

---

# *Double Parton Scattering with Herwig++*



*Miroslav Myška*

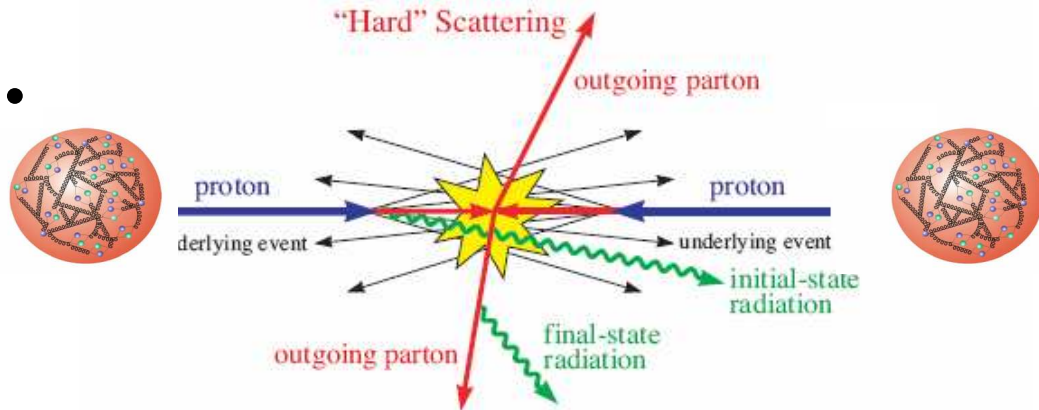
---



FNSPE CTU, Prague

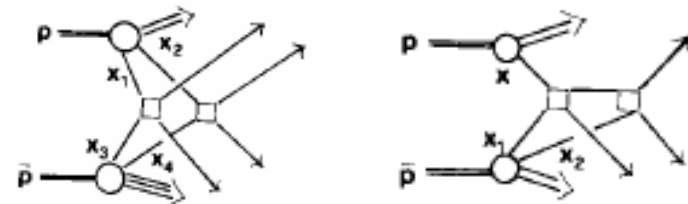
*8<sup>th</sup> MCnet meeting, 24 September 2010*

# Multiple Parton Scattering



Double (multiple) parton scattering (*DPS / MPS*)  
 = process, where two (more) parton-parton interactions occur in one hadron-hadron interaction

Disconnected scattering vs. Rescattering



DPS cross section :

$$\sigma_D = \frac{m}{2} \int_{p_T^c} \Gamma_A(x_1, x_2; b) \hat{\sigma}(x_1, x'_1) \hat{\sigma}(x_2, x'_2) \Gamma_B(x'_1, x'_2; b) dx_1 dx'_1 dx_2 dx'_2$$

# Multiple Parton Scattering

■ Poissonian model: 
$$\sigma_N = \int \frac{1}{N!} \left( \sigma_S F(b) \right)^N d^2b$$

- assumes no correlations for partons inside the hadron
- allows the simple factorization of complex distribution function:

$$\Gamma(x_1, x_2; b) = f(x_1) f(x_2) F(b)$$

- limited for low x region (mostly true)
- all the information about spatial distribution in transverse space is integrated into

DPS cross section is then simple: 
$$\sigma_D = \frac{1}{2} \frac{\sigma_S^2}{\sigma_{eff}}$$

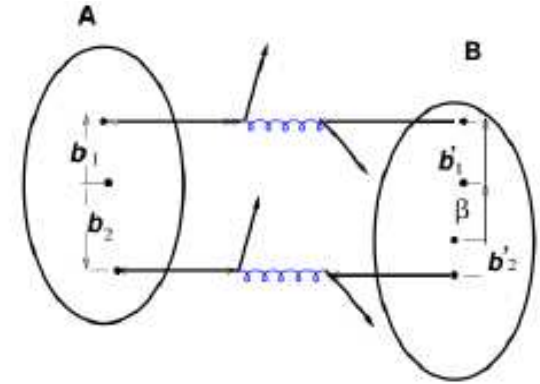
Effective cross section:

$$\sigma_{eff}^{-1} = \int [F(b)]^2 d^2b$$

- dPDF's: (e.g. GS08)
  - considering long. mom. correlations and conservation
  - obeying double DGLAP equation
  - the dPDFs and sPDFS are related by sum rules, e.g.

$$\sum_a \int_0^{1-x_2} dx_1 x_1 f_{ab}(x_1, x_2; Q^2) = (1-x_2) f_b(x_2, Q^2)$$

# Herwig++: Eikonal Model



- At fixed impact parameter  $\mathbf{b}$  partons interact independently
  - Poissonian distribution
- Mean number of scatters is (for general sub-processes):

$$\langle n(b = |\mathbf{b}|, s) \rangle = \int d^2\mathbf{b}' \int_{p_T^{\min 2}} dp_T^2 \sum_{ij} \frac{1}{1 + \delta_{ij}} \frac{d\hat{\sigma}_{ij}(x_1\sqrt{s}, x_2\sqrt{s}, p_T^2)}{dp_T^2} \\ \otimes G_{i/h_1}(x_1, \mathbf{b} - \mathbf{b}', \mu^2) \otimes G_{j/h_2}(x_2, \mathbf{b}', \mu^2) = A(b) \cdot \sigma^{\text{inc}}(s; p_T^{\min})$$

- $p_T^{\min}$  and  $\mu^2$  are important steering parameters for UE tune
  - There is a strong and constant correlation between  $p_T^{\min}$  and  $\mu^2$ : the smaller hadron radius always balance against a larger  $p_T$  cutoff (for UE activity)
  - Old tune (H++ 2.3.0):  $p_T^{\min} = 3.4 \text{ GeV}$ ,  $\mu^2 = 1.5 \text{ GeV}^2$
  - New tune (H++ 2.4.2):  $p_T^{\min} = 4.1 \text{ GeV}$ ,  $\mu^2 = 1.33 \text{ GeV}^2$
- Inclusive cross section (k int. of type 1, m int. of type 2):

$$\sigma_{k,m}(\sigma_a, \sigma_b) = \int d^2b \mathcal{P}_k(A(b)\sigma_a) \mathcal{P}_m(A(b)\sigma_b) = \int d^2b \frac{(A(b)\sigma_a)^k}{k!} e^{-A(b)\sigma_a} \frac{(A(b)\sigma_b)^m}{m!} e^{-A(b)\sigma_b}$$

# Herwig++: Eikonal Model

- Generation algorithm:

- factorization:

$$G(x, \mathbf{b}, \mu^2) = f(x, \mu^2) \cdot S(\mathbf{b})$$

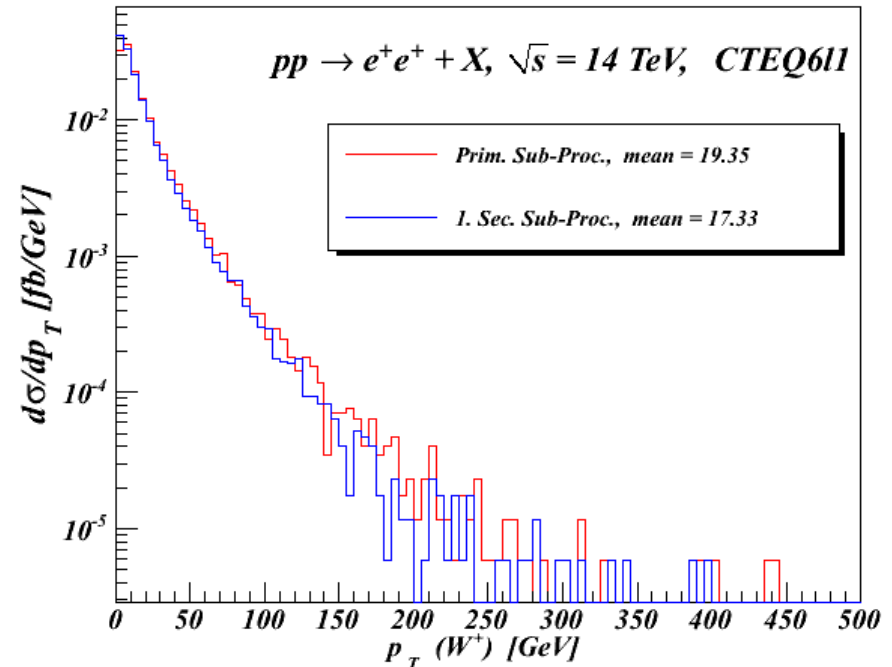
- elmag. form factor-like form:

$$S_{\bar{p}}(\mathbf{b}) = S_p(\mathbf{b}) = \int \frac{d^2\mathbf{k}}{2\pi} \frac{e^{i\mathbf{k}\cdot\mathbf{b}}}{(1 + \mathbf{k}^2/\mu^2)^2}$$

- Overlap function evaluation:

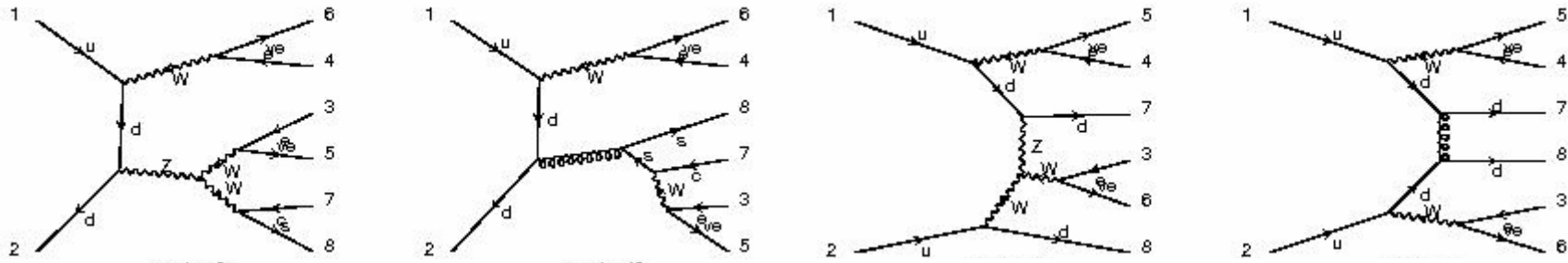
$$A(b) = \frac{\mu^2}{96\pi} (\mu b)^3 K_3(\mu b)$$

- Primary Sub-Process untouched
- 1<sup>st</sup> Secondary Sub-Process forced to be same ME
- Other Sub-Processes (till  $\langle n \rangle$ ) QCD 2→2
- Event veto for momentum conservation violation
- no  $p_T$  ordering (unlike e.g. Pythia)



# Process: $p+p \rightarrow W^+ + W^+ + X$

- Signal process: (DPS): 2 (hard) independent  $qq \rightarrow W^+ \rightarrow e^+ \nu_e$  sub-processes
- Background:
  - (SPS) SM  $e^+e^+ + jj$  production:



- SM  $W^+Z$  and  $ZZ$  production
  - while electron is outside the detector acceptance  $|\eta| < 2.5$  (Atlas ID)
  - no detector effects are considered (for now)
- heavy flavor quark production
- DPS and pile-up effects are assumed to be negligible

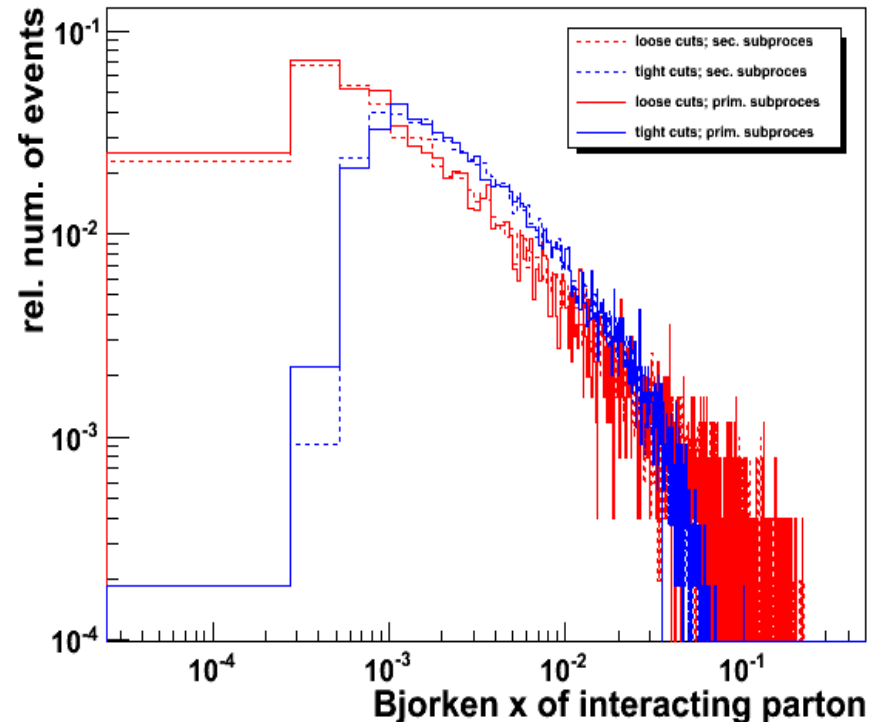
---

# Process: Generation

- DPS: Herwig++ 2.4.2
- SPS: Madgraph v.4 + Herwig++
- $\sqrt{s} = 7 \text{ \& } 14 \text{ TeV}$ ; PDF: cteq6l1, mrst2007lomod.lhgrid
- set MPIHandler:softInt Yes
- set MPIHandler:twoComp Yes
- (soft inv. radius squared: 7TeV(cteq): 0.51 GeV<sup>2</sup>, (mrst): 0.49 GeV<sup>2</sup>,  
14TeV(cteq): 0.47 GeV<sup>2</sup>, (mrst): 0.43 GeV<sup>2</sup>)
  
- Herwig++ (DPS generation):
  - insert SimpleQCD:MatrixElements[0] MEqq2W2ff
  - set MEqq2W2ff:Wcharge 1 #only W+
  - set MEqq2W2ff:Process 3 #only positrons
  - set LeptonKtCut:MinKT 0.001\*GeV
  - set LeptonKtCut:MaxEta 5
  - set LeptonKtCut:MinEta -5
  - set MassCut:MinM 0.002\*GeV

# Tight vs. loose generation cuts

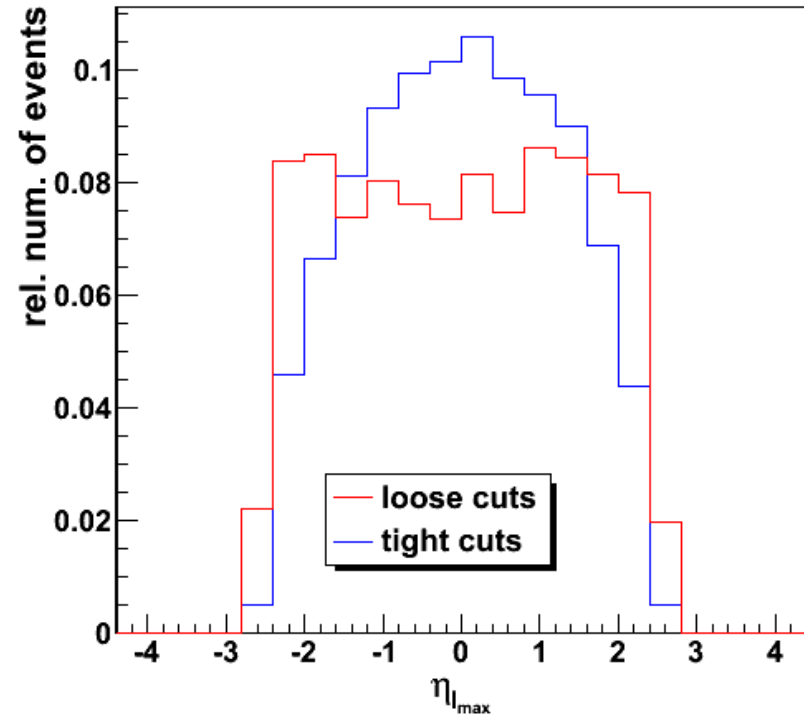
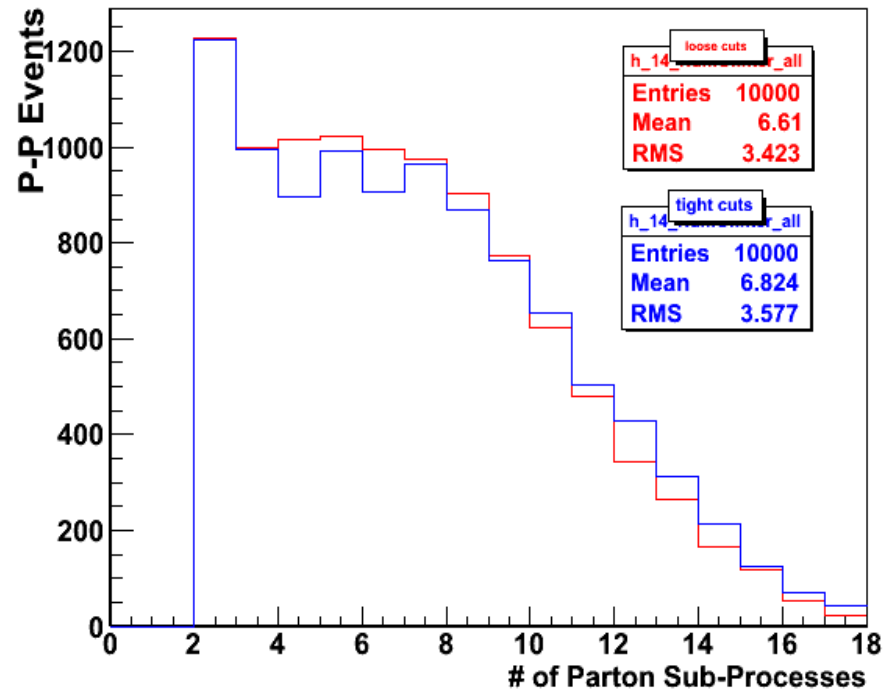
- cross section depends significantly on the cuts used for direct event filtering during the generation
- tight cuts:  $p_T(e^+) > 5\text{GeV}$ ,  $|\eta| < 2.5$
- loose cuts:  $p_T(e^+) > 1\text{MeV}$ ,  $|\eta| < 5$
- e.g.: 14 TeV & cteq6l1
  - $x$  – frac. of long. momentum dependence on cuts
  - no strong dependence on PDF's used





# Tight vs. loose generation cuts

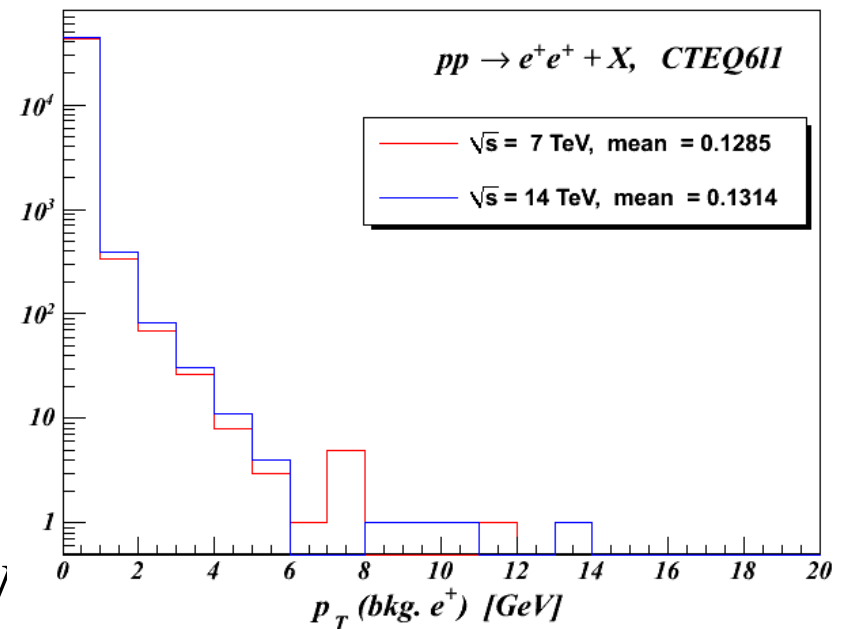
h\_14\_NumOfInter\_all



- plots for  $\sqrt{s} = 14$  TeV, PDF: cteq6l1
- filtering affects distribution of number of parton sub-processes as well as other kin. variables (e.g. pseudorapidity of hardest positron)

# Data Analysis

- final state requirements:
  - always at least two positrons with  $p_T > 5 \text{ GeV}$ ,  $|\eta| < 2.5$ 
    - 2 analyzed objects:  $e^+_{\text{max}}$ ,  $e^+_{\text{min}}$
  - if jet: then  $p_T > 5 \text{ GeV}$ ,  $|\eta| < 2.5$ 
    - anti- $k_T$  algorithm,  $R = 0.4$
    - analysed object:  $\text{jet}_{\text{max}}$
  - isolation requirements:  
 $\Delta R(e^+e^+) > 0.2$ ,  $\Delta R(e^+j) > 0.2$   
(57% survive isolation cuts)
  - veto on hard electrons  $p_T > 5 \text{ GeV}$

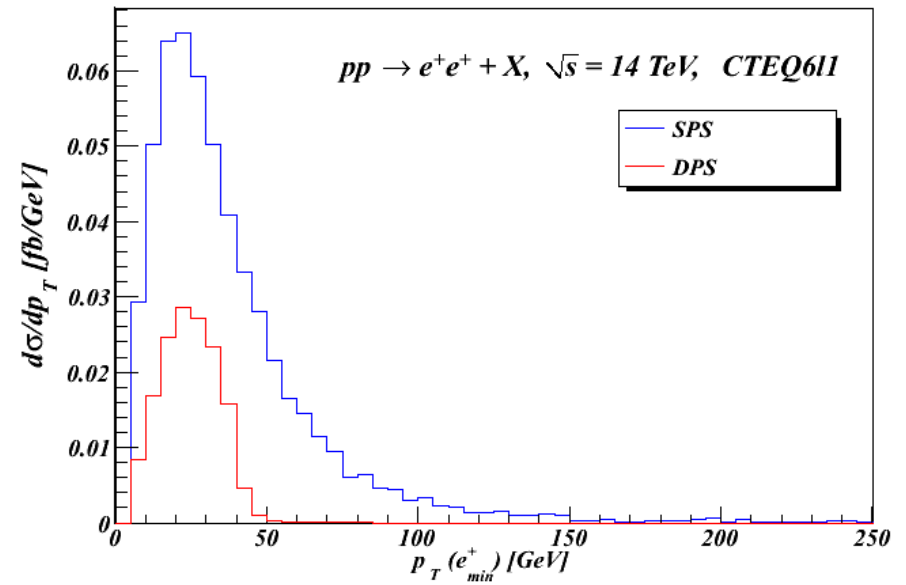
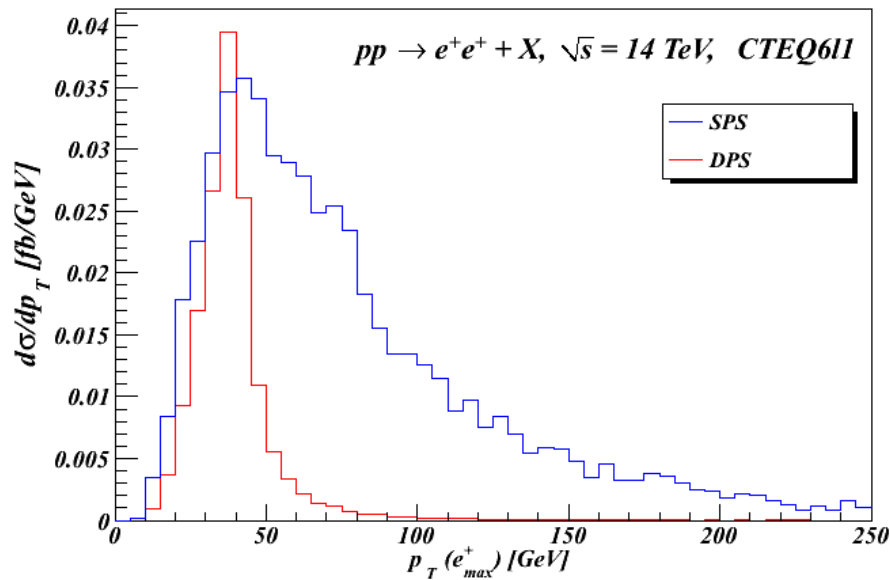


# Results:

## Transverse Momentum Distribution

Left: the hardest positron

Right: the 2<sup>nd</sup> hardest positron

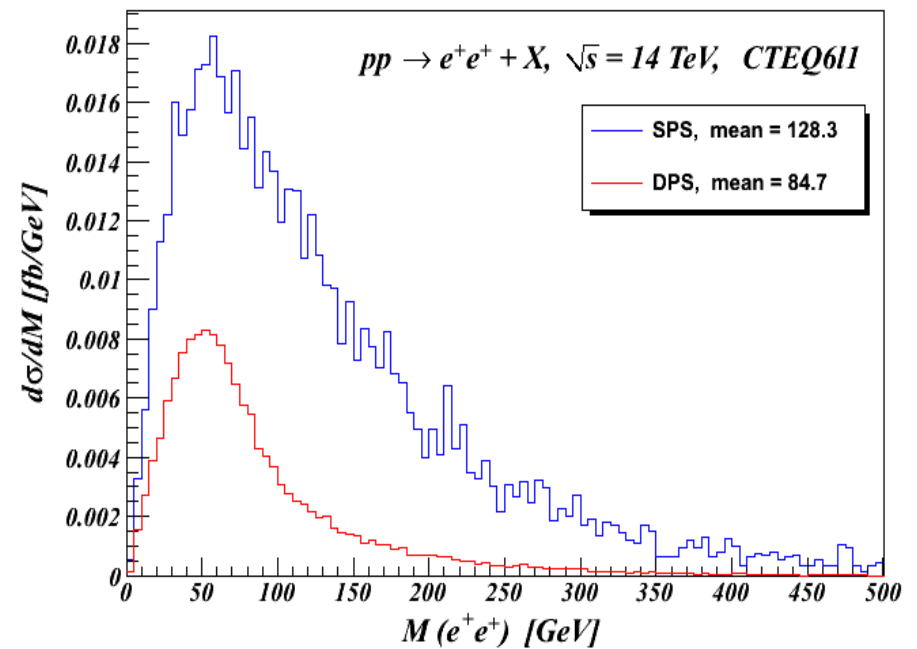
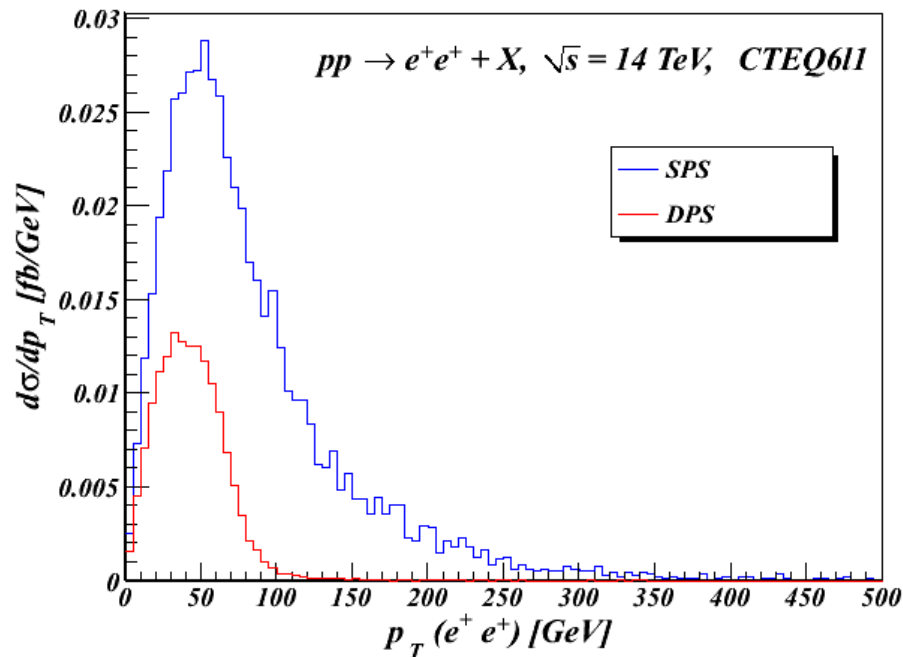


# Results:

## Positron-Positron Pair Characteristics

Left: transverse momentum

Right: invariant mass

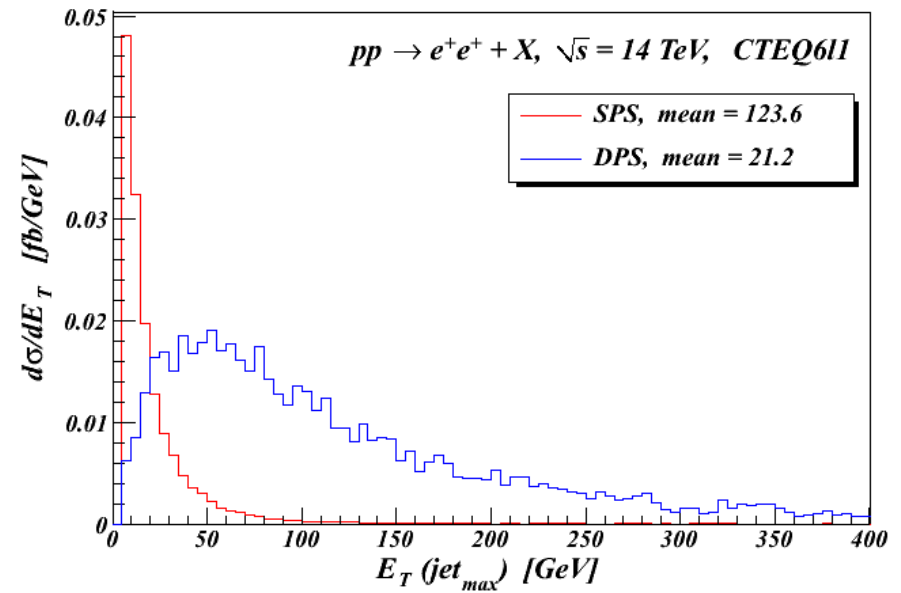
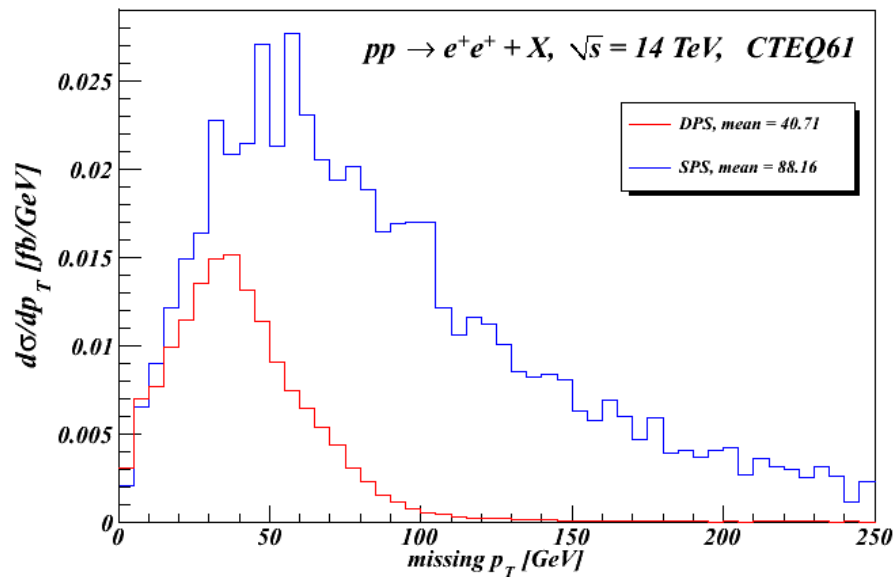


# Results:

## Background Characteristics

Left: missing transverse momentum

Right: transverse energy of the hardest jet

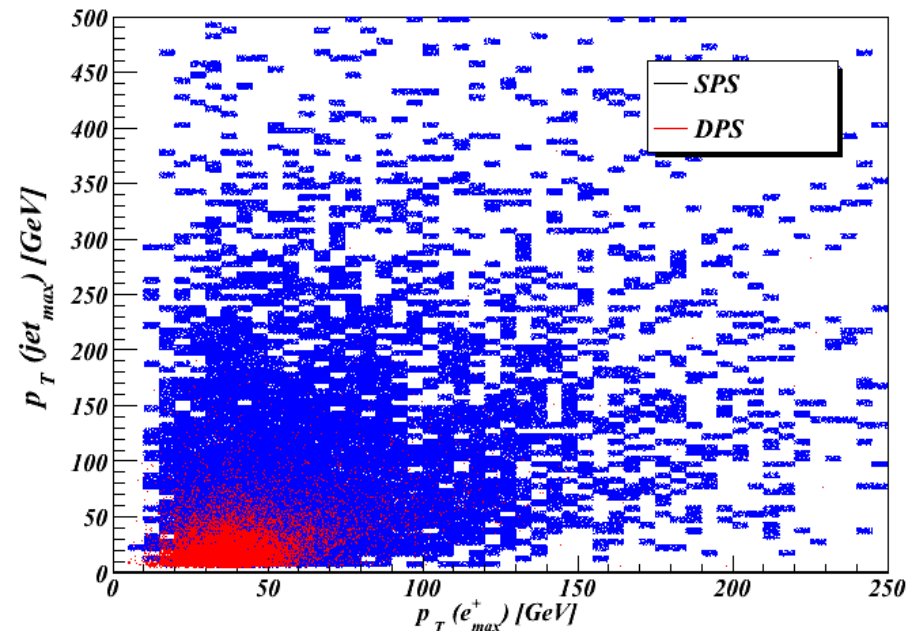
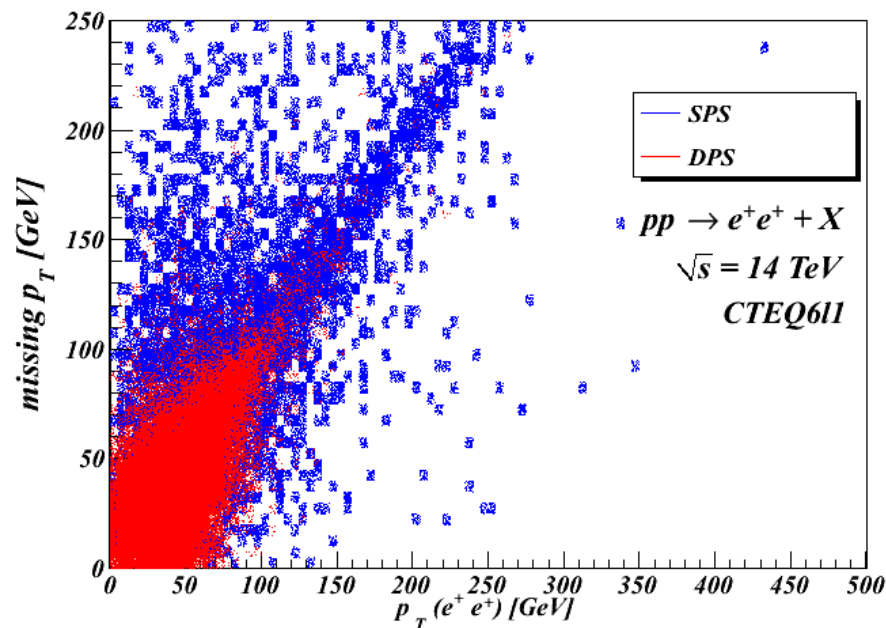


# Results:

## 2D histograms: demonstration of complementarity of selective vetoes

Left: transverse momentum: positron pair vs. missing part

Right: transverse momentum: hardest positron vs. hardest jet

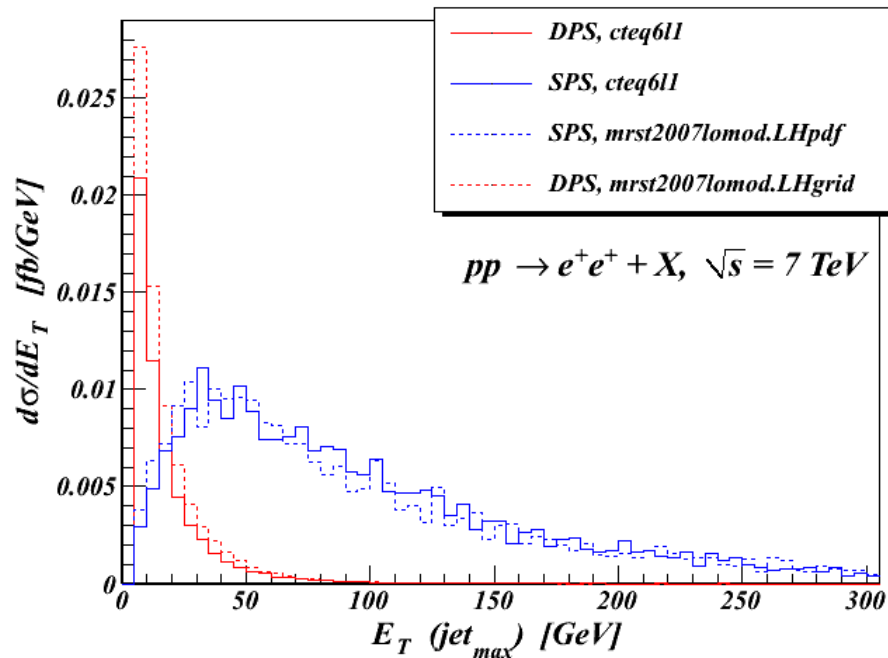


# Results: Cross sections; $\sigma_{\text{eff}} = 11.5 \text{ mb}$

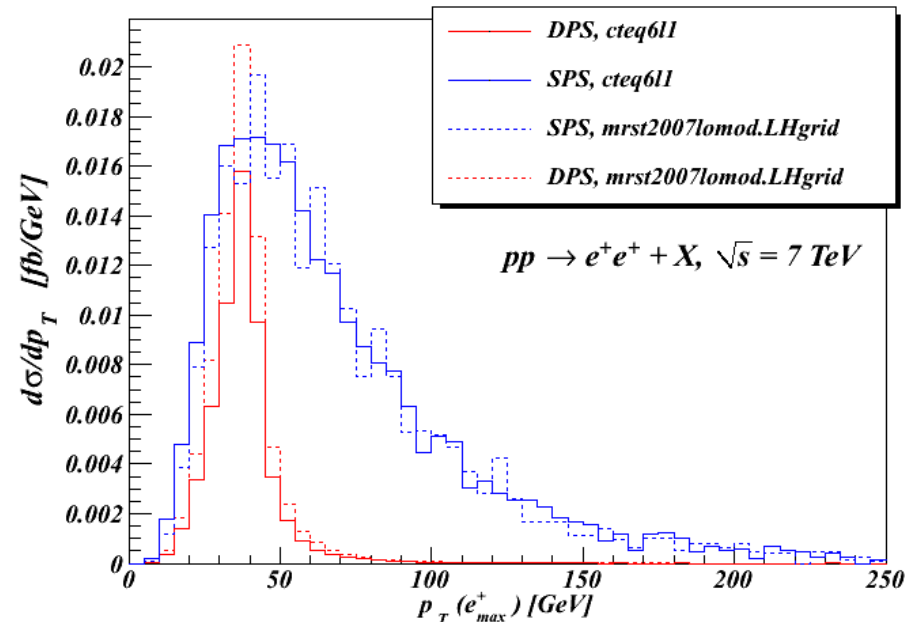
$\sigma$ [fb]	DPS				SPS: SM $e^+e^+ jj$			SPS: $W^+Z$
	14 TeV		7 TeV		14 TeV	7 TeV		14 TeV
	cteq	mrst	cteq	mrst	cteq	cteq	mrst	cteq
total	<b>0.75</b>	1.03	0.28	0.37	<b>2.69</b>	1.20	1.17	<b>6.45</b>
jet veto 20 GeV	<b>0.54</b> 72%	0.74 72%	0.22 79%	0.29 78%	<b>0.16</b> 6%	0.08 7%	0.09 8%	<b>4.45</b> 69%
jet veto 30 GeV	0.63 84%	0.85 83%	0.25 89%	0.33 88%	0.32 12%	0.17 14%	0.19 16%	5.16 80%

# Results: Comparison of PDF's -uncertainty in the cross section

Left: transverse momentum  
of the hardest jet



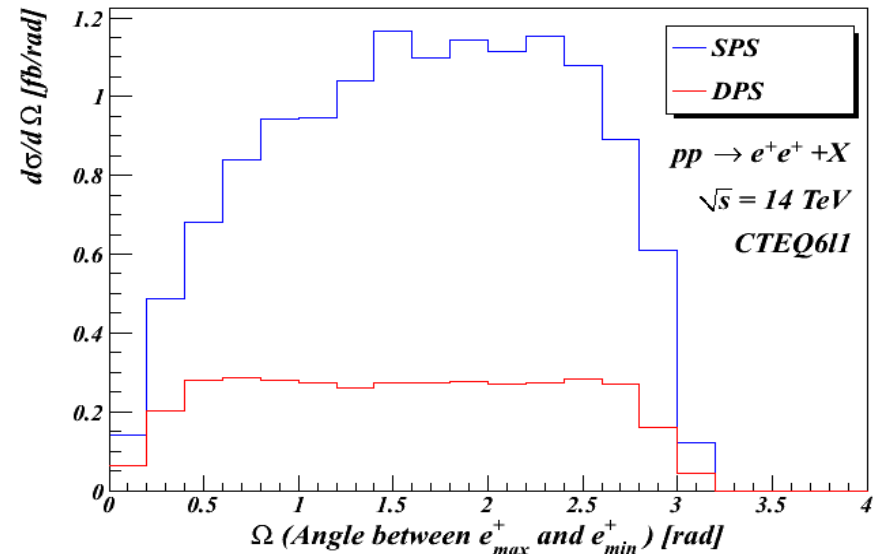
Right: transverse momentum  
of the hardest positron





# Another handlers to separate signal different from $p_T/E_T$ distributions

- distributions of  $\Phi$  and  $\eta$ , their difference and other functions could provide another helpful tool



- e.g. Pseudorapidity asymmetry.

$$a_{\eta_l} = \frac{\sigma(\eta_{l_1} \times \eta_{l_2} < 0) - \sigma(\eta_{l_1} \times \eta_{l_2} > 0)}{\sigma(\eta_{l_1} \times \eta_{l_2} < 0) + \sigma(\eta_{l_1} \times \eta_{l_2} > 0)}$$

---

# Summary and next steps

- Double parton production of same-sign W pair was performed for  $e^+ \nu_e$  decay channel
- SM background process  $pp \rightarrow e^+ e^+ jj$  was prepared and found to be strongly reduced by  $p_T$  vetoes
- SM background process  $pp \rightarrow W^+ Z$  under construction – preliminary results show it's dominance
- To do:
  - other sources of background need to be investigated as well as other final states (muon or multi-jets)
  - invisible component does not allow for usage of pair-to-pair variables characteristic for DPS events