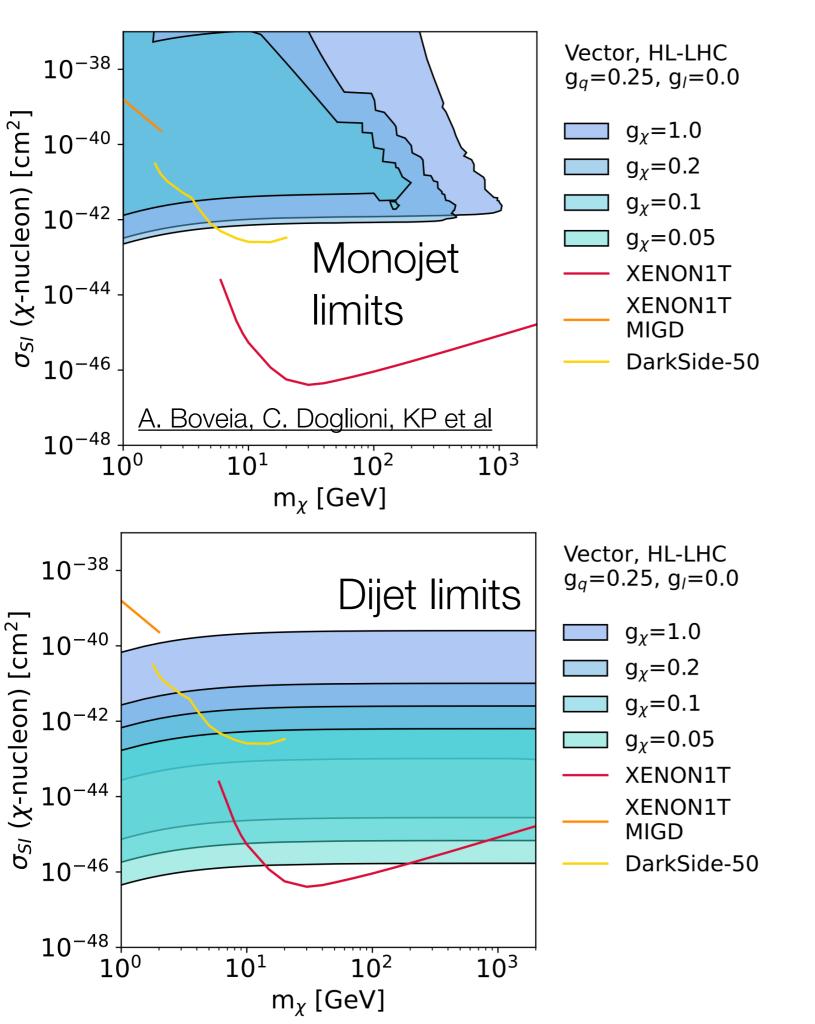
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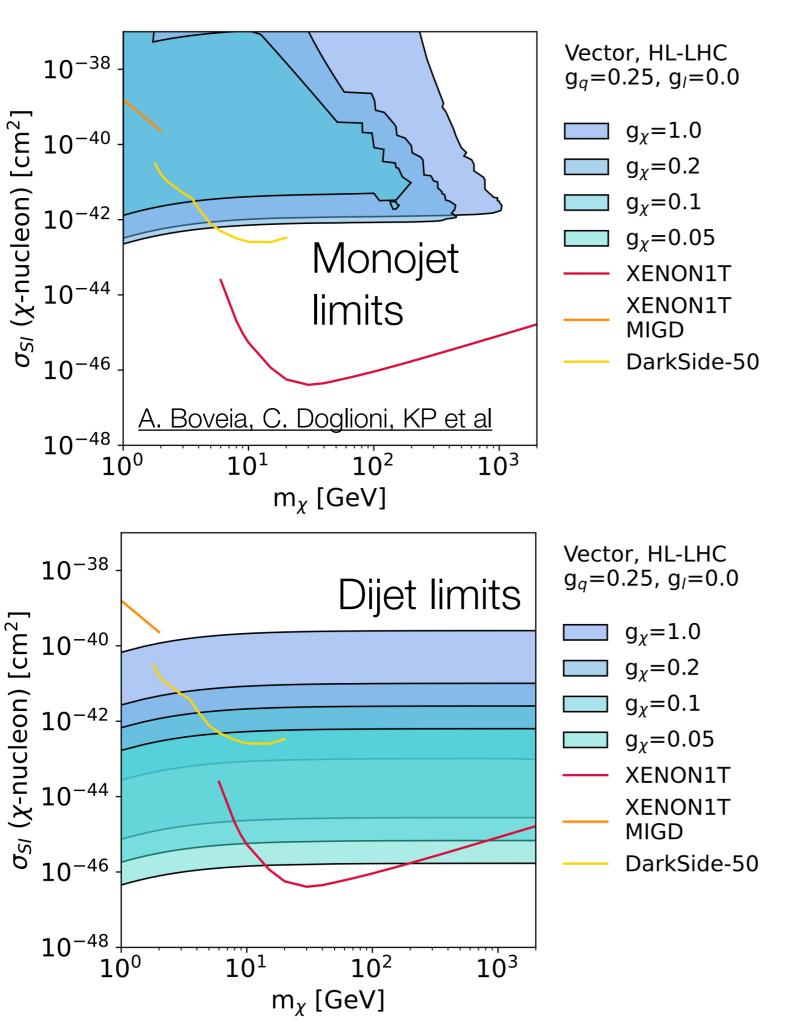
Show example I know best: LHC DMWG spin-1 simplified model

Must reduce 4-5 free parameters ($m_{\rm med}, m_{\chi}, g_{SM}, g_{\chi}$) to 2



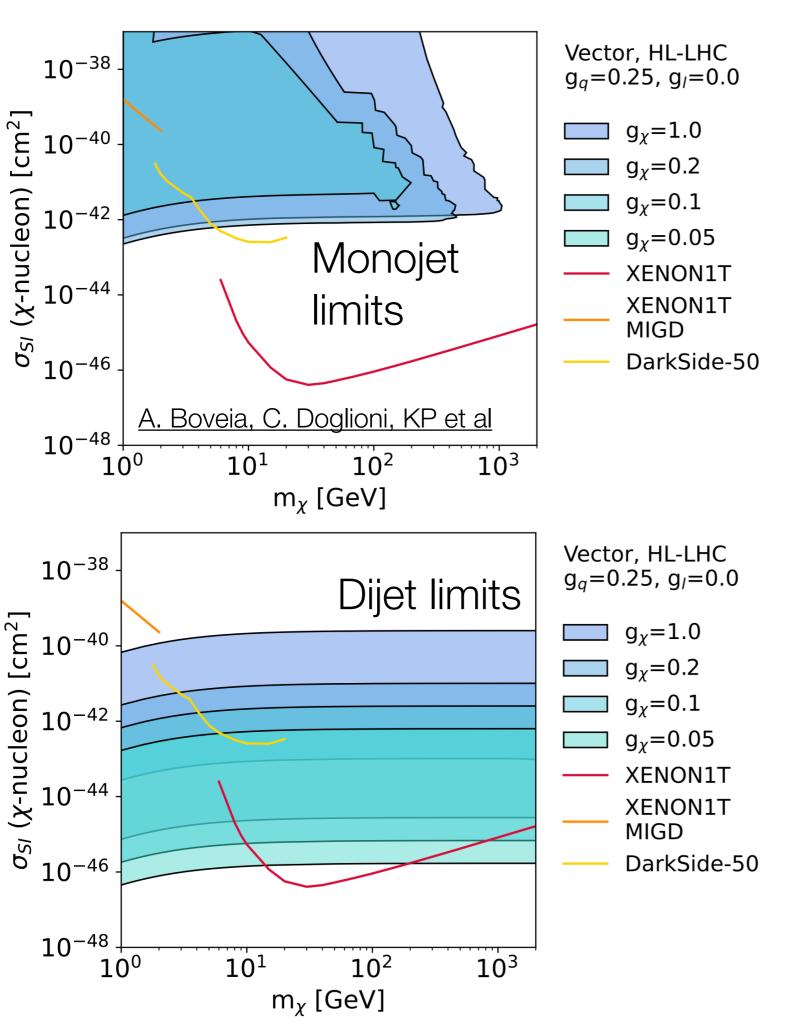


Couplings take explicit values



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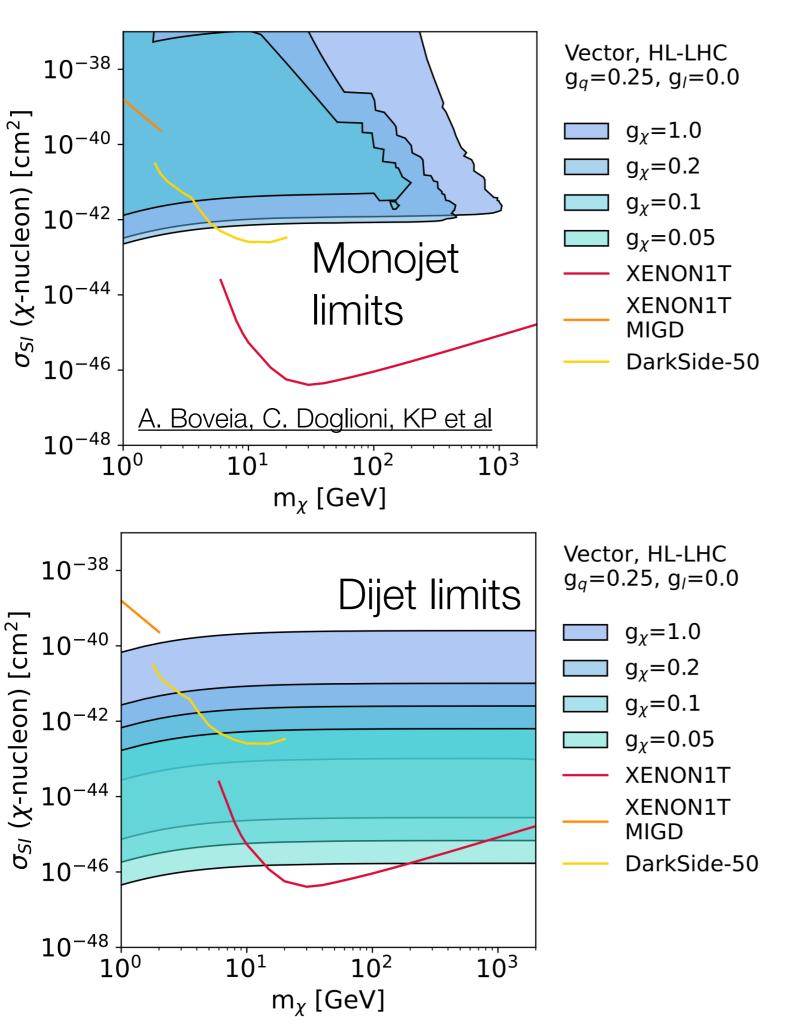
Mediator mass absorbed into y axis variable



Couplings take explicit values

Mediator mass absorbed into y axis variable

Implication: no constraint on mediator mass

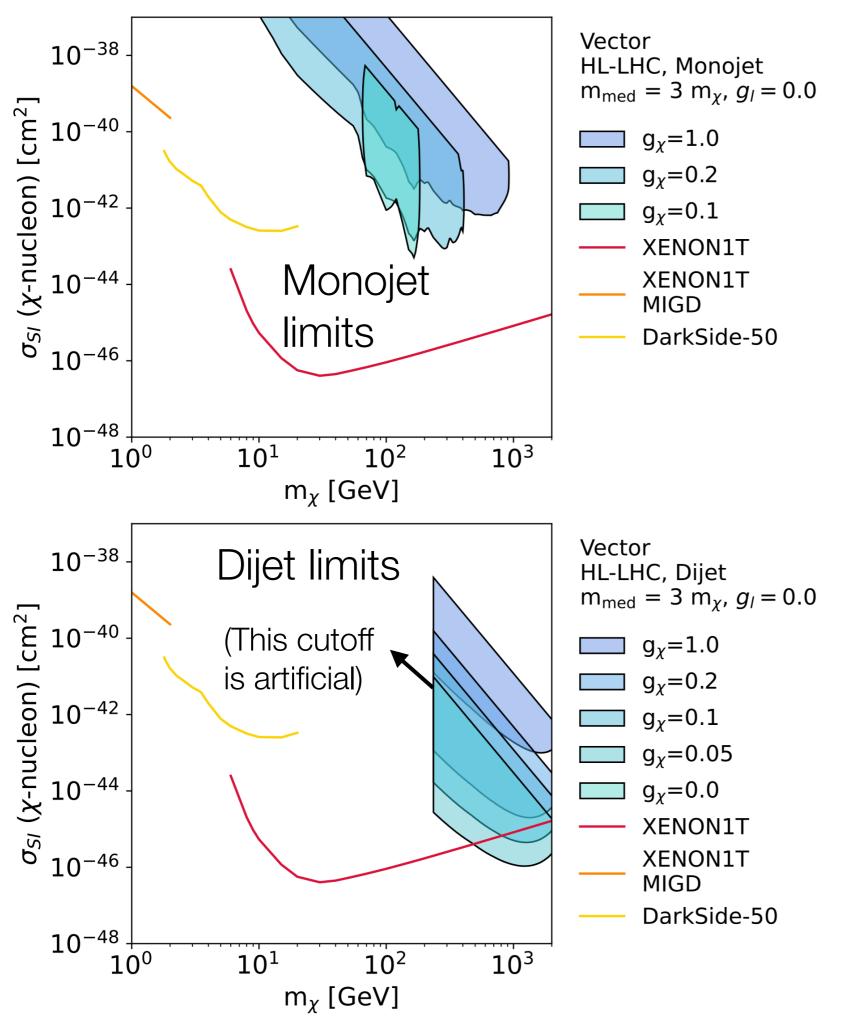


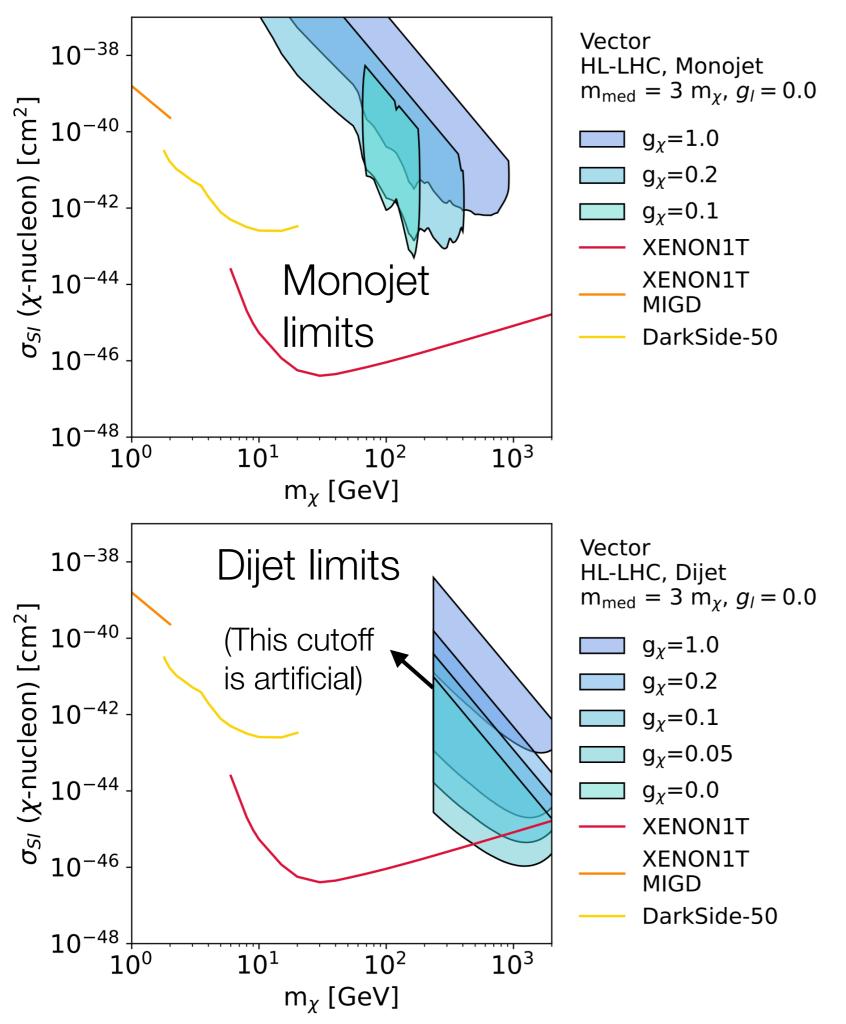
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Mediator mass absorbed into y axis variable

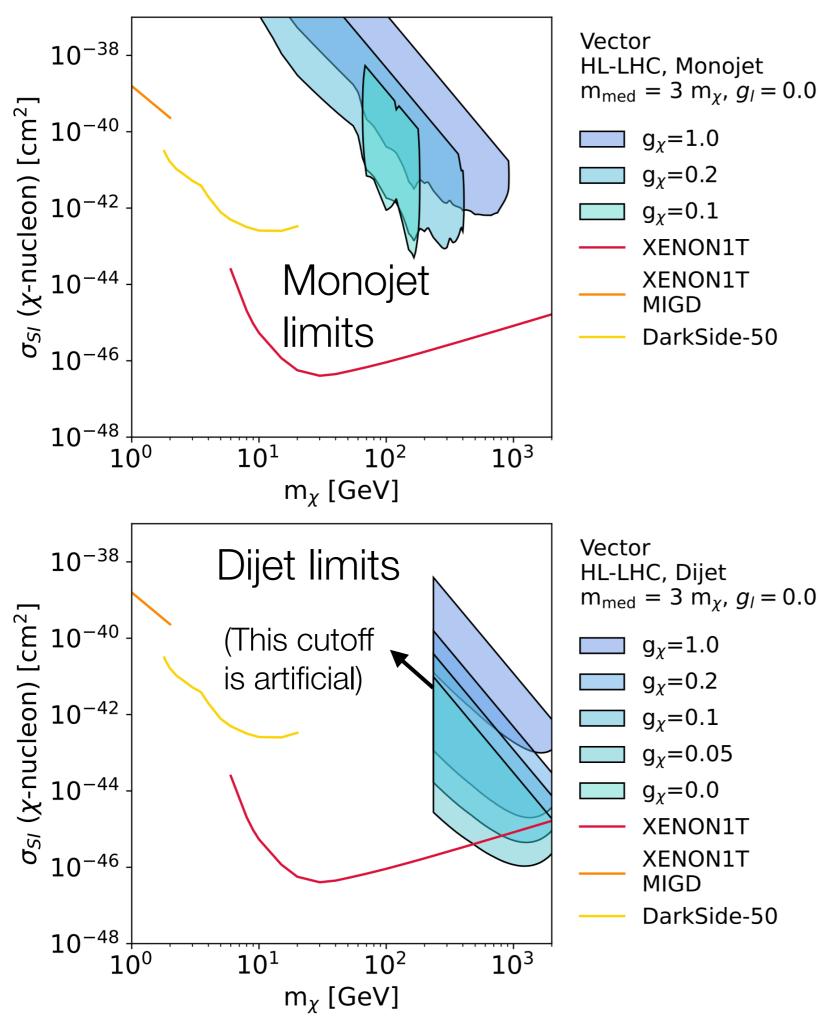
Implication: no constraint on mediator mass

Points with strong collider limits have high mediator mass to DM mass ratio



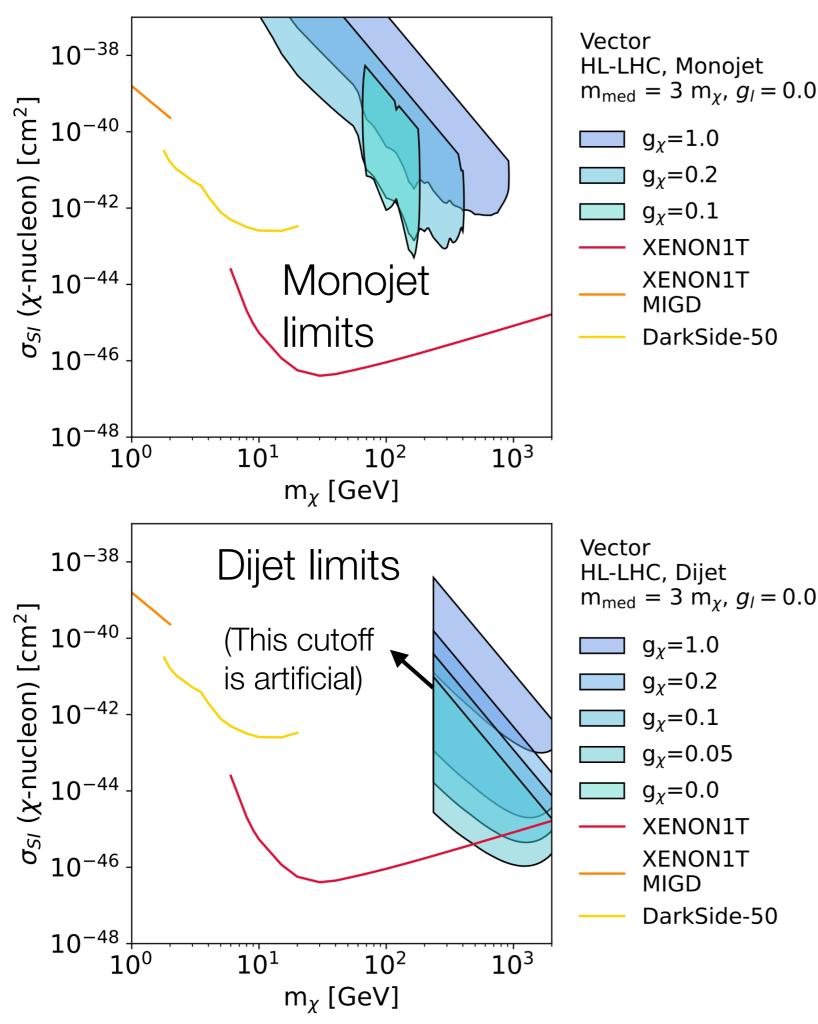


Now ratio between mediators is fixed and g_q is absorbed into y axis



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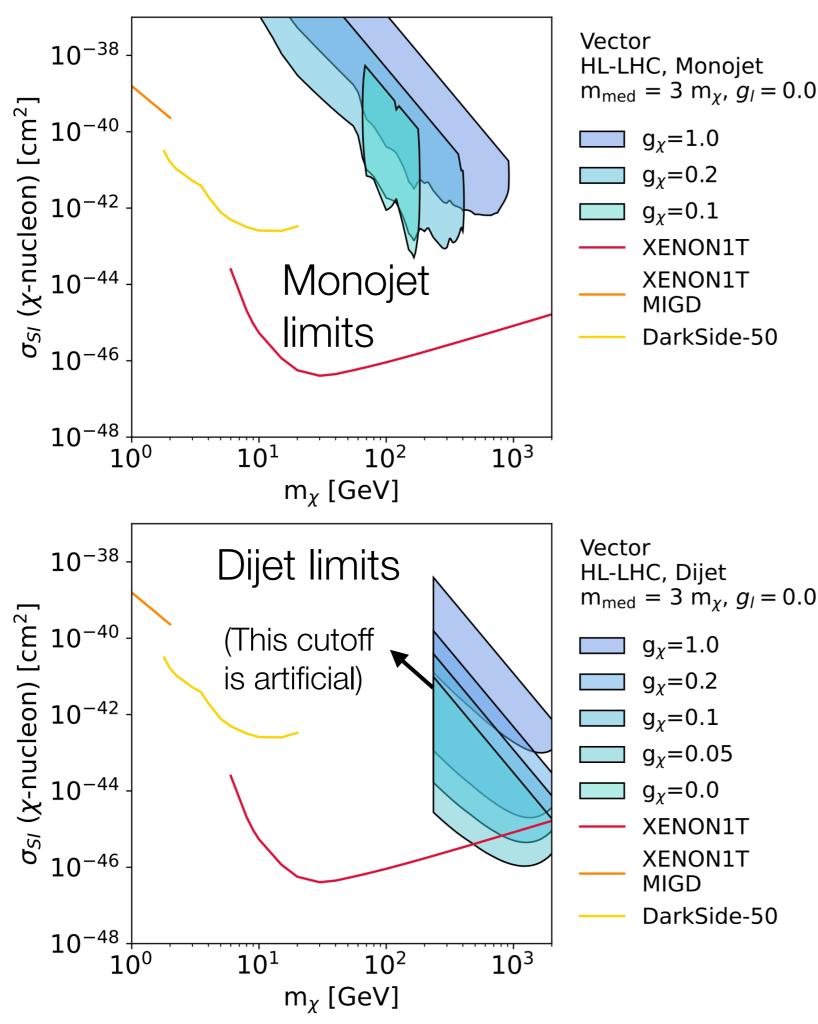
Colliders have unique strengths in accessing heavy mediators



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Colliders have unique strengths in accessing heavy mediators

Direct detection has unique strengths in accessing small couplings

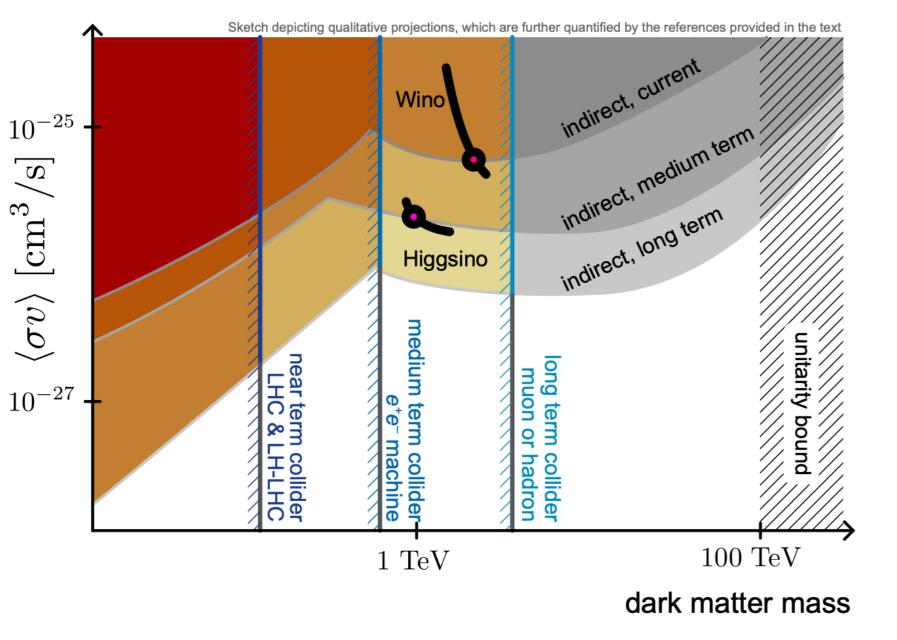


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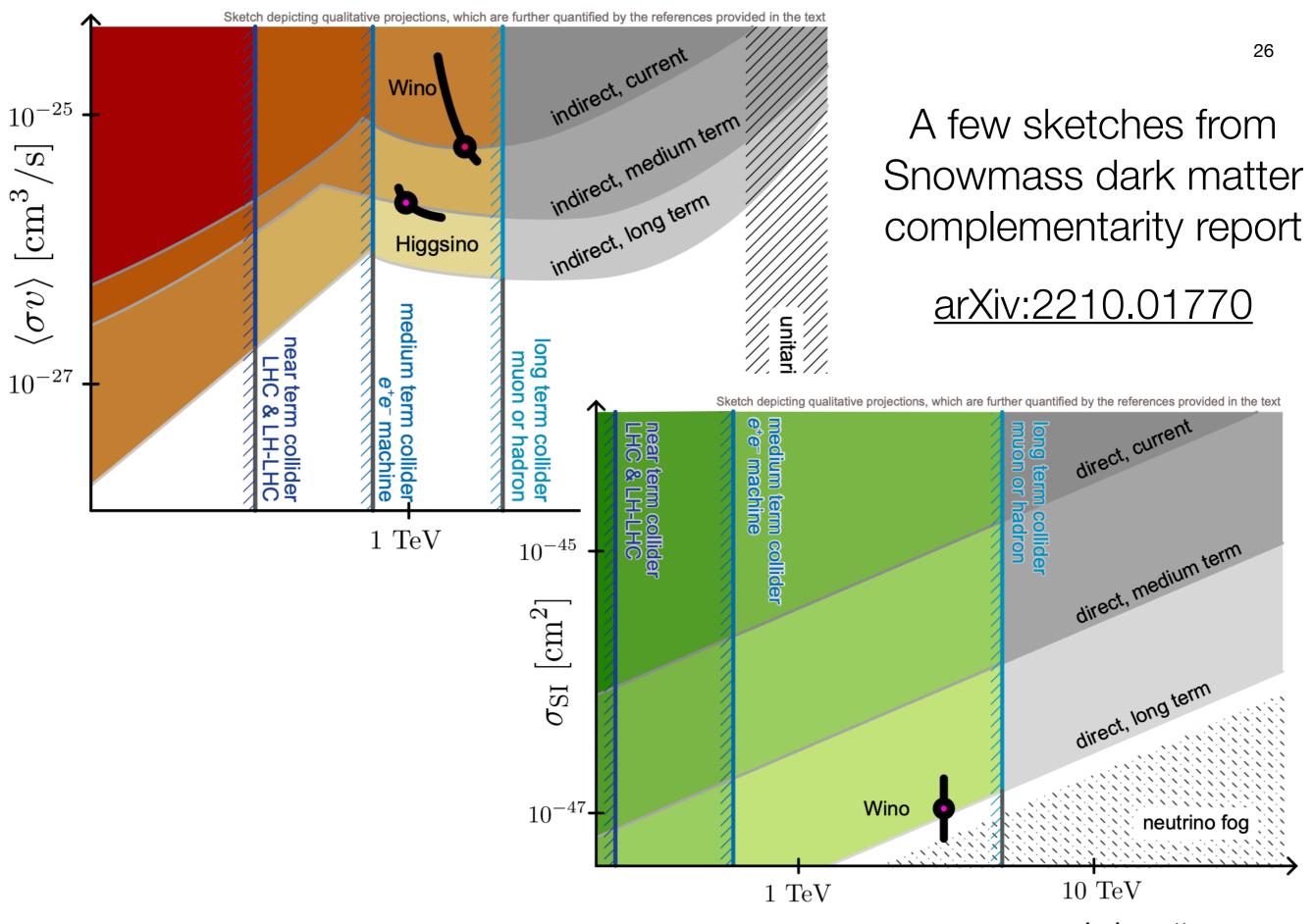
Direct detection has unique strengths in accessing small couplings

Must present both for complete picture

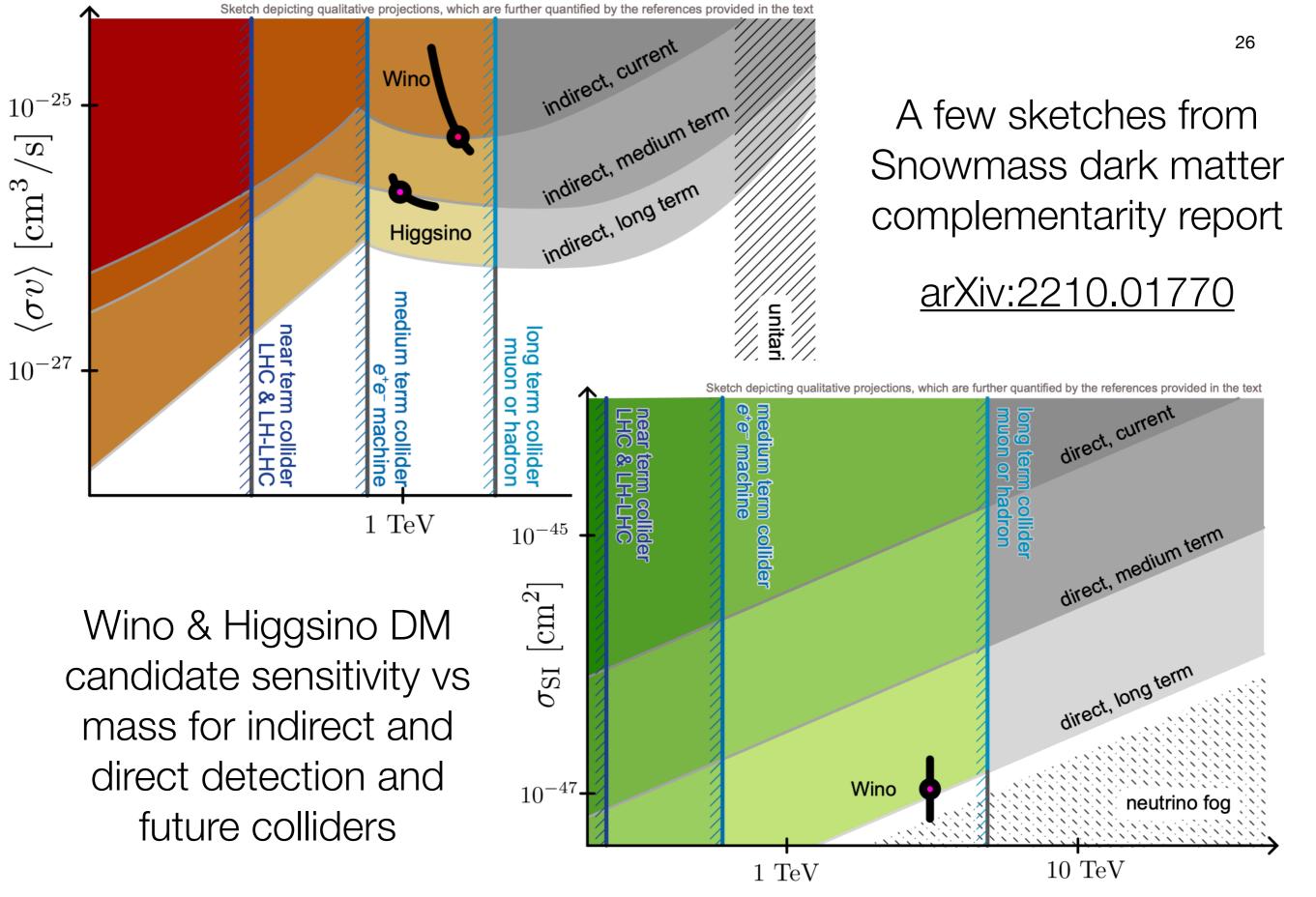


A few sketches from Snowmass dark matter complementarity report

arXiv:2210.01770



dark matter mass



dark matter mass



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There are also areas of DM phase space that only colliders can probe, just as there are areas that only direct or indirect detection experiments can probe

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There remains plenty of non-excluded space for cosmologically motivated particle dark matter above the ~GeV scale

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Complementarity, DM discovery potential, and the potential to exclude values aligning with cosmological observations should be thoroughly understood and included in future collider proposals

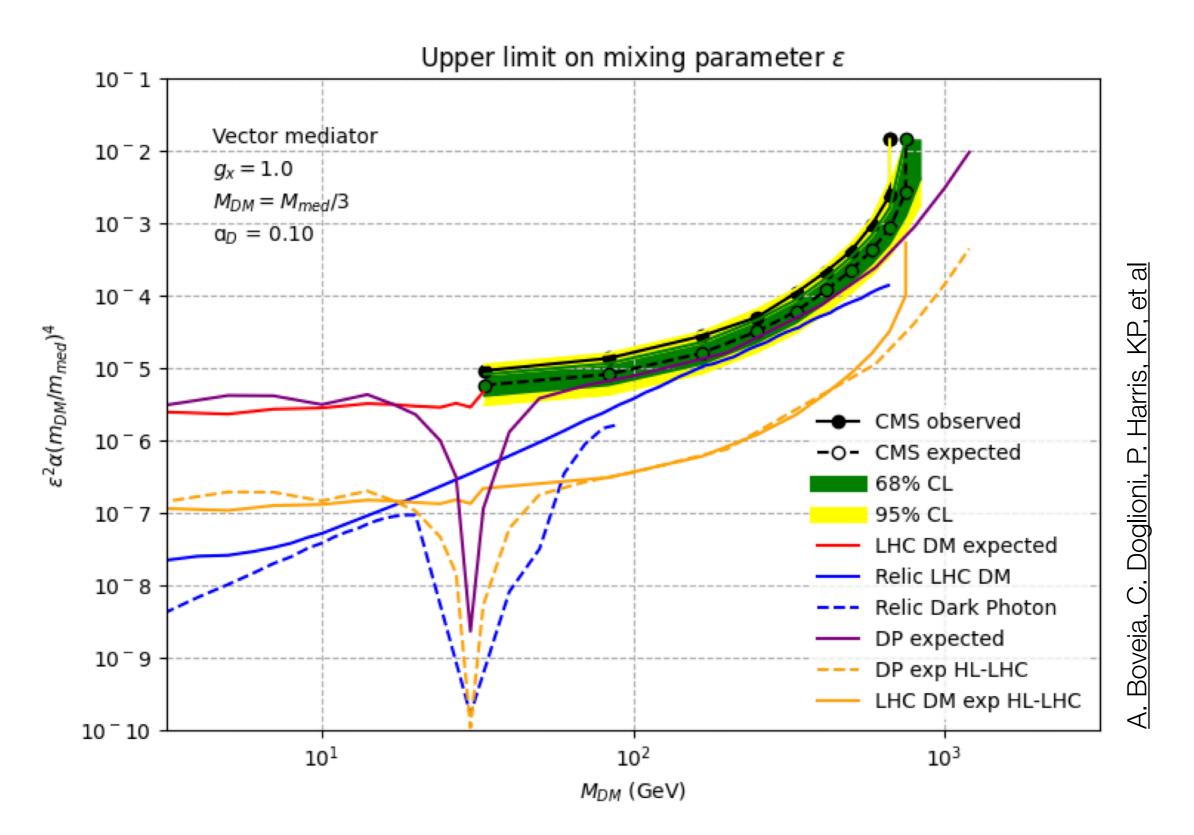
Additional materials

5

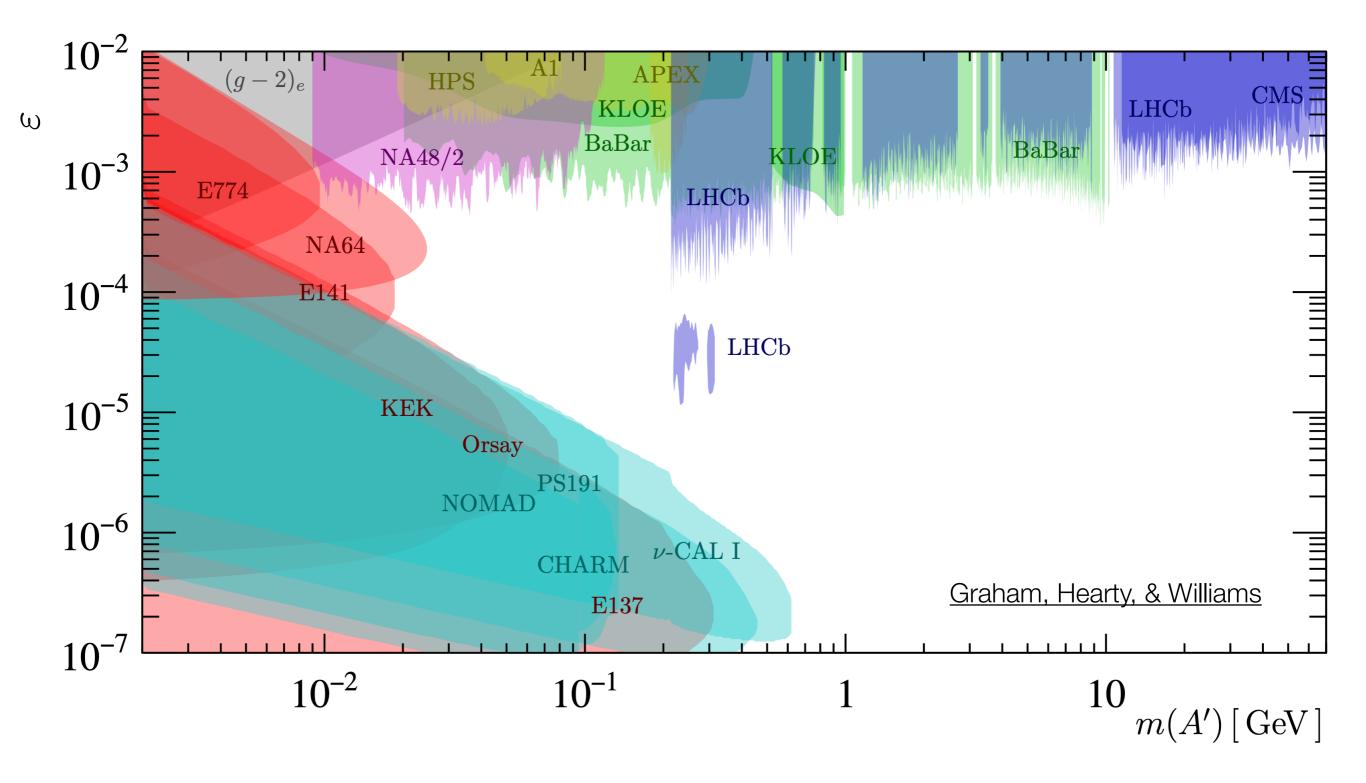
References

- LHC simplified models (s-channel mediators) <u>arXiv:1507.00966</u>
- LHC 2HDM+a model: <u>arXiv:1810.09420</u>
- Notes on Higgs portal: <u>arXiv:2001.10750</u>, <u>arXiv:1903.03616</u>
- Snowmass BSM topical group report <u>arXiv:2209.13128</u>
- Snowmass particle dark matter topical group report <u>arXiv:2209.07426</u>
- Snowmass DM complementarity report: <u>arXiv:2210.01770</u>
- Spin-1 projection comparisons for HL-LHC and FCC <u>arXiv:2206.03456</u>
- European Strategy briefing document: <u>cds link</u>

Comparison between true dark photon model and LHC simplified Z' mediator model, demonstrating good agreement above Z peak



Current limits on visible dark photon decays, by experiment

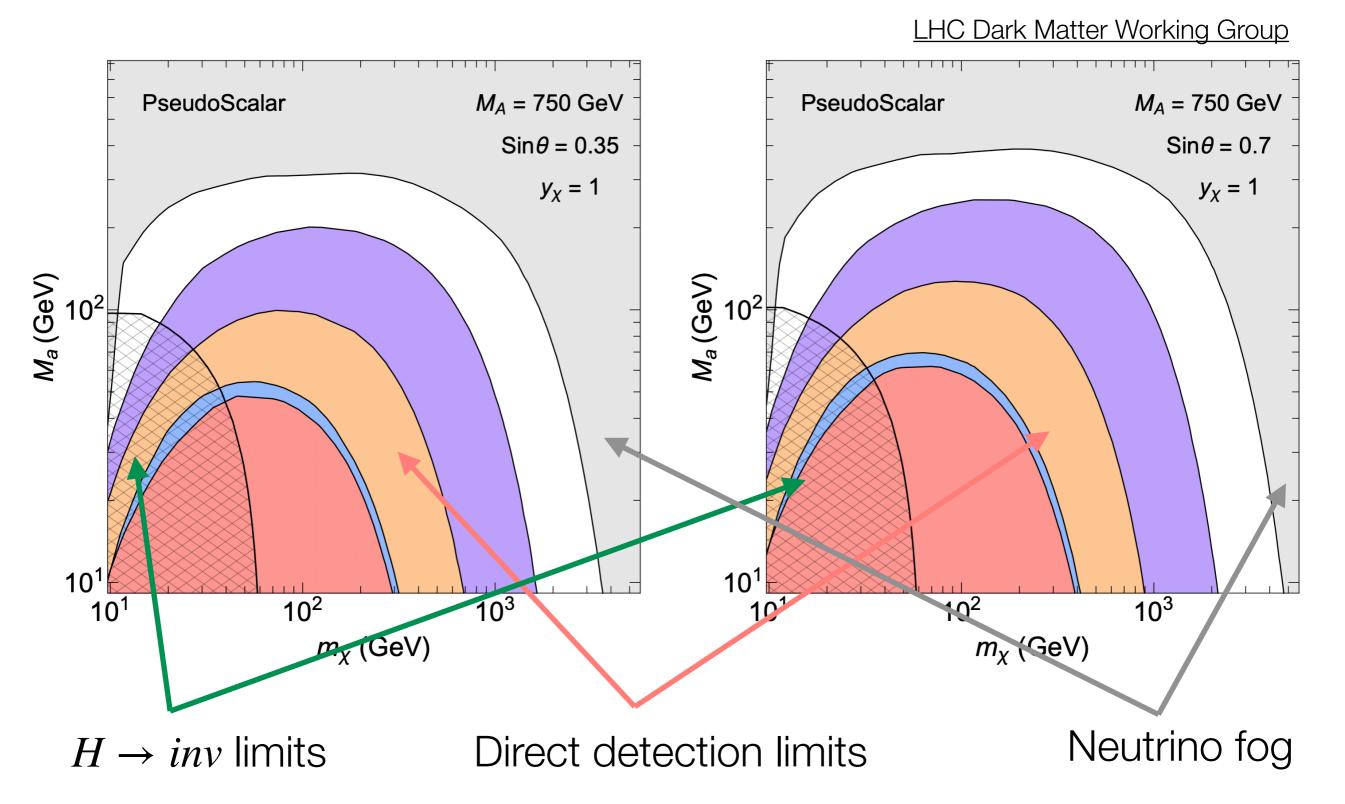


2HDM+a model and parameter choice description

The model considered here is the 2HDM+a model suggested by the LHC DM Working Group, which is the simplest gauge-invariant and renormalizable ultraviolet completion of the simplified pseudoscalar model initially recommended by the LHC DM Forum, which only contained the DM candidate and the mediator. This model is a type-II two-Higgs-doublet (2HDM) model to which an additional pseudoscalar a and a fermionic DM candidate χ are added. After electroweak symmetry breaking, the 2HDM contains five Higgs bosons: a lighter CP-even boson, h, a heavier CP-even boson, H, a CP-odd boson, A, and two charged bosons, $H\pm$. While the phenomenology of the model would be determined by 14 free parameters, some benchmark choices are made in order to match h with the observed SM Higgs boson, to ensure the stability of the Higgs potential, or to evade electroweak precision measurement constraints. In the end, the benchmarks are defined by five parameters: the mass of the heavy Higgs bosons, which are taken to be degenerate, $m_A = m_H = m_{H\pm}$; the mass of the pseudoscalar mediator, m_A ; the mass of the DM particle, m_{χ} ; the mixing angle θ between the two CP-odd states a and A; and the ratio of the vacuum expectation values of the two Higgs doublets, tan β .

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Shape of direct detection exclusions in 2HDM+a model, M_a vs m_x plane. Requires fixing of other three parameters



How spin-1 simplified model to DD plane conversion works

For details, see this talk

$$\sigma_{\rm SI} \simeq 6.9 \times 10^{-41} \text{ cm}^2 \cdot \left(\frac{g_q g_{\rm DM}}{0.25}\right)^2 \left(\frac{1 \text{ TeV}}{M_{\rm med}}\right)^4 \left(\frac{\mu_{n\chi}}{1 \text{ GeV}}\right)^2$$
1 variable 3 variables

Fix two and the other one becomes the thing that changes as σ_{SI} changes.

Implications and consequences can be very different, but can also be somewhat opaque when just looking at final 2D plot.